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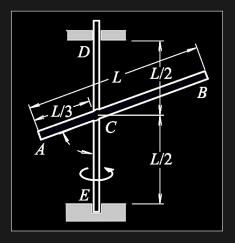
October 9, 2024

### MODULE 6 — Assignment

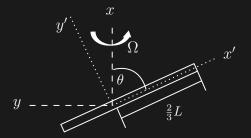
## EXERCISE 5.26

Thin bar ACB is welded to a shaft that rotates at the constant angular speed  $\Omega$ , so the angle  $\theta$  between the bar and the shaft is constant.

- (a) Derive expressions for the angular momentum  $\bar{H}_C$  and the kinetic energy of the bar. Draw a sketch of  $\bar{H}_C$ .
- (b) Based on an analysis of the manner in which  $\bar{H}_C$  in Part (a) rotates, derive an expression for  $\frac{\partial}{\partial t}\bar{H}_C$ .
- (c) Use Eq. (5.3.4) to evaluate  $\frac{\partial}{\partial t}\bar{H}_C$ , and compare it with the result of Part (b).



The following sketch shows the two frames of the system:



The rotation matrix that converts the  $\{xyz\}$  frame to the  $\{x'y'z'\}$  frame is:

$$R = \begin{bmatrix} \cos(-\theta) & \sin(-\theta) & 0 \\ -\sin(-\theta) & \cos(-\theta) & 0 \\ 0 & 0 & 1 \end{bmatrix}$$

Constructing the angular velocity vector:

$$\bar{\omega} = \Omega i$$

In the body frame is:

$$ar{\omega} = R \; \Omega \; i = \left[ egin{array}{ccc} \Omega \; \cos \left( heta 
ight) \; i' \ \Omega \; \sin \left( heta 
ight) \; j' \ 0 \; k' \end{array} 
ight]$$

From textbook Appendix, the centroidal inertia mass properties of the shaft are:

$$I_{xx} = 0$$

$$I_{yy} = \frac{1}{12}m L^2$$

$$I_{zz} = \frac{1}{12}m L^2$$

This expressed as the inertia tensor is:

$$I_{x'y'z'} = \left[ egin{array}{ccc} 0 & 0 & 0 \ 0 & rac{1}{12}m\,L^2 & 0 \ 0 & 0 & rac{1}{12}m\,L^2 \end{array} 
ight]$$

The distance from the body frame is:

$$d_{x'y'z'} = \left[ \begin{array}{ccc} \frac{1}{6} L & 0 & 0 \end{array} \right]$$

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Using the parallel axis theorem to get the parallel axis transformation of inertia matrix relative to the center of the frame of reference in question:

$$I_{pat} = m \left[ egin{array}{ccc} 0 & 0 & 0 \ 0 & rac{1}{36} L^2 & 0 \ 0 & 0 & rac{1}{36} L^2 \end{array} 
ight]$$

This makes the inertia tensor at point C:

$$I_C = I_{x'y'z'} + I_{pat} = m \begin{bmatrix} 0 & 0 & 0 \\ 0 & \frac{1}{9}L^2 & 0 \\ 0 & 0 & \frac{1}{9}L^2 \end{bmatrix}$$

From this we can compute the angular momentum as:

$$ar{H}_C = I_C \; ar{\omega} = rac{1}{9} m \, L^2 \, \Omega, \sin \left( heta 
ight) \; j'$$

 $\longrightarrow \mathcal{A}$ nswer

Since the bar rotates about the x axis, the rotation component in the terminal frame that is orthogonal to the j'-axis is  $\Omega \cos(\theta)$  i', therefore:

$$\dot{ar{H}}_C = rac{1}{9} m L^2 \Omega^2 \cos{( heta)} \sin{( heta)} k'$$

 $\longrightarrow \mathcal{A}$ nswer

If we use Eq. (5.3.4) to evaluate this, we get the same answer:

$$\dot{\bar{H}}_C = \frac{\partial}{\partial t} \ddot{\bar{H}}_C + \bar{\omega} \times \bar{H}_C$$

$$egin{aligned} \dot{ar{H}}_C &= \left[egin{array}{ccc} 0 \ i' \ 0 \ j' \ rac{1}{9} m \ L^2 \ \Omega^2 \cos \left( heta 
ight) \sin \left( heta 
ight) \ k' \ \end{array}
ight] \ &= rac{1}{9} m \ L^2 \ \Omega^2 \cos \left( heta 
ight) \sin \left( heta 
ight) \ k' \end{aligned}$$

 $o \mathcal{A}$ nswer

#### Matlab Code Used

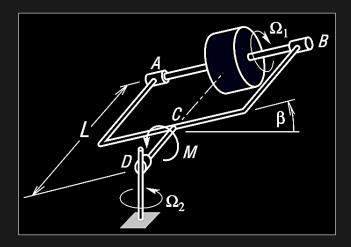
The following MATLAB code was used to solve this problem:

```
% EXERCISE 5.26
% Thin bar ACB is welded to a shaft that rotates at the constant
% angular speed, so the angle theta between the bar and the
% shaft is constant.
% (a) Derive expressions for the angular momentum H-C and the
% kinetic energy of the bar. Draw a sketch of H-C.
% (b) Based on an analysis of the manner in which H-C in Part (a)
% rotates, derive an expression for dH-C/dt.
% (c) Use Eq. (5.3.4) to evaluate dH -C/dt, and compare it
with the result of Part (b).

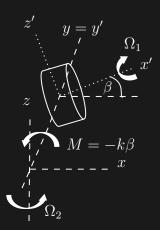
syms m L t
syms theta real
syms Omega real
R = Dynamics.Inertial2Body.zRot(-theta);
omega_bar = R*Omega*[1;0;0];
I = diag([0,(1/12)*m*L^2,(1/12)*m*L^2]);
I_pa = Dynamics.Inertia.parallelAxisTransform(m,[(2*L/3-L/2),0,0]);
H_C = (I+I_pa)*omega_bar;
H_C = diff(H_C,t) + cross(omega_bar,H_C);
```

# EXERCISE 6.8

The torque M acting on the gimbal of the gyroscopic turn indicator is exerted by a torsional spring, so  $M = -k\beta$ . The precession rate  $\Omega_2$  is a specified function of time, and the spin rate  $\Omega_1$  is held constant by a servomotor. Let  $I_1$  denote the moment of inertia of the flywheel about axis AB, and let  $I_2$  be the centroidal moment of inertia perpendicular to axis AB. Derive the differential equation of motion for  $\beta$ .



We will define the following frames for the system:



Construct the angular velocity vector  $\bar{\omega}$  of xyz by vectorially adding the simple rotation rates according to:

$$\bar{\omega} = \omega_1 \bar{e}_1 + \omega_2 \bar{e}_2 + \omega_3 \bar{e}_3$$

This gives us:

$$\bar{\omega} = \Omega_2 \ k + \dot{\beta}(-j') + \Omega_1(-i')$$

We can equate k, given the following rotation as:

$$R(-\beta) = \begin{bmatrix} \cos(\beta) & 0 & \sin(\beta) \\ 0 & 1 & 0 \\ -\sin(\beta) & 0 & \cos(\beta) \end{bmatrix}$$

$$k = R(-\beta) \begin{bmatrix} 0 \\ 0 \\ 1 \end{bmatrix} = \begin{bmatrix} \sin(\beta) & i' \\ 0 & j \\ \cos(\beta) & k' \end{bmatrix}$$

This allows us to express the angular velocity vector as:

$$\bar{\omega} = \Omega_2 \left( \sin(\beta) \ i' + \cos(\beta) \ k' \right) - \dot{\beta} \ j' - \Omega_1 \ i'$$

Or:

$$\bar{\omega} = \begin{bmatrix} \left\{ \sin \left( \beta \right) \, \Omega_2 - \Omega_1 \right\} \, i' \\ \left\{ -\dot{\beta} \right\} \, j' \\ \left\{ \cos \left( \beta \right) \, \Omega_2 \right\} \, k' \end{bmatrix}$$

Next solve for angular acceleration:

$$\bar{\alpha} = \sum_{n} \left( \dot{\omega}_n \bar{e}_n + \bar{\Omega}_n \times \omega_n \bar{e}_n \right)$$

Where:

$$\begin{split} & \omega_{n=1} = -\Omega_1 \ i', & \dot{\omega}_{n=1} = 0, & \Omega_{n=1} = \Omega_2 R \ k - \dot{\beta} \ j' \\ & \omega_{n=2} = -\dot{\beta} \ j', & \dot{\omega}_{n=2} = -\ddot{\beta} \ j', & \Omega_{n=2} = \Omega_2 R \ k \\ & \omega_{n=3} = \Omega_2 R \ k, & \dot{\omega}_{n=3} = \dot{\Omega}_2 R \ k, & \Omega_{n=3} = 0 \end{split}$$

Substituting these into  $\bar{\alpha}$ :

$$\bar{\alpha} = \begin{bmatrix} \left\{ \dot{\Omega}_2 \sin(\beta) +, \Omega_2 \dot{\beta} \cos(\beta) \right\} i' \\ \left\{ -\ddot{\beta} - \Omega_1 \Omega_2 \cos(\beta) \right\} j' \\ \left\{ \dot{\Omega}_2 \cos(\beta) - \Omega_2 \dot{\beta} \sin(\beta) - \Omega_1 \dot{\beta} \right\} k' \end{bmatrix}$$

Since the torque M is acting in the j' direction, we can refer to equation 6.1.6 in the textbook:

$$\sum \bar{M}_{cg} j' = k\beta = I_{yy}\alpha_y + (I_{zz} - I_{xx}) \omega_x \omega_z$$

Where the rotation of the frame, which is not rigidly fixed to the flywheel, is:

$$\omega = \Omega_2 \ k + \dot{\beta}(-j') = \begin{bmatrix} \Omega_2 \sin(\beta) \ i' \\ -\dot{\beta} \ j' \\ \Omega_2 \cos(\beta) \ k' \end{bmatrix}$$

And:

$$I_{xx} = I_1$$
$$I_{yy} = I_{zz} = I_2$$

This evaluates to:

$$egin{aligned} keta &= \cos{(eta)} \, \sin{(eta)} \, \, \Omega_2^{\,2} \, \left(I_1 - I_2
ight) - I_2\ddot{eta} - I_2\Omega_1 \, \cos{(eta)} \, \, \Omega_2 \ &\longrightarrow \mathcal{A}$$
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#### Matlab Code Used

The following MATLAB code was used to solve this problem:

```
%% EXERCISE 6.8
% The torque M acting on the gimbal of the gyroscopic turn indicator is
% exerted by a torsional spring, so M = -k beta. The precession rate
% Omega 2 is a specified function of time, and the spin rate Omega 1 is
% held constant by a servomotor. Let I1 denote the moment of inertia of
% the flywheel about axis AB, and let I2 be the centroidal moment of
% inertia perpendicular to axis AB. Derive the differential equation
```

```
clc,clear
syms k t beta(t)
syms Omega_1 Omega_2(t)
syms I_1 I_2
R = yRot(-beta);
omega_bar = Omega_2*R*[0;0;1] + diff(beta,t)*[0;-1;0]...
                + Omega_1*[-1;0;0];
Omega_N = Omega_2(t) *R(t) *[0;0;1];
alpha_bar = cross(omega_bar,Omega_1*[-1;0;0]) +...
            diff(Omega_2,t)*R*[0;0;1] +...
            simplify(cross(Omega_N,Omega_2*R*[0;0;1])) +...
            cross(Omega_N, diff(beta,t) *[0;-1;0]);
system = ([0;1;0]*k*beta(t) == I_2*alpha_bar(t)-...
    (I_2-I_1) *Omega_N(1) *Omega_N(3));
system(2)
R = [\cos(ang) \ 0 - \sin(ang) ; \% \text{ Neg body 2 in}
sin(ang) 0 cos(ang) ];
```

Submitted by Austin Barrilleaux on October 9, 2024.