TPG4190 Seismic data acquisition and processing Lecture 5: Deghosting

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Overview

- ► The wave equation
- ► Up- and downgoing waves
- ► Removal of the ghost effect

The acoustic wave equation descries wave motion in a material with no shear forces, i.e. only P-waves can exist.

$$\rho(\mathbf{x})\ddot{u}_i(\mathbf{x},t) = \partial_i \sigma(\mathbf{x},t) + f_i(\mathbf{x},t), \tag{1}$$

$$\sigma(\mathbf{x},t) = \kappa(\mathbf{x})\partial_i u_i(\mathbf{x},t) + q(\mathbf{x},t). \tag{2}$$

- ightharpoonup
 ho: Density
- \triangleright u_i : Component i of particle dispalcement
- $ightharpoonup \sigma: -p$ where p is the pressure.
- \triangleright κ : Bulk modulus
- ► *f_i*: component *i* of source body force
- q: source of volume injection type

Divide equation (2) with density and differentiate equation (2) two times w.r.t. time:

$$\ddot{u}_i(\mathbf{x},t) = \frac{1}{\rho(\mathbf{x})} \partial_i \sigma(\mathbf{x},t) + \frac{f_i(\mathbf{x},t)}{\rho(\mathbf{x})},\tag{3}$$

$$\ddot{\sigma}(\mathbf{x},t) = \kappa(\mathbf{x})\partial_i \ddot{u}_i(\mathbf{x},t) + \ddot{q}(\mathbf{x},t). \tag{4}$$

Insert equation (3) into equation (4) to obtain:

$$\ddot{\sigma}(\mathbf{x},t) = \kappa(\mathbf{x})\partial_i \left[\frac{1}{\rho(\mathbf{x})} \partial_i \sigma(\mathbf{x},t) \right] + \partial_i \left[\frac{f_i(\mathbf{x},t)}{\rho(\mathbf{x})} \right] + \ddot{q}(\mathbf{x},t). \quad (5)$$

Assume f=0 and constant density $\rho=\rho_0$

$$\ddot{\sigma}(\mathbf{x},t) = c^2(\mathbf{x})\partial_i\partial_i\sigma(\mathbf{x},t) + \ddot{q}(\mathbf{x},t). \tag{6}$$

where $c(\mathbf{x}) = \sqrt{\kappa(\mathbf{x})/\rho_0}$ or

$$\nabla^2 \sigma(\mathbf{x}, t) - \frac{1}{c^2(\mathbf{x})} \ddot{\sigma} = s(\mathbf{x}, t)$$
 (7)

where $s(\mathbf{x},t) = -q(\mathbf{x},t)/c^2(\mathbf{x})$

We now assume that the velocity is independent of x and y coordinates and only depends on depth; c(x) = c(z)

$$\nabla^2 \sigma(\mathbf{x}, t) - \frac{1}{c^2(z)} \ddot{\sigma} = s(\mathbf{x}, t)$$
 (8)

 $abla^2$ operator is

$$\nabla^2 \sigma(\mathbf{x}, t) = \frac{\partial^2 \sigma}{\partial_x^2}(\mathbf{x}, t) + \frac{\partial^2 \sigma}{\partial_y^2}(\mathbf{x}, t) + \frac{\partial^2 \sigma}{\partial_z^2}(\mathbf{x}, t)$$
(9)

The Fourier transform over x- and y is obtained by replacing $\partial_x^2 \to (-ik_x)^2 = -k_x^2$ and $\partial_y^2 \to (-ik_y)^2 = -k_y^2$.

$$\nabla^2 \sigma(\mathbf{x}, t) \to (-k_x^2 - k_y^2) \sigma(k_x, k_y, z, t) + \frac{\partial^2 \sigma}{\partial z^2} (k_x, k_y, z, t) \quad (10)$$

We also have for the fourier transform over time

$$\frac{\partial^2 \sigma}{\partial_t^2}(\mathbf{x}, t) \to -\omega^2 \sigma(k_x, k_y, \omega) \tag{11}$$

Putting this together we get for the wave equation

$$-(k_x^2 + k_y^2)\sigma(\mathbf{x}, t) + \frac{d^2\sigma}{d^2z}(k_x, k_y, z, \omega) + \frac{\omega^2}{c^2(z)}\sigma(k_x, k_y, z, \omega)$$
(12)
= $s(k_x.k_y, z, \omega)$ (13)

Using

$$k_z^2 = \omega^2/c^2(z) - (k_x^2 + k_y^2) = \omega^2/c^2(z) - k_r^2,$$
 (14)

where $k_r^2 = k_x^2 + k_y^2$.

we have

$$\frac{d^2\sigma(k_r,z,\omega)}{d^2z} + k_z^2(k_r,z,\omega)\sigma(k_r,z,\omega) = s(k_r,z,\omega)$$
 (15)

The solution of the this wave equation is

$$\sigma(k_r, z, \omega) = U(k_r, \omega) \exp(ik_z z) + D(k_r, \omega) \exp(-ik_z z), \quad (16)$$

Where the constants U and D are interpreted as the amplitude of the upgoing wave and the downgoing wave, respectively. We also have from the equation of motion

$$\rho_0 \ddot{u}_z(\mathbf{x}, t) = \rho_0 \dot{v}_z(\mathbf{x}, t) = \frac{\partial \sigma(\mathbf{x}, t)}{\partial z}.$$
 (17)

Or in the wave-number frequency domain

$$-i\omega\rho_0 v_z(k_r,\omega,z) = \frac{\partial\sigma(k_r,\omega,z)}{\partial z},$$
 (18)

$$v_z(k_r, \omega, z) = \left(\frac{-1}{i\omega\rho_o}\right) \frac{\partial\sigma(k_r, \omega, z)}{\partial z}.$$
 (19)

Using the last equation together with equation (16)

$$v_z(k_r, \omega, z) = \left(\frac{-ik_z}{i\omega\rho_o}\right) \left[U(k_r, \omega) \exp(ik_z z) - D(k_r, \omega) \exp(-ik_z z)\right], (20)$$

For simplicity we set z=0 at the receiver depth. Equation (16) and (20) then gives

$$\left(\frac{i\omega\rho}{-ik_z}\right)v_z(k_r,\omega) = U(k_r,\omega) - D(k_r,\omega), \qquad (21)$$

$$\sigma(k_r,\omega) = U(k_r,\omega) + D(k_r,\omega). \qquad (22)$$

U-D separation

Adding and subtracting equations (22) and (22) gives

$$U(k_r,\omega) = \left(\frac{i\omega\rho}{-2ik_z}\right)v_z(k_r,\omega) + \sigma(k_r,\omega), \qquad (23)$$

$$D(k_r,\omega) = -\left(\frac{i\omega\rho}{-2ik_z}\right)v_z(k_r,\omega) + \sigma(k_r,\omega). \tag{24}$$









