

*Local variation of the velocity magnitude*

Some simple but useful results follow immediately from the expression for local vorticity in terms of rectangular co-ordinates with axes parallel to the location directions of  $\mathbf{u}$ , of the principal normal to the streamline (directed towards the centre of curvature), and of the binomial to the streamline. If  $(s, n, b)$  represent co-ordinates in these three dimensions respectively, and  $(u, v, w)$  are the corresponding velocity components, we have

$$v = w = 0, \quad u = q, \quad \frac{\partial v}{\partial s} = \frac{q}{R}, \quad \frac{\partial w}{\partial s} = 0$$

locally, where  $R$  is the local radius of curvature of the streamline. The components of the vorticity locally are then

$$\omega_s = \frac{\partial w}{\partial n} - \frac{\partial v}{\partial b}, \quad \omega_n = \frac{\partial u}{\partial b}, \quad \omega_b = \frac{u}{R} - \frac{\partial u}{\partial n}.$$

Moreover,  $\left( \frac{\partial}{\partial s}, \frac{\partial}{\partial n}, \frac{\partial}{\partial b} \right) u = \left( \frac{\partial}{\partial s}, \frac{\partial}{\partial n}, \frac{\partial}{\partial b} \right) q$

locally. Hence, in irrotational flow we have

$$\frac{\partial q}{\partial n} = \frac{q}{R}, \quad \frac{\partial q}{\partial b} = 0. \quad (6.2.13)$$

The first of the relations (6.2.13) shows that, when the streamlines . . .