

Assignment 7

Quantum Information & Computing

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Theory // Quantum Ising Model in Transverse Field

The **Quantum Ising model** describes a many-body quantum system, which consists of a linear chain of N interacting $1/2$ -spins in presence of an external field of intensity λ .

$$\hat{\mathcal{H}} = \sum_i^{N-1} \sigma_x^{(i)} \sigma_x^{(i+1)} + \lambda \sum_i^N \sigma_z^{(i)}$$

the notation $\sigma_x^{(i+1)}$ contracts $\mathbb{1}_1 \otimes \dots \otimes \mathbb{1}_{i-1} \otimes \sigma_x \otimes \mathbb{1}_{i+1} \otimes \dots \otimes \mathbb{1}_N$.

Being it simple but non-trivial, the Ising model is used as a benchmark test for many numerical algorithms.

Using a Mean Field (MF) approach - therefore on the ansatz $|\Psi_{MF}\rangle = \otimes_i^N |\psi_i\rangle$ - the energy density $\varepsilon_0 = E_0/N$ of the ground state of Ising model would be roughly, at the limit $N \rightarrow \infty$,

$$\varepsilon_0 = \begin{cases} -1 - \lambda^2/4 & \text{if } |\lambda| < 2 \\ -|\lambda| & \text{else} \end{cases} \quad (1)$$

Even though this formula is not exact, it encodes a meaningful physical behaviour: the discontinuity of the second derivative in $\lambda = 2, -2$ implies a **quantum phase transition**.

In this homework, we tackle the Quantum Ising model using a "direct" approach.



1. We explicitly build $\hat{\mathcal{H}}$ for finite N and given λ
2. We **diagonalize** $\hat{\mathcal{H}}$, taking the first k eigenvalues and eigenvectors
3. We repeat for different values of λ , sampling in the interval $\lambda \in [0, 3]$


The main computational problems are:

- $\hat{\mathcal{H}}$ is a symmetric matrix of real (double precision) values, of size $2^N \cdot 2^N$. Therefore, the number of elements scales exponentially in N . For instance, in the case of $N = 15$ it would require ~ 8 GB of RAM to store the matrix as `real*8` values. Scaling up to $N = 16$ would require ~ 32 GB of RAM.
- Diagonalization is the most time-consuming step, as it scales exponentially in N , too.

Optimizations // Packed storage and inplace tensor products

These optimizations came to my mind:

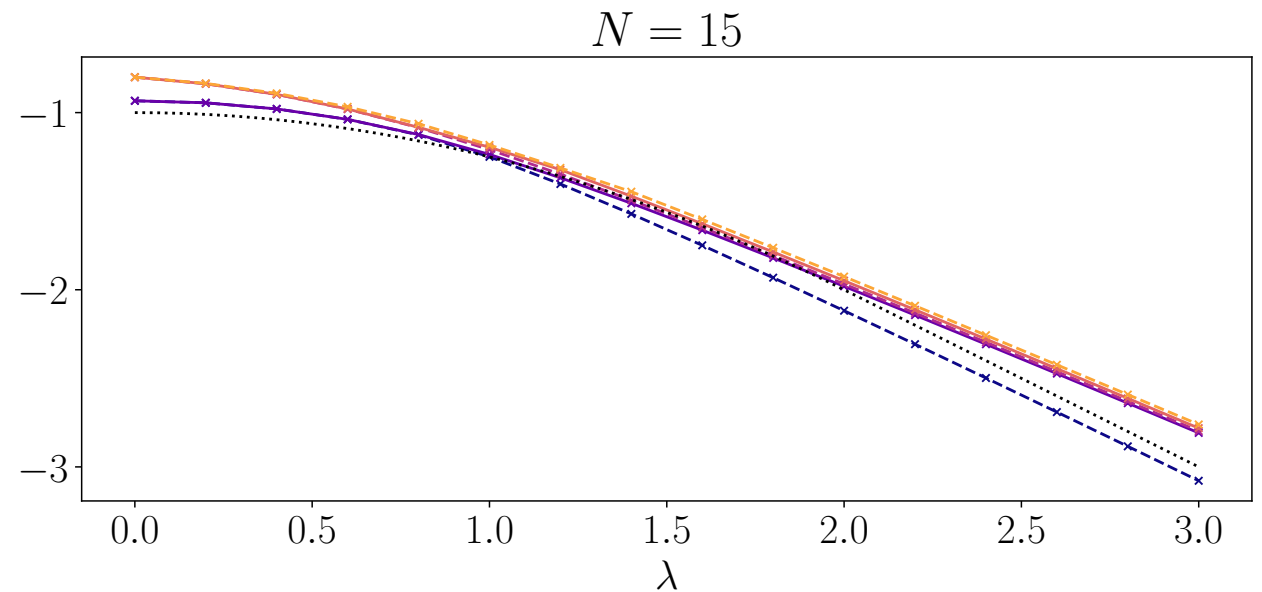
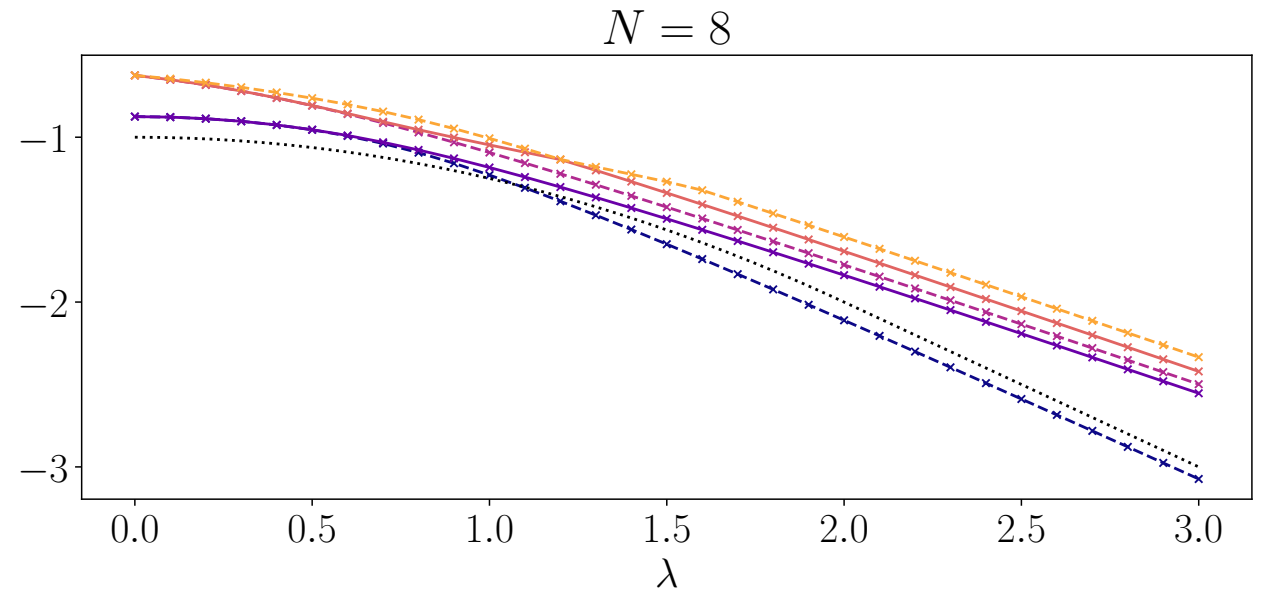
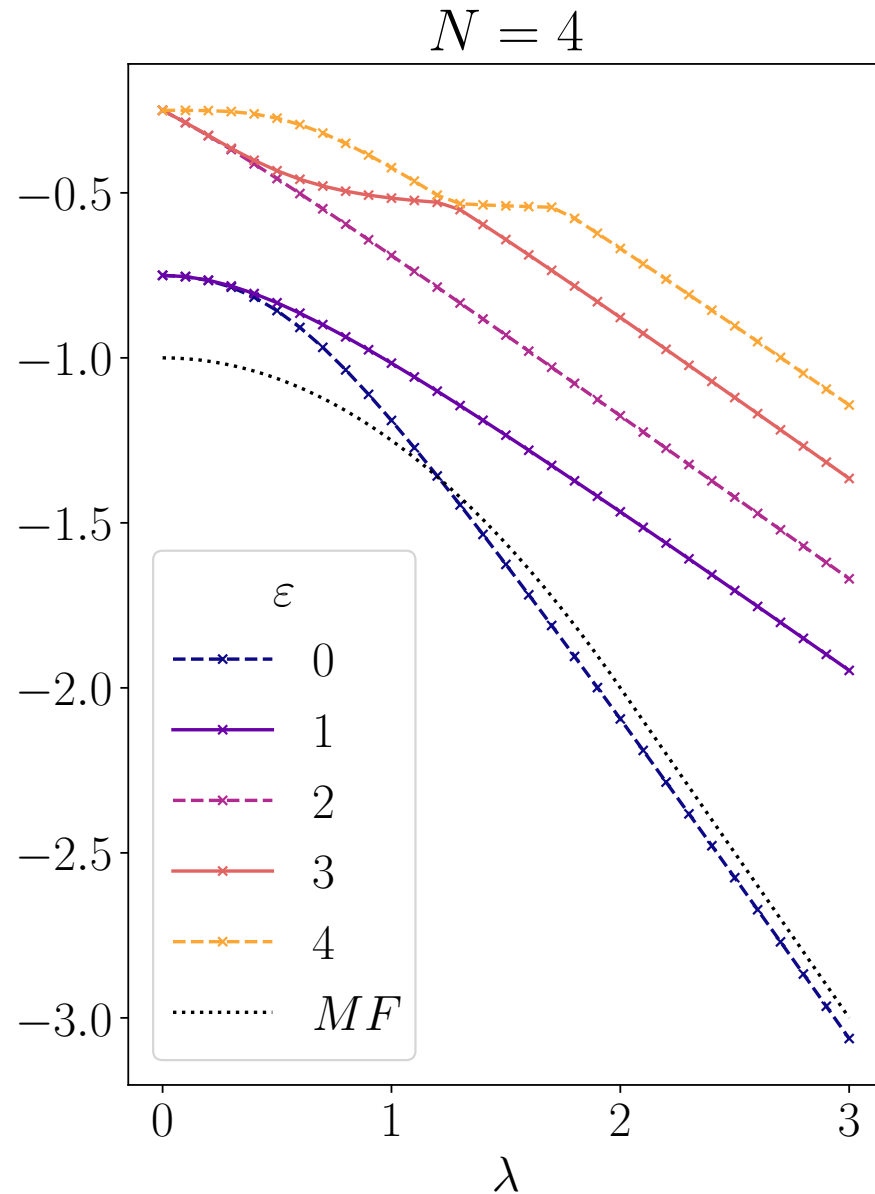
- In order to gain a factor 2 on the storage requirement, we can store the Hamiltonian in (upper triangular) **packed format**, as it is symmetric.
 ⇒  The scaling is still exponential, but reduced by a factor 2. An Hamiltonian for $N = 15$ would require ~ 4 GB of RAM.
-  All the mathematical operations have to be implemented for upper triangular matrices. For instance, what about the tensor product operations \otimes required to build $\hat{\mathcal{H}}$?
- To make the diagonalization virtually faster, I have used LAPACK with optimized multicore routines from **OpenBLAS**. Even though the algorithm still scales exponentially, the new backend makes the computer work in parallel with more cores...

To solve the issue () , I have implemented routines for **tensor product of upper triangular matrices** ($A \otimes B \leftrightarrow \text{kronp}(A, B)$) with **inplace sum** as further optimization; i.e. the chain of operations

$$C \leftarrow \text{kronp}(A, B) + C$$

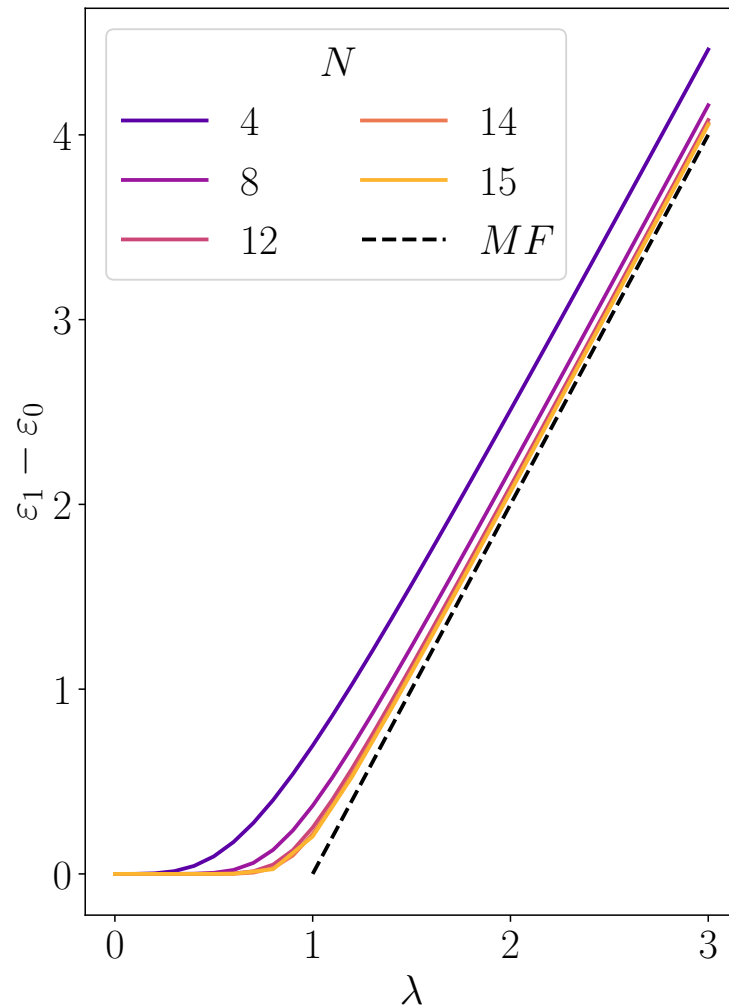
will not require the intermediate storage of the matrix $\text{kronp}(A, B)$. Eventually, I optimize with specific routines the tensor products with **left and right identities**: $\text{kronp_l1}(n, X)$ (for $\mathbb{1}_n \otimes X$) and $\text{kronp_r1}(X, n)$ (for $X \otimes \mathbb{1}_n$).

Results // Energy density spectrum for N spins

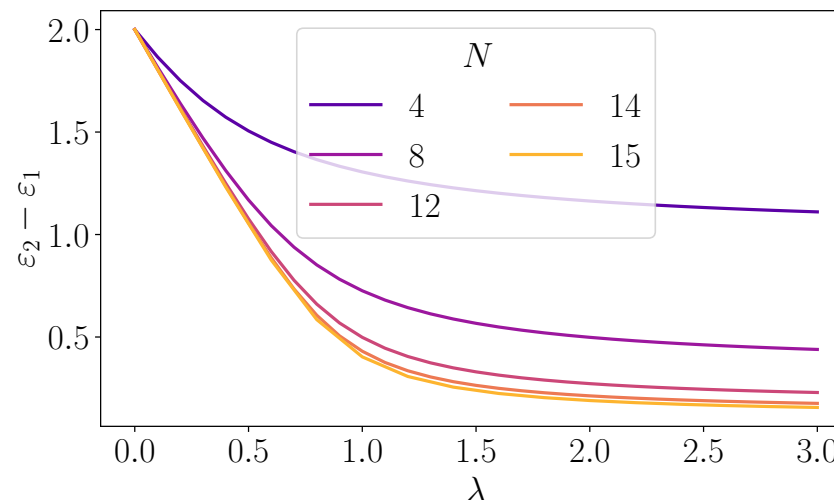


Results // Energy gap

Another interesting perspective on the Ising problem is the **energy gap between the ground state and the first two excited levels** ε_1 and ε_2 .



- In the area of stronger external field $|\lambda| > 1$, known in **literature** as **disordered phase**, we observe that systems with more spins (larger N) have ground-state energy closer to the theoretical limit in MF theory.
- In the regime $0 < |\lambda| < 1$ the system is said to be in **ordered phase**, and the gs breaks the spin-flip symmetry.
- For $\lambda \rightarrow 0$, all the levels are degenerate, as expected.



The gap between the second excited state and the first excited state

$$\varepsilon_2 - \varepsilon_1 \rightarrow 2$$

for $\lambda \rightarrow 0$.