New Results Concerning Collective Motion in Triaxial Nuclei

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Nuclear Deformation

Most of the nuclei are either *spherical* or *axially symmetric* in their ground-state.

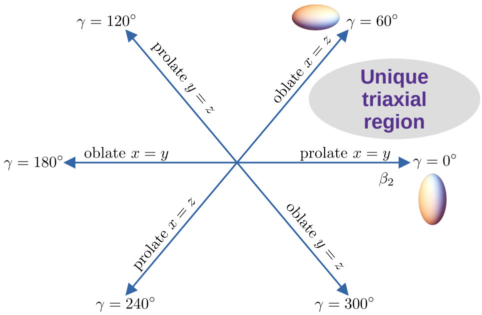
Deformation parameter β (*Bohr, 1969*): preserves axial symmetry

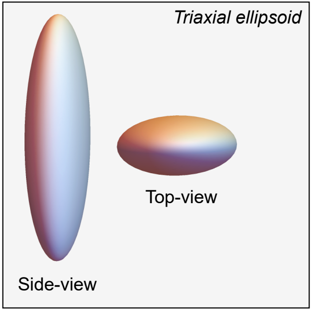
Figure 1: **spherical:** $\beta = 0$ **prolate:** $\beta > 0$ **oblate:** $\beta < 0$

Nuclear Triaxiality

Non-axial shape

Deviations from symmetric shapes can occur across the chart of nuclides → **triaxial nuclei**.
The triaxiality parameter γ (*Bohr, 1969*): departure from axial symmetry





Fingerprints for Triaxiality

- Experimentally, stable triaxial nuclei represent a real challenge
- Clear signatures for confirming stable triaxiality in nuclei
  1. Chiral symmetry breaking (*Frauendorf, 1997*)
  2. **Wobbling motion** (*Bohr & Mottelson, 1975*)

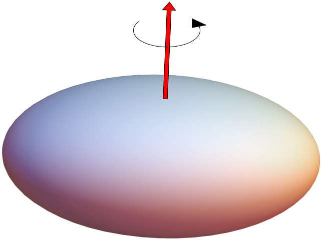
Wobbling Motion (WM)

- Unique to non-axial nuclei
- Predicted 50 years ago for even-A nuclei
- First experimental evidence for ^{163}Lu (*Ødegård*, 2001)
- Currently: confirmed wobblers within the mass regions $A \approx [100, 130, 160, 180]$.

Energy of Deformed Nuclei

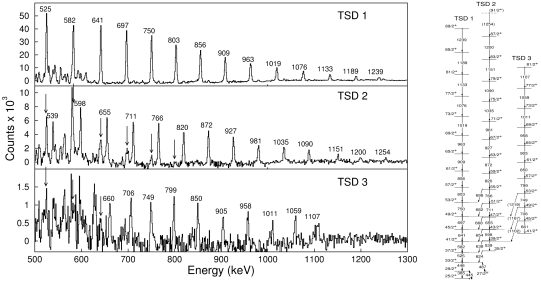
Collective Motion

- A nucleus - *droplet* - can generate angular momentum from the rotation and vibration of the droplet itself
- Each individual nucleon contributes to the total angular momentum → *collectiveness*
- ⚠Rotation can occur only if the nuclear potential is *deformed*



Triaxial Rotor Energy

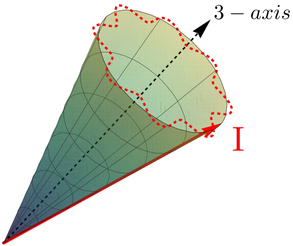
- A triaxial nucleus can rotate about any of the three axes
- Rotation about the axis with **the largest moment of inertia** (MOI) is energetically the most favorable: $E_{\text{rot}} \propto \frac{\hbar^2}{2\mathcal{I}_{\text{max}}} I(I+1)$
- MOI anisotropy $\rightarrow$ the *main rotation* around $\mathcal{J}_{\text{max}}$ is disturbed by the other two axes $\rightarrow$ *total motion of the rotating nucleus has an oscillating behavior*

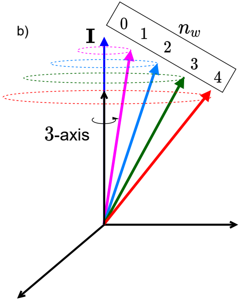


Figures from Ødegård et al., 2001

Wobbling Motion

- Total angular momentum **I** disaligned w.r.t. body-fixed axes
- The a.m. **precesses** and **wobbles** around the axis with $\mathcal{J}_{\mathrm{max}}$
- The precession of **I** can increase by **tilting**
- Tilting by an energy quanta $\sim$ *vibrational character* $\rightarrow$ **wobbling phonon** $n_w = 0, 1, 2...$



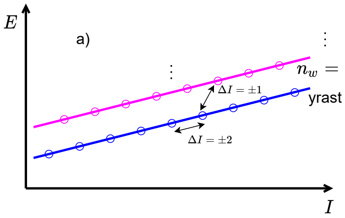


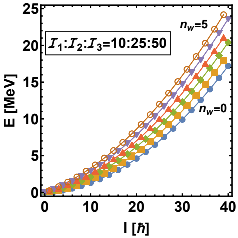
Wobbling Spectrum

Even-A Nuclei

- Employing the Harmonic Approximation *(Bohr, 1969)*
- $\hat{H}$ composed of a *rotational* part and *harmonic oscillation* (i.e., wobbling) part:

$$\hat{H} = \frac{\hbar^2}{2\mathcal{J}_{\max}} I(I+1) + \hbar\omega_{\text{wob}}\left(n_w + \frac{1}{2}\right), n = 0, 1, 2, \ldots \quad (1)$$





New Results for A=130

Recent findings for even-even nuclei

- Two wobbling bands have been identified experimentally in ^{130}Ba (*Petrache et al., 2019*)
- DFT+PRM description of the wobbling motion described the excited spectra (*Chen et al, 2019*)
- Stable triaxiality for $\beta = 0.24$ and $\gamma = 21.5^{\circ}$
- Infer spin-dependence for $\mathcal{J}_{1,2,3}$

Harmonic Approximation

- Employed an energy spectrum of harmonic type as Eq. 1
- Reproduced the excited spectra $\{B1, B2\}$

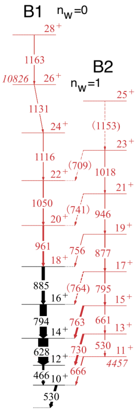


Figure from Petrache et al., 2019

New Results for A=130 II

Harmonic Approximation

- Employed an energy spectrum of harmonic type according to Eq. 1:

$$E_{I,n_w} = \frac{\hbar^2}{2\mathcal{J}_3} I(I+1) + \hbar\omega_{\text{wob}}\left(n_w + \frac{1}{2}\right) \qquad (2)$$

- wobbling frequency is a function of the three MOI (fixed ordering $\mathcal{J}_3 > \mathcal{J}_{1,2}$)

$$\hbar\omega_{\text{wob}} = 2I \times f(\mathcal{J}_1, \mathcal{J}_2, \mathcal{J}_3) \qquad (3)$$

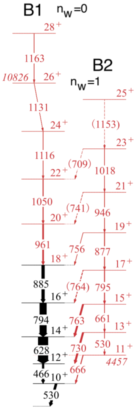
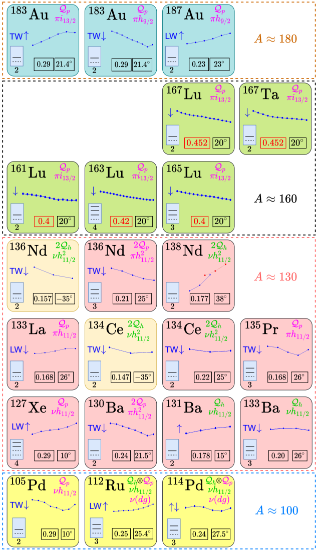


Figure from Petrache et al., 2019

Experimental Evidence



Wobbling nuclei (up to date)

Poenaru, 2022, in progress