

Modeling Stability of Magnetic Fields in Magnetars

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Background

Neutron Stars - Remnant core of massive star

- Composed almost completely of neutrons
- Core densities $\sim 10^{15} \text{ g cm}^{-3}$

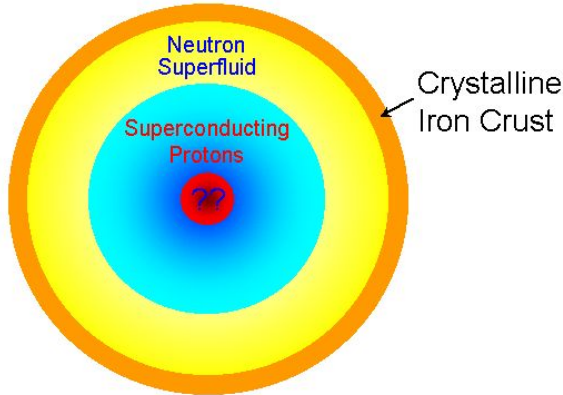


Image credit: Ohio State University

Pulsars - Rotating neutron star with radio beam and B-field

- B-field $\sim 10^{12} - 10^{13}$ gauss
- Average rotational period: $\sim 0.7 \text{ s}$

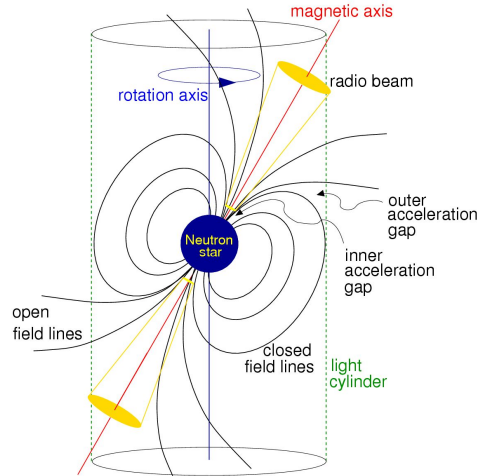


Image credit: NRAO

Magnetars - Highly magnetic pulsars

- B-field $\sim 10^{15}$ gauss, *strongest in universe*
- Fields alter structure of star via Lorentz force.

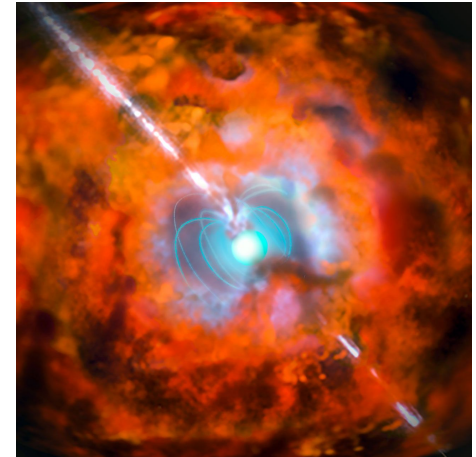
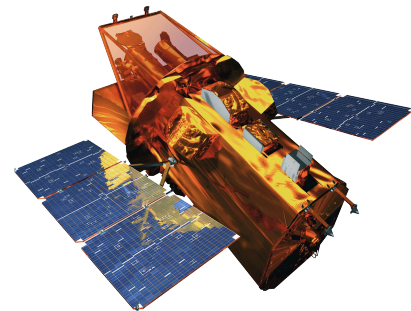


Image credit: ESO

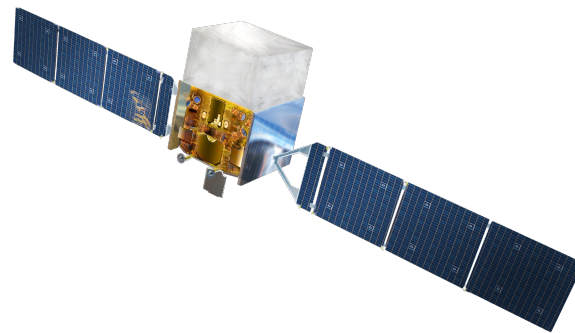
Why Study Magnetars?

- Rapidly growing field: Rate of detection dramatically increased following launch of Swift and Fermi space telescopes.
 - ◆ First detection in 1979.
- Potential sources of gravitational waves (GWs):
 - ◆ If the magnetic axis is not aligned with the rotational axis, deformation due to B-field will not rotate symmetrically, causing a time-varying quadrupole, forming GWs.
- Detection of GWs from Magnetars would inform, set boundaries for models of neutron star composition.



Swift Space Telescope: Launched November, 2004

NASA E/PO, Sonoma State University/Aurore Simonnet



Fermi Gamma-ray Space Telescope:
Launched June, 2008

NASA E/PO, Sonoma State University/Aurore Simonnet

Research Goals

- Model stellar structure:
 - Describe how state variables such as pressure vary
 - Achieve stable equilibrium between stellar pressure and force of gravity
- Model magnetic field structure:
 - Verify stability of particular model configurations
 - Show stability is maintained through evolution in time
- Verify that additional phenomena within simulations have physical justifications.

Stellar Model: Structural Equations

→ Hydrostatic Equilibrium: $\frac{dP}{dr} = -G \frac{M_r \rho}{r^2}$ *

→ Polytropic equation of state (EOS): $P(\rho) = K \rho^\gamma$

◆ Determine γ via choice of polytropic index 'n' via the relationship: $\gamma = \frac{n+1}{n}$

→ The choice of $n = 1$ suits density profile of neutron stars well.

Based on our choice of EOS, we can completely describe stellar structure via the following equations

$$P_{n=1}(r) = K \rho(r)^2 \quad \rho(r) = \rho_c \frac{\sin(\pi r/R) R}{r \pi}$$

$$r < R$$

ρ_c is core density $\sim 10^{15} \text{ g cm}^{-3}$

* M_r , known as the *interior mass*, is the mass interior to the sphere of radius r

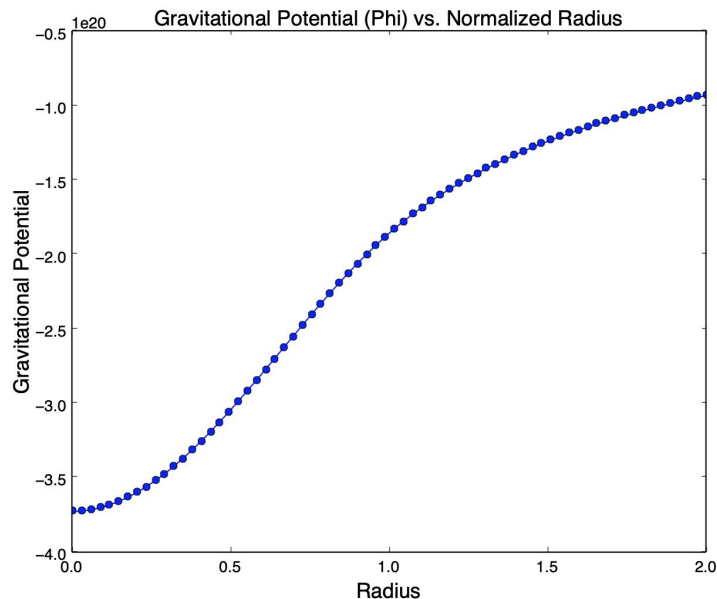
Stellar Model: Gravitational Potential

- Hydrostatic equilibrium requires a balancing force opposing gravity. We seek solutions to the spherically symmetric form of Poisson's equation for gravitational potential energy:

$$\frac{1}{r^2} \frac{\partial}{\partial r} \left(r^2 \frac{\partial \varphi}{\partial r} \right) = 4\pi G \rho$$

- $\varphi(r)$ in three equations:
 - $\varphi(r = 0)$
 - $\varphi(r)$ for $0 < r < 1.0$
 - $\varphi(r)$ for $r \geq 1.0$

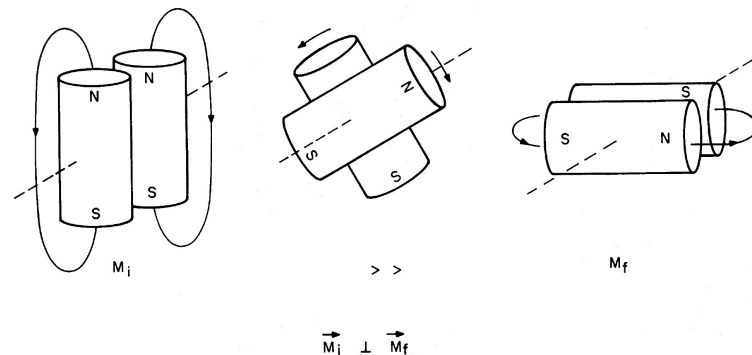
- We require $\varphi(r)$ to be piecewise continuous across the computational domain:



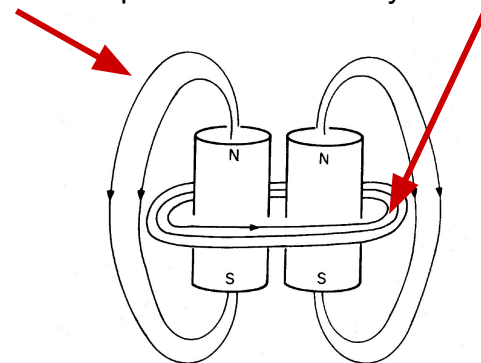
Magnetic Field Structure

- Flowers & Ruderman (1977):
Stable field configurations require poloidal and toroidal field components.
 - Poloidal and Toroidal components act in tandem; when one wanes the other strengthens, allowing dynamic stability despite perturbations.
- Toroidal component is non-zero only interior to $r = R_{\text{star}}$, Poloidal component exists everywhere.

Unstable poloidal configuration*



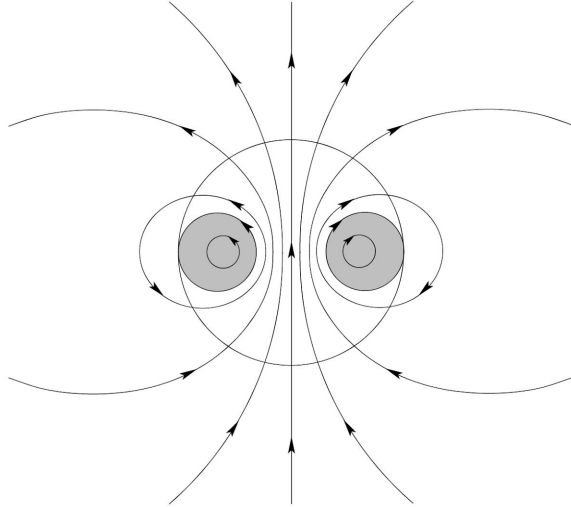
Poloidal component stabilized by toroidal field*



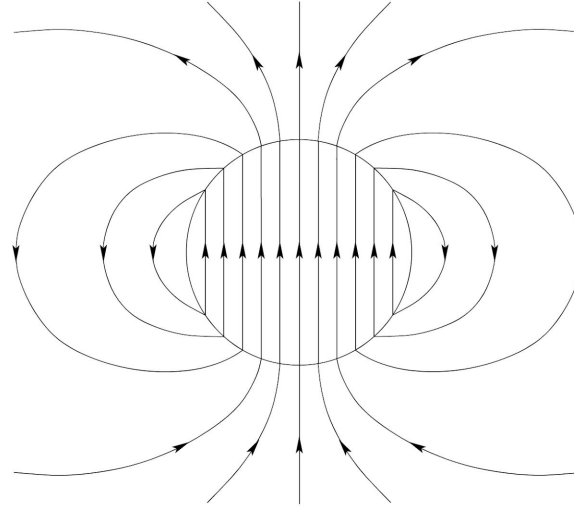
Magnetic Field Model

- Via Haskell et al. 2008, specific to the $n = 1$ polytrope model
- Stable as mixed poloidal-toroidal field

Stable poloidal-toroidal configuration



Unstable pure poloidal configuration



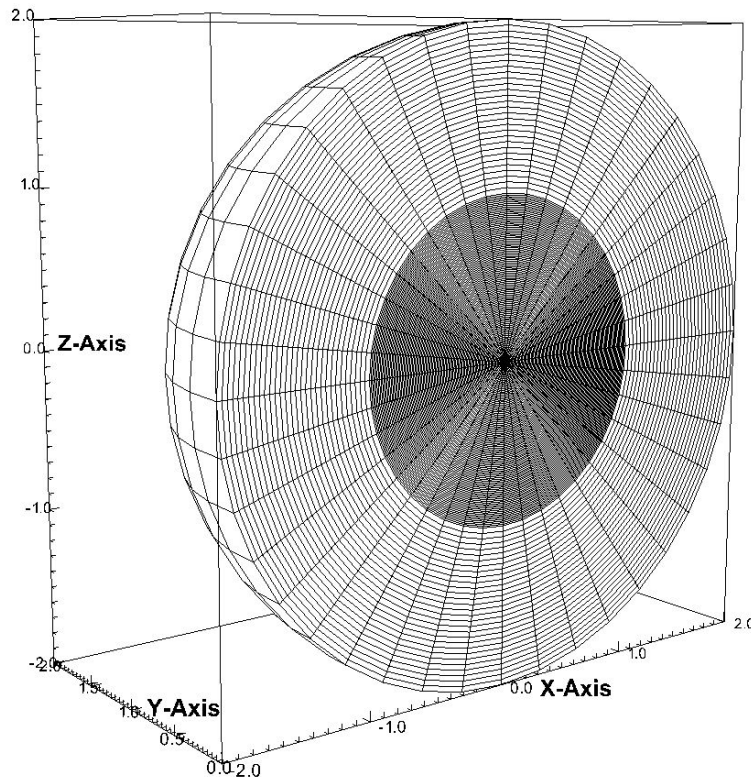
Methods

Software:

- PLUTO - Computational package for evolving computation
- VisIt - Open source visualization and graphical analysis software

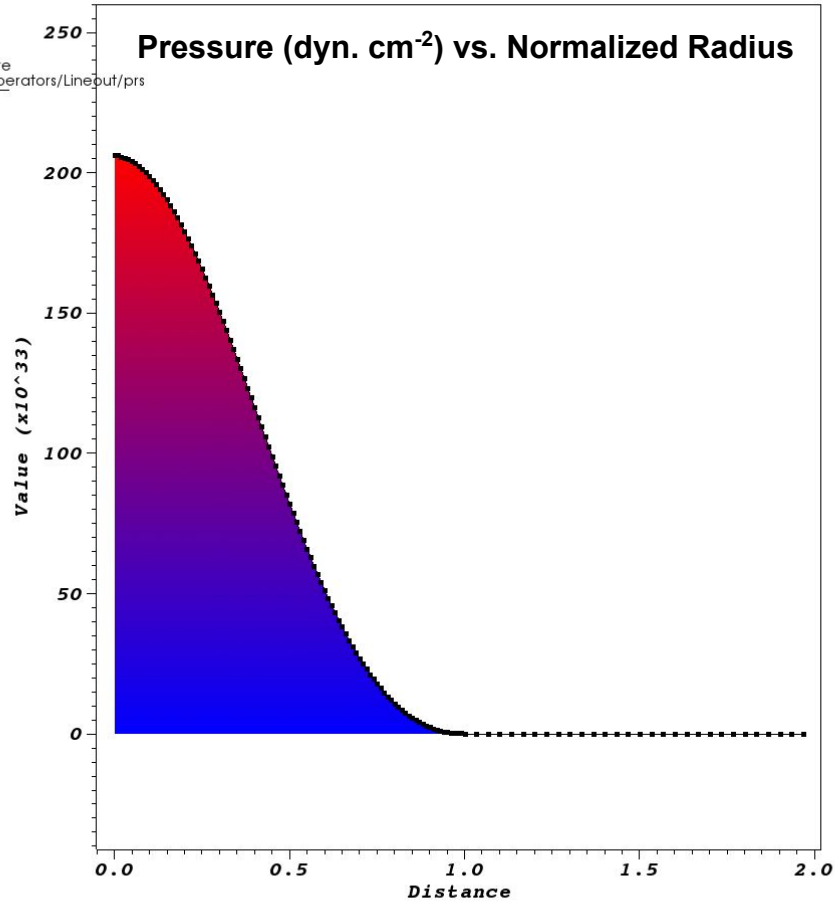
The Computational Domain:

- Spherical domain for evolving simulations.
- 100 equally spaced points for $0 \leq r \leq 1.0$, 30 equally spaced points for $1.0 < r \leq 2.0$
- 20 equally spaced points in both the domains of Θ and ϕ .

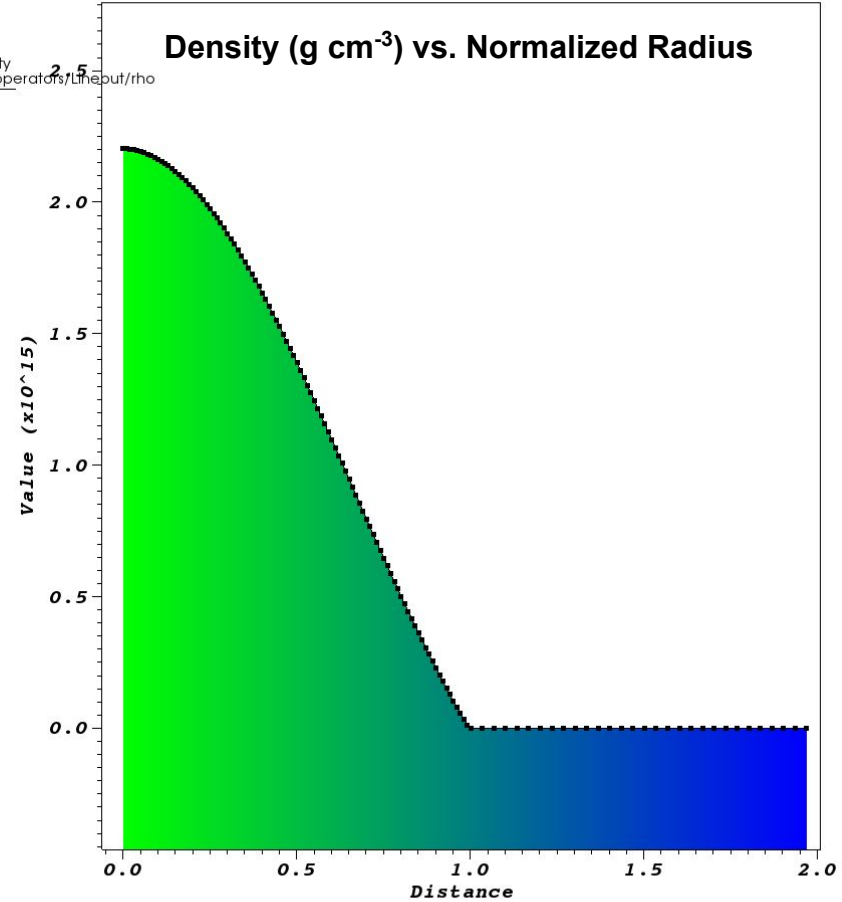


A hemisphere of the computational domain

Pressure (dyn. cm⁻²) vs. Normalized Radius

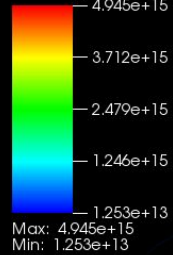


Density (g cm⁻³) vs. Normalized Radius



DB: data.0049.vtk
Cycle: 49 Time: 0.0489639

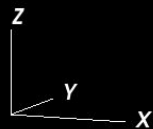
Pseudocolor
Var: Magnetic_Field - Speed



Toroidal Field

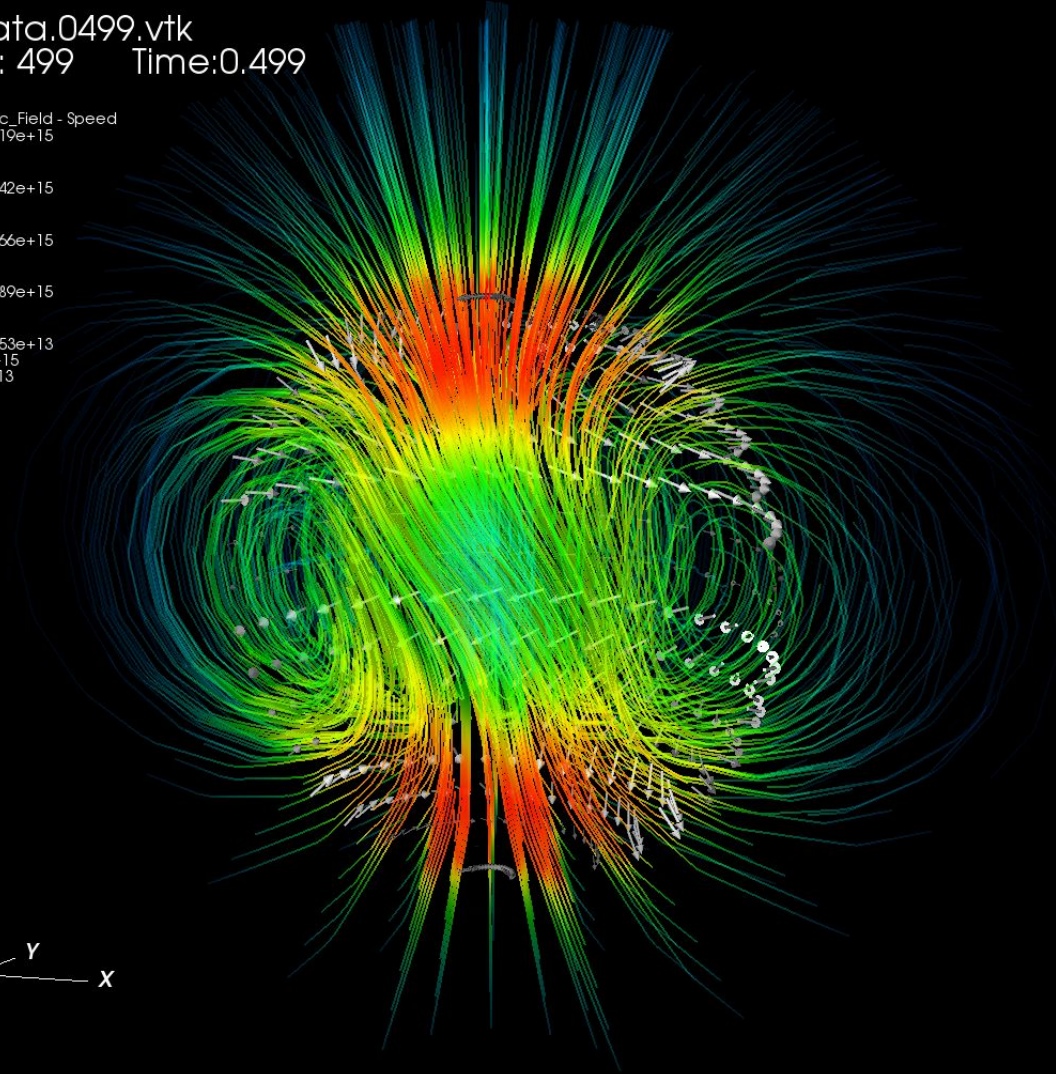
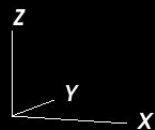
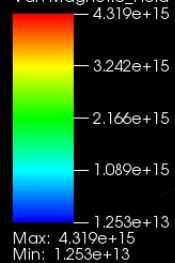
Poloidal Field

Velocity Vector
Field



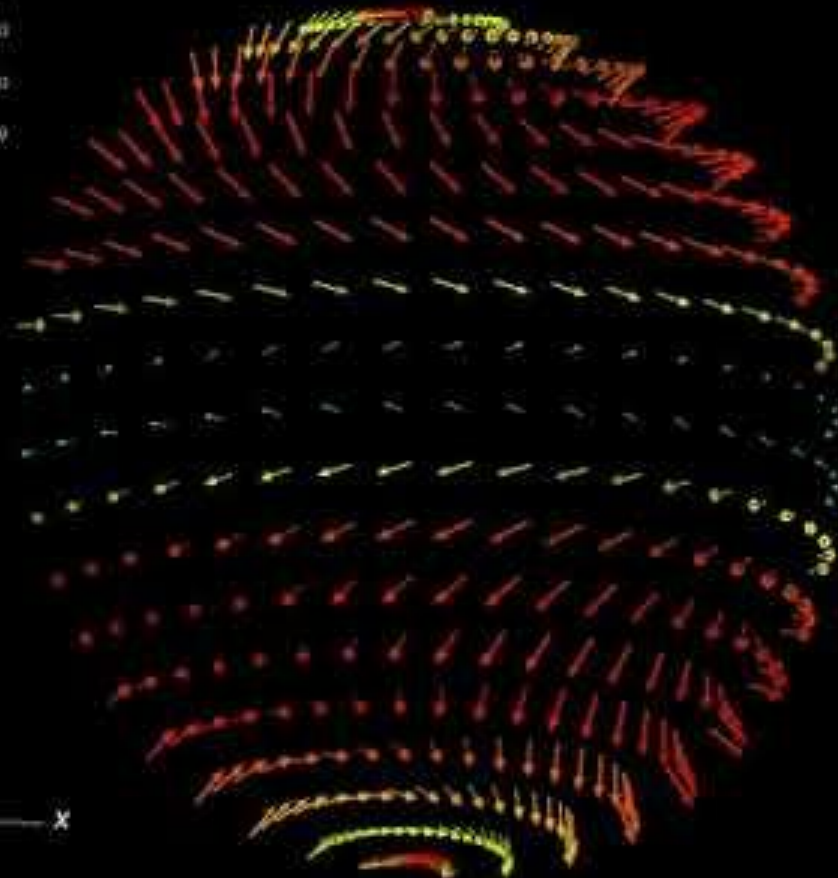
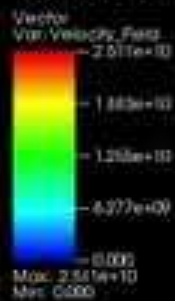
DB: data.0499.vtk
Cycle: 499 Time:0.499

Pseudocolor
Var: Magnetic_Field - Speed



DB: data.0316.vtk
Cycle: 316 Time: 0.315979

Variable:
Velocity Vector
Field



Verification and Validation

- Verify model in hydrostatic equilibrium:
 - Gravitational potential balances pressure gradient.
 - Mass does not spontaneously eject from stellar surface.
- Confirm stability of magnetic field model:
 - Poloidal and toroidal field components self-regulate, evolve with stability under time.
- Velocity vector field suggests over time, velocity perturbations settle near equatorial plane of star.
 - This plane is also referred to as the neutral line, where the magnetic field is minimized.
 - This behavior is validated given the Lorentz force in this region is also minimized

Future Work: Improving our Model

- Relativistic Equation of State:

- Tolman-Oppenheimer-Volkoff (TOV) equation (replaces newtonian formulation of hydrostatic equilibrium:

$$\frac{dP}{dr} = -\frac{Gm(r)}{r^2}\rho\left(1 + \frac{P}{\rho c^2}\right)\left(1 + \frac{4\pi r^3 P}{m(r)c^2}\right)\left(1 - \frac{2Gm(r)}{rc^2}\right)^{-1}$$

- EOS as a function of many variables:

$$P = P(\rho, T, x_p, \dots)$$

- As we change our EOS, we must choose a new magnetic field model:

- Fujisawa & Kisaka (Sept. 2014) via SLy EOS.

Acknowledgements

- I wish to extend appreciation and thanks to my co-mentors Dr. Michelle Kuchera and Dr. Kristen Thompson.
- I also thank former Duke undergraduate-student Emily Kuhn for providing an important starting point in informing my research.

Additional Photo Credits

- Slide 2
 - Neutron star structure: <http://www.astronomy.ohio-state.edu/~pogge/Ast162/Unit3/extreme.htm>
 - Pulsar schematic: <https://www.cv.nrao.edu/course/ast534/Pulsars.html>
 - Artists rendition of magnetar: <http://www.eso.org/public/images/eso1527a/>
- Slide 3
 - Swift Telescope: <http://swift.sonoma.edu/resources/multimedia/images/>
 - Fermi Telescope: <https://science.nasa.gov/toolkits/spacecraft-icons>

Requirements & Assumptions

- Our assumption of a polytropic EOS neglects the influence of temperature and composition on pressure, and a more realistic EOS will be a function of many variables:

$$P = P(\rho, T, x_p, \dots)$$

- Temperature plays an important role in young neutron stars, adding a buoyancy force. Additionally, young neutron stars develop separate superfluids for neutrons and protons, which can be accounted for by proton fraction ' x_p ' dependence in the EOS.
- Our polytropic EOS follows from the ideal gas thermal EOS given we require an adiabatic relationship between pressure and density.
- Model does not account for rotation, justified by relatively slow rotation of magnetars: mean rotational period of known magnetars ~ 6.7 s