

# DESCRIPTION OF THE WOBBLING MOTION THROUGH A BOSON METHOD

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# Outline

1 Introduction

2 Triaxial Shapes

## Nuclear Radius

The **shape** of the nucleus is most generally described in terms of the *nuclear radius*:

$$R(\theta, \varphi; t) = R_0 \left( 1 + \sum_{\lambda=0}^{\infty} \sum_{\mu=-\lambda}^{\lambda} \alpha_{\lambda\mu}(t) Y_{\lambda}^{\mu}(\theta, \varphi) \right) \quad (1)$$

- The  $\alpha_{\lambda\mu}$  are collective coordinates  $\implies$  *vibrations of the nucleus*.
- $Y_{\lambda}^{\mu}$  are the spherical harmonics.

# Nuclear shapes

Most nuclei are spherical or axially symmetric in the ground state.



Figure 1: **Spherical:**  $\beta_2 = 0$  ; **Prolate:**  $\beta_2 > 0$  ; **Oblate:**  $\beta_2 < 0$

# Quadrupole deformations

- Most relevant vibrational degrees of freedom in nuclei.
- Play a crucial role in the rotational spectra of nuclei.

## Quadrupole radius

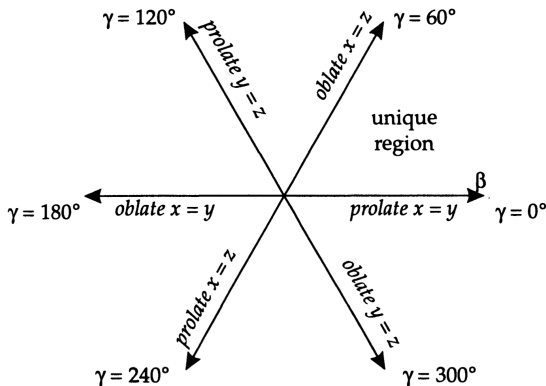
For pure quadrupole deformations:

$$R(\theta, \varphi) = R_0 \left( 1 + \sum_{\mu} \alpha_{2\mu} Y_2^{\mu}(\theta, \varphi) \right) , \quad (2)$$

Using A. Bohr's description, the coordinates  $\alpha_{2\mu}$  can be reduced to only two *deformation parameters*:  $\beta_2$  (*eccentricity*) and  $\gamma$  (**triaxiality**).

# Nuclear triaxiality

- Besides the axially symmetric shapes (i.e., spherical, prolate, and oblate), nuclei can be **triaxial**  $\implies$  lack of symmetry along any of the principal axes.
- The asymmetry is given by the non-zero value of  $\gamma$ .



# Triaxial ellipsoid

Schematic example with a triaxial ellipsoid ( $\gamma \neq 0$ )  $\beta_2 > 0$ .

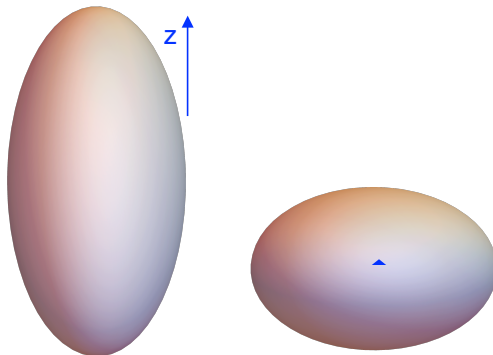


Figure 2: **Left:** side-view. **Right:** top view.

- Probing triaxiality experimentally is a real challenge (e.g., large and complex detector setups).
- Only two fingerprints known so far: **chiral motion** (Frauendorf, 1997) and **wobbling motion** (Bohr and Mottelson, 1975).

## Wobbling Motion (WM)

- Collective effect → *unique* to triaxial nuclei.
- Predicted almost 50 years ago, first experimental confirmation: in 2001 (Odegard et al.) for  $^{163}\text{Lu}$ .
- In present, few wobblers are experimentally confirmed in the mass regions:  $A \approx 130, 160, 180 \rightarrow$  A list of all known wobblers will be available in my PhD thesis (Chapter 3).



# Wobbling motion

## Triaxial nuclei

A triaxial nucleus can rotate about any of the three axes.

The rotational angular momentum (a.m.) is NOT aligned along any of the body-fixed axes  $\Rightarrow$  **precesses** and **wobbles** around the axes with the largest MOI.

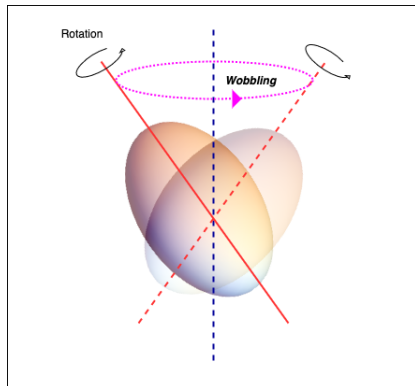


Figure 3: Schematic representation for the nuclear wobbling motion.

## Wobbling bands

Sequences of  $\Delta I = 2\hbar$  rotational bands that are built on different *wobbling phonon excitations*.

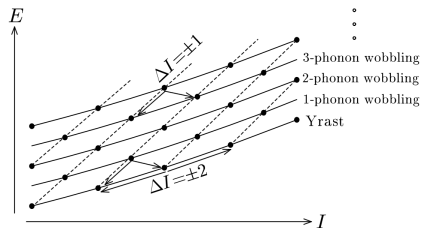


Figure 4: Rotational-band structures of the wobbling motion.

- For  $^{163}\text{Lu}$ :  $n_w = 0, 1, 2, 3$  wobbling phonon numbers, respectively.
- Nuclei have **large** quadrupole moments
- **Strong E2** character for the electro-magnetic transitions.

# Even-Even vs. Even-Odd Nuclei

## Theoretical frameworks for **even-mass** nuclei

- 1 *Harmonic Approximation(s)* (Bohr and Mottelson, 1975)
- 2 Triaxial-Rotor-Model (Davydov and Filippov, 1958)
- 3 Boson-approximations (Tanabe, 1971)

## Theoretical frameworks for **odd-mass** nuclei

- 1 Particle Rotor Model (Hamamoto, 2002)
- 2 Tilted-axis wobbling (Frauendorf and Meng, 1997)
- 3 RPA, Mean-Field Theories, GCM+AMP...

## Recent work on wobbling motion

- RPA for  $^{163}\text{Lu}$ , Raduta et al (PRC, 2017)
- Tilted-axis wobbling for  $^{135}\text{Pr}$ , R. Budaca (PRC, 2018)
- PRM for  $^{163}\text{Lu}$ , R. Poenaru (IJMPE, 2021)

## Description of WM for an even-odd nucleus

- one single-particle (nucleon) *coupled* to an *even-even* triaxial core.
- the nucleon is moving in a quadrupole deformed mean-field generated by the core
- particle + rotor coupling drives the entire system to large (and stable) deformations ( $\epsilon \sim 0.2 - 0.4$ ).

Hamiltonian:

$$\hat{H}_{\text{rot}} = \sum_{k=1}^3 A_k \left( \hat{I}_k - \hat{j}_k \right)^2 . \quad (3)$$

$A_k \rightarrow$  inertia parameters:  $A_k = (2\mathcal{I}_k)^{-1}$ .

# Theoretical Formalism - Rotational Hamiltonian