INERTIAL FRAME DRAGGING

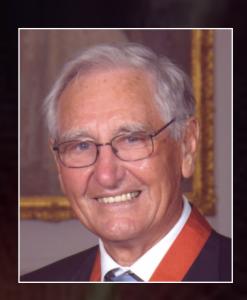
NEAR KERR BLACK-HOLE

KERR BLACK-HOLE

• Roy Kerr (1934 -)

Discovered exact solutions of Einstein's equation describing a spinning (uncharged)

Black-Hole, which is known as Kerr Solution (1960).



- Any (isolated) Black-Hole can be described by just three quantities:
- 1. Mass
- 2. Spin
- 3. Electric Charge

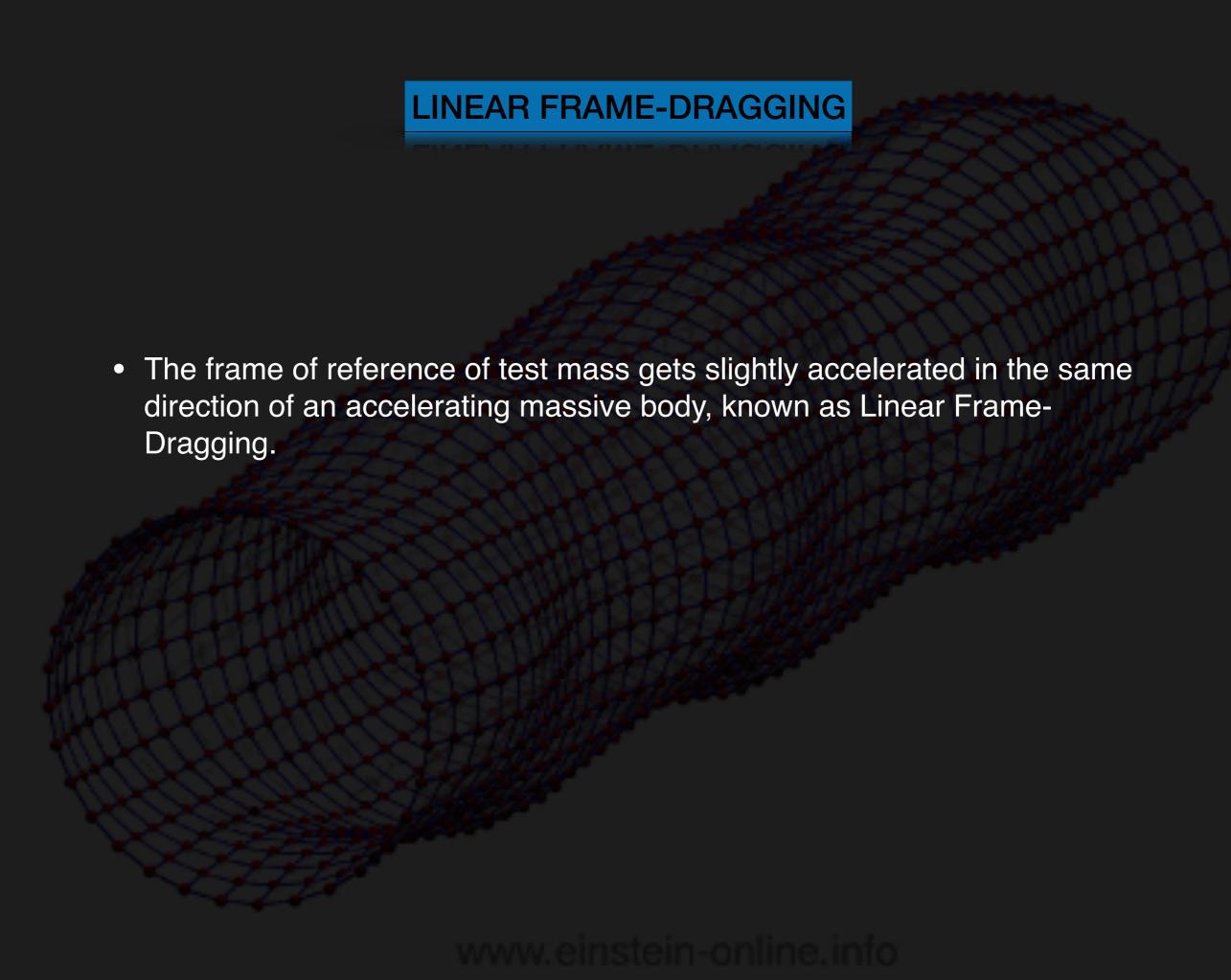
FRAME-DRAGGING

The non-static stationary distributions of mass-energy distorts the space-time. This effect is known as Frame-Dragging.

Einstein's General Theory of Relativity predicted this effect.

There are two basic types of Frame- Dragging:

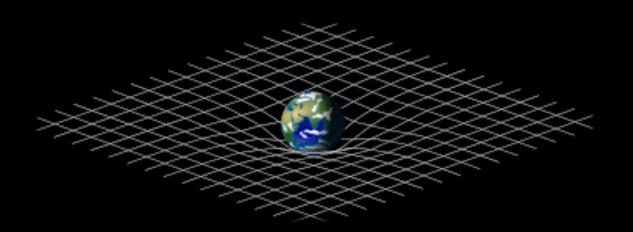
- 1. Linear Frame-Dragging
- 2. Rotational Frame- Dragging



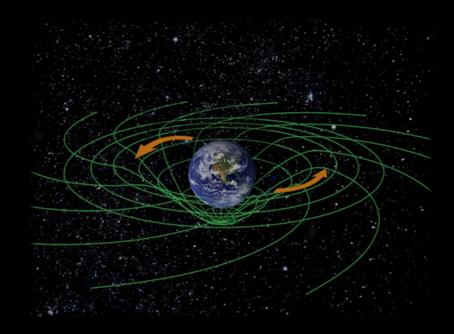
ROTATIONAL FRAME-DRAGGING

• The rotation of a massive object distorts the space-time metric locally, such that, it makes the orbit of a nearby test particle to precess around, this effect is known as Rotation Frame-Dragging.

Non-Rotating Massive Object



Rotating Massive Object



KERR SOLUTION & AZIMUTHAL VELOCITY

Kerr Solution line element :

$$ds^{2} = c^{2}\partial\tau^{2} = \left(1 - \frac{2mr}{\rho^{2}}\right)c^{2}\partial t^{2} + \frac{4mcraSin^{2}\theta}{\rho^{2}}\partial t\partial\phi - \frac{\rho^{2}}{\Delta}\partial r^{2} - \rho^{2}\partial\theta^{2} - \left((r^{2} + a^{2})Sin^{2}\theta + \frac{2mra^{2}Sin^{4}\theta}{\rho^{2}}\right)\partial\phi^{2}$$

$$\rho^2 = r^2 + a^2 Cos^2 \theta$$

$$\Delta = r^2 + a^2 - 2mr$$

$$a = \frac{J}{Mc}$$

$$m = \frac{GM}{c^2}$$

 ϕ and t are Cyclic Coordinates

Therefore,
$$\frac{\partial L}{\partial \dot{\phi}}$$
 = constant

where,
$$L = \frac{1}{2} \left[\left(1 - \frac{2mr}{\rho^2} \right) c^2 \dot{t}^2 + \frac{4mcraSin^2 \theta}{\rho^2} \dot{t} \dot{\phi} - \frac{\rho^2}{\Delta} \dot{r}^2 - \rho^2 \dot{\theta}^2 - \left((r^2 + a^2)Sin^2 \theta + \frac{2mra^2Sin^4 \theta}{\rho^2} \right) \dot{\phi}^2 \right]$$

$$\frac{4mcraSin^2\theta}{\rho^2}\dot{t} - 2\left((r^2 + a^2)Sin^2\theta + \frac{2mra^2Sin^4\theta}{\rho^2}\right)\dot{\phi} = constant$$

Take constant = 0

$$\frac{4mcraSin^2\theta}{\rho^2}\dot{t} - 2\left((r^2 + a^2)Sin^2\theta + \frac{2mra^2Sin^4\theta}{\rho^2}\right)\dot{\phi} = 0$$

$$\dot{\phi} \to 0 \ as \ r \to \infty$$

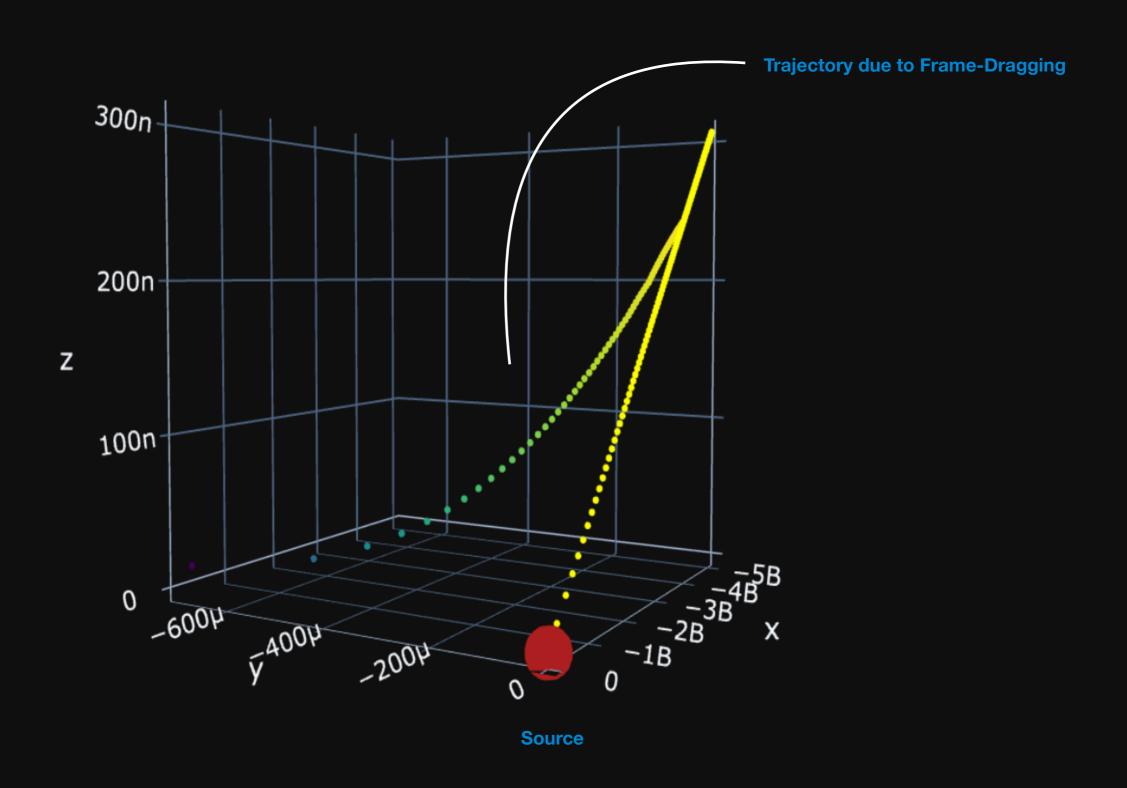
$$\Omega = \frac{d\phi}{dt} = \frac{\dot{\phi}}{\dot{t}} = \frac{2mcra}{(r^2 + a^2)\rho^2 + 2mra^2Sin^2\theta}$$

where,
$$\rho^2 = r^2 + a^2 Cos^2 \theta$$

$$a = \frac{J}{Mc}$$

$$m = \frac{GM}{c^2}$$

Inertial Frame Dragging In Kerr Space-Time



GRAVITY PROBE-B EXPERIMENT

- It was a satellite-based mission (2004) by a Stanford group and NASA.
- Experimentally measure the Frame-Dragging effect due to Gravity.
- Four perfectly spherical gyroscope balls were sent to polar orbit of Earth.

$$\Omega \approx \frac{2mca}{r^3} = 7.64e - 15 \text{ rad/s}$$

PYTHON CODE FOR SIMULATING ROTATIONAL FRAME DRAGGING

```
import numpy as np
import astropy.units as u
from plotly.offline import init_notebook_mode
from einsteinpy.coordinates import BoyerLindquistDifferential
from einsteinpy.bodies import Body
from einsteinpy.geodesic import Geodesic
from einsteinpy.metric.kerrnewman import KerrNewman
init_notebook_mode(connected=True)
## Source ##
m = 2e30 # kg
spin_factor = 0.3 # spin_factor = J/(Mc) (metres)
## Test Particle ##
# initial position cordinates
r = 5e9
                   # (metres)
theta = np.pi * 0.50
                       # (radians)
                    # (radians)
phi = np.pi
# Source
Attractor = Body(name="attractor", mass = m * u.kg, R=0 * u.m, differential = None, a = spin_factor * u.m, q = 0 * u.C, parent=None)
# Coordinate System
sph_obj = BoyerLindquistDifferential(r * u.m, theta * u.rad, phi * u.rad, 0 * u.m/u.s, 0 * u.rad/u.s, 0 * u.rad/u.s, spin_factor * u.m)
Object = Body(name="testparticle", mass = 0 * u.kg, R = 0 * u.m, differential = sph_obj, a = 0 * u.m, q = 0 * (u.C/u.kg), parent=Attractor)
# geodesic simulation
geodesic = Geodesic(body = Object, time=0 * u.s, end_lambda= ((1 * u.year).to(u.s)).value/930, step_size=((0.5 * u.min).to(u.s)).value, metric=KerrNewman)
## Each element here is w.r.t. proper-time (lambda) ##
vals = geodesic.trajectory
t = np.array(vals[:, 0])
x = np.array(vals[:, 1])
y = np.array(vals[:, 2])
z = np.array(vals[:, 3])
```

```
spin_factor=0
# Source
Attractor = Body(name="attractor", mass = m * u.kg, R=0 * u.m, differential = None, a = spin_factor * u.m, q = 0 * u.C, parent=None)
# Coordinate System
sph_obj = BoyerLindquistDifferential(r * u.m, theta * u.rad, phi * u.rad, 0 * u.m/u.s, 0 * u.rad/u.s, 0 * u.rad/u.s, spin_factor * u.m)
# Test Particle
Object = Body(name="testparticle", mass = 0 * u.kg, R = 0 * u.m, differential = sph_obj, a = 0 * u.m, q = 0 * (u.C/u.kg), parent=Attractor)
# geodesic simulation
geodesic = Geodesic(body = Object, time=0 * u.s, end_lambda= ((1 * u.year).to(u.s)).value/930, step_size=((0.5 * u.min).to(u.s)).value, metric=KerrNewman)
## Each element here is w.r.t. proper-time (lambda) ##
vals = geodesic.trajectory
t1 = np.array(vals[:, 0])
x1 = np.array(vals[:, 1])
y1 = np.array(vals[:, 2])
z1 = np.array(vals[:, 3])
import plotly.graph_objects as go
fig = go.Figure()
fig.add trace(
  go.Scatter3d(
    x=x,
    y=y,
    z=z,
    mode="markers",
    marker=dict(size=2, color=y, colorscale='Viridis', opacity=1),
    name="Frame Dragged Trajectory"))
fig.add_trace(
  go.Scatter3d(
    x=x1,
    y=y1,
    z=z1,
    mode="markers",
    marker=dict(size=2, color='yellow', colorscale='Viridis', opacity=1),
    name="Ordinary Trajectory"))
fig.add_trace(
  go.Scatter3d(
    x = [0],
    y=[0],
    z=[0],
    mode="markers",
    marker=dict(size=15, color='firebrick', colorscale='Viridis', opacity=1),
    name="Source"))
fig.update_layout(title={
     'text': "Inertial Frame Dragging In Kerr Space-Time",
     'y':0.9,
     'x':0.4,
     'xanchor': 'center',
     'yanchor': 'top'},
    margin=dict(l=0, r=0, b=0, t=0),
    template="plotly_dark")
fig.show()
```