

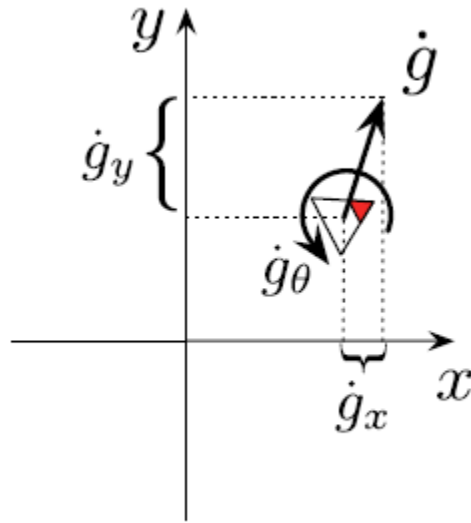
# Geometry of Locomotion

## Chapter 2.x – Rigid Body Velocities

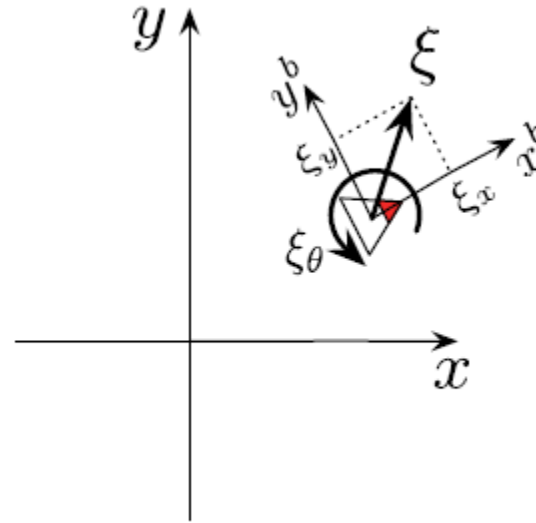
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# Body Velocity

A system's *body velocity*  $\xi$  is its velocity expressed in the instantaneous local coordinate frame,



(a) World velocity.



(b) Body velocity.

why

$$\begin{bmatrix} \xi^x \\ \xi^y \\ \xi^\theta \end{bmatrix} = \begin{bmatrix} \cos \theta & \sin \theta & 0 \\ -\sin \theta & \cos \theta & 0 \\ 0 & 0 & 1 \end{bmatrix} \begin{bmatrix} \dot{g}_x \\ \dot{g}_y \\ \dot{g}_\theta \end{bmatrix}$$

rotating the translational component by  $-\theta$

[equivalent to rotating the reference frame by  $\theta$ ]

$$\begin{aligned} \dot{g} &= \begin{bmatrix} \cos \theta & \sin \theta & 0 \\ -\sin \theta & \cos \theta & 0 \\ 0 & 0 & 1 \end{bmatrix}^{-1} \xi \\ &= \begin{bmatrix} \cos \theta & -\sin \theta & 0 \\ \sin \theta & \cos \theta & 0 \\ 0 & 0 & 1 \end{bmatrix} \xi \end{aligned}$$

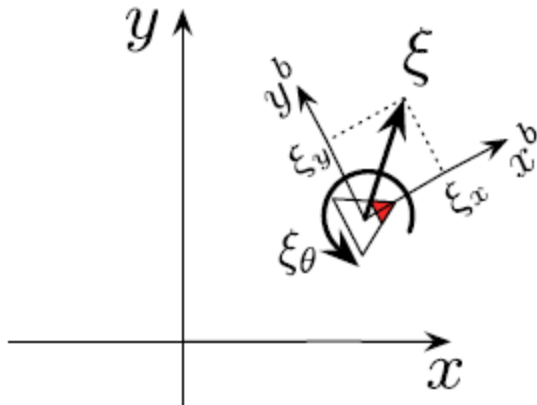
# Left Lifted Action

$$\begin{aligned}
 g &= (x, y, \theta) & T_g L_h &= \frac{\partial(hg)}{\partial g} \\
 h &= (u, v, \beta) & &= \begin{bmatrix} \frac{\partial(x \cos \beta - y \sin \beta + u)}{\partial x} & \frac{\partial(x \cos \beta - y \sin \beta + u)}{\partial y} & \frac{\partial(x \cos \beta - y \sin \beta + u)}{\partial \theta} \\ \frac{\partial(x \sin \beta + y \cos \beta + v)}{\partial x} & \frac{\partial(x \sin \beta + y \cos \beta + v)}{\partial y} & \frac{\partial(x \sin \beta + y \cos \beta + v)}{\partial \theta} \\ \frac{\partial(\theta + \beta)}{\partial x} & \frac{\partial(\theta + \beta)}{\partial y} & \frac{\partial(\theta + \beta)}{\partial \theta} \end{bmatrix} \\
 & & &= \begin{bmatrix} \cos \beta & -\sin \beta & 0 \\ \sin \beta & \cos \beta & 0 \\ 0 & 0 & 1 \end{bmatrix}.
 \end{aligned}$$

$T_g L_h$  preserves body velocities:

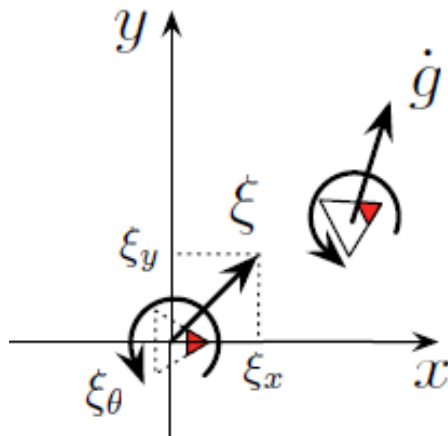
For any left action  $L_h$  that rotates the body frame by an angle  $\beta$ , the accompanying lifted action  $T_g L_h$  rotates the velocity vector by the same amount, leaving its expression in the body frame unchanged.

# More left lift



$h$  is set equal to  $g^{-1}$

$L_{g^{-1}}$  and  $T_g L_{g^{-1}}$  take a rigid body from  $g$  place it at the origin with equivalent body velocity,



$$\xi = T_g L_{g^{-1}} \dot{g} \in T_e G,$$

interchangeably treat the body velocity either in its standard interpretation as the set of forward, lateral, and rotational velocities of the body, or as a vector in the tangent space at the origin

$$\dot{g} = (T_g L_{g^{-1}})^{-1} \xi = T_e L_g \xi.$$

# Right Action

$R_h g$  on  $SE(2)$  finds the frame at position and orientation  $h$  with respect to  $g$ .

$$(\dot{g}h) = \begin{bmatrix} 1 & 0 & -(u \sin \theta + v \cos \theta) \\ 0 & 1 & u \cos \theta - v \sin \theta \\ 0 & 0 & 1 \end{bmatrix} \dot{g},$$

$$\begin{aligned} T_g R_h &= \frac{\partial(g h)}{\partial g} \\ &= \begin{bmatrix} \frac{\partial(x+u \cos \theta-v \sin \theta)}{\frac{\partial x}{\partial(y+u \sin \theta+v \cos \theta)}} & \frac{\partial(x+u \cos \theta-v \sin \theta)}{\frac{\partial y}{\partial(y+u \sin \theta+v \cos \theta)}} & \frac{\partial(x+u \cos \theta-v \sin \theta)}{\frac{\partial \theta}{\partial(y+u \sin \theta+v \cos \theta)}} \\ \frac{\partial(\theta+\beta)}{\frac{\partial x}{\partial x}} & \frac{\partial(\theta+\beta)}{\frac{\partial y}{\partial y}} & \frac{\partial(\theta+\beta)}{\frac{\partial \theta}{\partial \theta}} \end{bmatrix} \\ &= \begin{bmatrix} 1 & 0 & -(u \sin \theta + v \cos \theta) \\ 0 & 1 & u \cos \theta - v \sin \theta \\ 0 & 0 & 1 \end{bmatrix}, \end{aligned}$$

# Right Action

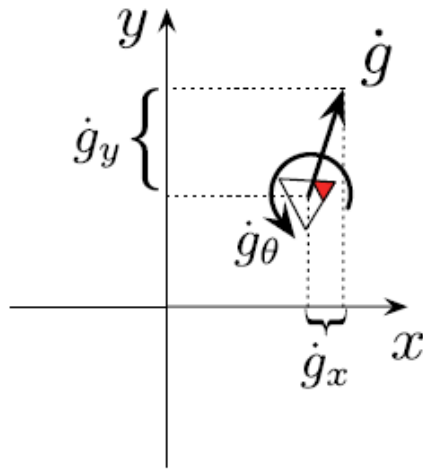
- an action  $R_h$  maps between rigidly attached frames
- the “equivalent” vector selected by the accompanying lifted action  $T_g R_h$  ensures that the new frame’s velocity is compatible with the rigid attachment.
- right lifted action preserves their *spatial velocity*
- velocity of the (possibly imaginary) point on that body that is instantaneously over the origin, calculated as

$$\xi^s = T_g R_{g^{-1}} \dot{g} = \begin{bmatrix} 1 & 0 & y \\ 0 & 1 & -x \\ 0 & 0 & 1 \end{bmatrix} \dot{g}$$

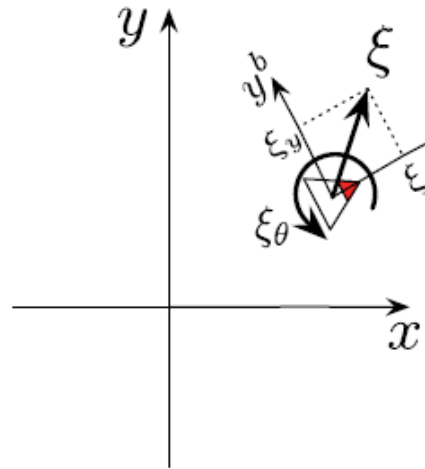
Lie algebra

$$\dot{g} = (T_g R_{g^{-1}})^{-1} \xi^s = T_e R_g \xi^s = \begin{bmatrix} 1 & 0 & -y \\ 0 & 1 & x \\ 0 & 0 & 1 \end{bmatrix} \xi^s$$

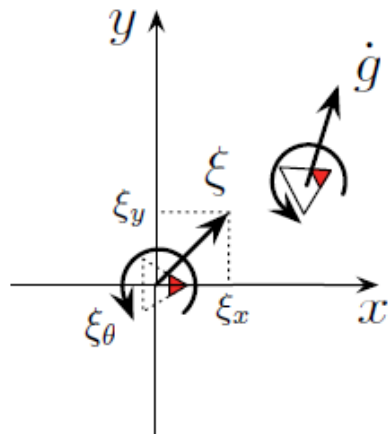
# Velocities Summarized



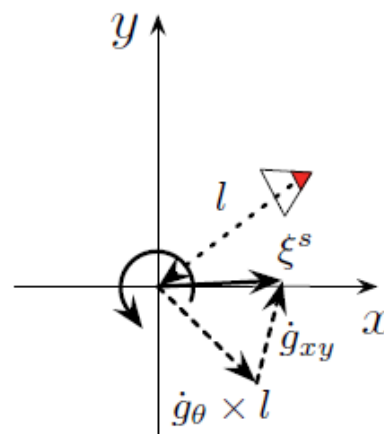
(a) World velocity.



(b) Body velocity.



(c) Alternate interpretation of body velocity



(d) Spatial velocity.

# Adjoint Operation

$$\xi^s = \overbrace{(T_g R_{g^{-1}})(T_e L_g)}^{Ad_g} \xi,$$

$$\xi = \overbrace{(T_g L_{g^{-1}})(T_e R_g)}^{Ad_g^{-1}} \xi^s \quad \text{adjoint inverse action } Ad_g^{-1} = Ad_{g^{-1}}$$



# Adjoint on SE(2)

$$\xi^s = \underbrace{\begin{bmatrix} 1 & 0 & y \\ 0 & 1 & -x \\ 0 & 0 & 1 \end{bmatrix}}_{T_g R_{g^{-1}}} \underbrace{\begin{bmatrix} \cos \theta & -\sin \theta & 0 \\ \sin \theta & \cos \theta & 0 \\ 0 & 0 & 1 \end{bmatrix}}_{T_e L_g} \xi$$

$$= \underbrace{\begin{bmatrix} \cos \theta & -\sin \theta & y \\ \sin \theta & \cos \theta & -x \\ 0 & 0 & 1 \end{bmatrix}}_{Ad_g} \xi.$$

$$\xi = \underbrace{\begin{bmatrix} \cos \theta & \sin \theta & 0 \\ -\sin \theta & \cos \theta & 0 \\ 0 & 0 & 1 \end{bmatrix}}_{T_g L_{g^{-1}}} \underbrace{\begin{bmatrix} 1 & 0 & -y \\ 0 & 1 & x \\ 0 & 0 & 1 \end{bmatrix}}_{T_e R_g} \xi^s$$

$$= \underbrace{\begin{bmatrix} \cos \theta & \sin \theta & x \sin \theta - y \cos \theta \\ -\sin \theta & \cos \theta & x \cos \theta + y \sin \theta \\ 0 & 0 & 1 \end{bmatrix}}_{Ad_g^{-1}} \xi^s.$$