Modern Physics

Luis A. Anchordoqui

Department of Physics and Astronomy Lehman College, City University of New York

> Lesson V October 1, 2015

Table of Contents

Particle Dynamics

- Particle Decay and Collisions
 - Conservation of 4-momentum and all that...
 - Two-body decay of unstable particles





Advanced Search
Google Search I'm Feeling Lucky
Language Tools

- Newtons first law of motion holds in special relativistic mechanics as well as nonrelativistic mechanics
- In absence of forces body is at rest or moves in straight line at constant speed

$$\frac{d\mathbf{U}}{d\tau} = 0 \tag{1}$$

 $U(\gamma c, \gamma \vec{u})$ (1) implies that \vec{u} is constant in any inertial frame

• Objective of relativistic mechanics:

introduce the analog of Newtons 2nd law

$$\vec{F} = m\vec{a} \tag{2}$$

- There is nothing from which this law can be derived but plausibly it must satisfy certain properties:
 - It must satisfy the principle of relativity
 i.e.
 □ take the same form in every inertial frame
 - 2 It must reduce to (1) when the force is zero
 - It must reduce to (2) in any inertial frame
 when speed of particle is much less than speed of light

$$m\frac{d\mathbf{U}}{d\tau} = f \tag{3}$$

- m respective
 m res
- f ➡ 4-force.
- Using $d\mathbf{U}/d\tau = A$ (3) can be rewritten in evocative form

$$f = mA \tag{4}$$

- This represents 4-equations
 but they are not all independent
- Normalization of the 4-velocity $\bowtie U_{\mu}U^{\mu}=c^2$ implies

$$m\frac{d(\mathbf{U}\cdot\mathbf{U})}{d\tau} = 0 \Rightarrow f\cdot\mathbf{U} = 0 \tag{5}$$

(5) shows there are only 3 independent equations of motion
 – same number as in Newtonian mechanics –

4-momentum defined by

$$p = m U (6)$$

equation of motion can be rewritten as

$$\frac{dp}{d\tau} = f \tag{7}$$

important property of 4-momentum rinvariant mass

$$p_{\mu} p^{\mu} = m^2 c^2 \tag{8}$$

ullet Components of 4-momentum related \vec{u} according to

$$p^0 = \frac{mc}{\sqrt{1 - u^2/c^2}}$$
 and $\vec{p} = \frac{m\vec{u}}{\sqrt{1 - u^2/c^2}}$ (9)

• For small speeds $u \ll c$

$$p^{0} = mc + \frac{1}{2}m\frac{u^{2}}{c} + \cdots$$
 and $\vec{p} = m\vec{u} + \cdots$ (10)

- \vec{p} reduces to usual 3-momentum
- p⁰ reduces to kinetic energy per units of c plus mass in units of c
- p

 also called energy-momentum 4-vector

$$p^{\mu} = (E/c, \vec{p}) = (m\gamma c, m\gamma \vec{u}) \tag{11}$$

• mass is part of energy of relativistic particle

$$p_{\mu} p^{\mu} = m^2 c^2 \Rightarrow E = (m^2 c^4 + \vec{p}^2 c^2)^{1/2}$$
 (12)

• For particle at rest (12) reduces to $\mathbb{E} E = mc^2$

In particular inertial frame...

 Connection between relativistic equation of motion and Newton's laws can be made more explicit by defining 3-force

$$\frac{d\vec{p}}{dt} \equiv \vec{F} \tag{13}$$

- ullet Same form as Newton's law but with relativistic expression for $ec{p}$
- Only difference relation momentum to velocity

$$\vec{f} = \frac{d\vec{p}}{d\tau} = \frac{d\vec{p}/dt}{dt/d\tau} = \gamma \vec{F}$$
 (14)

4-force acting on particle can be written in terms of 3-force

$$f = (\gamma \vec{F} \cdot \vec{u}, \gamma \vec{F}) \tag{15}$$

• Time component of equation of motion

$$\frac{dE}{dt} = \vec{F} \cdot \vec{u} \tag{16}$$

familiar relation from Newtonian mechanics

- Time component of equation of motion consequence of other 3
- In terms of three force equations of motion take same form as they do in usual Newtonian mechanics but with relativistic expressions for energy and momentum
- ullet For $v\ll c$ relativistic version of Newton's second law reduces to the familiar nonrelativistic form
- Newtonian mechanics is low-velocity approximation of relativistic mechanics

Light rays

- Massless particles move at speed of light along null trajectories
- Proper time interval between any two points is zero
- Curve x = ct could be written parametrically in term of λ

$$x^{\mu} = U^{\mu}\lambda$$
 with $U^{\mu} = (c, c, 0, 0)$ (17)

U is a null vector

$$\mathbf{U} \cdot \mathbf{U} = 0 \tag{18}$$

- Different parametrizations give different tangent 4-vectors but all have zero length
- With choice (17)

$$\frac{d\mathbf{U}}{d\lambda} = 0 \tag{19}$$

light ray equation of motion is same as for particle

 Basic law of collision mechanics conservation of 4-momentum: Sum of 4-momenta of all particles going into point-collision is same as sum of 4-momenta of all those coming out

$$\sum^* p_i = 0 \tag{20}$$

∑* sum that counts pre-collision terms positively and post-collision terms negatively

- For closed system conservation of total 4-momentum can be shown to be result of spacetime homogeneity
- Whether law is actually true must be decided by experiment
- Countless experiments have shown that total 4-momentum of isolated system is constant

CM frame

- If we have a system of particles with 4-momenta p_i subject to no forces except mutual collisions
- ullet total 4-momentum $p_{\mathrm{tot}} = \sum p_i$ is timelike and future-pointing
- ullet there exists an inertial frame S in which spatial components of p_{tot} vanish

S should be called center-of-momentum frame but is called CM frame

Invariant mass

- Invariant mass of two particles with 4-momenta p_a and p_b $m_{ab}^2c^2=(p_a+p_b)^2$
- Invariant mass useful to find mass of short-live unsatble particles from momenta of their observed decay products
- Consider $X \rightarrow a + b \bowtie p_X = p_a + p_b$

$$m_X^2 c^2 = (p_a + p_b)^2 = p_a^2 + p_b^2 + 2p_a \cdot p_b$$

= $m_a^2 c^2 + m_b^2 c^2 + 2E_a E_b / c^2 - 2\vec{p}_a \cdot \vec{p}_b$ (21)

- In high energy experiment
 3-momenta and masses of particles a and b must be measured
- For charged particles this requires a magnetic field and tracking of trajectory to measure bending as well as some means of particle identification
- One must also identify the vertex and measure the opening angle

- Consider decay process $X \rightarrow ab$
- In CM frame for a and b mother particle X is at rest
- 4-momenta

$$p_X = (Mc, 0, 0, 0)$$
 $p_a = (E_a/c, \vec{p}_a)$ $p_b = (E_b/c, \vec{p}_b)$ (22)

Conservation of 4-momentum requires:

$$p_X = p_a + p_b \ \vec{p}_a = -\vec{p}_b$$

Omitting subscript on 4-momenta
energy conservation reads

$$E_a + E_b = \sqrt{m_a^2 c^4 + p^2 c^2} + \sqrt{m_b^2 c^4 + p^2 c^2} = Mc^2$$
 (23)

Solving (23) for p

$$p = c \frac{\sqrt{[M^2 - (m_a - m_b)^2][M^2 - (m_a + m_b)^2]}}{2M}$$
 (24)

Immediate consequence

$$M \ge m_a + m_b \tag{25}$$

X decays only if mass exceeds sum of decay products masses

- Conversely
 - if particle has mass exceeding masses of two other particles particle is unstable and decays unless decay is forbidden by some conservation law
- e.g. conservation of charge, momentum, and angular momentum Momenta of daughter particles and energies fixed by 3 masses: from energy conservation (23) $E_b = \sqrt{E_a^2 m_a^2 c^4 + m_b^2 c^4}$

solve to get

$$E_a = \frac{1}{2M}(M^2 + m_a^2 - m_b^2)c^2 \tag{26}$$

similarly

$$E_b = \frac{1}{2M}(M^2 + m_b^2 - m_a^2)c^2 \tag{27}$$

 No preferred direction in which the daughter particles travel decay is said to be isotropic

□ daughter particles travelling back-to-back in X rest frame

- Of interest is also two-body decay of unstable particles in flight
- In-flight decays
 only way to measure mass of neutral particle
- Take z-axis along direction of flight of mother particle

$$p_X = (E/c, 0, 0, p)$$
 $p_a = (E_a/c, \vec{p}_{a\perp}, p_{az})$ $p_b = (E_b/c, \vec{p}_{b\perp}, p_{bz})$

By momentum conservation
 transverse momentum vectors

$$\vec{p}_{\perp} \equiv \vec{p}_{a\perp} = -\vec{p}_{b\perp} \tag{28}$$

 Energies and z components of particle momenta related to those in the CM frame by a Lorentz boost with a boost velocity equal to the speed of the mother particle

$$E_a/c = \gamma(E_a^*/c + \beta p_{az}^*)$$

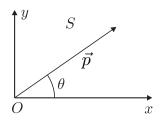
$$p_{az} = \gamma(p_{az}^* + \beta E_a^*/c)$$

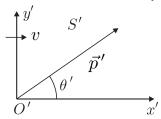
$$\vec{p}_{a\perp} = \vec{p}_{a\perp}^*$$

$$\beta = pc/E \text{ and } \gamma = E/(Mc^2)$$

L. A. Anchordogui (CUNY)

This completely solves problem
 e.g. we can find angles which daughter particles
 make with z-axis and with each other as functions of p_X





In S we have

$$p^{\mu} = (E/c, p\cos\theta, p\sin\theta, 0) \tag{29}$$

in S' it follows that

$$p'^{\mu} = (E'/c, p'\cos\theta', p'\sin\theta', 0)$$
 (30)

Applying $S \to S'$ Lorentz transformation

$$p'\cos\theta' = \gamma^*(p\cos\theta - \beta^*E/c)$$

$$p'\sin\theta' = p\sin\theta$$
 (31)

SO

$$\tan \theta' = \frac{p \sin \theta}{\gamma^* (p \cos \theta - \beta^* E/c)}$$
 (32)

or

$$\tan \theta' = \frac{\sin \theta}{\gamma^* (\cos \theta - \beta^* / \beta)} \tag{33}$$

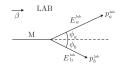
 $\beta^* = v/c$ represented velocity of S' wrt S and $\beta = pc/E$ represented velocity of particle in S

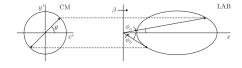
Inverse relation is found to be

$$an \theta = \frac{\sin \theta'}{\gamma^* (\cos \theta' + \beta^* / \beta')}$$
 (34)

 $\beta' = p'c/E'$ we velocity of particle in S'







Example

- Mother particle of mass M is traveling with velocity $\beta = |pc|/E$
- In the CM frame \blacksquare particle a has energy E_a and momentum \vec{q} @ angle θ with respect to x'-axis
- Particle momenta in the lab frame

$$p_a^{\text{lab}}\cos\phi_a = \gamma(q\cos\theta + \beta E_a/c)$$

$$p_b^{\text{lab}}\cos\phi_b = \gamma(-q\cos\theta + \beta E_b/c)$$
(35)

 Use inverse Lorentz transformation to obtain variables in CM frame from measured parameters in lab 2nd approach

start from energy-momentum conservation

$$E = E_a + E_b = \sqrt{m_a^2 c^4 + p_a^2 c^2} + \sqrt{m_b^2 c^4 + p_b^2 c^2}$$
 (36)

$$\vec{p} = \vec{p}_a + \vec{p}_b \tag{37}$$

• Substituting in (36) p_b^2 by $(\vec{p} - \vec{p}_a)^2$

$$p_a = \frac{(M^2 + m_a^2 - m_b^2)c^2 \ p\cos\theta_a \pm 2E \ \sqrt{M^2 p^{*2} - m_a^2 p^2 \sin^2\theta_a}}{2(M^2 c^2 + p^2 \sin^2\theta_a)}$$

- By demanding p_a to be real $M^2 p^{*2} m_a^2 p^2 \sin^2 \theta_a \ge 0$
- This condition is satisfied for all angles θ_a if $Mp^*/(m_ap)>1$ negative sign must be rejected \bowtie unphysical $p_a<0$ for $\theta_a>\pi/2$
- If $Mp^*/(m_ap) < 1$ region of parameter space in which both signs must be kept: for each value of $\theta_a < \theta_{a,\max}$ there are two values of p_a and correspondingly also two values of p_b