

November 11, 2024

MODULE 11 — Assignment

Problem 1: Solve Ginsberg 9.28

The absolute velocity of a particle may be represented by the components v_x , v_y , and v_z relative to the axes of a moving reference system xyz . Suppose that the angular velocity $\bar{\omega}$ of xyz and the velocity \bar{v}_O of the origin of xyz are known as functions of time. Derive the Gibbs-Appell equations of motion relating the quasi-velocities $\dot{\gamma}_1 = v_x$, $\dot{\gamma}_2 = v_y$, and $\dot{\gamma}_3 = v_z$ to the resultant force acting on the particle.

Where:

$$\bar{\omega} = \left\langle \begin{array}{c} \omega_x \\ \omega_y \\ \omega_z \end{array} \right\rangle$$

Given that:

$$\bar{v} = \left\langle \begin{array}{c} v_x \\ v_y \\ v_z \end{array} \right\rangle = \left\langle \begin{array}{c} \dot{\gamma}_1 \\ \dot{\gamma}_2 \\ \dot{\gamma}_3 \end{array} \right\rangle$$

Solving for acceleration:

$$\begin{aligned} \bar{a} &= \frac{\partial \bar{v}}{\partial t} + \bar{\omega} \times \bar{v} \\ &= \frac{\partial \dot{\gamma}}{\partial t} + \bar{\omega} \times \dot{\gamma} \\ &= \left\langle \begin{array}{c} \ddot{\gamma}_1 - \dot{\gamma}_2 \omega_z + \dot{\gamma}_3 \omega_y \\ \ddot{\gamma}_2 + \dot{\gamma}_1 \omega_z - \dot{\gamma}_3 \omega_x \\ \ddot{\gamma}_3 - \dot{\gamma}_1 \omega_y + \dot{\gamma}_2 \omega_x \end{array} \right\rangle \end{aligned}$$

Given that the Gibbs-Appell function for a system of particles is:

$$S = \sum_p \frac{1}{2} m \bar{a}_p \cdot \bar{a}_p$$

For this single particle case:

$$\begin{aligned} S &= \frac{1}{2} m (\bar{a} \cdot \bar{a}) \\ &= \frac{1}{2} m \left[(\ddot{\gamma}_3 - \dot{\gamma}_1 \omega_y + \dot{\gamma}_2 \omega_x)^2 + (\ddot{\gamma}_2 + \dot{\gamma}_1 \omega_z - \dot{\gamma}_3 \omega_x)^2 + (\ddot{\gamma}_1 - \dot{\gamma}_2 \omega_z + \dot{\gamma}_3 \omega_y)^2 \right] \end{aligned}$$

Where the equations of motion are calculated as:

$$\frac{\partial S}{\partial \ddot{\gamma}_j} = \Gamma_j = \Gamma_1$$

The virtual work associated with the forces applied to the particle is:

$$\delta W = \sum \bar{F} \cdot \delta \bar{r} = \sum_{j=1}^K \Gamma_j \delta \gamma_j = \sum \bar{F} \cdot \left\langle \begin{matrix} \delta \gamma_1 \\ \delta \gamma_2 \\ \delta \gamma_3 \end{matrix} \right\rangle$$

The equation of motion is solved for as:

$$\frac{\partial S}{\partial \ddot{\gamma}} = m \left\langle \begin{matrix} \ddot{\gamma}_1 - \dot{\gamma}_2 \omega_z + \dot{\gamma}_3 \omega_y \\ \ddot{\gamma}_2 + \dot{\gamma}_1 \omega_z - \dot{\gamma}_3 \omega_x \\ \ddot{\gamma}_3 - \dot{\gamma}_1 \omega_y + \dot{\gamma}_2 \omega_x \end{matrix} \right\rangle = \Gamma_1 = \left\langle \begin{matrix} \sum F_x \\ \sum F_y \\ \sum F_z \end{matrix} \right\rangle$$

Or:

$$\begin{aligned} m (\ddot{\gamma}_1 - \dot{\gamma}_2 \omega_z + \dot{\gamma}_3 \omega_y) &= \sum F_x \\ m (\ddot{\gamma}_2 + \dot{\gamma}_1 \omega_z - \dot{\gamma}_3 \omega_x) &= \sum F_y \\ m (\ddot{\gamma}_3 - \dot{\gamma}_1 \omega_y + \dot{\gamma}_2 \omega_x) &= \sum F_z \end{aligned}$$

→ Answer

Where:

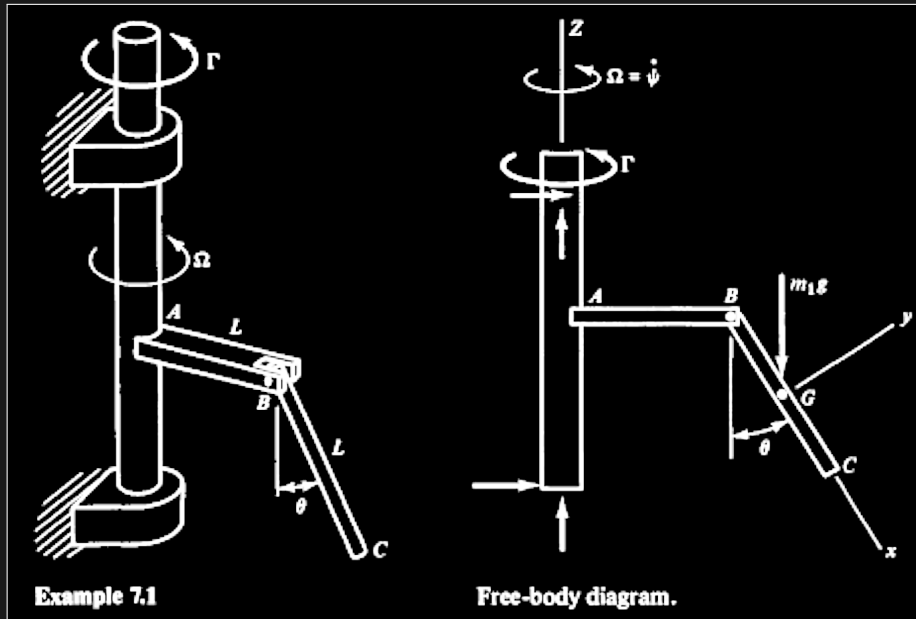
$$\dot{\gamma} = \left\langle \begin{matrix} \dot{\gamma}_1 \\ \dot{\gamma}_2 \\ \dot{\gamma}_3 \end{matrix} \right\rangle = v = \left\langle \begin{matrix} v_x \\ v_y \\ v_z \end{matrix} \right\rangle = v_0 + \bar{\omega} \times \bar{r} = \left\langle \begin{matrix} v_{0x} - \omega_z y + \omega_y z \\ v_{0y} + \omega_z x - \omega_x z \\ v_{0z} - \omega_y x + \omega_x y \end{matrix} \right\rangle$$

Where:

Problem 2:

Use the Gibbs-Appell approach to find the equations of motion for this problem.

A torque Γ applied to the vertical shaft of the T-bar causes the rotation rate Ω about the vertical axis to increase in proportion to the angle θ by which bar BC swings outward, that is, $\Omega = c\theta$. The mass of bar BC is m_1 and the moment of inertia of the T-bar about its axis of rotation is I_2 . Determine the equations of motion for the system, and for the torque Γ .



$$\left(\begin{array}{c} \frac{L \cos(\psi) \cos(\theta)}{2} \frac{\partial}{\partial t} \dot{\gamma}_1 - \frac{L \cos(\psi) \sin(\theta)}{2} \dot{\gamma}_1^2 - \frac{L c^2 \cos(\psi) \theta^2 (\sin(\theta)+2)}{2} - \frac{L c \sin(\psi) \dot{\gamma}_1 (\sin(\theta)+2)}{2} - L c \cos(\theta) \sin(\psi) \dot{\gamma}_1 \theta \\ \frac{L \cos(\theta) \sin(\psi)}{2} \frac{\partial}{\partial t} \dot{\gamma}_1 - \frac{L \sin(\psi) \sin(\theta)}{2} \dot{\gamma}_1^2 - \frac{L c^2 \sin(\psi) \theta^2 (\sin(\theta)+2)}{2} + \frac{L c \cos(\psi) \dot{\gamma}_1 (\sin(\theta)+2)}{2} + L c \cos(\psi) \cos(\theta) \dot{\gamma}_1 \theta \\ \frac{L \sin(\theta)}{2} \frac{\partial}{\partial t} \dot{\gamma}_1 + \frac{L \cos(\theta)}{2} \dot{\gamma}_1^2 \end{array} \right)$$

Submitted by Austin Barrilleaux on November 11, 2024.