

$$0 = G_i(\underline{\varrho}) = \frac{4}{\varrho_i} + \frac{4\alpha + 2}{a^2 + \varrho_i} + \frac{4\beta + 2}{b^2 + \varrho_i} + \frac{4\gamma + 2}{c^2 + \varrho_i} + \sum_{\substack{j=1 \\ i \neq j}}^n \frac{8}{\varrho_i - \varrho_j}$$

$$\frac{\omega^2}{c_0^2} = 2 \left( \frac{\alpha}{a^2} + \frac{\beta}{b^2} + \frac{\gamma}{c^2} \right) - \sum_{i=1}^n \frac{4}{\varrho_i} \quad i = 1, \dots, n$$

$$\delta \ln(z) = \text{const.} \cdot x_1^\alpha x_2^\beta x_3^\gamma \prod_{i=1}^n \left( 1 - \frac{x_1^2}{a^2 + \varrho_i} - \frac{x_2^2}{b^2 + \varrho_i} - \frac{x_3^2}{c^2 + \varrho_i} \right) \quad n \geq 1$$

$$1 \quad n = 0$$

$$G_i(\underline{\varrho}) = - \frac{\partial V(\underline{\varrho})}{\partial \varrho_i}$$

~ we can find the  $\varrho$ -s with a modified Newton-Raphson method  
 • it should be stabilized so it does not diverge.

$$G(\underline{\varrho}^{(k)}) \neq 0 \quad \text{the } k^{\text{th}} \text{ approx.}$$

$$G(\underline{\varrho}^{(k)} + \delta \underline{\varrho}) = G(\underline{\varrho}^{(k)}) + \underbrace{\left( \frac{\partial G}{\partial \underline{\varrho}} \right)_{\underline{\varrho}^{(k)}}}_{\text{derivative matrix}} \cdot \delta \underline{\varrho} = 0$$

~ derivative matrix.

$$\delta \underline{\varrho} = - \left( \frac{\partial G}{\partial \underline{\varrho}} \right)_{\underline{\varrho}^{(k)}}^{-1} G(\underline{\varrho}^{(k)})$$

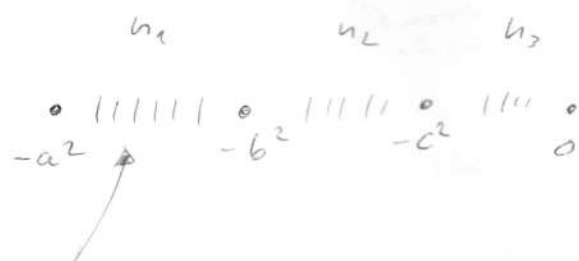
~ computationally expensive

$$\left( \frac{\partial G}{\partial \underline{\varrho}} \right)_{\underline{\varrho}^{(k)}}^{-1} \delta \underline{\varrho} = G(\underline{\varrho}^{(k)}) \quad \text{is better.}$$

some numerical trickery + Gauss-elimination

$$\underline{Q}^{(k+1)} = \underline{Q}^{(k)} + \delta \underline{Q}$$

• this can be iterated to get a good  $\underline{Q}$



we put  $n_i$  changes here equidistantly  $\rightarrow$  "lin space"

• this will be  $\underline{Q}^{(0)}$

• components are ordered  $\mathcal{Q}_1 < \mathcal{Q}_2 < \dots$

• after adding  $\delta \underline{Q}$  the relation between  $\mathcal{Q}_i$ -s should stay the same.

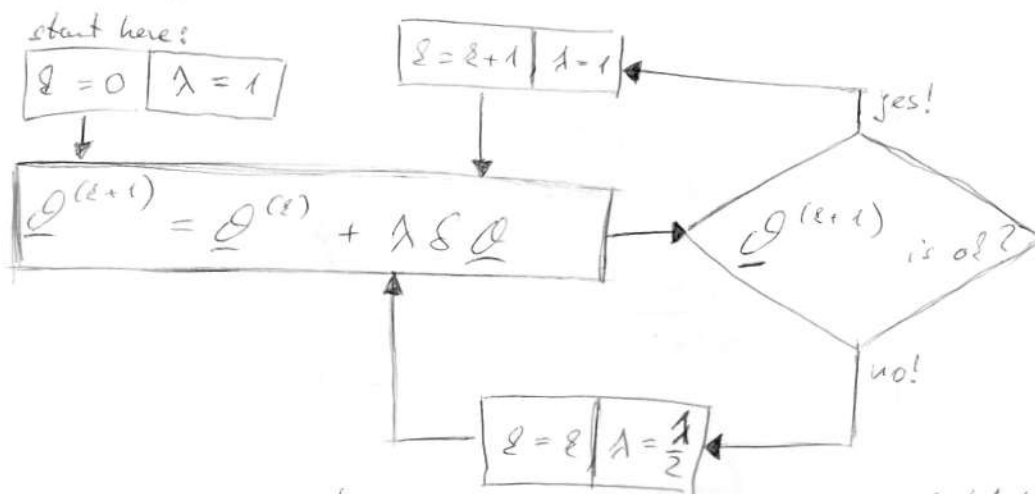
$\downarrow$

$$-a^2 < \mathcal{Q}_1 \dots < \mathcal{Q}_{n_1} < -b^2 < \dots < \mathcal{Q}_{n_1+n_2+n_3} < 0$$

• the "walls" stay at the same place after every iteration!

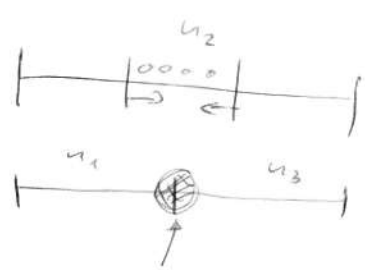
• if the config is okay Newton-Raphson can be used to refine the solution.

• we can try  $\frac{1}{2} \cdot \delta \underline{Q}$  if the previous trial was ~~not~~ bad.



• only a few starting steps are needed to stabilize the process.

$a, b, c \rightarrow a, b$



$\omega_1 = \omega_x, \omega_y$   
 $\omega_z$

cylindrical case

new, renormalized charge in the middle with no conserved quantity.

it is easier to start from scratch:  $x_1^\alpha x_2^\beta x_3^\gamma$

$$\delta u = \rho^{lm} e^{in\phi} z^\beta \left\{ \prod_{i=1}^n \left( 1 - \frac{\rho^2}{a^2 + \mathcal{Q}_i} - \frac{z^2}{b^2 + \mathcal{Q}_i} \right) \right.$$

the new pre-factor  $\dots 1$

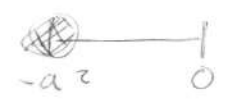
same calc., same equations. can be repeated...

$$\frac{\omega^2}{c_0^2} = 2 \left( \frac{|lm|}{a^2} + \frac{\beta}{b^2} \right) - \sum_{i=1}^n \frac{4}{\mathcal{Q}_i} \quad \text{trap freq.}$$

$$0 = \frac{4|lm| + 4}{a^2 + \mathcal{Q}_i} + \frac{4\beta + 2}{b^2 + \mathcal{Q}_i} + \frac{4}{\mathcal{Q}_i} + \sum_{\substack{j=1 \\ i \neq j}}^n \frac{8}{\mathcal{Q}_i - \mathcal{Q}_j} \quad \text{force - eq.}$$

$a, b \rightarrow a$

$\omega_0 = \omega_x = \omega_y = \omega_z$   
 spherical case



renormalized value of the fixed charge

$$r^l Y_{lm}'(\theta, \phi) \left\{ \prod_{i=1}^n \left( 1 - \frac{r^2}{a^2 + \mathcal{Q}_i} \right) \right\} = \delta u$$

$$\frac{\omega^2}{c_0^2} = \frac{2l}{a^2} - \sum_{i=1}^n \frac{4}{\mathcal{Q}_i} \quad \text{trap freq.}$$

$$0 = \frac{4}{d_i} + \frac{4\ell + 6}{a^2 + d_i} + \sum_{\substack{j=1 \\ i \neq j}}^n \frac{8}{d_i - d_j} \quad \text{face - eq.}$$

→ this form gives a new evaluation of a hypergeometric polynomial (which is the known solution for this prob.)

→ it also gives the roots ...

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