

State	T_e	ω_e	$\omega_e x_e$	B_e	α_e	D_e (10^{-4}cm^{-1})	r_e (Å)	Observed Transitions		References	
								Design.	ν_{00}		
$^4\text{He}_2$		$\mu = 2.0013016_3$		$D_e^0 = 0.00090 \text{ eV}^a$ $D_0(a^3\Sigma_u^+) = 1.850 \text{ eV}^b$ $D_0(A^1\Sigma_u^+) = 2.355 \text{ eV}^b$		$\text{I.P.} = 22.223 \text{ eV}^c$				NOV 1976 A	
				Rydberg series of $np\pi$ $^3\Pi_g$ levels (b, e, i, l, p, r, s, t, u, ...)	}	$\nu^d = 34302.20 - R/(n - 0.071)^2, n = 2 \dots 17.$				(23)(39)	
				Rydberg series of $np\sigma$ $^3\Sigma_g^+$ levels (c, g, k', n, p', r', s', t', ...)			}	$\nu^d = 34302.20 - R/(n - 0.777)^2, n = 3 \dots 12.$			
u	$^3\Pi_g$ 10pπ 177291	[1628.6 ₉] ^e Z (35.2 ₅)		7.21 ₂ ^e	0.22 ₂ ^e	[5.03] ^e			1.081	u→a, R	33189.16 Z
t'	$^3\Sigma_g^+$ 10pσ [177969]			[8.27] ^f				t'→a, V	33026.6 Z	(23)	
t	$^3\Pi_g$ 9pπ 177027	[1629.1 ₅] ^g Z (35.2 ₅)		7.21 ₂ ^g	0.23 ₀ ^g	[5.07] ^g	1.081	t→a, R	32925.96 Z	(23)	
s'	$^3\Sigma_g^+$ 9pσ [177636]			[8.04] ^f				s'→a, V	32693.2 Z	(23)	
s	$^3\Pi_g$ 8pπ 176658	[1629.30] ^h Z (35.2 ₅)		7.21 ₃ ^h	0.223 ₄ ^h	5.1 ^h	1.080 ₆	s→a, R	32556.66 Z	(23)	
r'	$^3\Sigma_g^+$ 8pσ [177154]			[7.78] ^f				r'→a, V	32211.7 Z	(23)	
r	$^3\Pi_g$ 7pπ 176117	(1700.5 ₆) ⁱ (35.2 ₅)		[7.104 ₄] ⁱ	i	[5.08] ⁱ	[1.0889]	r→d, R	11625.3 Z	(39)	
								r→a, R	32016.56 Z	(23)	
								p'→a,	31478.6 Z	(23)	
p'	$^3\Sigma_g^+$ 7pσ [176421]			[7.543] ^f							
q	$^3\Delta_u$ 6dδ [176195]	}				[7.092] ^j	[5.0] ^j	[1.089 ₈]	q→c, {	V 20365	(1)(39)
									R (20330)		
									R 20288		
	$^3\Pi_u$ 6dπ [(176169)]								{	V 26483	(1)(39)
	$^3\Sigma_u^+$ 6dσ [176120]								R (26466)		
									R 26409		
p	$^3\Pi_g$ 6pπ 175281	1701.18 Z 35.35		7.220 ^k	0.224	5.14	1.080 ₁	p→d, R	10788.6 ₈ Z	(39)	
								p→a, R	31179.93 Z	(23)	
o	$^3\Sigma_u^+$ 6sσ [176001]			[7.109]		[5.1]	[1.088 ₅]	o→c, V	20168.8 Z	(2)(39)	
								o→b, R	26290.3 Z	(2)(39)	

State	T_e	ω_e	$\omega_e x_e$	B_e	α_e	D_e (10^{-4} cm^{-1})	r_e (\AA)	Observed Transitions		References
								Design.	ν_{00}	
$^4\text{He}_2$ (continued)										
n $^3\Sigma_g^+$	6p6	174389	[1619.52]	Z (36.5)	7.475 ₄ ^l	0.249 ₀	[7.2 ₁] ^m	1.0615	n → a,	30283.26 Z (3)(67)
m $\left\{ \begin{array}{l} ^3\Delta_u \\ ^3\Pi_u \\ ^3\Sigma_u^+ \end{array} \right.$	5d8	[174863]	}		[7.09] ⁿ [7.07] ^o		[1.09 ₁]		m → c,	$\left\{ \begin{array}{l} V (19039) \\ R (18944) \\ R 18899 \end{array} \right.$ (39)
	5dπ	[(174778)]							m → b,	$\left\{ \begin{array}{l} V 25152 \\ R (25070) \\ R 25019 \end{array} \right.$ (39)
	5dσ	[174730]								
l $^3\Pi_g$	5pπ	173884	[1633.96] ^p	Z (35.2 ₅)	7.226 ^p	0.222 ^p	[5.12] ^p	1.079 ₇	l → d, R 9393.9 Z (39) l → a, R 29785.31 Z (23)	

He_2 : ^aAverage of two independent values for the well depth obtained from measurements of the total (44) and differential (45)(66) elastic scattering cross sections (0.88₈ and 0.90₅ meV, resp.). Ab initio values range from 0.78 to 1.04 meV (33)(43)(51)(53)(56)(64). Both experiment and theory agree that no bound vibrational level exists in the potential well, i.e. $D_0^0 = 0$; see (27)(63). A somewhat higher D_e value (0.99 meV) was derived (61) from the temperature dependence of the relaxation time in dilute ^3He . For measurements of the short-range potential ($0.49 < r(\text{\AA}) < 1.56$) see (57); the long-range potential is discussed by (28)(29).

^bBased on $D_0^0(\text{He}_2^+)$. From a detailed interpretation of the Hopfield continuum and the 600 \AA absorption and emission bands (49) derives $D_e(A^1\Sigma_u^+) = 2.50 \text{ eV}$.

^cRelative to $\text{He}(^1\text{S}) + \text{He}(^1\text{S})$, i.e. $\text{I.P.}(\text{He}_2) = \text{I.P.}(\text{He}) - D_0^0(\text{He}_2^+)$. The I.P. for the lowest stable state (a $^3\Sigma_u^+$) is 4.25297 eV.

^dGiving the $v=0, N=0$ levels (real or hypothetical) of $\text{np}6$

$^3\Sigma_g^+$ and $\text{np}\pi$ $^3\Pi_g^-$ relative to $2s6$ $^3\Sigma_u^+$, $v=0, N=0$.

^eRefers to Π^- . $B_0(^3\Pi^+) \approx 5.6$, strongly affected by l -uncoupling.

^fEffective value, strongly affected by l -uncoupling.

^gRefers to Π^- . $B_0(^3\Pi^+) = 5.87$, strongly affected by l -uncoupling.

^hRefers to Π^- . $B_0(^3\Pi^+) = 6.13$, strongly affected by l -uncoupling.

ⁱConstants refer to Π^- ; $B_0(^3\Pi^+) = 6.37_5$ affected by l -uncoupling. $v=1$ is perturbed; approximate deperturbed constants for Π^- : $B_1 = 6.88_6$, $\Delta G(\frac{1}{2}) = 1629.7$.

^jStrong l -uncoupling. The rotational constants (1) refer to the average of Π^- and Δ^- .

^kRefers to Π^- . $B_0(^3\Pi^+) = 6.630$ affected by l -uncoupling.

^lSeveral small accidental perturbations.

^m $D_1 = 6.8_2 \times 10^{-4}$, $H_0 = 14.2 \times 10^{-8}$, $H_1 \approx -26 \times 10^{-8}$.

ⁿAverage of Π^- and Δ^- as given by (39).

^oAverage of Σ^+ , Π^+ , Δ^+ as given by (39).

^pRefers to Π^- .

State	T_e	w_e	$w_e x_e$	B_e	α_e	D_e (10^{-4}cm^{-1})	r_e (\AA)	Observed Transitions		References
								Design.	ν_{00}	
$^4\text{He}_2$ (continued)										
k $3\Sigma_u^+$ 5s σ (173698)		[1635.3]	Z		7.23 ₂	0.23		1.079	k \rightarrow c, V 18683.5 Z (39) k \rightarrow b, R 24804.8 Z (39)	
k' $3\Sigma_g^+$ 5p σ 172236		1686.90	Z 38.10		7.379 ^a	0.349 ^b	[5.8]	1.068 ₄	k' \rightarrow a, R 28127.58 Z (67)	
j $\left\{ \begin{array}{l} 3\Delta_u \\ 3\Pi_u \end{array} \right.$ 4d δ 171573		1702.2 ₄ ^c	Z 35.0 ₇		7.208 ₈ ^c	0.224 ₈ ^d	5.2 ^c	1.0810	j \rightarrow c, $\left\{ \begin{array}{l} V \\ R \end{array} \right.$ 16583.18 ^c Z (50) 16400.69 ^c Z	
j $\left\{ \begin{array}{l} 3\Pi_u \\ 3\Sigma_u^+ \end{array} \right.$ 4d π 171402		1680.9 ₄ ^c	Z 40.8 ₁		7.186 ₀ ^c	0.229 ₆ ^e	5.4 ^c	1.0827	j \rightarrow b, $\left\{ \begin{array}{l} V \\ R \end{array} \right.$ 22704.5 ^c Z (50) 22522.0 ^c Z 22437.8 ^f	
j $\left\{ \begin{array}{l} 3\Sigma_u^+ \\ 3\Sigma_u^+ \end{array} \right.$ 4d σ 171323		1669.7 ₉ ^f	39.0 ₉	Strongly perturbed by ℓ -uncoupling and by interaction with the h state. ^h						
i $3\Pi_g$ 4p π 171290		[1637.94] ^g	Z (35.2 ₅)		7.242 ^g	0.223 ^g	[5.1 ₄] ^g	1.078 ₅	i \rightarrow a, R 27193.01 Z (23)(39)	
h $3\Sigma_u^+$ 4s σ (170884)		[1637.9]	Z		7.26 ₄ ^h	0.23	(5.2 ₄)	1.077	h \rightarrow c, V 15870.7 Z (2)(39)(50) h \rightarrow b, R 21992.2 Z (2)(39)(50)*	
g $3\Sigma_g^+$ 4p σ 167714		[1589.92]	Z (41)		7.220 ₇	0.247 ₈	[5.38] ⁱ	1.0801	g \rightarrow a, R 23597.00 Z (67)	
f $\left\{ \begin{array}{l} 3\Delta_u \\ 3\Pi_u \end{array} \right.$ 3d δ 166303		1706.82	Z 35.10		7.230 ^j	0.227 ^k	[5.26] ^l	1.079 ₄	f \rightarrow c, $\left\{ \begin{array}{l} V \\ R \end{array} \right.$ 11316.06 Z (17) 10864.53 Z 10659.33 Z	
f $\left\{ \begin{array}{l} 3\Pi_u \\ 3\Sigma_u^+ \end{array} \right.$ 3d π 165877 ^m		1661.48	Z 44.79		7.136 ^j	0.235 ⁿ	[5.34] ^o	1.086 ₅	f \rightarrow b, $\left\{ \begin{array}{l} V \\ R \end{array} \right.$ 17437.3 Z (13)* 16985.8 Z (13)(39) 16780.6 Z	
f $\left\{ \begin{array}{l} 3\Sigma_u^+ \\ 3\Sigma_u^+ \end{array} \right.$ 3d σ 165685		1635.77	Z 44.41		7.071 ^j	0.246 ^p	[5.31] ^q	1.091 ₄		
e $3\Pi_g$ 3p π 165598		1721.22	Z 34.97 ₀ ^s		7.283 ₈ ^r	0.221 ₅ ^t	5.22	1.0754	e \leftrightarrow a, ^u R 21507.26 Z (36)	
d $3\Sigma_u^+$ 3s σ 164479		1728.01	Z 36.13 ^v		7.341 ₂	0.224 ₄ ^w	5.32	1.0712	d \rightarrow c, V 9502.7 Z (13) d \rightarrow b, R _V 15623.1 Z (13)*	
c $3\Sigma_g^+$ 3p σ 155053		1583.85	Z 52.7 ₄ ^x		7.004 ₈	0.310 ₅ ^y	[5.56] ^z	1.0966	c \rightarrow a, R 10889.48 Z (11)	
b $3\Pi_g$ 2p π 148835		1769.07	Z 35.02 ^{a'}		7.447 ₃ ^{b'}	0.219 ₆ ^{c'}	[5.30] ^{d'}	1.0635	b \rightarrow a, R 4768.2 Z (5)(14)	
a $3\Sigma_u^+$ 2s σ 144048		1808.56	Z 38.2 ₁ ^{e'f'}		7.703 ₆ ^{g'}	0.228 ₁ ^{h'}	5.56 ^{i'}	1.0457	(a-X) 144935 ^{j'}	

He₂:

- ^aSeveral small accidental perturbations.
- ^b $\gamma_e = + 0.030$.
- ^cThe vibrational and rotational constants refer to Π^- and Δ^- which are less affected by ℓ -uncoupling.
- ^d $\gamma_e = - 0.0038$.
- ^e $\gamma_e = - 0.0046$.
- ^fConstants refer to $N'=1$.
- ^gConstants refer to $3\Pi^-$.
- ^h(50) give average effective constants for the four interacting components $j(3\Delta_u^+, 3\Pi_u^+, 3\Sigma_u^+)$ and $h\ 3\Sigma_u^+$.
- ⁱ $H_0 = 2.0 \times 10^{-8}$.
- ^jThese constants are corrected for ℓ -uncoupling effects.
- ^k $\gamma_e = - 0.0046$.
- ^l $D_1 = 5.28 \times 10^{-4}$, $D_2 = 5.50 \times 10^{-4}$, $H_0 = 2.5 \times 10^{-8}$.
- ^mAb initio calculations of $f\ 3\Pi_u$ and $F\ 1\Pi_u$ (10)(26) yield excellent agreement with the observed constants and confirm the presence of substantial potential maxima; see also (9).
- ⁿ $\gamma_e = - 0.0069$.
- ^o $D_1 = 5.34 \times 10^{-4}$, $D_2 = 5.74 \times 10^{-4}$, $H_0 = 2.9 \times 10^{-8}$.
- ^p $\gamma_e = - 0.0095$.
- ^q $D_1 = 5.34 \times 10^{-4}$, $D_2 = 5.45 \times 10^{-4}$, $H_0 = 0.91 \times 10^{-8}$.
- ^rThe rotational constants refer to Π^- (36); $B(3\Pi^-) - B(3\Pi^+) \approx + 0.072$. Slightly different constants were given by (4) who also derived constants for 3He_2 .
- ^s $w_e y_e = - 0.038$. See also ^r.
- ^t $\gamma_e = - 0.00138$.

- ^uObserved in absorption in a pulsed discharge (19).
- ^v $w_e y_e = - 0.1267$.
- ^w $\gamma_e = - 0.00273$.
- ^x $w_e y_e = - 1.2566$, $w_e z_e = - 0.4875$.
- ^y $+ 0.1629(v+\frac{1}{2})^2 - 0.0655(v+\frac{1}{2})^3$.
- ^z $D_1 = 5.76 \times 10^{-4}$, $D_2 = 6.11 \times 10^{-4}$, ...; $H_0 = 1.32 \times 10^{-8}$, $H_1 = -0.40 \times 10^{-8}$, $H_2 = -5.3 \times 10^{-8}$, ...
- ^a $w_e y_e = - 0.0483$.
- ^bThe constants refer to $3\Pi^-$ and were derived by (13) from the $d, f \rightarrow b$ bands; $B(3\Pi^-) - B(3\Pi^+) \approx + 0.026$. The triplet splitting is partially resolved in the $d \rightarrow b$ bands (13).
- ^c $\gamma_e = - 0.00170$.
- ^d $D_1 = 5.34 \times 10^{-4}$, ...; $H_0 = 2.80 \times 10^{-8}$, $H_1 = 3.60 \times 10^{-8}$, ...
- ^e $w_e y_e = - 0.38$ (36).
- ^fThere is good experimental (20)(30) and theoretical (7) (21)(48) evidence for a potential maximum in this state. (20) place the maximum at 0.067 eV above the asymptote; the net dissociation energy is 1.850 eV.
- ^gFrom molecular beam magnetic resonance experiments (58) (65) have determined the triplet splitting for $N=1$ and 3. The splitting constants (extrapolated to $N=0$) are $\lambda = - 0.03666$, $\gamma = - 0.0000808\text{ cm}^{-1}$. An ab initio calculation (55) gives $\lambda = - 0.04089$.
- ^h $\gamma_e = - 0.00462$ (36).
- ⁱ $\beta_e = + 0.07 \times 10^{-4}$; $H_e \approx 2.7 \times 10^{-8}$ (36).
- ^jEnergy of a $3\Sigma_u^+, v=0, N=0$ above $\text{He}(^1S) + \text{He}(^1S)$, based on $D_0^0(\text{He}_2^+) = 19073\text{ cm}^{-1}$. See also ^c p. 297.

State	T_e	w_e	$w_e x_e$	B_e	α_e	D_e (10^{-4}cm^{-1})	r_e (\AA)	Observed Transitions		References
								Design.	ν_{00}	
$^4\text{He}_2$ (continued)										
S $1\pi_g$	8p π [177515]			(7.21)	(0.22)		(1.08 ₁)	S \rightarrow A, R	30228.6 Z	(39)
R $1\pi_g$	7p π [176983]			(7.22)	(0.22)		(1.08 ₀)	R \rightarrow A, R	29696.4 Z	(39)
P $1\pi_g$	6p π [176160]			(7.22)	(0.22)		(1.08 ₀)	P \rightarrow A, R	28873.9 Z	(39)
M $\left\{ \begin{array}{l} 1\Delta_u \\ 1\pi_u \\ 1\Sigma_u^+ \end{array} \right.$	5d δ [(174838)]			$\left[\begin{array}{l} 7.09 \\ 7.07 \end{array} \right]^a_b$			[1.09 ₁]	M \rightarrow B, $\left\{ \begin{array}{l} (24050) \\ (24000) \\ (23960) \end{array} \right.$	(39)	
	5d π [(174788)]									
	5d σ [(174748)]									
L $1\pi_g$	5p π [174794]			(7.23)	(0.22 ₂)		(1.07 ₉)	L \rightarrow A, R	27507.8 Z	(39)
J $\left\{ \begin{array}{l} 1\Delta_u \\ 1\pi_u \\ 1\Sigma_u^+ \end{array} \right.$	4d δ [172416]			[7.097] ^c		[5.0]	[1.089 ₄]	J \rightarrow C, $\left\{ \begin{array}{l} V \\ R \end{array} \right.$	14183.90 ^c Z 14058.37 ^c Z 13990.32 ^d Z	(50)
	4d π [172290]			[7.080] ^c		[5.4]	[1.090 ₈]			
	4d σ [172222] ^d			Strongly perturbed by l -uncoupling and interaction with H $1\Sigma_u^+$. ^e						
I $1\pi_g$	4p π [172266]			(7.24 ₂)	(0.22 ₃)		(1.078)	I \rightarrow A, R	24979.6 Z	(39)
H $1\Sigma_u^+$	4s σ [171951]			(7.26) ^e	(0.23)		(1.07 ₇)	H \rightarrow C, V	13719.5 Z	(39)(50)
								H \rightarrow B, R	21163.5 Z	(39)(50)
F $\left\{ \begin{array}{l} 1\Delta_u \\ 1\pi_u \\ 1\Sigma_u^+ \end{array} \right.$	3d δ 166304	1706.59 Z	35.06	7.230 ^g	0.225 ^h	[5.20] ⁱ	1.079 ₄	F \rightarrow B, $\left\{ \begin{array}{l} V \\ R \end{array} \right.$	16360.9 Z 16008.3 Z 15837.5 Z	(12)(17)
	3d π 165971 ^f	1670.57 Z	40.03	7.156 ^g	0.235	[5.24] ^j	1.084 ₉			
	3d σ (165813)	[1564.25] Z	(40)	7.098 ^g	0.246	[5.21] ^k	1.089 ₄			
E $1\pi_g$	3p π 165911	1721.1 ₉ Z	34.76 ^l	7.270 ₅ ^m	0.215 ₆ ⁿ	5.20	1.076 ₄	E \rightarrow A, R	19476.61 Z	(36)
D $1\Sigma_u^+$	3s σ 165085	1746.43 Z	35.54	7.365	0.218 ₀ ^o	5.24 ^p	1.069 ₄	D \rightarrow B, ^q R _V	15161.81 Z	(12)
								D \rightarrow X, ^r	continuum	
C $1\Sigma_g^+$	3p σ 157415	1653.43 Z	41.0 ₄ ^s	7.052	0.215 ^t	5.08 ^u	1.092 ₉	C \rightarrow A, R	10945.50 Z	(11)(39)
B $1\pi_g$	2p π 149914	1765.76 Z	34.39 ^v	7.403 ^w	0.216 ^w	5.02 ^w	1.066 ₇	(B - A)	3501.5 ^x	
A $1\Sigma_u^+$	2s σ 146365 ^y	1861.3 ₃ Z	35.28 ^z	7.778 ₉	0.216 ₆ ^{a'}	5.44	1.0406	A \leftrightarrow X, ^{b'}	147279 ^{c'}	(8)(18)(22)
								Hopfield continuum		(38)

State	T_e	ω_e	$\omega_e x_e$	B_e	α_e	D_e (10^{-4} cm^{-1})	r_e (\AA)	Observed Transitions		References
								Design.	v_{00}	
$^4\text{He}_2$ (continued)										
X $1\Sigma_g^+$	0	Repulsive potential with very small well ($D_e = 0.90 \text{ meV}$).					2.97 ^d			

- He₂: ^aAverage of Π^- and Δ^- as given by (39).
^bAverage of Σ^+ , Π^+ , Δ^+ as given by (39).
^cThese constants refer to Π^- and Δ^- which are less affected by ℓ -uncoupling.
^dRefers to $N=1$.
^e(50) give average effective constants for the four interacting components $J(1\Delta_u^+, 1\Pi_u^+, 1\Sigma_u^+)$ and $H(1\Sigma_u^+)$.
^fSee ^mp. 295.
^gThese constants are corrected for ℓ -uncoupling effects.
^h $\gamma_e = -0.0045$.
ⁱ $D_1 = 5.26 \times 10^{-4}$, $H_0 = 2.18 \times 10^{-8}$.
^j $D_1 = 5.26 \times 10^{-4}$, $H_0 = 1.97 \times 10^{-8}$.
^k $D_1 = 5.29 \times 10^{-4}$.
^l $\omega_e y_e = -0.023$.
^mThe rotational constants refer to Π^- (36); $B(\Pi^-) - B(\Pi^+) \approx +0.044$.
ⁿ $\gamma_e = -0.00227$.
^o $\gamma_e = -0.0059$.
^p $\beta_e = +0.19 \times 10^{-4}$.
^qFranck-Condon factors (25).
^rThe weak maximum near 676 \AA in the Hopfield continuum is ascribed by (37) to the transition $D \rightarrow X$.
^s $\omega_e y_e = +0.3549$, $\omega_e z_e = -0.1315$. Calculations of (47) give a potential hump of 0.22 eV at 2.09 \AA ; see also (59).
^t $\gamma_e = -0.0111$. | and (60).

- ^u $\beta_e = +0.10 \times 10^{-4}$; $H_0 = 1.72 \times 10^{-8}$, ...
^v $\omega_e y_e = -0.0267$ (12).
^w B_e refers to Π^- ; $B(\Pi^-) - B(\Pi^+) = +0.019$; $\gamma_e = -0.0015$,
^xFrom (39). | $\beta_e = +0.05 \times 10^{-4}$ (12).
^yRKR potential curve (34)[see also (39)]; ab initio potential (42). The latter gives vibrational and rotational levels in good agreement with the experimental values. (8) have established, from the absorption and emission bands near 600 \AA , a potential maximum of 0.059 eV in the A $1\Sigma_u^+$ state. Theoretical work by (15)(16)(47) gives maxima of 0.084, 0.153, 0.061 eV, respectively; see also (59)(60).
^z $\omega_e y_e = -0.136$ (36).
^a $\gamma_e = -0.00273$ (36).
^bTransitions from the low vibrational levels of A $1\Sigma_u^+$ to X $1\Sigma_g^+$ give rise to the Hopfield continuum; see MOISPEC 1, 404. Transitions from the high vibrational levels as well as the continuous range of energy levels of A $1\Sigma_u^+$ to X $1\Sigma_g^+$ give rise to diffuse bands near 600 \AA observed in emission (18) (22) and absorption (8)(38) with quite different intensity distribution. See also (6)(37) and (31)(52)(54). Observed absorption coefficients near 600 \AA (38) agree fairly well with those predicted by (41).
^cEnergy of the $v=0, N=0$ level of A $1\Sigma_u^+$ relative to He(¹S) + He(¹S), calculated from the corresponding value for a $3\Sigma_u^+$ by (continued p. 299)