CALCULATION OF LLL BERRY CURVATURE IN SPACE OF METRICS

D.B.

We calculate the adiabatic curvature resulting from varying the LLL wavefunctions with respect to the torus geometry, as parametrized by flat metrics

$$g(V,\tau) = \frac{V}{\tau_2} \begin{pmatrix} 1 & \tau_1 \\ \tau_1 & |\tau|^2 \end{pmatrix}$$

Starting from the Landau level wavefunctions on the torus (ASZ 14), we can calculate the Berry curvature using (ASZ 6),

$$F_{ij} = \operatorname{Im} \sum_{\ell=1}^{N} \left\langle \partial_{i} \phi_{\ell} \middle| \partial_{j} \phi_{\ell} \right\rangle,$$

where ℓ indexes the single particle wavefunctions, and N is the number of particles. We will focus on obtaining this curvature for a particular index ℓ . We will find that F does not depend on ℓ , so the result for multiple particles filling N levels follows trivially.

The normalized lowest Landau level wavefunction for a given metric $g(V, \tau)$ is

$$\phi_{\ell}(\mathbf{x}) = \frac{(2\tau_2 B)^{1/4}}{\sqrt{V}} \sum_{n=-\infty}^{\infty} e^{i\pi\tau B(\widetilde{y}+n)^2} e^{-2i\pi \left(\varphi_2 + B/2\right)(\widetilde{y}+n\right)} e^{2i\pi (nB+\ell)x}$$

with $\tilde{y} = y + y_{\ell} = y + \ell/B + \varphi_1/B + 1/2$.

To compute components of F involving τ_1 , τ_2 , we note

$$\phi_{\ell}(\mathbf{x}) = \frac{(2\tau_2 B)^{1/4}}{V^{1/2}} \Phi_{\ell}(\mathbf{x})$$

with Φ holomorphic in τ . Therefore

$$\partial_{\tau_1} \phi_{\ell}(\mathbf{x}) = \frac{(2\tau_2 B)^{1/4}}{V^{1/2}} \partial_{\tau} \Phi_{\ell}(\mathbf{x})$$

and

$$\partial_{\tau_2} \phi_{\ell}(x, y) = \frac{(2\tau_2 B)^{1/4}}{V^{1/2}} \left[\frac{1}{4\tau_2} \Phi_{\ell}(\mathbf{x}) + i \partial_{\tau} \Phi_{\ell}(\mathbf{x}) \right].$$

Explicitly,

$$\partial_{\tau}\Phi_{\ell}(\mathbf{x}) = \sum_{n=-\infty}^{\infty} \left(i\pi B(\widetilde{y}+n)^2 \right) e^{i\pi\tau B(\widetilde{y}+n)^2} e^{-2i\pi \left(\varphi_2 + B/2 \right) (\widetilde{y}+n)} e^{2i\pi (nB+\ell)x}.$$

For F_{V,τ_2}^{ℓ} ,

$$\begin{split} F_{V,\tau_2}^\ell &= -F_{\tau_2,V}^\ell = \operatorname{Im} \left[\left\langle \partial_V \phi_\ell \middle| \partial_{\tau_2} \phi_\ell \right\rangle \right] \\ &= \operatorname{Im} \left[\int\limits_Q \sqrt{\det g} \, \mathrm{d}x \, \mathrm{d}y \, \partial_V \phi_\ell^*(\mathbf{x}) \partial_{\tau_2} \phi_\ell(\mathbf{x}) \right] \\ &= \operatorname{Im} \left[-\frac{1}{2V} \int\limits_Q \mathrm{d}x \, \mathrm{d}y \, \phi_\ell^*(\mathbf{x}) \left(\frac{1}{4\tau_2} \phi_\ell(\mathbf{x}) + i \sqrt{2\tau_2 B} \, \partial_\tau \Phi_\ell(\mathbf{x}) \right) \right], \end{split}$$

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where the integral is over $Q = [0, 1] \times [0, 1]$. The first term in the integral is clearly real, so

$$F_{V,\tau_2}^{\ell} = \frac{\sqrt{2\tau_2 B}}{2V} \operatorname{Im} \left[i \int_{Q} dx \, dy \, \Phi_{\ell}^{*}(\mathbf{x}) \partial_{\tau} \Phi_{\ell}(\mathbf{x}) \right]$$
$$\equiv \frac{\sqrt{2\tau_2 B}}{2V} \operatorname{Im} \left[iI_1 \right].$$

The integral I_1 is

$$\begin{split} I_{1} &= \int\limits_{Q} \mathrm{d}x \, \mathrm{d}y \, \Phi_{\ell}^{*}(\mathbf{x}) \partial_{\tau} \Phi_{\ell}(\mathbf{x}) = \int\limits_{Q} \mathrm{d}x \, \mathrm{d}y \, \left[\sum_{m=-\infty}^{\infty} e^{-i\pi\tau^{*}B(\widetilde{y}+m)^{2}} e^{2i\pi\left(\phi_{2}+B/2\right)(\widetilde{y}+m)} e^{-2i\pi(mB+\ell)x} \right] \\ &\times \left[\sum_{n=-\infty}^{\infty} \left(i\pi B(\widetilde{y}+n)^{2} \right) e^{i\pi\tau B(\widetilde{y}+n)^{2}} e^{-2i\pi\left(\phi_{2}+B/2\right)(\widetilde{y}+n)} e^{2i\pi(nB+\ell)x} \right] \\ &= \int\limits_{Q} \mathrm{d}x \, \mathrm{d}y \, \sum_{m,n} i\pi B(\widetilde{y}+n)^{2} e^{i\pi B\left(\tau(\widetilde{y}+n)^{2}-\tau^{*}(\widetilde{y}+m)^{2}\right)} e^{2i\pi(\phi_{2}+B/2)(m-n)} e^{2i\pi Bx(n-m)} \end{split}$$

Performing the *x* integration, we have, since *B* is an integer

$$\int_{0}^{1} \mathrm{d}x \, e^{2i\pi Bx(n-m)} = \delta_{n,m},$$

and so

$$I_{1} = \int_{0}^{1} dy \sum_{n} i\pi B(\tilde{y} + n)^{2} e^{i\pi B(\tau - \tau^{*})(\tilde{y} + n)^{2}}$$

$$= i\pi B \sum_{n = -\infty}^{\infty} \int_{0}^{1} dy (\tilde{y} + n)^{2} e^{-2\pi B \tau_{2}(\tilde{y} + n)^{2}}$$

$$= i\pi B \sum_{n = -\infty}^{\infty} \int_{y_{\ell} + n}^{y_{\ell} + n + 1} du u^{2} e^{-2\pi B \tau_{2} u^{2}}$$

$$= i\pi B \int_{-\infty}^{\infty} du u^{2} e^{-2\pi B \tau_{2} u^{2}}$$

$$= i\pi B \left(\frac{1}{4\sqrt{2}\pi (B\tau_{2})^{3/2}}\right)$$

$$= \frac{i}{4\tau_{2}\sqrt{2\tau_{2}B}}$$

Since I_1 is pure imaginary,

$$F_{V,\tau_2}^{\ell} = \frac{\sqrt{2\tau_2 B}}{2V} \text{Im} [iI_1] = 0.$$

For F_{V,τ_1}^{ℓ} ,

$$F_{V,\tau_1}^{\ell} = \operatorname{Im} \left[-\frac{\sqrt{2\tau_2 B}}{2V} \int_{Q} dx \, dy \, \Phi_{\ell}^*(\mathbf{x}) \, \partial_{\tau} \Phi_{\ell}(\mathbf{x}) \right]$$
$$= -\frac{\sqrt{2\tau_2 B}}{2V} \operatorname{Im} [I_1]$$
$$= -\frac{1}{8\tau_2 V} \neq 0 \, [???]$$

For F_{τ_1,τ_2}^{ℓ} ,

$$\begin{split} F_{\tau_{1},\tau_{2}}^{\ell} &= -F_{\tau_{2},\tau_{1}}^{\ell} = \operatorname{Im}\left[\left\langle\partial_{\tau_{1}}\phi_{\ell}\middle|\partial_{\tau_{2}}\phi_{\ell}\right\rangle\right] \\ &= \operatorname{Im}\left[\int_{Q} \sqrt{\det g} \; \mathrm{d}x \; \mathrm{d}y \; \partial_{\tau_{1}}\phi_{\ell}^{*}(\mathbf{x})\partial_{\tau_{2}}\phi_{\ell}(\mathbf{x})\right] \\ &= \operatorname{Im}\left[\sqrt{2\tau_{2}B}\int_{Q} \mathrm{d}x \; \mathrm{d}y \; \partial_{\tau}\Phi_{\ell}^{*}(\mathbf{x}) \left(\frac{1}{4\tau_{2}}\Phi_{\ell}(\mathbf{x}) + i\partial_{\tau}\Phi_{\ell}(\mathbf{x})\right)\right] \\ &= \operatorname{Im}\left[\sqrt{2\tau_{2}B}\left(\frac{1}{4\tau_{2}}\int_{Q} \mathrm{d}x \; \mathrm{d}y \; \partial_{\tau}\Phi_{\ell}^{*}(\mathbf{x})\Phi_{\ell}(\mathbf{x}) + i\int_{Q} \mathrm{d}x \; \mathrm{d}y \; \partial_{\tau}\Phi_{\ell}^{*}(\mathbf{x})\partial_{\tau}\Phi_{\ell}(\mathbf{x})\right)\right] \\ &= \operatorname{Im}\left[\sqrt{2\tau_{2}B}\left(\frac{1}{4\tau_{2}}I_{1}^{*} + i\int_{Q} \mathrm{d}x \; \mathrm{d}y \; \partial_{\tau}\Phi_{\ell}^{*}(\mathbf{x})\partial_{\tau}\Phi_{\ell}(\mathbf{x})\right)\right] \\ &= \sqrt{2\tau_{2}B}\operatorname{Im}\left[\frac{1}{4\tau_{2}}I_{1}^{*} + iI_{2}\right] \end{split}$$

The integral I_2 is

$$\begin{split} I_2 &= \int\limits_{Q} \mathrm{d}x \, \mathrm{d}y \, \partial_{\tau} \Phi_{\ell}^*(\mathbf{x}) \partial_{\tau} \Phi_{\ell}(\mathbf{x}) \\ &= \int\limits_{Q} \mathrm{d}x \, \mathrm{d}y \, \sum_{m,n} \left(-i\pi B (\widetilde{y} + m)^2 \right) \left(i\pi B (\widetilde{y} + n)^2 \right) e^{i\pi B \left(\tau (\widetilde{y} + n)^2 - \tau^* (\widetilde{y} + m)^2 \right)} e^{2i\pi (\phi_2 + B/2)(m - n)} e^{2i\pi B x (n - m)} \\ &= \pi^2 B^2 \sum_{n = -\infty}^{\infty} \int\limits_{0}^{1} \mathrm{d}y \, (\widetilde{y} + n)^4 e^{-2\pi B \tau_2 (\widetilde{y} + n)^2} \\ &= \pi^2 B^2 \int\limits_{-\infty}^{\infty} \mathrm{d}u \, u^4 e^{-2\pi B \tau_2 (u)^2} \\ &= \pi^2 B^2 \frac{3}{16\sqrt{2}\pi^2 (B\tau_2)^{5/2}} \\ &= \frac{1}{\sqrt{2\tau_2 B}} \frac{3}{16\tau_2^2}. \end{split}$$

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So that

$$\begin{split} F_{\tau_1,\tau_2} &= \sqrt{2\tau_2 B} \, \mathrm{Im} \, \left[\frac{1}{4\tau_2} \left(\frac{-i}{4\tau_2 \sqrt{2\tau_2 B}} \right) + i \frac{1}{\sqrt{2\tau_2 B}} \frac{3}{16\tau_2^2} \right] \\ &= \mathrm{Im} \, \left[i \left(-\frac{1}{16\tau_2} + \frac{3}{16\tau_2^2} \right) \right] \\ &= \frac{1}{8\tau_2^2} \end{split}$$