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# FERMI-LAT OBSERVATIONS OF EXTENDED GAMMA-RAY EMISSION IN THE DIRECTION OF SNR $\mathrm{G}150.3+4.5$

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#### ABSTRACT

We report here a dedicated analysis of the  $\gamma$ -ray emission around supernova remnant (SNR) G150.3+4.5, observed with the Large Area Telescope (LAT) on board the Fermi Gamma-Ray Space Telescope. The Second Catalog of Hard Fermi-LAT Sources reported detection of a hard spectrum, spatially extended source from 50 GeV - 2 TeV, partially overlapping G150.3+4.5. Lowering the energy threshold to 1 GeV, we significantly detect a large ( $\sigma = 1.40^{\circ} \pm 0.03^{\circ}$ ) extended  $\gamma$ -ray source consistent with the entirety of the radio shell and displaying a power law spectral index of  $1.82 \pm 0.04$ . An obtained HI spectrum toward the SNR suggests that the remnant could be one of the closest to us. Estimates of its age, within the context of other LAT observed SNRs, indicate that G150.3+4.5 is in the Sedov-Taylor phase, and is compatible with a dynamically-young remnant [JAM: modfly this with low density in mind?]. Despite the spectral similarities with other unevolved SNRs, ROSAT all-sky survey observations show no prominent X-ray emission in the region. We model the broadband non-thermal radiation from G150.3+4.5 using a published radio spectrum of the SNR and the GeV results presented here. We find that the emission is best described by [JAM: or, can be described equally well by blah. End with othe naima results]

Keywords: Supernova Remnants,  $\gamma$ -rays, Cosmic rays, Radio

#### 1. INTRODUCTION

SNRs have long been thought to be the most-likely accelerators of cosmic rays (CRs) up to the knee of the CR energy spectrum, with diffusive shock acceleration being the primary mechanism accelerating the charged particles to  $\gamma$ -ray emitting energies (see Reynolds (2008) for a review of SNRs from X-rays to  $\gamma$ -rays ). Fermi-LAT was instrumental [JAM: pun!] in demonstrating that CR protons can indeed be accelerated by SNR shock fronts (through detection of the characteristic "pion bump" feature), and are capable of generating the observed  $\gamma$ -ray emission in SNRs (Ackermann et al. 2013; Jogler & Funk 2016). In addition, observations of SNRs with the LAT have proven to be vital in uncovering a large swath of the  $\gamma$ -ray SNR population; both evolved SNRs interacting with dense surrounding material, as well as dynamically young remnants useful for probing acceleration directly at the shock (Acero et al. 2016).

The recently updated Pass 8 LAT event reconstruction provides a significantly improved angular resolution, acceptance, and background event rejection (Atwood et al. 2013a,b), all of which lead to an increase in the effective energy range and sensitivity of the LAT. Leveraging the increased sensitivity afforded by Pass 8 data, Ackermann et al. (2016) performed an all-sky analysis from 50 GeV to 2 TeV (referred to as the second catalog of hard Fermi-LAT sources, or 2FHL), directly connecting GeV LAT observations with those of ground-based Chernkov telescopes at higher energies. While it's troublesome for Cherenkov telescopes operating under pointed observations to detect broadly extended sources on the sky (i.e. sources larger than the telescopes field of view (FOV). the LAT, with its all-sky survey mode and wide FOV, is well suited for this task. The 2FHL catalog detected significant spatial extension from 31 sources above 50 GeV, 5 of which had not previously been detected as extended. Of particular interest, one of the 5 blindly detected sources, 2FHL J0431.2+5553e, was a large extended source (modeled as a uniform disk with radius.  $\sigma = 1.27^{\circ} \pm 0.04^{\circ}$ ), exhibiting a hard power-law spectral index ( $\Gamma = 1.66 \pm 0.20$ ). This 2FHL source was found to be coincident with a recently detected radio SNR, G150.3+4.5. Faint emission from the eastern portion of the shell of G150.3+4.5 was first reported in Gerbrandt et al. (2014) (called G150.8+3.8), and considered a strong SNR candidate due to the semi-circular shape of the emission, clearly non-thermal spectrum, and the presence of red optical filamentary structures. Gao & Han (2014) performed follow-up observations of the region using Urumqi 6 cm survey data (as well as Effelbserg 11cm and 21cm data and CGPS 1420 MHz and 408 MHz observations), taking advantage of the survey's extended Galactic latitude range, up to b=20°. They reported clear detection of a 2.5° wide by 3° high, synchrotron emitting, shell-like object (G150.3+4.5), bolstering an SNR origin for the radio emission.

2FHL J0431.2+5553e only partially overlaps the northern region of G150.3+4.5, so the nature of the extended source is uncertain. In this paper, we perform an in depth study of the  $\gamma$ -ray emission in the direction of SNR G150.3+4.5, extending the energy from 50 GeV in 2FHL, down to 1 GeV. We report here detection of a significantly extended source whose extent matches well with that of G150.3+4.5. We describe the LAT observations and explore the spectral and spatial properties of the extended  $\gamma$ -ray source in §2. In §3 we employ archival HI and X-ray observations to assess the properties of the environment G150.3+4.5 resides in. Finally, in §4 we discuss potential  $\gamma$ -ray emission scenarios and model the broadband emission from the source to constrain the origin of the GeV emission and understand the connection between the radio detected source G150.3+4.5 and the  $\gamma$ -ray one.

#### 2. Fermi-LAT OBSERVATIONS AND ANALYSIS

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## 2.1. Data Set and Reduction

Fermi-LAT is a pair conversion telescope sensitive to 156 high energy  $\gamma$ -rays from 20 MeV to greater than 1 TeV 157 (Ackermann et al. 2016), operating primarily in a sky-158 survey mode which views the entire sky every 3 hours. 159 The LAT has a wide field of view ( $\sim$ 2.4 sr), a large 160 effective area of  $\sim$ 8200 cm<sup>2</sup> above 1 GeV for on axis 161 events and a 68% containment radius angular resolution 162 of  $\sim$ 0.8° at 1 GeV. For further details on the instrument 163 and its performance see Atwood et al. (2009) and Ackermann et al. (2012).

In this analysis, we analyzed 7 years of Pass 8 data, <sup>166</sup> from August 2nd 2008 to August 2nd 2015. Source class <sup>167</sup> events were analyzed within a  $14^{\circ}x14^{\circ}$  region centered <sup>168</sup> on SNR G150.3+4.5 using the P8R2\_SOURCE\_V6 in- <sup>169</sup> strument response functions, with a pixel size of  $0.1^{\circ}$ . <sup>170</sup> To reduce contamination from earth limb  $\gamma$ -rays, only <sup>171</sup> events with zenith angle less than  $100^{\circ}$  were included. <sup>172</sup>

For spectral and spatial analysis we utilized both 173 the standard Fermi Science Tools (version 10-01-174 01)<sup>1</sup>, and the binned maximum likelihood package <sup>175</sup> pointlike (Kerr 2010). pointlike provides methods 176 for simultaneously fitting the spectrum, position, and 177 spatial extension of a source, and was extensively vali- 178 dated in Lande et al. (2012). Both packages fit a source 179 model, the Galactic diffuse emission, and an isotropic 180 component (which accounts for the background of mis- 181 classified charged particles and the extragalactic diffuse 182  $\gamma$ -ray background)<sup>2</sup> to the observations. In this anal- <sup>183</sup> ysis, we used the standard Galactic diffuse ring-hybrid 184 model scaled for Pass 8 analysis, gll\_iem\_v06.fits (mod-185 ulated by a power law function with free index and 186 normalization), and for the isotropic emission, we used iso\_P8R2\_SOURCE\_V6\_v06.txt, extrapolated to 2 TeV as in Ackermann et al. (2016).

In our source model for the region, we included sources from the third Fermi-LAT catalog (Acero et al. 2015, 3FGL) within 15° of the center of our region of interest (RoI). We replaced the position and spectrum of any 3FGL pulsars in the region with their corresponding counterpart from the LAT 2nd pulsar catalog (Abdo Residual emission unaccounted for by et al. 2013). 3FGL sources is present in the RoI due to the increased time range and different energy selection with respect to that in 3FGL. We added to the RoI several significant (TS > 16) point sources to account for this unmodeled emission and minimize the global residuals. The closest of these sources added was over 1° away from the edge of the best fit GeV disk. Considering the size of the PSF at 1 GeV, the affect of these sources on the disk fit was assumed to be negligible and we don't discuss them further. The normalization and spectral index of sources within 5° of the center of the RoI were free to vary, whereas all other source parameters were fixed. A preliminary maximum likelihood fit of the RoI was performed, and sources with a test statistic (TS) < 9 (TS) is defined as,  $TS = 2 \operatorname{Log}(\mathcal{L}_1/\mathcal{L}_0)$  where  $\mathcal{L}_1$  is the likelihood of source plus background and  $\mathcal{L}_0$  that of just the background) were removed from the model.

## 2.2. Morphological Analysis

Studying the spatial extension of sources with the LAT is non-trivial due to the energy-dependent point spread function (PSF) and strong diffuse emission present in the Galactic plane. Soft spectrum point sources and uncertainties in the diffuse model can act as sources of systematic error when not accurately modeling extended emission as such, particularly at low energies where the PSF is broad. To strike a balance between the best angular resolution and minimal source and diffuse contamination, we restrict our morphological analysis to energies between 1 GeV and 1 TeV. We divide this energy range into 12 logarithmically spaced bins for both pointlike and gtlike binned likelihood analyses.

Three unidentified 3FGL sources are located within the extent of G150.3+4.5. 3FGL J0425.8+5600, located approximately 0.6° from the center of the SNR, is the closest of the three sources and is described with a power law spectrum of index  $\Gamma = 2.35 \pm 0.17$  in the 3FGL catalog. The closest radio source to 3FGL J0425.8+5600 is NVSS J042719+560823, at 0.25 away (Condon et al. 1998). 3FGL J0423.5+5442, exhibits a power law spectral index,  $\Gamma = 2.63 \pm 0.15$ , with no clear multiwavelength source association. Finally, 3FGL J0426.7+5437 has a pulsarlike spectrum, yet in a timing survey performed with the 100-m Effelsberg radio telescope, Barr et al. (2013) were unable to detect pulsations from the source down to a limiting flux density of  $\sim 0.1$  mJy. This source is located about 0.84° from the center of the SNR. We discuss 3FGL J0426.7+5437 and potential association with G150.3+4.5 further in §4.2. Figure 1 is a counts map of the region, showing the location of the 3FGL sources.

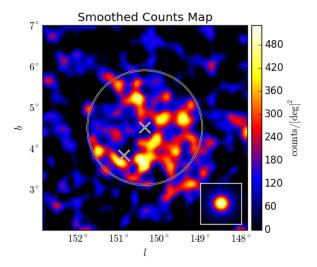


Figure 1. Smoothed background subtracted residual counts map above 1 GeV with  $0.1^{\circ}$ x  $0.1^{\circ}$  pixels, centered on SNR G150.3+4.5. 3FGL J0426.7+5437 and the diffuse backgrounds are included in the region model, 3FGL J0425.8+5600 and 3FGL J0423.5+5442 are not. Inset shows the size of the PSF above 1 GeV, demonstrating the  $\gamma$ -ray emission is extended well beyond the extent of the inset PSF. [JAM: add more description after i remake the map. Add the other 2 3FGL sources? The point of this is to give an sense of what our model thinks the GeV emission from the SNR looks like remove the extended source? 3FGL removed in a different color? Remove titles, use same cmap as figure 1, bigger bolder font.

http://fermi.gsfc.nasa.gov/ssc/
ttp://fermi.gsfc.nasa.gov/ssc/data/access/lat/
BackgroundModels.html

In our analysis, we removed 3FGL J0425.8+5600  $_{206}$  and 3FGL J0423.5+544 from the RoI, but kept 3FGL  $_{207}$  J0426.7+5437 in the model since preliminary analyses  $_{208}$  showed clear positive residual emission at the position of  $_{209}$  the source if it was removed from the RoI. Figure 2 shows  $_{210}$  a residual TS map for the region around G150.3+4.5.  $_{211}$  This point source detection-significance map was created  $_{212}$  by placing a point source modeled with a power law of  $_{213}$  photon index  $\Gamma$  = 2 at each pixel and gives the significance of detecting a point source at each location above  $_{215}$  the background.

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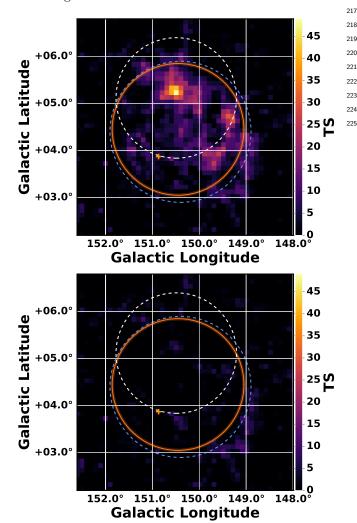


Figure 2. Background subtracted residual TS map above 1 GeV with  $0.1^{\circ}$ x  $0.1^{\circ}$  pixels, centered on SNR G150.3+4.5. The orange circle and translucent shading show the fit disk radius and  $1\sigma$  errors, respectively, for the extended source, the orange cross shows the position of 3FGL J0426.7+5437 (included in the background model), blue dashed circle is the extent of the radio SNR, and white dashed circle depicts 2FHL J0431.2+5553e. Bottom map includes G150.3+4.5 in the background model, top does not.

We modeled the excess emission in the direction of <sup>234</sup> G150.3+4.5 with a uniform intensity, radially-symmetric <sup>235</sup> disk, simultaneously fitting the spatial and spectral com- <sup>236</sup> ponents of the model via pointlike. The extension of <sup>237</sup> the disk was initialized with a seed radius of  $\sigma = 0.1^{\circ}$  and <sup>238</sup> position centered on the radio position of G150.3+4.5. <sup>239</sup> We define the significance of extension as in Lande <sup>240</sup> et al. (2012); TS<sub>ext</sub> =  $2 \log(\mathcal{L}_{ext}/\mathcal{L}_{ps})$ , with  $\mathcal{L}_{ext}$  be- <sup>241</sup>

ing the likelihood of the model with the extended source and  $\mathcal{L}_{ps}$  that of a point source located at the peak of emission interior to the extended source. For the disk model we found that  $TS_{ext} = 298$ , for the best fit radius,  $\sigma = 1.40^{\circ} \pm 0.03^{\circ}$ , and position, R.A. = 55.46° ± 0.03°, DEC. =  $66.91^{\circ} \pm 0.03^{\circ}$ , all in excellent agreement with the radio SNR size and centroid determined in Gao & Han (2014). Figure 3 shows radially integrated counts for the region as a function of angular radius squared. It's clear from this figure that there is significant excess of counts above the Galactic diffuse radiation in this region that is adequately modeled by a symmetric disk. We tried adding back in to our model the two removed 3FGL sources but both were insignificant when fit on top of the best fit disk. The bottom map in Figure 2 is a residual TS map of the same region as the top map, but with the disk source included in the background model, demonstrating that the disk can account well for the emission in the region and justifying the exclusion of the two aforementioned 3FGL sources.

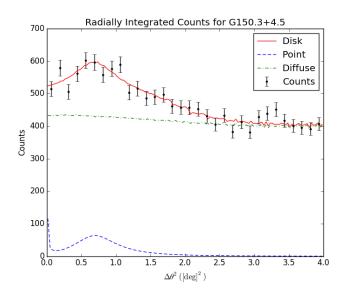


Figure 3. Radially integrated counts map centered on the GeV emission coincident with G150.3+4.5. Red line shows the expected counts for a uniform intensity disk with radius,  $\sigma = 1.40^{\circ}$ , green line is that of the Galactic diffuse background. Error bars on data points are statistical. [JAM: Replot without the point model? I feel like it's distracting and doesn't really add anything.]

The morphology of the radio emission is suggestive of an elliptical or ring morphology, so both of these spatial models were tested as well. For the ring model, the fit reduced to a disk with parameters matching those stated above. Using the elliptical model showed a weak improvement over the radially symmetric model at the  $2.6\sigma$  level ( $\Delta TS = 9$  with two additional degrees of freedom), which we did not consider significant enough to say the GeV emission had an elliptical morphology (see Table 1). For the remainder of this study, we only considered the disk spatial model.

2FHL J0431.2+5553e is the extended source detected in the 2FHL catalog found to be overlapping the northern region of G150.3+4.5 Ackermann et al. (2016). The source has a power law spectral index  $\Gamma = 1.66 \pm 0.2$ , and disk radius  $\sigma = 1.27^{\circ} \pm 0.04^{\circ}$  (see Figure 2). When

	Spatial Model	$TS_{ext}$	$TS^a$	$\sigma$ [°]	R.A. [°]	DEC [°] <sup>287</sup>
-	Disk Elliptical Disk 2FHL (free) <sup>b</sup> 2FHL (fixed)	278.843 189.048 260.317 260.317	-32.850 3.033 -3.277 -3.277	49.80 398.64 63.87 63.87	LMC IC 443 Puppis A Puppis A	gal <sub>289</sub> snr <sub>290</sub> snr <sub>291</sub>
	Disk & 2FHL <sup>b</sup>	260.317	-3.277	63.87	Puppis A	snr <sup>292</sup>

Note. — 2FHL (free) corresonds to the model where a disk match ing 2FHL J0431.2+5538e, was included in the likelihood model and the spectral and spatial parameters we free to vary. For the 2FHL (fixed) model, 2FHL J0431.2+5553e was included with spatial parameters fixed In the Disk & 2FHL model, we included both the best-fit disk determined in §2.2, fixed in position and size, and added a source resembling the 2FHL J0431.2+5553e with free spectral and spatial parameters. This model reports the fit values of 2FHL J0431.2+5553e [JAM: Hayen't filled in real numbers for G150 yet just copied from 2FHL table.] [JAM: Other things to add  $N_{dof}$ , LL, spectral params?]

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comparing the best fit extension of the 2FHL source with the result from this paper, factoring in the uncertainty in both extension and position, we see that the  $> 50~{\rm GeV}$  and  $> 1~{\rm GeV}$  results are not incompatible. It is likely that the paucity of events above 50  ${\rm GeV}$  is the cause of the smaller fit radius, as opposed to the difference arising from the effects of an energy dependent morphology. To explore the connection between the 2FHL and above 1  ${\rm GeV}$  emission, we tested a few other spatial hypotheses.

First, we replaced the  $\sigma = 1.40^{\circ}$  disk with an another disk matching the spectral and spatial parameters of 2FHL J0431.2+5553e and calculated the likelihood with this new source's position and extension fixed. For this hypothesis, we find  $TS_{ext} = 165$ , and TS = 226, demonstrating that the fixed disk matching the 2FHL source is clearly disfavored over the previously determined best fit disk at this energy. Our next test consisted of placing a second extended source on top of the best fit disk detected above 1 GeV. We added a source, initially matching the spatial and spectral parameters of 2FHL J0431.2+5553e, to our source model of the region (in addition to the  $\sigma = 1.40^{\circ}$  disk), and fit its spectrum and extension. Fitting a second extended source in this region serves two purposes: 1. it acts as a check on whether there was residual emission unaccounted for by the pre- 297 viously best-fit disk, and 2. it allows us to determine if 298 the best fit disk can be split into two spectrally distinct, components. This fit resulted in the source wandering north (but still partially overlapping G150.3+4.5) and having an insignificant extension,  $TS_{ext} = 4$ . Details on the spatial parameters are given in Table 1.

#### 2.3. Spectral Analysis

After determining the best fit morphology with  $^{303}$  pointlike for the GeV emission coincident with SNR  $^{304}$  G150.3+4.5, we used those results as a starting point for our gtlike maximum-likelihood fit of the region to es- $^{306}$  timate the best spectral parameters for our model. The  $^{307}$  LAT data is well described by a power law from 1 GeV to  $^{308}$  1 TeV with a photon index,  $\Gamma = 1.82 \pm 0.04$ , and energy  $^{309}$  flux above 1 GeV of  $(7.3 \pm 0.72) \times 10^{-11}$  erg cm<sup>-2</sup> s<sup>-1</sup> and TS = 389 [JAM: pointlike results were index = 1.80  $^{311}$  flux =  $(7.17 \pm 0.73 \times 10^{-11})$  erg cm<sup>-2</sup> s<sup>-1</sup>]. We tested  $^{312}$ 

the  $\gamma$ -ray spectrum of the extended disk for spectral curvature using a log-normal model (Log Parabola), and find no significant deviation from a power law ( $\Delta$ TS  $\sim$  1). Figure 4 shows the best-fit power law spectral energy distribution for the GeV source whose morphology was described in Section 2.2. Spectral data points were obtained by dividing the energy range into 12 logarithmically spaced bins and modeling the source with a power law of fixed spectra index,  $\Gamma=2$ . We over plotted the SED of 3FGL J0426.7+5437, to demonstrate how the spectra of the two sources would grow more confused at energies below 1 GeV [JAM: Need to remake Fig 1 with J0426].

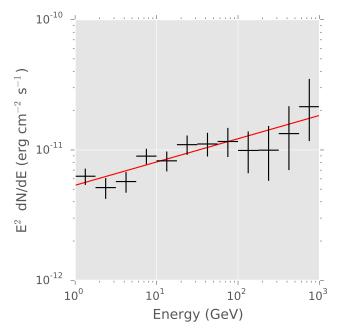


Figure 4. Spectral energy distribution for the extended source coincident with SNR G150.3+4.5 from 1 GeV to 1 TeV. Red line corresponds to the best fit power law model. Crosses are shown with with statistical error bars. Grey dashed line is the SED of 3FGL J0426.7+5437, modeled with an exponential cut-off power law. [JAM: need to remake this and actually add J0426]]

## 3. MULTIWAVELENGTH OBSERVATIONS AND ANALYSIS

 $3.1.~H_2$ 

[JAM: Jack is working on this]

3.2. *X-ray* 

[JAM: Dan is working on this.]

#### 4. DISCUSSION AND RESULTS

## 4.1. SNR or PWN?

The follow-up observations of the  $\gamma$ -ray emission in the direction of G150.3+4.5, presented here, of the source detected above 50 GeV in 2FHL have led to the detection of an extended  $\gamma$ -ray source whose centroid and radius match extremely well with those of the radio detected SNR. The broad size of the extended source and correlation with the radio shell leave few plausible scenarios for the nature of the GeV emission. Namely, the GeV emission can arise from the wind nebula of the putative

<sup>&</sup>lt;sup>a</sup> Calculated in gtlike

<sup>&</sup>lt;sup>b</sup> Started with disk matching the spectral/spatial parameters of 2FHL J0431.2+5553e, then left them free to fit in the likelihood model.

pulsar of G150.3+4.5 or the GeV emission corresponds  $_{376}$  to  $\gamma$ -rays produced in the SNR. We argue here that the  $_{377}$  SNR is favored over a pulsar wind nebulae (PWN) as the  $_{378}$  generator of the observed  $\gamma$ -rays.

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The first problem with the PWN hypothesis is that 380 there is no pulsar candidate detected near the centroid of the SNR to power a PWN. While 3FGL J0425.8+5600 is the closest  $\gamma$ -ray source to the cen-  $^{381}$ ter of the remnant, it does not have a pulsar-esque spectrum, it lies about 0.25° away, and we showed in §2.2 that with the best-fit disk hypothesis, neither 383 3FGL J0425.8+5600 nor 3FGL J0423.5+5442 are sig- 384 nificant in the likelihood model of the region. 3FGL 385 J0426.7+5437, with a spectrum reminiscent of a pul- 386 sar, may actually be one, but as discussed previously, 387 Barr et al. (2013) detect no pulsations from the source 388 [JAM: something to say about SWIFT nondetection?]). 389 Furthermore, the source is 0.84° away from the cen- 390 troid of G150.3+4.5. Typical pulsar ballistic velocities 391 range from  $V_{\rm PSR} \sim 400-500~{\rm km~s^{-1}}$ , with extreme velocities exceeding 1000 km s<sup>-1</sup> (Gaensler & Slane 2006). 393 If 3FGL J0426.7+5437 was the compact remnant of the 394 progenitor star that birthed G150.3+4.5, it would have 395 to be traveling with a velocity,  $V_{PSR} = 1125 \text{ km s}^{-1}$  (assuming an age of 5 kyr, which we derive in the following 397 section,  $\S4.2$ ), and would make it one of the fastest known 398 pulsars (Chatterjee et al. 2005). While possible, this scenario is unlikely without further evidence to support such 400 a high velocity. [JAM: Fastest pulsar (till 2011 at least) 401 1100 km/s, more recent ref?

Another argument disfavoring the PWN scenario is 403 that, despite the hard  $\gamma$ -ray spectral index extending to 404 TeV energies, ROSAT X-ray observations detect no significant emission suggestive of a PWN in the direction 406 of G150.3+4.5 (see §3.2). Typical PWNe spectral indices range from about  $-0.3 \lesssim \alpha \lesssim 0$  (Gaensler & Slane 408 2006). The radio spectral index as determined in Gao & 409 Han (2014) ( $\alpha = 0.4 \pm 0.17$  for part of the eastern shell, 410  $\alpha = 0.69 \pm 0.24$  for a region in western shell) suggests 411 that the radio object is likely not a PWN.

Many of the arguments disfavoring the PWN hypothesis in fact bolster that of SNR. First and foremost in favor of an SNR origin for the  $\gamma$ -ray emission is the excellent agreement between the GeV best-fit disk radius and centroid with that of the radio shell. The radio shell-like appearance, non-thermal radio spectrum, and strands of red optical filamentary structures led both Gao & Han (2014) and Gerbrandt et al. (2014) to regard the radio spectral source an SNR as opposed to a PWN. The radio spectral londex, while not quite in line with typical PWN spectra, list actually common of SNRs.

While the above factors lend credence to an SNR ori-  $_{423}$  gin for the GeV  $\gamma$ -rays the PWN scenario can not be  $_{424}$  ruled out due to the lack of an associated pulsar. Re-  $_{425}$  gardless, for the remainder of this study, we assumed the  $_{426}$  observed  $\gamma$ -rays were produced in the shock front of SNR  $_{427}$  G150.3+4.5

## 4.2. G150.3+4.5 in Context

Having associated the  $\gamma$ -ray emission with G150.3+4.5,  $_{431}$  next, we assessed the evolutionary state of the remnant  $_{432}$  to place it in context within the current population of  $_{433}$  LAT SNRs. Using the most viable HI kinematic dis-  $_{434}$  tance, d  $\approx 0.38$  kpc derived in §3.1, we showed that the  $_{435}$ 

projected radius of G150.3+4.5 is  $R \approx 9.4$  pc. Employing a standard Sedov-Taylor solution for the expansion of a blast wave, we estimated the age of G150.3+4.5. In the Sedov phase, the radius of the shock front is given by.

$$R_{ST} = 0.314 \left(\frac{E_{51}}{n_0}\right)^{1/5} t_{\rm yr}^{2/5} \text{pc}$$
 (1)

[JAM: None of this is taking into account the low density upper limit. give Age for a few n's] Where  $E_{51}$  is the kinetic energy output of the supernova in units of  $10^{51}$  erg, and  $n_0$  the ambient density the shock is expanding into in units of cm<sup>-3</sup>. Assuming standard values of 1 for  $E_{51}$  and  $n_0$  we solved equation 1 for  $t_{yr}$  (the current age of the remnant in years) and used the value of R derived for G150.3+4.5 to estimate the age of the SNR as  $t \approx 4.9$  kyr.

Figure 5 shows the SED of G150.3+4.5 overlaid on the spectra of a selection of other LAT observed SNRs with ages ranging from  $\sim 10^3 - 10^4 \text{yr}$ . G150.3+4.5 exhibits a hard spectrum extending to TeV energies with no spectral break (breaks are commonly seen in LAT SNRs interacting with nearby molecular material (Hewitt & Lemoine-Goumard 2015)) and appears spectrally similar to the younger SNRs like RX J1713.7-3946 and RX J0852.0-4622. In figure 6, we plotted the luminosity of several LAT SNRs against their squared diameters (a proxy for age, as evident from equation 1). [JAM: should be using SNR cat fig 18?. need to reword things if I'm just taking their figure and putting my points on]. Similarly, with it's low luminosity [JAM: give L here. If I use the W41 LvsD2 plot, 100MeV-100 GeV, if SNR cat, 1-100 GeV, G150.3+4.5 appears to correlate well with the younger sect of LAT SNRs. Our age estimate alone does not unambiguously determine the evolutionary state of G150.3+4.5. However, when combined with the results of Figures 5 and 6 comparing G150.3+4.5 to the population of other LAT SNRs, it indicates that G150.3+4.5 is more compatible with a dynamically unevolved, noninteracting (with the surrounding interstellar medium) stage of expansion.

## 4.3. Nonthermal Modeling

SNR shock fronts are known accelerators of cosmic rays to very high energies. There are potentially multiple radiation mechanisms operating at the shock that produce GeV  $\gamma$ -rays. Accelerated electrons can give rise to inverse Compton (IC) emission via upscattering of ambient cosmic microwave background (CMB), stellar, and IR photon fields, as well as non-thermal bremsstrahlung radiation. Energetic protons can collide with ambient protons in the surrounding, producing neutral pions which decay into  $\gamma$ -ray photons.

To infer the properties of the underlying relativistic particle populations in the SNR environment, it is vital to understand the origin of the observed  $\gamma$ -ray emission detected from G150.3+4.5. To do so, we employ the naima Python package. naima is an open-source code base that computes the non-thermal radiation from a relativistic particle population (Zabalza 2015). It utilizes known parameterizations and analytic approximations to the various non-thermal processes (i.e., synchrotron, IC, bremsstrahlung, and pion decay emission), which results

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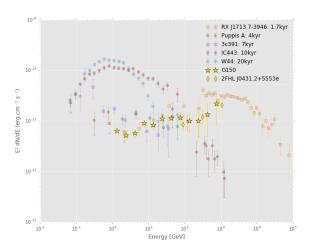
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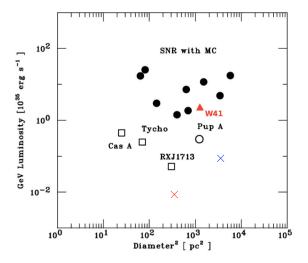
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**Figure 5.** SEDs for several LAT observed SNRs with ages spanning  $\sim 10^3 - 10^4 \mathrm{yr}$ . The GeV spectrum of G150.3+4.5 is shown as stars. [JAM: Replot with white bkg, bigger font, lines, change colors/shapes? get rid of 2FHL, less, different SNRs? I need refs for each]



**Figure 6.** Luminosity of several LAT SNRs plotted against their <sup>471</sup> [JAM: radio? GeV?] diameter square. [JAM: taken from W41 paper, overplotted G150] [JAM: Should I actually remake this myself or is it ok to just use the one from the other paper with my points on it? Add more text when I settle on a plot.]

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in the calculations being computationally inexpensive. 476 naima also makes use of emcee, a Markov chain Monte 477 Carlo (MCMC) ensemble sampler for Bayesian parameter estimation (Foreman-Mackey et al. 2013). The sam- 479 pler is used to find the best-fit parameters of the radiative 480 models to the observed photon SED for a given particle 481 distribution function.

To determine the best fit parameters, naima calls 483 emcee to sample the log-likelihood function (i.e., the 484 likelihood of the observed data given the assumed spectrum) of the radiative model. [JAM: should I include what 486 the likelihood function looks like here?] The radiative 487 models require as input a particle distribution function 488 to model the present-age electron or proton spectrum. 489 We used a one-zone, homogeneous particle distribution 490 model (which naima inherently assumes) and scaled the 491

Table 2 Naima Model Best Fit Parameters

s	$K_{\mathrm{ep}}$	$A_{p}$	Ba	$E_{cutoff(e)}$	$E_{cutoff(p)}$			
Fixed $K_{ep} = 0.01$								
$1.5 \pm 0.2$	0.01	-32.850	49.80	$_{ m LMC}$	2			
$\frac{\text{Fixed K}_{\text{ep}}}{1.5 \pm 0.2}$		-3.277	63.87	Puppis A	2			
$\frac{\text{Fixed K}_{\text{ep}}}{1.5 \pm 0.2}$		-3.277	63.87	Puppis A	2			
$\frac{\text{Fixed s}}{2 \pm 0.2}$	1	-3.277	63.87	Puppis A	2			

Note. — [JAM: add better caption] Results from naima model? Right now the free params are index, kep, eleccut, protcut, B. Fixed are nh, all the IC photon field values distance (this is just for determining flux) [JAM: Correct values aren't in yet][JAM: ]add units to params

likelihood function by a uniform prior probability distribution. For this work, we model the separate proton and electron and spectra as power laws with an exponential cut off,

$$\frac{dN}{dE_{(e,p)}} = A_{(e,p)} (E/E_0)^{-s} \exp\left(\frac{-E}{E_{\text{cutoff }(e,p)}}\right) \qquad (2)$$

where E is the particle energy,  $E_0$  the reference energy, s the spectral index, and  $E_{\rm cutoff}$  the cutoff energy. The electron distribution's normalization is related to the proton normalization through the electron-to-proton ratio scaling factor,  $A_e = K_{ep}A_p$ . We also assumed that the electron and proton distributions have the same spectral shape.

For our radiation models, we assumed a gas density,  $n_0=1~{\rm cm}^{-3}$  for proton-proton [JAM: and bremss when I get it working] interactions For IC emission, we include CMB ( Talk about free/ fixed params of the model, reference the table, and figure to show best fit, discuss results and what the fits imply regarding lep/had dom and energy in e- p.

[JAM: Used radio SED from (Gerbrandt et al. 2014)] [JAM: estimates min of negative loglike through sambling]

[JAM: Say something about what the radio index is and the connection to the gev index]

[JAM: Discuss implications of the naima fits. Do they show preference for lep/had? suggest something about total energy content in particles?]

[JAM: Another section here? can I just put a paragraph about what future observations will elucidate on in the conclusions?]

[JAM: looks like young SNRs, but no x-ray and lower luminosity than most other LAT SNRs. It's possibly just that ROSAT is no sensitive enough.]

[JAM: Deeper x-ray observations to search for the compact stellar remnant. ROSAT all-sky not that sensitive, dedicated X-ray to search for thermal/nonthermal components]

[JAM: What can be done in TeV? The difficulty pointed TeV observations are that it might be difficulty for them to detect such broadly extended emission (why?

a Calculated in gtlike

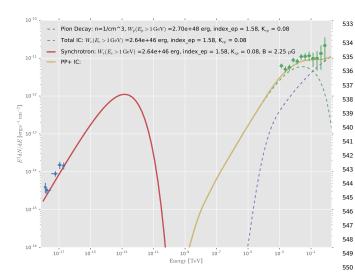


Figure 7. G150.3+4.5. Naima SED [JAM: bigger font, get rid of text for each line, better colors

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I know they'd have to tile their observations, but there's 554 something inherently difficulty for them about observing large extended sources. Is it just that the emission is spread out so it might be faint and below the detection threshold?). What about HAWC? Why does it not de- 558 tect the emission that Fermi clearly shows extending to 559 VHE energies? Is it not as sensitive at this energy for 560 some reason? HAWC energy range extends to 100 GeV.]

#### 5. CONCLUSIONS

We analyzed 7 years of Fermi-LAT data in the direction of SNR G150.3+4.5, lowering the energy threshold from that previously reported in the 2FHL catalog, and report detection of significantly extended  $\gamma$ -ray  $^{567}$ emission coincident with the entirety of the radio remnant's shell. We find the emission from 1 GeV to 1 TeV  $^{569}\,$ to be well described by a power law of spectral index  $^{570}$  $\Gamma = 1.82 \pm 0.04$ , with morphology consistent with a uniform disk with best-fit radius,  $\sigma = 1.40^{\circ} \pm 0.03^{\circ}$ . Based on radio and  $\gamma$ -ray properties of emission in the direction of G150.3+4.5, within the context of the current  $^{574}$ LAT SNR population, we argued that the GeV emission  $^{575}$ likely originates in the shock of G150.3+4.5, and disfavor a PWN origin. To estimate the distance to the SNR, we obtained an HI spectrum toward G150.3+4.5 from the Leiden/Argentine/Bonn survey of Galactic HI. Cal-  $^{576}\,$ culating distances from the derived HI velocity peaks, 578 we showed that the most reasonable distance estimate 579 places G150.3+4.5 at a distance of d = 0.4 kpc, making 580 it one of the closest known SNRs detected by the LAT. Using this distance and a standard Sedov-Taylor SNR  $^{582}$ evolution model, we estimate the age of the G150.3+4.5  $^{\circ\circ}_{584}$ to be t  $\sim 5$  kyr. Say something about X-ray once I know <sub>585</sub> itTo assess the underlying particle population acting in 586 G150.3+4.5 we use the naima Python package to fit the 587 observed radio and  $\gamma$ -ray SED to non-thermal electron 588  $_{589}^{589}$ and proton radiation models. We find that blah, which 590 suggests more blah. Something about how G150 fits in 591 with other LAT detected SNRs based on age, spectrum 592 luminosity, spectral modeling. End with what further 593 observations can get us. 595

[JAM: thanks?]

#### 6. SCRATCH

 $L_{\gamma} = 1.3 \times 10^{33} \text{ erg s}^{-1} \text{ from 1 GeV to 1 TeV for best}$ d and flux above

energy flux from 100 MeV100 to GeV:

 $4.84 \times 10^{-11} {
m erg \ cm^{-2} \ s^{-1}}$   $L_{\gamma} = 8.6 \times 10^{32} {
m erg \ s^{-1}}$  from 100 MeV to 100 GeV for best d and flux in same range

energy flux from 1 GeV

 $3.83 \times 10^{-11} \text{ erg cm}^{-2} \text{ s}^{-1}$   $L_{\gamma} = 6.8 \times 10^{32} \text{ erg s}^{-1} \text{ from 1 geV to 100 GeV for}$ best d and flux in same range

For diamMax = 60pc, dmax = 1.22kpc, and Lmax (100 mev-100 GeV) = 8.7 e+33

JAM: One potential scenario is that the  $\gamma$ -rays in this region are produced by a pulsar wind nebula (PWN) generated by the putative pulsar of SNR G150.3+4.5.

[JAM: for synchrotron  $\alpha = (1-s)/2$ , where  $\alpha$  is the radio spectral index, and s the electron distribution power law index. Same for IC below break?]

JAM: SNR cat figure 8 suggests there are only 4(ish) SNRs with an index less than 2

[JAM: eastern shell (Jack called this overall) radio in $dex \alpha = 0.4 \pm 0.17$  Gao & Han (2014), but  $\alpha = 0.69 \pm 0.24$ for the western

JAM: For energies below the high energy break, For pion and bremss  $\Gamma = 2\alpha + 1$  (says SNR cat) For IC,  $\Gamma = \alpha + 1$  ) for positive  $\alpha$ 

[JAM: From Gaensler & Slane (2006) Typical indices for PWNe are  $\sim -0.3 \lesssim \alpha \lesssim 0$  in the radio band, and  $(\Gamma \approx 2)$  in the X-ray band. So  $\alpha$  is not inconsistent, but at the boundary.

JAM: For puppis A paper, why did they use particle index = gam photon index?

JAM: cr abundances at earth kep = 0.01 (Hillas 2005). I think in the puppis A paper the use 0.02?

[JAM: Sooo, my index is consistent with either?]

JAM: how off can dist be? SNR catalog says 8 SNRs are withing 1.5 kpc and have some kind of classification in the catalog. These are (closest first) Vela, cygnus loop , Vela Jr, RX J1713, G073, S147, IC443, Monoceros loop. if 0.4 kpc is correct for G150, it's the second closest LAT detected SNR. Even at 1.5 kpc it would be within top

# REFERENCES

Abdo, A. A., et al. 2013, ApJS, 208, 17 2.1 Acero, F., et al. 2015, ApJS, 218, 23 2.1 2016, ApJS, 224, 8 1 Ackermann, M., et al. 2012, ApJS, 203, 4 2.1 2013, Science, 339, 807 1 -. 2016, ApJS, 222, 5 1, 2.1, 2.2 Atwood, W., et al. 2013a, ArXiv:1303.3514 1 Atwood, W. B., et al. 2009, ApJ, 697, 1071 2.1 -. 2013b, ApJ, 774, 76 Barr, E. D., et al. 2013, MNRAS, 429, 1633 2.2, 4.1 Chatterjee, S., et al. 2005, ApJL, 630, L61 4.1 Condon, J. J., Cotton, W. D., Greisen, E. W., Yin, Q. F., Perley, R. A., Taylor, G. B., & Broderick, J. J. 1998, AJ, 115, 1693 2.2 Foreman-Mackey, D., Hogg, D. W., Lang, D., & Goodman, J. 2013, PASP, 125, 306 4.3 Gaensler, B. M., & Slane, P. O. 2006, ARA&A, 44, 17 4.1, 6 Gao, X. Y., & Han, J. L. 2014, A&A, 567, A59 1, 2.2, 4.1, 6 Gerbrandt, S., Foster, T. J., Kothes, R., Geisbüsch, J., & Tung, A. 2014, A&A, 566, A76 1, 4.1, 4.3

Hewitt, J. W., & Lemoine-Goumard, M. 2015, Comptes Rendus Physique, 16, 674 4.2

Jogler, T., & Funk, S. 2016, ApJ, 816, 100 1
 Kerr, M. 2010, PhD thesis, University of Washington,
 arXiv:1101.6072 2.1

Lande, J., et al. 2012, ApJ, 756, 5
Reynolds, S. P. 2008, ARA&A, 46, 89
Zabalza, V. 2015, ArXiv e-prints
4.3