

# Modern Particle Physics Experiments

## Particle Identification Detectors

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**Lecture 06**  
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## Detector concepts

Depending on the particle type and application, particle detectors can be divided into three main classes:

- **Tracking detectors**

Measure position/trajectory of charged particles,  
based on energy losses due to ionization or activation of material.

We try to minimize particle interactions

⇒ gaseous detectors or **thin semiconductor layers**

- **Calorimeters**

Measure particle energy by absorbing it in the dense medium

Interactions of high energy incident particle

⇒ electromagnetic or hadronic cascade

- **Particle identification detectors** (today)

Use different processes to improve particle identification capabilities

Cherenkov detectors, Transition radiation detectors, Time-Of-Flight ...

## References

- Particle Physics Reference Library (vol.2)  
Review of the state of the art in detector physics and related data-taking technology (open access)
- PDG reviews:
  - Passage of particles through matter
  - Particle detectors at accelerators
  - Particle detectors for non-accelerator physics

## Particle Identification Detectors

- 1 General concept of collider experiments
- 2 Time-of-Flight measurement
- 3 Cherenkov detectors
- 4 Transition radiation
- 5 Large detector systems

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## General concept - layer structure

Modern **universal detectors** at particle colliders are build from many **different sub-components**.

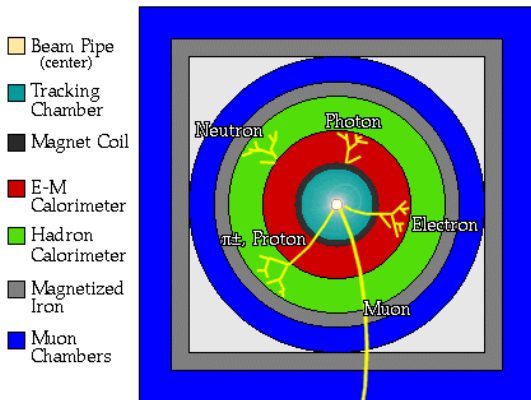
We arrange detectors in such a way as to obtain **best measurement** for all possible particles, as well as their **(at least partial) identification**.

⇒ detectors which **interact least** with produced particles are placed **closest** to the interaction point **gas detectors, thin silicon sensors**

⇒ detectors which **absorb** particles are placed at **largest distances** from interaction point **calorimeters, muon detectors**

## General concept - layer structure

Modern **universal detectors** at particle colliders are build from many different **sub-components**.



## Universal detector

Layout describing most of recent, present-day and future experiments at colliders (LEP, HERA, Tevatron, LHC, ILC, CLIC, FCCee, FCChh...):

Starting from the center of the detector:

- vertex detector
  - as close to the beam line as possible, used to measure the exact position of the interaction point, allows for identification of short-lived particles (secondary vertexes)
  - silicon pixel detectors



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Starting from the center of the detector:

- vertex detector  
as close to the beam line as possible, used to measure the exact position of the interaction point, allows for identification of short-lived particles (secondary vertexes)  
silicon pixel detectors
- tracking detectors  
measure tracks of charged particles, allows to determine their momentum from bending in magnetic field  
gas detectors or silicon strip detectors

## Universal detector

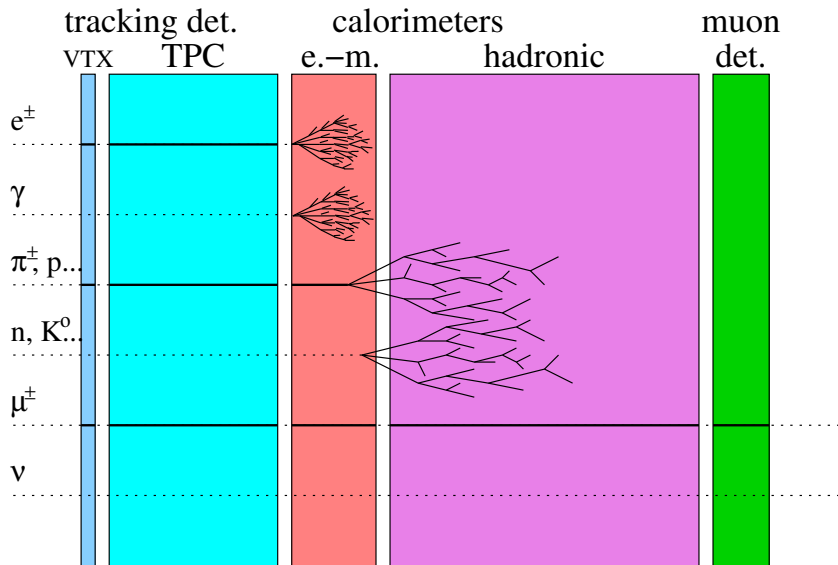
- electromagnetic calorimeter  
electron and photon energy measurement  
dense material absorbing EM cascade  
(copper, lead, tungsten)

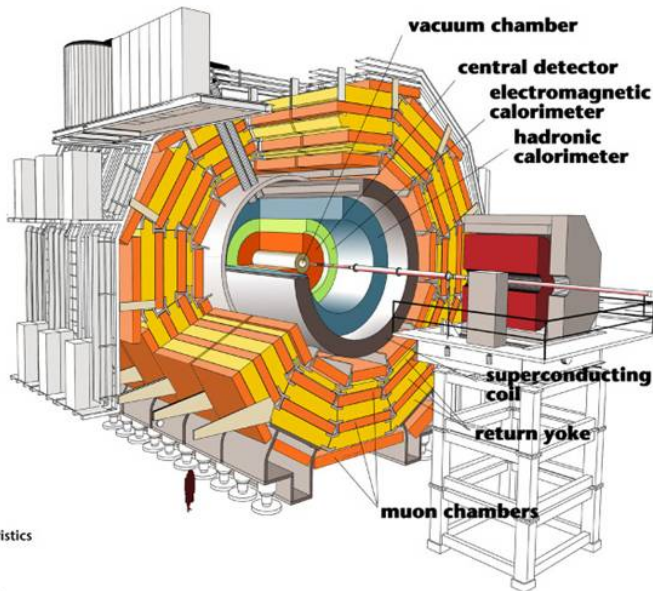
## Universal detector

- electromagnetic calorimeter  
electron and photon energy measurement  
dense material absorbing EM cascade  
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- hadron calorimeter  
hadron energy measurement (protons, neutrons, pions, kaons)  
dense material absorbing hadronic cascade (iron, lead, uranium)  
hadronic cascade is many times longer than EM one

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hadronic cascade is many times longer than EM one
- muon detectors  
identify muons - only charged particles which can pass both calorimeters with small energy losses

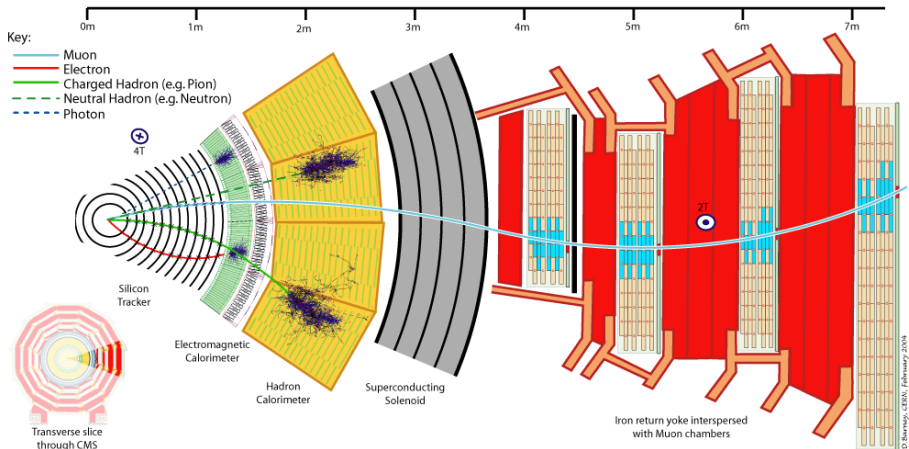




## Detector characteristics

Width: 22m  
Diameter: 15m  
Weight: 14'500t

## Compact Muon Solenoid - CMS



## Particle Identification

General concept, layout of universal collider experiment allows for partial particle identification. Following particles can be directly identified:

- electrons, positrons and photons  
dense shower in EM calorimeter, with or without matching track
- muons  
charged track + MIP signal in calorimeters + signal in muon detectors
- charged and neutral hadrons  
all other particles (with and without matched track);  
for small fraction of events, more detailed identification is possible  
based on the decay reconstruction...

The general problem is that universal detector (tracking+calorimeters) can not differentiate between different hadron types.

**Can we do it with some dedicated instruments?**

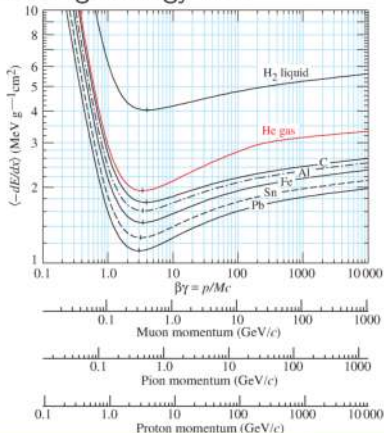


## Ionisation losses Lecture 4

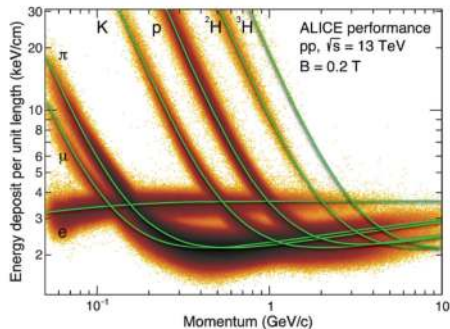
Universal shape of the dependence for different particles!

Scales with  $\beta\gamma = p/M$ .

### Average energy loss



### Ionization losses in ALICE TPC



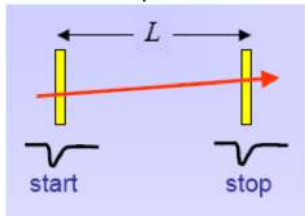
Only partial identification possible...

## Particle Identification Detectors

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## Principle of operation

Time of flight measurement can be used to estimate particle mass



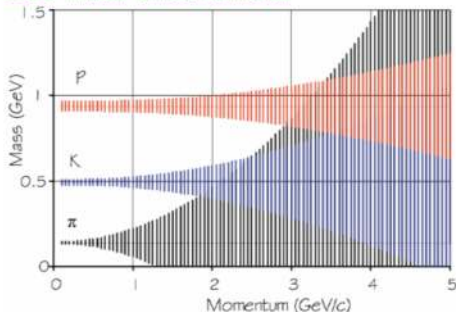
$$p = \beta\gamma m$$

$$l = \beta ct \Rightarrow m^2 = \frac{p^2}{c^2} (c^2 t^2 - l^2)$$

Example:

$$l = 12m, \sigma_t = 150ps, \frac{\sigma_p}{p} = 1\%$$

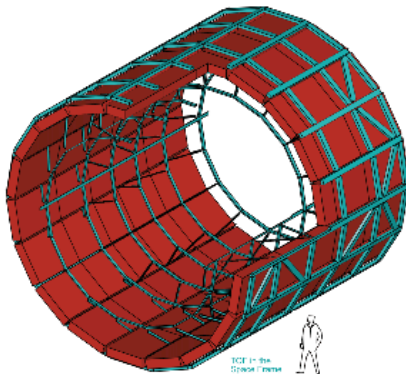
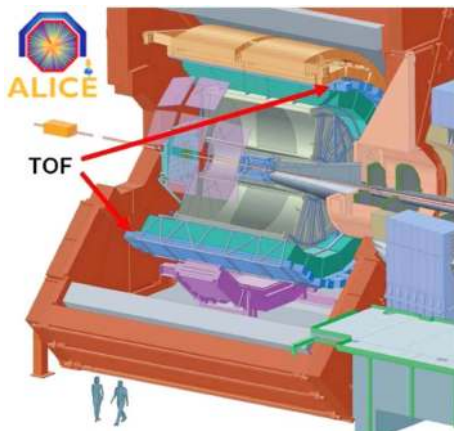
$\pm 1\sigma$  mass uncertainties



Efficient identification only in low momentum...

## Example: ALICE ToF Detector

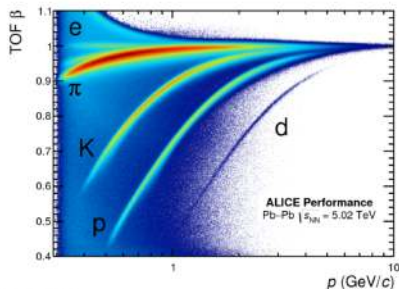
Placed between tracking detector system and calorimeters



## ALICE ToF: performance

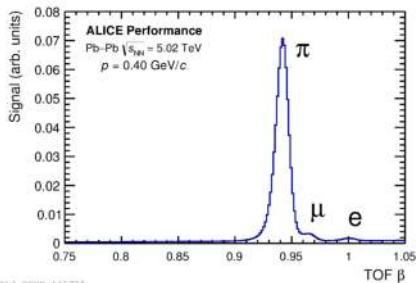
arXiv:1809.00574

2D event distribution



ALI-PRF-141119

Reconstructed  $\beta$  for  $p = 0.4$  GeV

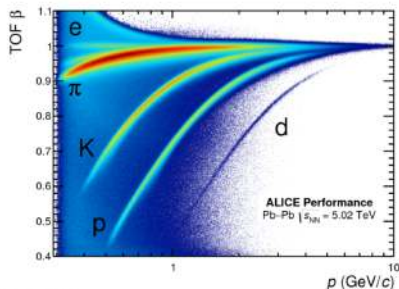


ALI-PRF-141735

## ALICE ToF: performance

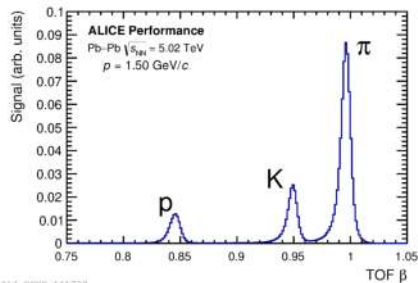
arXiv:1809.00574

2D event distribution



ALI-PPHF-104119

Reconstructed  $\beta$  for  $p = 1.5$  GeV

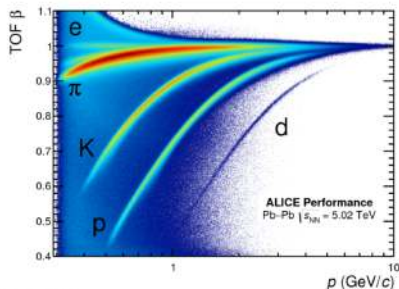


ALI-PPHF-141739

## ALICE ToF: performance

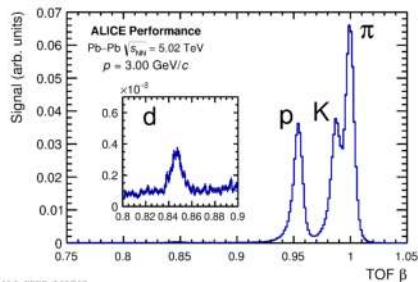
arXiv:1809.00574

### 2D event distribution



ALI-PPHF-141116

### Reconstructed $\beta$ for $p = 3$ GeV

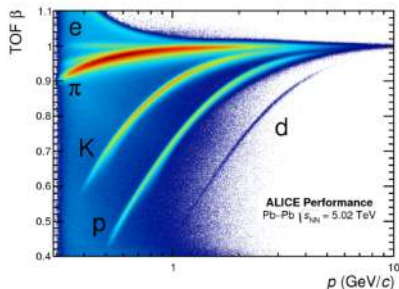


ALI-PPHF-141743

## ALICE ToF: performance

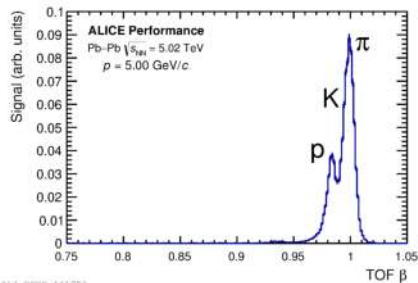
arXiv:1809.00574

2D event distribution



ALI-PRF-104119

Reconstructed  $\beta$  for  $p = 5$  GeV

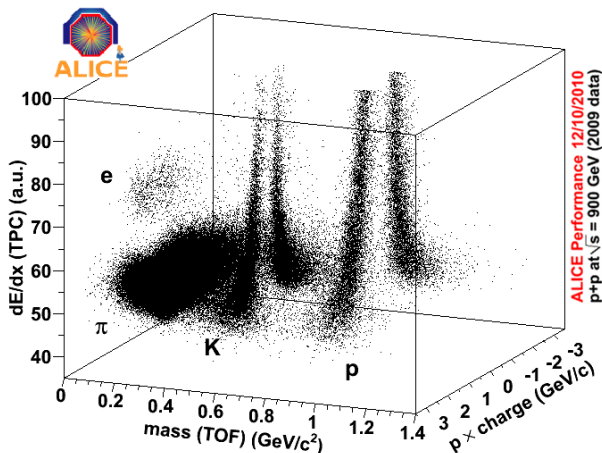


ALI-PRF-141751



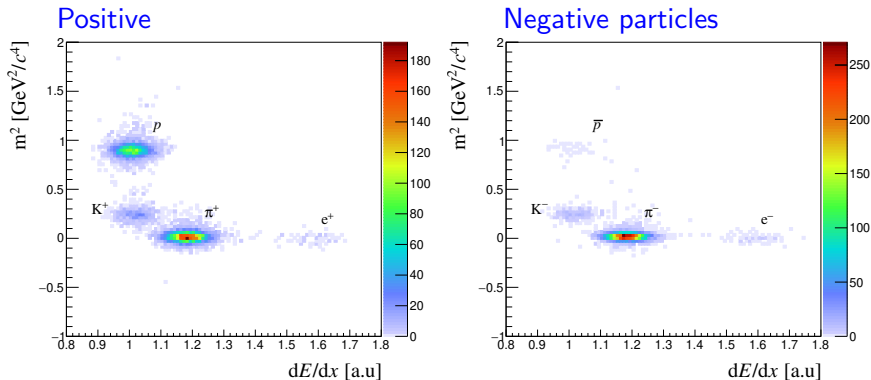
## ALICE ToF vs $dE/dx$ [arXiv:1101.3276](https://arxiv.org/abs/1101.3276)

Improved PID capabilities by combining the ionization measurements in the TPC and the mass calculated using the TOF signal



## NA69/SHINE ToF vs $dE/dx$ P.Podlaski, PhD Thesis

Improved PID capabilities by combining the ionization measurements in the TPC and the mass calculated using the TOF signal



Example distributions for data bin:  $2 \text{ GeV}/c < p < 3 \text{ GeV}/c$  &  $0.5 \text{ GeV}/c < p_T < 0.6 \text{ GeV}/c$

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- 1 General concept of collider experiments
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- 3 Cherenkov detectors**
- 4 Transition radiation
- 5 Large detector systems

## Cherenkov radiation

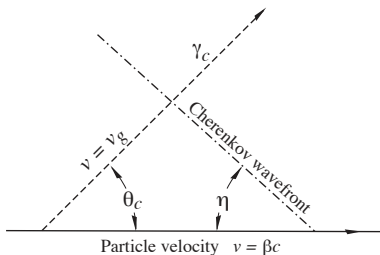
When speed of charged particle is greater than the local phase speed of light ( $\beta > \frac{1}{n}$ ), small part of its energy is lost in form of Cherenkov radiation

Angle of radiation:

$$\cos \theta_c = \frac{1}{n\beta}$$

Wide spectra of radiation

$$\begin{aligned} \frac{d^2 N_\gamma}{dE_\gamma dx} &= \frac{\alpha z^2}{hc} \sin^2 \theta_c(E_\gamma) \\ &\approx 370 \frac{1}{\text{eV} \cdot \text{cm}} \cdot \sin^2 \theta_c \end{aligned}$$

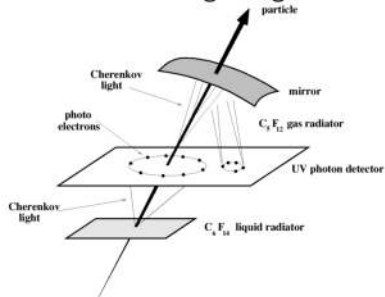


Measurement of the radiation opening angle  $\Rightarrow$  particle velocity

First applications: threshold counters

## RHIC Ring Imaging CHerenkov detector

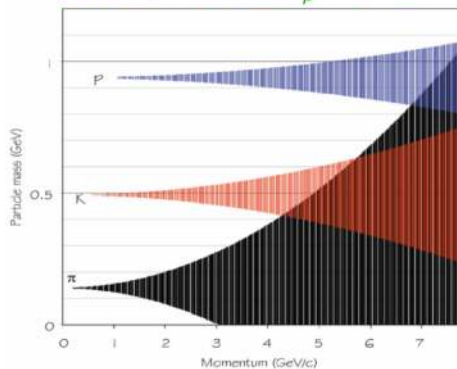
When the light is emitted towards the spherical mirror, it can be focused to form ring images



Ring size allows for direct emission angle  $\Rightarrow$  speed determination

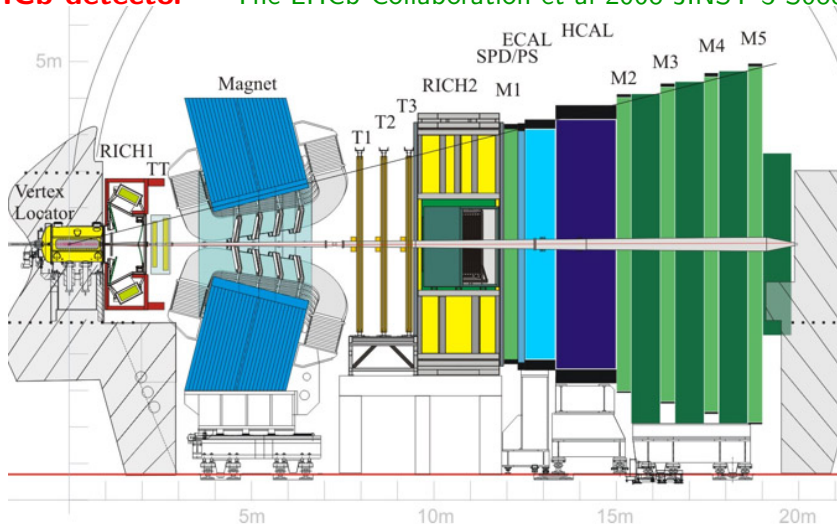
Example:

$$n = 1.333, \sigma_\theta = 15 \text{ mrad}, \frac{\sigma_p}{p^2} = 5 \cdot 10^{-5}$$



## LHCb detector

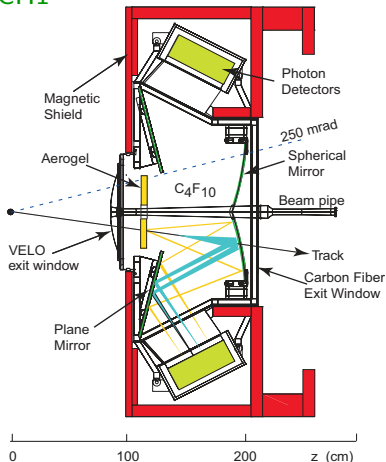
The LHCb Collaboration et al 2008 JINST 3 S08005



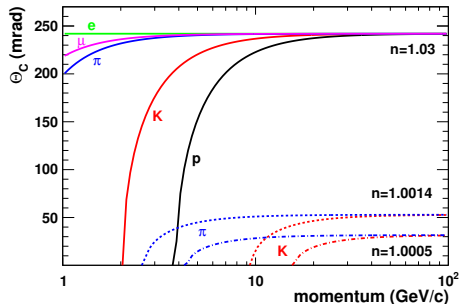
## LHCb RICH

arXiv:1101.3276

### RICH1



Three radiators to cover large momentum range

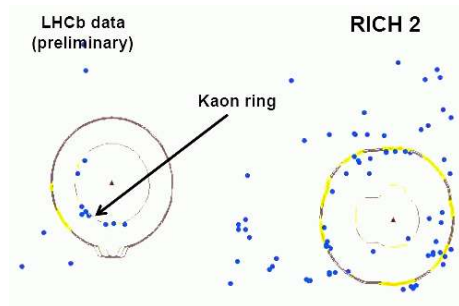
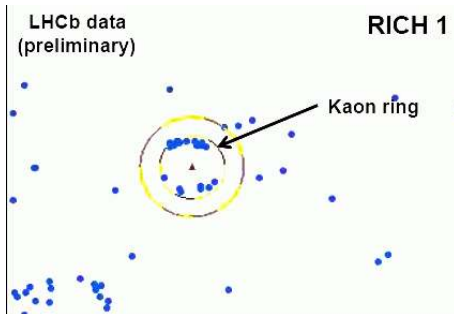


Physics goal:

K/ $\pi$  separation up to 100 GeV

## LHCb RICH

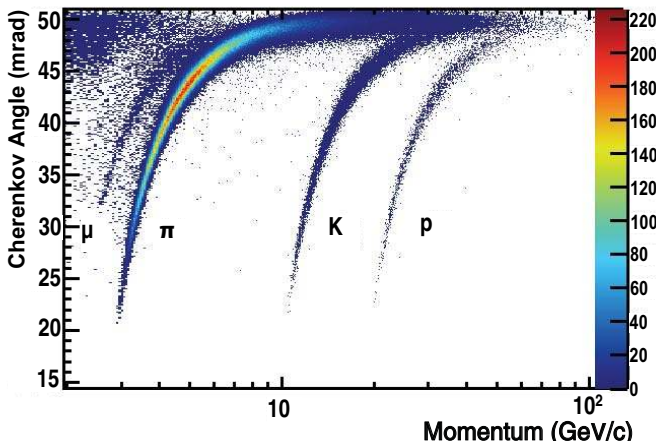
RICH images from first LHCb data





## LHCb RICH Carla Marin @ ICHEP'2018

Combining light rings and momentum measurements in LHC Run 2:



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## Transition radiation

Charged particle radiates also when it crosses suddenly from one medium to another with different optical properties

Average radiated energy:

$$W \approx \frac{\alpha}{3} \hbar \omega_p \gamma$$

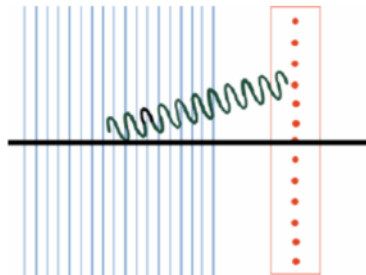
$\omega_p$  - plasma frequency ( $\hbar \omega_p \sim 20 \text{ eV}$ )

Photon energies:  $\hbar \omega \approx \frac{1}{4} \hbar \omega_p \gamma$

$\Rightarrow$  emission probability  $\sim \alpha = \frac{1}{137}$

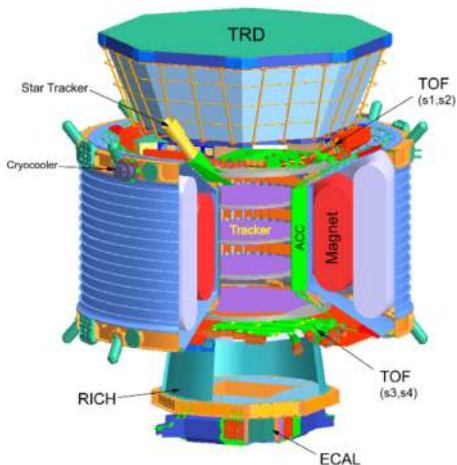
Signal enhanced by using a stack of N foil radiators

Used mainly for  $e^\pm$  identification at large energies ( $\gamma = E/m$ )



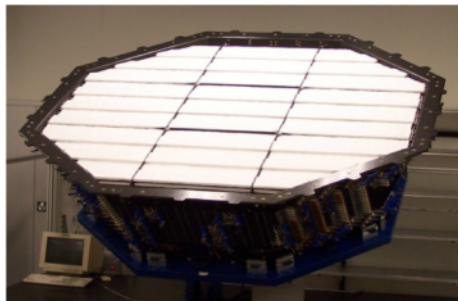
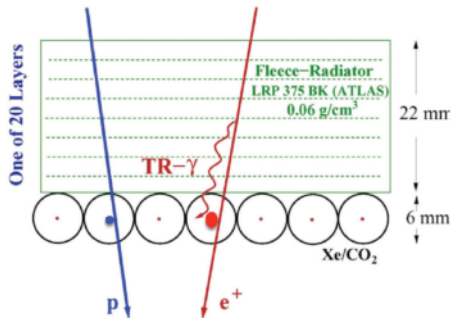
## AMS 02

Particle detector at ISS. Multiple particle identification capabilities



## AMS 02 TRD

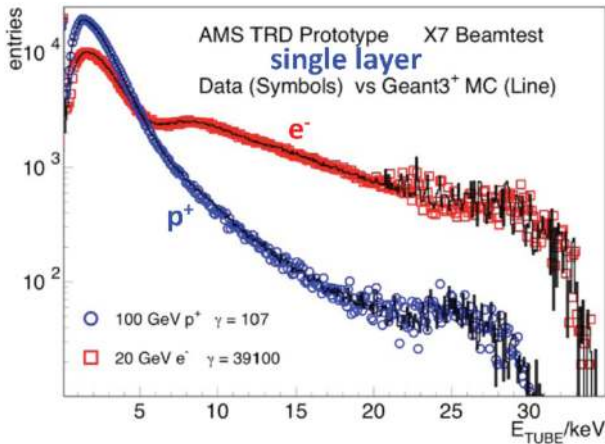
Detector consists of 20 detection layers (radiator + straw tubes)



Focus on  $e^+/p$  discrimination (radiation  $\sim \gamma$ )

## AMS 02

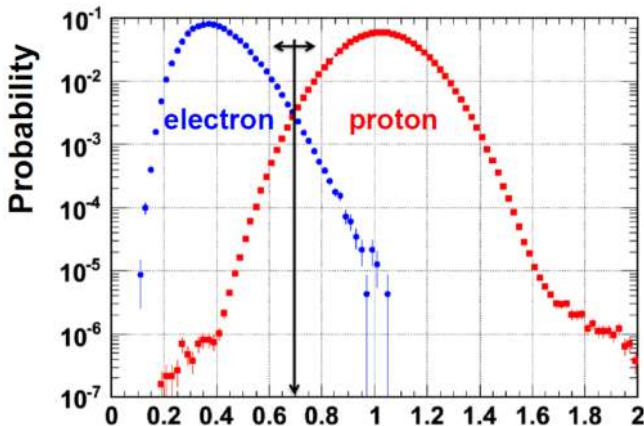
Results of TRD prototype tests at CERN SPS



Electrons radiate much more, but still not enough in single layer...

## AMS 02

Detector performance (20 layers) as measured at ISS

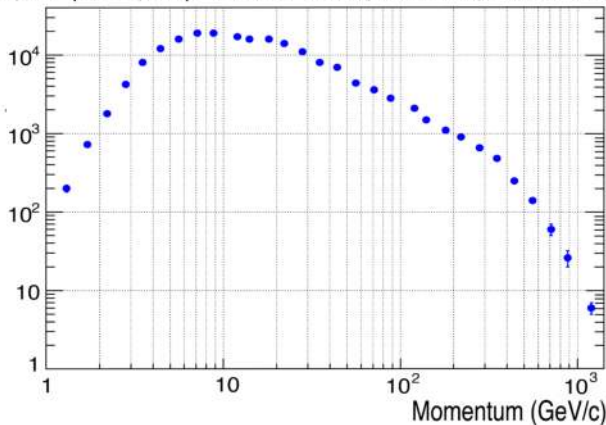


$$\text{TRD estimator} = -\log\left(\frac{P_e}{P_e + P_p}\right)$$

TRD estimator

## AMS 02

ISS data analysis (20 layers): proton background rejection factor

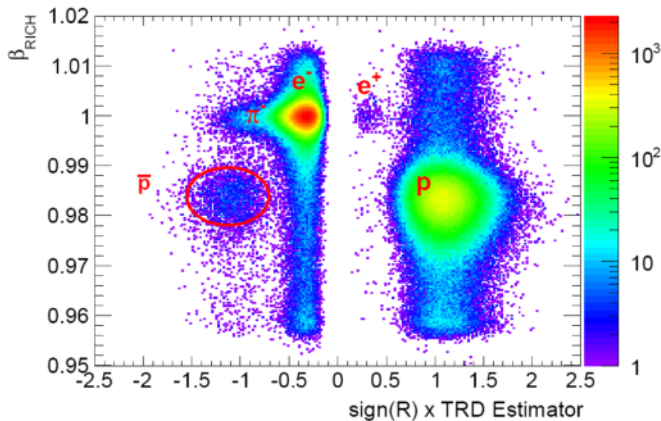


Proton background reduction in  $e^+$  flux measurement by factor  $\sim 10^4$



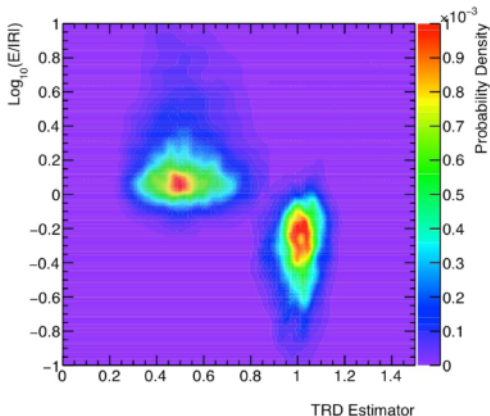
## AMS 02

Correlation of TRD estimator and RICH detector response



## AMS 02

EM calorimeter response (relative to track momentum) vs TRD estimator



For protons, only small fraction of energy is deposited in EM calorimeter

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## Primary cosmic rays

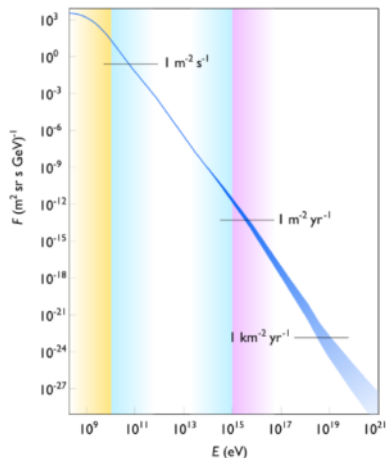
Observed in the cosmic space, outside Earth's atmosphere

Composition:

- protons ( ${}^1\text{H}$ )  $\sim 86\%$
- $\alpha$  particles ( ${}^4\text{He}$ )  $\sim 13\%$
- heavier nuclei  $\sim 1\%$
- neutrons, electrons  $\ll 1\%$

(neglecting neutrinos and gammas)

Same as the “composition of the Universe...”

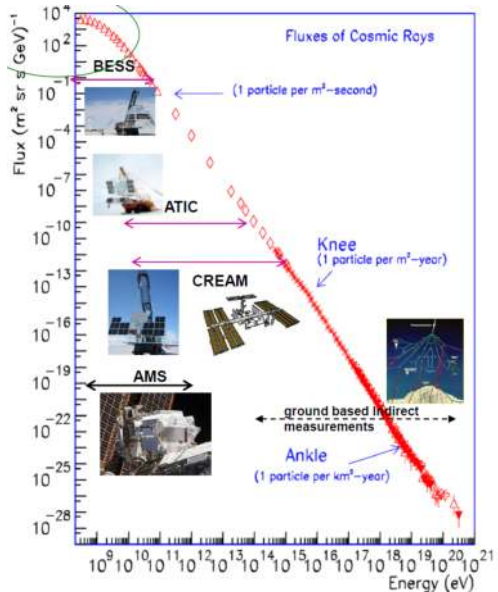


## Primary cosmic rays

Earth's atmosphere shields us from the primary cosmic rays.

They can be measured directly only with stratospheric balloon or **satellite experiments**

For high energies, we can measure cosmic rays indirectly, by observing development of the **Extensive Air Showers** (EAS) in the atmosphere or **Secondary Cosmic Rays** at the surface



## Secondary cosmic rays

Result of the primary cosmic ray interactions in the Earth's atmosphere. Secondary particles, mainly pions and kaons, are copiously produced in these interactions. Both pions and kaons are unstable, produced in their decay chain are muons and electrons.

At the sea level:

- muons  $\mu^\pm \sim 70\%$
- electrons  $e^\pm \sim 25\%$
- protons and pions  $\pi^\pm \sim 3\%$

Average flux:

180 charged particles per  $m^2 \cdot s$



## Extensive air showers

Proton  $10^{15}$  eV:

at ground

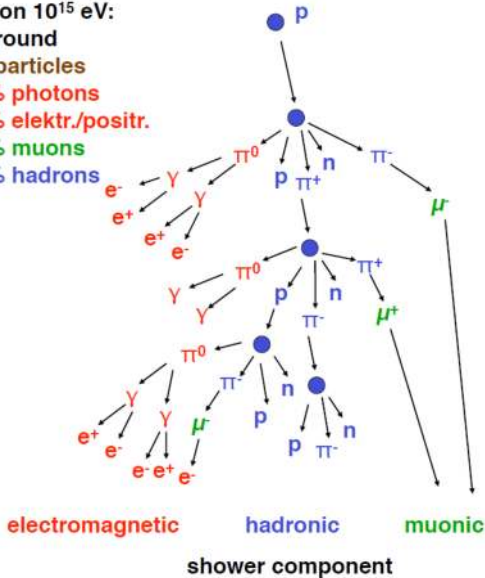
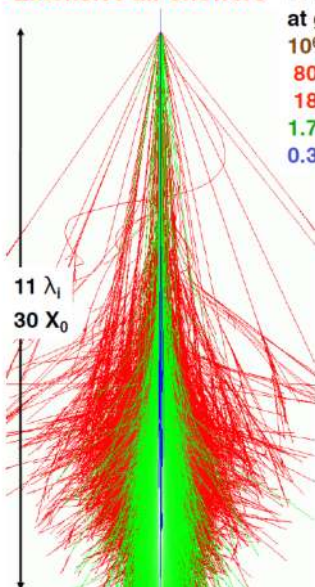
$10^6$  particles

80% photons

18% elektr./positr.

1.7% muons

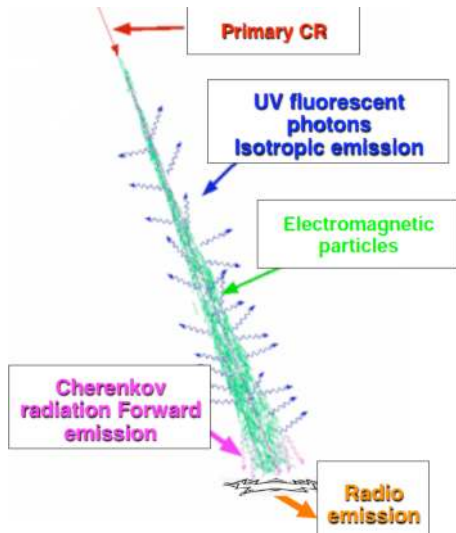
0.3% hadrons



## Extensive Air Showers

Detection channels:

- charged particles reaching Earth's surface, mainly electrons and muons (secondary cosmic rays)
- Cherenkov radiation from electrons in EM core of the cascade
- fluorescent light from excited nitrogen particles in the atmosphere
- radio emission

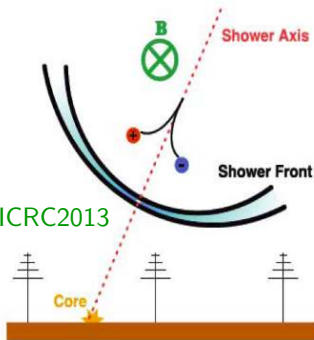




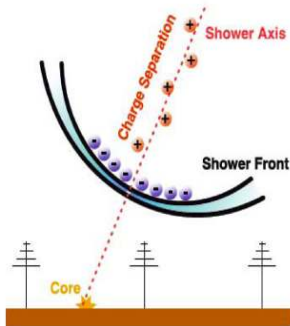
## EAS radio emission

Radio emission is due to the space separation of negative and positive particles in the cascade. Two separation mechanisms:

### Geomagnetic



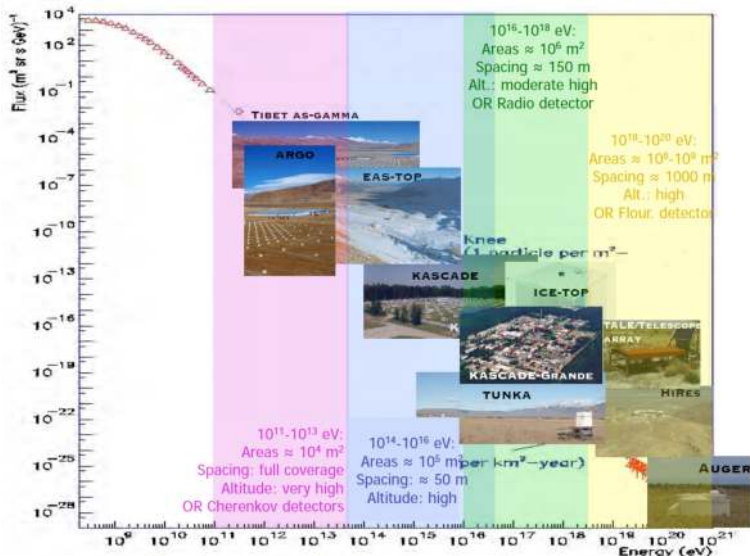
### Askaryan



T. Huege, ICRC2013

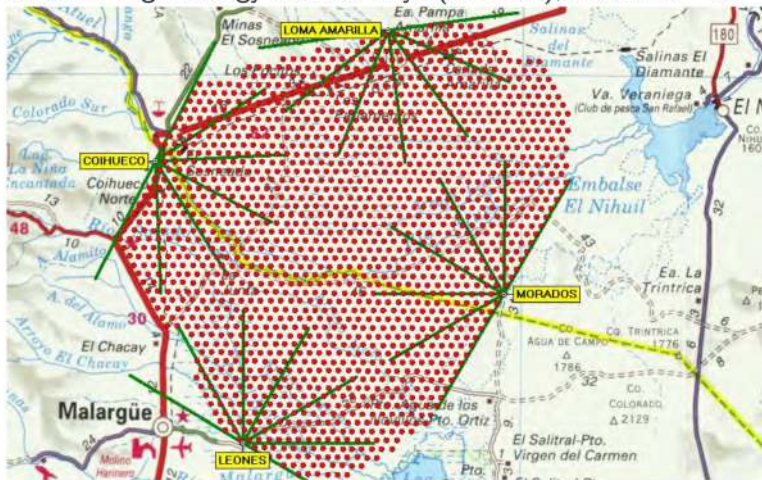
Different polarisation of radio emission for the two components

## Extensive Air Shower experiment



## AUGER

Largest particle physics experiment built so far (3000 km<sup>2</sup>, 1600 units)  
focus on Ultra High Energy Cosmic Rays (UHECR), above 10<sup>18</sup> eV

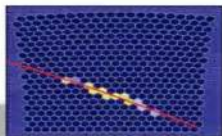


## AUGER

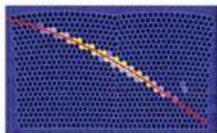
Hybrid EAS measurement: fluorescent telescopes and surface detectors

Event: 1364365

Los Morados

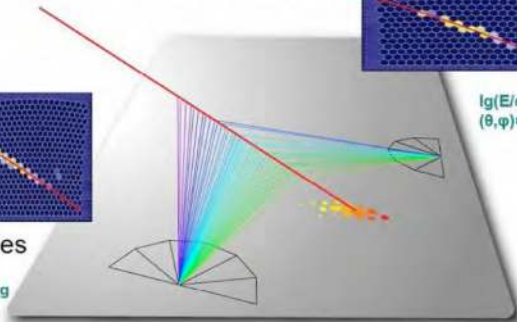


$\lg(E/eV) \sim 19.2$   
 $(\theta, \phi) = (63.7, 148.4)$  deg



Los Leones

$\lg(E/eV) \sim 19.3$   
 $(\theta, \phi) = (63.7, 148.3)$  deg



SD array:  $\lg(E/eV) \sim 19.1$   
 $(\theta, \phi) = (63.3, 148.9)$  deg

⇒ validation of energy reconstruction and energy calibration

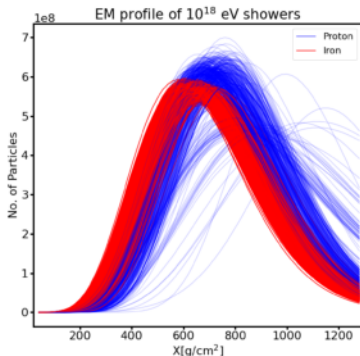
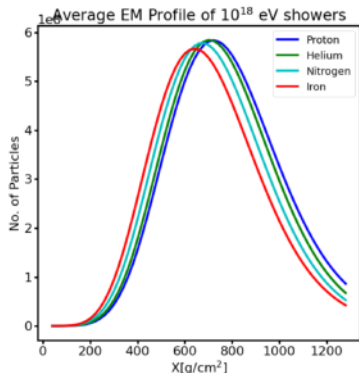
## Primary Cosmic Ray identification simulations by P.Mitra

Shower development in the atmosphere depends on the initial particle mass

⇒ we can discriminate between single protons and heavier nuclei

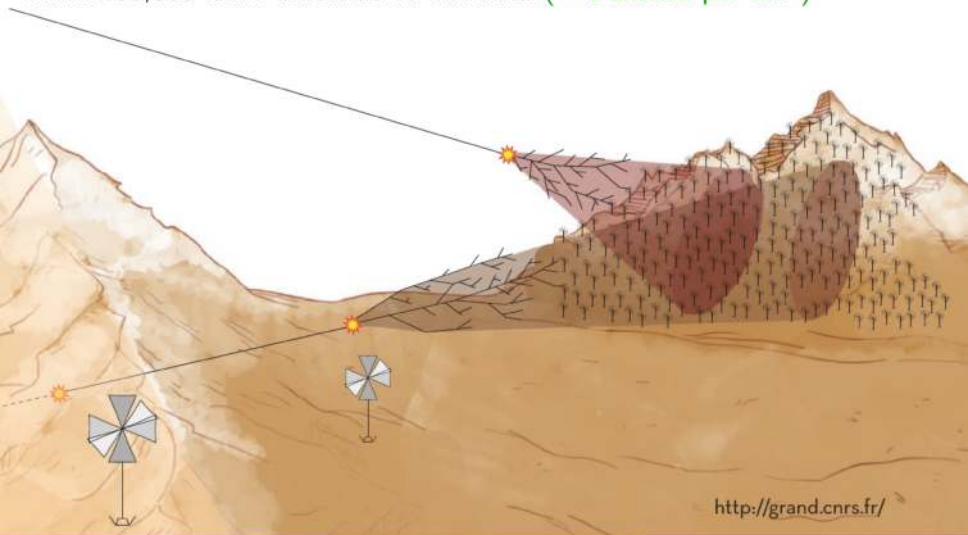
Base on average depth of shower  
maximum for given energy

Based on maximum depth variation  
for given energy



## Giant Radio Array for Neutrino Detection (GRAND)

Aim: 200,000 radio antennas on hill sides ( $\sim 1$  antenna per  $\text{km}^2$ )

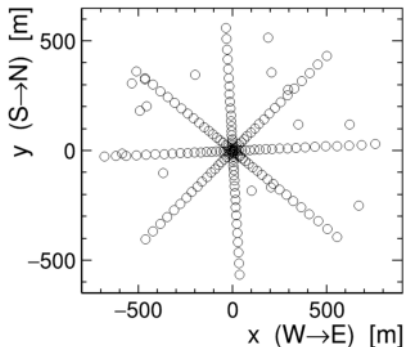


<http://grand.cnrs.fr/>

## Shower reconstruction `06_uhcr_detection.ipynb`

Principle: shower development is measured on Earth surface with very sparse grid of detectors.

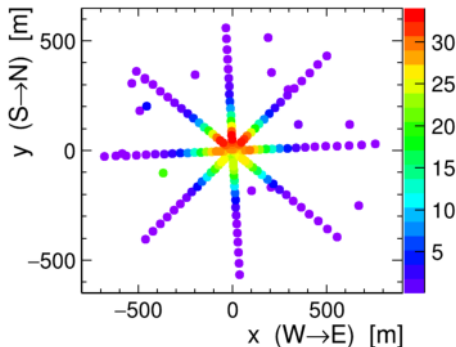
Example configuration used for simulation studies



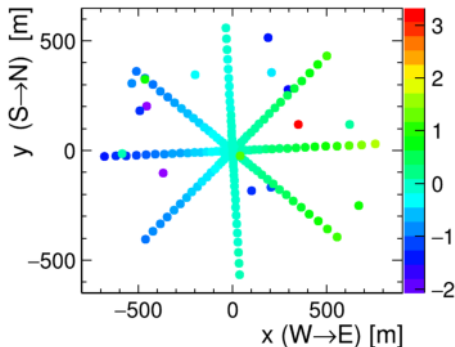
## Shower reconstruction `06_uhcr_detection.ipynb`

Principle: shower development is measured on Earth surface with very sparse grid of detectors.

Radio pulse amplitude



Pulse arrival time



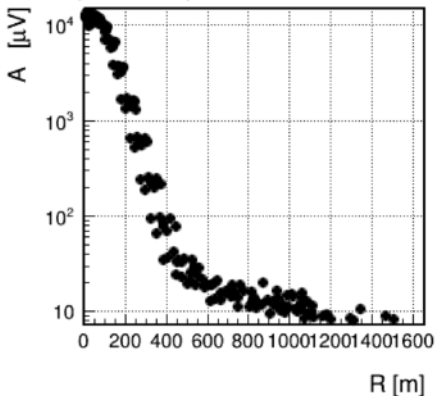


## Shower reconstruction

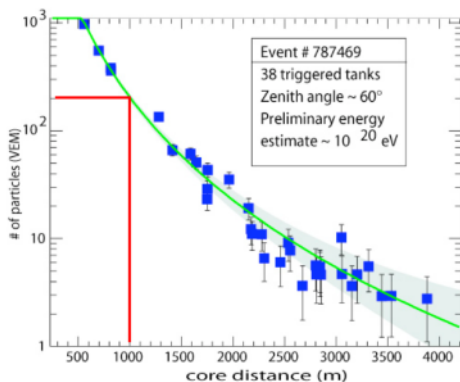
Reconstructed radial profile can be used to extract shower energy

### GRAND simulation

Radio pulse radial profile



### AUGER data

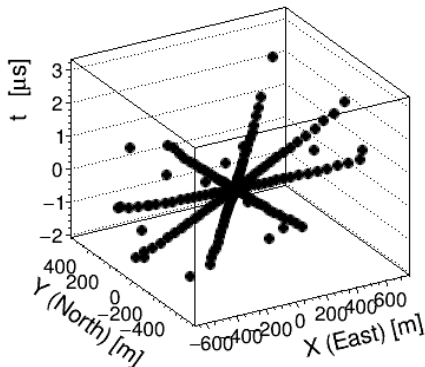


AUGER energy measurement based on S(1000)

## Reconstruction of shower arrival direction

Implement the procedure for fitting the shower arrival direction (zenith and azimuth angles) based on the radio pulse arrival times measured with ground-based antennas.

Radio pulse arrival time



Reconstruct shower direction for the six example data files available with the lecture notebook.

Try to estimate effective time reconstruction precision