

# Detailed Proof and Analysis of Grover's Search Algorithm

Raunak Kumar Singh

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## 1 Introduction and Problem Statement

The unstructured search problem is a fundamental challenge in computer science. Classically, searching an unsorted database of  $N$  items requires  $O(N)$  queries in the worst case. Grover's algorithm provides a quadratic speedup, solving the problem in  $O(\sqrt{N})$  queries.

### 1.1 Formal Setup

We are given a search space of size  $N = 2^n$ . We assume the existence of a black box function (oracle)  $f : \{0, 1\}^n \rightarrow \{0, 1\}$  defined as:

$$f(x) = \begin{cases} 1 & \text{if } x \text{ is a solution (marked item)} \\ 0 & \text{if } x \text{ is not a solution} \end{cases} \quad (1)$$

Let  $M$  be the number of solutions such that  $f(x) = 1$ . Our objective is to find a state  $x^*$  such that  $f(x^*) = 1$  with high probability.

## 2 Geometric Formalism

To understand the algorithm, we analyze the evolution of the state vector within a specific 2-dimensional subspace of the Hilbert space  $\mathcal{H} = \mathbb{C}^N$ .

### 2.1 Defining the Basis States

We define two normalized superposition states:

**Definition 1** (Good and Bad States). *The uniform superposition of all solution states ( $|good\rangle$ ) and non-solution states ( $|bad\rangle$ ) are:*

$$|good\rangle = \frac{1}{\sqrt{M}} \sum_{x:f(x)=1} |x\rangle, \quad |bad\rangle = \frac{1}{\sqrt{N-M}} \sum_{x:f(x)=0} |x\rangle \quad (2)$$

### 2.2 The Initial State

The algorithm begins by initializing  $n$  qubits to  $|0\rangle$  and applying Hadamard gates ( $H^{\otimes n}$ ). This creates the uniform superposition  $|\psi\rangle$ , which lies entirely within the plane spanned by  $|good\rangle$  and  $|bad\rangle$ :

$$|\psi\rangle = H^{\otimes n} |0\rangle^{\otimes n} = \frac{1}{\sqrt{N}} \sum_{x=0}^{N-1} |x\rangle \quad (3)$$

We can rewrite  $|\psi\rangle$  using an angle  $\theta$ :

$$|\psi\rangle = \sin \theta |\text{good}\rangle + \cos \theta |\text{bad}\rangle \quad (4)$$

where the angle  $\theta$  is determined by the ratio of solutions to the total space:

$$\sin \theta = \sqrt{\frac{M}{N}} \quad (5)$$

### 2.3 The Orthogonal State

For the algebraic derivation of the rotation, it is useful to define a state  $|\bar{\psi}\rangle$  that is orthogonal to  $|\psi\rangle$  in this 2D plane:

$$|\bar{\psi}\rangle = \cos \theta |\text{good}\rangle - \sin \theta |\text{bad}\rangle \quad (6)$$

Note that  $\langle \psi | \bar{\psi} \rangle = 0$ .

## 3 The Grover Iteration

The core of the algorithm is the Grover Iterator  $G$ , which is applied  $k$  times. The iterator consists of two reflections:

$$G = U_s U_f \quad (7)$$

where  $U_f$  is the Oracle and  $U_s$  is the Diffusion Operator.

### 3.1 Step 1: The Oracle ( $U_f$ )

The oracle marks solutions by flipping their phase:  $|x\rangle \rightarrow (-1)^{f(x)} |x\rangle$ . Geometrically, this is a reflection about the  $|\text{bad}\rangle$  axis.

$$U_f |\psi\rangle = -\sin \theta |\text{good}\rangle + \cos \theta |\text{bad}\rangle \quad (8)$$

To understand the rotation, let us express this reflection in terms of the  $|\psi\rangle, |\bar{\psi}\rangle$  basis. Using the substitutions  $|\text{good}\rangle = \sin \theta |\psi\rangle + \cos \theta |\bar{\psi}\rangle$  and  $|\text{bad}\rangle = \cos \theta |\psi\rangle - \sin \theta |\bar{\psi}\rangle$ :

$$\begin{aligned} U_f |\psi\rangle &= -\sin \theta (\sin \theta |\psi\rangle + \cos \theta |\bar{\psi}\rangle) + \cos \theta (\cos \theta |\psi\rangle - \sin \theta |\bar{\psi}\rangle) \\ &= (-\sin^2 \theta + \cos^2 \theta) |\psi\rangle - 2 \sin \theta \cos \theta |\bar{\psi}\rangle \end{aligned}$$

Using the double-angle trigonometric identities ( $\cos 2\theta = \cos^2 \theta - \sin^2 \theta$  and  $\sin 2\theta = 2 \sin \theta \cos \theta$ ):

$$U_f |\psi\rangle = \cos 2\theta |\psi\rangle - \sin 2\theta |\bar{\psi}\rangle \quad (9)$$

### 3.2 Step 2: The Diffusion Operator ( $U_s$ )

The diffusion operator is defined as  $U_s = 2|\psi\rangle\langle\psi| - I$ . Geometrically, this is a reflection about the vector  $|\psi\rangle$ . When we apply  $U_s$  to Equation (9):

- The component parallel to  $|\psi\rangle$  remains unchanged.
- The component orthogonal to  $|\psi\rangle$  (i.e.,  $|\bar{\psi}\rangle$ ) gets flipped.

$$\begin{aligned} G |\psi\rangle &= U_s (\cos 2\theta |\psi\rangle - \sin 2\theta |\bar{\psi}\rangle) \\ &= \cos 2\theta |\psi\rangle - \sin 2\theta (-|\bar{\psi}\rangle) \\ &= \cos 2\theta |\psi\rangle + \sin 2\theta |\bar{\psi}\rangle \end{aligned}$$

### 3.3 Step 3: Calculating the Total Rotation (The $3\theta$ Term)

We have established that one Grover iteration rotates the state by  $2\theta$  relative to the start state. Let us convert this back to the  $|\text{good}\rangle, |\text{bad}\rangle$  basis to see the new amplitude.

Substituting  $|\psi\rangle$  and  $|\bar{\psi}\rangle$  back into the equation:

$$\begin{aligned} |\psi_1\rangle &= \cos 2\theta(\sin \theta |\text{good}\rangle + \cos \theta |\text{bad}\rangle) + \sin 2\theta(\cos \theta |\text{good}\rangle - \sin \theta |\text{bad}\rangle) \\ &= (\sin \theta \cos 2\theta + \cos \theta \sin 2\theta) |\text{good}\rangle + (\cos \theta \cos 2\theta - \sin \theta \sin 2\theta) |\text{bad}\rangle \end{aligned}$$

We use the angle addition formulas:

- $\sin(A + B) = \sin A \cos B + \cos A \sin B$
- $\cos(A + B) = \cos A \cos B - \sin A \sin B$

Setting  $A = \theta$  and  $B = 2\theta$ , we obtain:

$$|\psi_1\rangle = \sin(3\theta) |\text{good}\rangle + \cos(3\theta) |\text{bad}\rangle \quad (10)$$

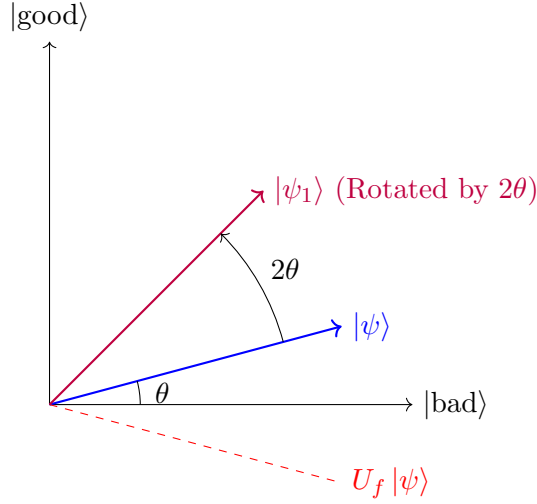


Figure 1: Geometric interpretation of one Grover Iteration. The state is rotated by  $2\theta$  towards  $|\text{good}\rangle$ .

## 4 Complexity and Convergence

By induction, applying the Grover iterator  $k$  times results in the state:

$$|\psi_k\rangle = \sin((2k + 1)\theta) |\text{good}\rangle + \cos((2k + 1)\theta) |\text{bad}\rangle \quad (11)$$

The probability of measuring a solution is  $P_k = \sin^2((2k + 1)\theta)$ .

### 4.1 Optimal Iterations

We wish to stop when the probability is maximal, i.e.,  $(2k + 1)\theta \approx \pi/2$ .

$$k \approx \frac{\pi}{4\theta} - \frac{1}{2} \quad (12)$$

Assuming  $N \gg M$ ,  $\theta \approx \sin \theta = \sqrt{M/N}$ . This yields the complexity:

$$k \approx \frac{\pi}{4} \sqrt{\frac{N}{M}} = O(\sqrt{N}) \quad (13)$$

## 5 Error Analysis and Bounds

A crucial question is what happens if the optimal  $k$  is not an integer. Let  $\tilde{k}$  be the exact real number that satisfies  $(2\tilde{k} + 1)\theta = \pi/2$ . Let  $\bar{k}$  be the closest integer to  $\tilde{k}$ .

$$|\bar{k} - \tilde{k}| \leq \frac{1}{2} \quad (14)$$

### 5.1 Probability of Failure

If we iterate  $\bar{k}$  times, the probability of *not* observing a solution (failure) is given by the cosine component squared:

$$P(\text{fail}) = \cos^2((2\bar{k} + 1)\theta) \quad (15)$$

We can substitute  $(2\bar{k} + 1)\theta$  by utilizing the deviation from the optimal angle:

$$\begin{aligned} (2\bar{k} + 1)\theta &= (2\tilde{k} + 2(\bar{k} - \tilde{k}) + 1)\theta \\ &= (2\tilde{k} + 1)\theta + 2(\bar{k} - \tilde{k})\theta \\ &= \frac{\pi}{2} + 2(\bar{k} - \tilde{k})\theta \end{aligned}$$

Substituting this back into the failure probability:

$$\begin{aligned} P(\text{fail}) &= \cos^2\left(\frac{\pi}{2} + 2(\bar{k} - \tilde{k})\theta\right) \\ &= \sin^2(2(\bar{k} - \tilde{k})\theta) \end{aligned}$$

Since  $|\bar{k} - \tilde{k}| \leq 1/2$ , the argument of the sine function is bounded by  $\theta$ :

$$|2(\bar{k} - \tilde{k})\theta| \leq \theta \quad (16)$$

Using the inequality  $\sin x \leq x$  for small  $x$ , and specifically here  $\sin^2(\theta) = M/N$ :

$$P(\text{fail}) \leq \sin^2(\theta) = \frac{M}{N} \quad (17)$$

## 6 Conclusion

Grover's algorithm successfully rotates the initial state vector from the superposition of all states towards the solution states.

1. It provides a quadratic speedup over classical search.
2. Even with discrete iteration steps, the error is bounded by  $M/N$ , which is negligible for large  $N$ .