#### **EECS 16B Midterm 1 Review Session**

Presented by George Higgins Hutchinson, Ben Liao, Rohin Tangriala October 9, 2020

#### **Disclaimer**

Although some of the presenters may be course staff, the material covered in the review session may not be an accurate representation of the topics covered in and difficulty of the exam.

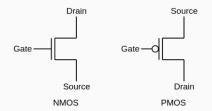
Slides are posted at 691 on Piazza.

### **HKN Drop-In Tutoring**

• These details should be written.

# **CMOS Transistors and Logic**

#### **Transistors**

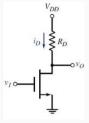


Two varieties of MOSFETs: P-type and N-type Any MOSFET has a characteristic threshold voltage  $V_{th}$  NMOS "turns on" (connects drain to source) when  $V_{GS} > V_{th}$  PMOS turns on when  $V_{GS} < V_{th}$ 

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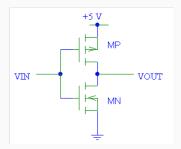
#### **NMOS Logic**

We can build an inverter (a circuit that flips a 1 to a 0 and vice versa) with a single N-type MOSFET!



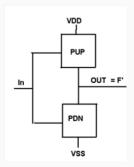
What is  $v_O$  when  $v_I > V_{th}$ ? When  $v_I < V_{th}$ ? Key disadvantage: what is the power dissipated when  $v_I > V_{th}$ ? How can we rectify this?

#### **CMOS Logic**



Now, we have the same logical function as the NMOS inverter, but we're using more transistors and much less power is dissipated. Why?

### **CMOS Logic**



This is broadly called Complementary Metal Oxide Semiconductor (CMOS) logic, using a Pull-Up Network of P-type and a Pull-Down Network of N-type MOSFETs.

Now we can build circuits that perform logical functions out of MOSFETs!

#### True or False:

- 1. Power (in the EE16B model) is dissipated in a CMOS circuit only when there is switching
- 2. NMOS devices turn on with a large VGS and off with low VGS
- For a given choice of inputs in a CMOS inverter circuit, we can create a low resistance path from ground to high voltage

#### True or False:

1. Power (in the EE16B model) is dissipated in a CMOS circuit only when there is switching  $\frac{1}{2}$ 

#### True

- 2. NMOS devices turn on with a large VGS and off with low VGS
- For a given choice of inputs in a CMOS inverter circuit, we can create a low resistance path from ground to high voltage

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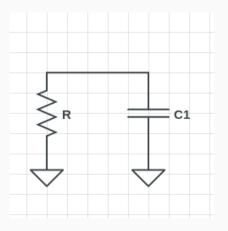
- 2. NMOS devices turn on with a large VGS and off with low VGS **True**
- For a given choice of inputs in a CMOS inverter circuit, we can create a low resistance path from ground to high voltage False

## **RC Circuits**

#### **RC Circuits**

The capacitor and resistor in a NOT circuit form the most basic RC circuit:

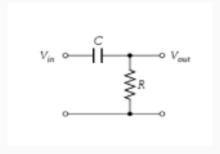
Write down a differential equation describing the circuit below:



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Write down a differential equation describing the circuit below:



We have a differential equation describing  $V_{out}$  in terms of  $V_c$ .

$$\frac{dV_c}{dt} = -\frac{1}{RC}V_c$$

How do we actually solve it?

We have a differential equation describing  $V_{out}$  in terms of  $V_c$ .

$$\frac{dV_c}{dt} = -\frac{1}{RC}V_c$$

How do we actually solve it?

Think of the differential operation as a linear operator that scales  $V_c$ , since  $V_c$  is one of its eigenfunctions:

$$\left[\frac{d}{dt}\right]V_c = \lambda V_c$$

Which are the eigenfunctions of differentiation?

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$$Ae^{\lambda t}$$

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$$\frac{dV_c}{dt} = -\frac{1}{RC}V_c$$

The solution to our first order differential equations is therefore:

$$V_c(t) = V_c(0)e^{\frac{-1}{RC}t}$$

### RC Differential Equation: Non-homogenous case

How do you solve a RC circuit with a voltage source?



Applying KCL at the top right node, along with Ohm's law and the capacitor relationship, we get:

$$\frac{dV_c}{dt} = \frac{1}{VC}(V_s - V_c)$$

We can't easily solve this equation, so we change variables to

$$x = V_c - V_s$$

### RC Differential Equation: Non-homogenous case

Now, we have  $\frac{dx}{dt} = -\frac{x}{RC}$ 

We already know how to solve this differential equation, and we get

$$x(t) = x_0 e^{-\frac{t}{RC}}$$

Finally, change back to the original variables by substituting  $V_c$  -  $V_s$  for x

$$V_c(t) = V_c(0)e^{-\frac{t}{RC}} + V_s(1 - e^{-\frac{t}{RC}})$$

# Change of Basis

### **Change of Basis**

In the standard basis, we write vectors as a linear combination of the standard basis vectors.

$$\vec{x} = \begin{bmatrix} x_1 \\ x_2 \end{bmatrix} = x_1 \begin{bmatrix} 1 \\ 0 \end{bmatrix} + x_2 \begin{bmatrix} 0 \\ 1 \end{bmatrix}$$

Likewise, we can write our x vector as a linear combination of some other basis vectors. For example, if our V-basis has basis vectors  $\vec{v_1}$  and  $\vec{v_2}$ :

$$\vec{x} = \tilde{x_1} \cdot \vec{v_1} + \tilde{x_2} \cdot \vec{v_2} = V\vec{\tilde{x}}$$

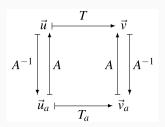
We can go from the V-basis to the standard basis by applying the V-matrix to x, and go from the standard basis to the V-basis by applying  $V^{-1}$ .

### Change of Basis Diagram

- When you are doing change of basis for systems of equations, it is useful to use a diagram mapping transformations between basis
- Given T, to find T<sub>a</sub> you would trace the path from u<sub>a</sub> to v<sub>a</sub>, applying each transformation to the left of the previous:

$$A \Longrightarrow TA \Longrightarrow A^{-1}TA$$
 $T_a = A^{-1}TA$ 

• Step by step, we have  $Au_a = u \implies TAu_a = v \implies A^{-1}TAu_a = v_a$ 



# Diagonalization

### Diagonalization

- The idea is that we want to change into a basis in which the system  $A\vec{x} = \vec{y}$  is represented by a diagonal matrix. So how do we find such basis?
- Remember that for all eigenvalue-eigenvector pairs we have:  $A\vec{v}=\lambda\vec{v}$
- Let's use our eigenvectors as our basis. Doing so we obtain:

$$\vec{x} = V\tilde{\vec{x}}, \vec{y} = V\tilde{\vec{y}}, \text{and} \Lambda \tilde{\vec{x}} = \tilde{\vec{y}}$$

Where the upper-case lambda represents the diagonal matrix with eigenvalues on the diagonals.

Transforming back to the standard basis, we get:

$$A = V \Lambda V^{-1}$$

### Diagonalization Cont.

- Let's analyze this a bit further, why is it important to be able to do this?
- We see that diagonalizing the matrix makes the system much easier to solve, why?
- Also, we see that there is a "home state" for every system of linearly independent equations, i.e. the space in which the system's components are uncoupled.

2 Minute Break!

# **Differential Equations**

#### **Solving Differential Equations**

Differentiation is linear!

$$\frac{d(c_1x(t)+c_2y(t))}{dt}=c_1\frac{dx(t)}{dt}+c_2\frac{dy(t)}{dt}$$

### Solving systems of differential equations

Write the system in the following form:

$$\frac{dx(t)}{dt} = Ax(t)$$

Find the eigenvalues/eigenvectors of A and transform the system into the eigenbasis. If A has distinct eigenvalues, the general solution is:

$$x(t) = c_1 e^{\lambda_1 t} \vec{v_1} + c_2 e^{\lambda_2 t} \vec{v_2}$$

Then, change back to the standard basis.

Solve the following differential equation:

$$\ddot{x} - \dot{x} - 2x = 0$$

$$x(0) = 2, \dot{x}(0) = 1$$

$$\ddot{x} - \dot{x} - 2x = 0$$
,  $x(0) = 2, \dot{x}(0) = 1$ 

Solution:

Write in matrix form and find eigenvectors:

$$\begin{bmatrix} \dot{x} \\ \ddot{x} \end{bmatrix} = \begin{bmatrix} 0 & 1 \\ 2 & 1 \end{bmatrix} \begin{bmatrix} x \\ \dot{x} \end{bmatrix}$$

$$\lambda_1 = 2, v_1 = \begin{bmatrix} 1 \\ 2 \end{bmatrix}, \lambda_2 = -1, v_2 = \begin{bmatrix} -1 \\ 1 \end{bmatrix}$$

Use general form to solve:

$$\begin{bmatrix} x \\ \dot{x} \end{bmatrix} = c_1 e^{2t} v_1 + c_2 e^{-t} v_2$$

$$\ddot{x} - \dot{x} - 2x = 0$$
,  $x(0) = 2$ ,  $\dot{x}(0) = 1$ 

Solution:

Now plug in initial conditions:

$$\begin{bmatrix} 2 \\ 1 \end{bmatrix} = \begin{bmatrix} x(0) \\ \dot{x}(0) \end{bmatrix} = c_1 v_1 + c_2 v_2 = \begin{bmatrix} c_1 - c_2 \\ 2c_1 + c_2 \end{bmatrix}$$

Solving, we get 
$$c_1 = 1, c_2 = -1$$

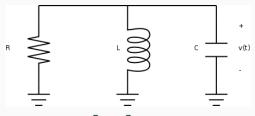
so 
$$x(t) = e^{2t} + e^{-t}$$

$$\ddot{x} - \dot{x} - 2x = 0$$
,  $x(0) = 2$ ,  $\dot{x}(0) = 1$ 

Solution:

Sanity check: compute derivatives and check that the original system is satisfied

$$x(t) = e^{2t} + e^{-t}, \dot{x}(t) = 2e^{2t} - e^{-t}, \ddot{x}(t) = 4e^{2t} + e^{-t}$$
$$\implies \ddot{x} - \dot{x} - 2x = 0, x(0) = 2, \dot{x}(0) = 1$$



Given 
$$x(t) = \begin{bmatrix} v(t) \\ i_L(t) \end{bmatrix}$$

find matrix A such that  $\frac{dx}{dt} = Ax$ 

Direct all currents into ground.

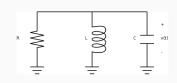
$$i_R + i_L + i_C = 0$$

$$\frac{V}{R} + i_L = -C \frac{dV}{dt}$$

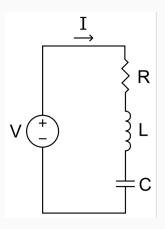
$$-\frac{1}{CR}V + -\frac{1}{C}i_L = \frac{dV}{dt}$$

Also

$$L\frac{di_L}{dt} = V \implies A = \begin{bmatrix} -\frac{1}{RC} & -\frac{1}{C} \\ \frac{1}{L} & 0 \end{bmatrix}$$



At t < 0, the circuit is at steady state. At  $t \ge 0$ , the voltage source is set to 0. Find a differential equation for  $i_L$  for t > 0.

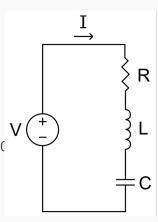


Solution: use KCL/KVL to get

$$i_R = i_L = i_C, V_C + V_R + V_L = 0$$

$$\frac{V_R}{R} = i_L = C \frac{dV_C}{dt}, \frac{1}{C} i_L + R \frac{di_L}{dt} + L \frac{d^2 i_L}{dt} = 0$$

$$\begin{bmatrix} \frac{di_L}{dt} \\ \frac{d^2i_L}{dt} \end{bmatrix} = \begin{bmatrix} 0 & 1 \\ -\frac{1}{LC} & -\frac{R}{L} \end{bmatrix} \begin{bmatrix} i_L \\ \frac{di_L}{dt} \end{bmatrix}$$

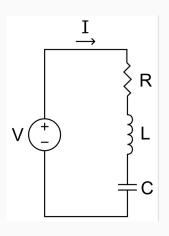


If R = 0, L = C = 1, solve the differential equation:

Initial conditions:  $i_L(0) = 0$ ,  $\frac{di_L}{dt}(0) = -V_C = -V$ 

$$\begin{bmatrix} \frac{di_{l}}{dt} \\ \frac{d^{2}i_{l}}{dt} \end{bmatrix} = \begin{bmatrix} 0 & 1 \\ -\frac{1}{LC} & -\frac{R}{L} \end{bmatrix} \begin{bmatrix} i_{L} \\ \frac{di_{L}}{dt} \end{bmatrix}$$
$$\lambda_{1} = i, v_{1} = \begin{bmatrix} i \\ -1 \end{bmatrix}$$

$$\lambda_2 = -i, v_2 = \begin{bmatrix} i \\ 1 \end{bmatrix}$$



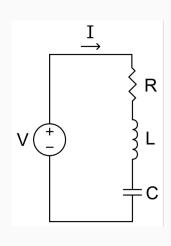
$$\begin{bmatrix} i_{L} \\ \frac{di_{L}}{dt} \end{bmatrix} = c_{1}e^{it}v_{1} + c_{2}e^{-it}v_{2}$$

$$ic_1 + ic_2 = 0, -c_1 + c_2 = -V$$

$$\implies c_1 = \frac{V}{2}, c_2 = -\frac{V}{2}$$

SO

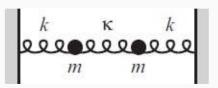
$$i_L = \frac{V}{2}ie^{it} - \frac{V}{2}e^{-it}$$



#### **Example: 2 coupled diff eqs**

Consider a system of 2 coupled harmonic oscillators, described by

$$m_i\ddot{x}_i = -(k+\kappa)x_i + \kappa x_j$$



Write this system in matrix form

$$m_i\ddot{x}_i = -(k+\kappa)x_i + \kappa x_j$$

Write this system in matrix form

$$\vec{x} = \begin{bmatrix} -\frac{k+\kappa}{m} & \frac{\kappa}{m} \\ \frac{\kappa}{m} & -\frac{k+\kappa}{m} \end{bmatrix} \vec{x}$$

Now, diagonalize the system

$$\begin{bmatrix} -\frac{k+\kappa}{m} & \frac{\kappa}{m} \\ \frac{\kappa}{m} & -\frac{k+\kappa}{m} \end{bmatrix} = PDP'$$

- $P = \hat{a}$
- D = ?

Now, diagonalize the system

$$\begin{bmatrix} -\frac{k+\kappa}{m} & \frac{\kappa}{m} \\ \frac{\kappa}{m} & -\frac{k+\kappa}{m} \end{bmatrix} = PDP'$$

$$P = \begin{bmatrix} 1 & 1 \\ 1 & -1 \end{bmatrix}$$

$$D = \begin{bmatrix} k & 0 \\ 0 & 2\kappa + k \end{bmatrix}$$

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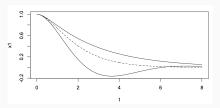
$$D = \begin{bmatrix} k & 0 \\ 0 & 2\kappa + k \end{bmatrix}$$

Since our original matrix was linearly taking a SECOND derivative in time, with only real eigenvalues, we expect solutions to be in the form of real sinusoids, optionally phaseshifted.

$$x(t) = A_s \begin{bmatrix} 1 \\ 1 \end{bmatrix} \cos(\sqrt{\frac{k}{m}}t + \phi_s) + A_f \begin{bmatrix} 1 \\ -1 \end{bmatrix} \sin(\sqrt{\frac{k+2\kappa}{m}}t + \phi_f)$$

# What do we do with different types of eigenvalues?

Each interesting "case" of eigenvalues (real, imaginary, and repeated) is complementary to a case of resonance in 2d ODEs, so let's discuss them together:



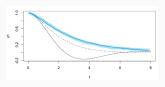
## Overdamping: two real eigenvalues

 When a system has two real eigenvalues, it will have two real eigenvectors, so the solution will be a combination of exponentials in the form

$$ae^{\lambda_1 t} + be^{\lambda_2 t}$$

For negative lambdas, this produces exponential decay.

At large t, greater eigenvalue dominates behavior, so how fast the system approaches to 0 is determined by the larger eigenvalue.



## Underdamping: imaginary or complex eigenvalues

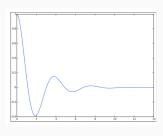
 You can still solve the system of differential equations as usual, but your solutions will be complex:

$$(a_1j + b_1)e^{c_1j+d_1} + (a_2j + b_2)e^{c_2j+d_2}$$

 You can rearrange terms and apply Euler's formula to write each term as a real exponential multiplied by a sinusoid

$$e^{d_1}(\alpha_1 sin(...) + \beta_1 cos(...)) + e^{d_2}(\alpha_2 sin(...) + \beta_2 cos(...))$$

If  $d_1$  and  $d_2$  are negative, then you can think of this as a sinusoid where the amplitude decays to 0



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## Critical Damping: repeated eigenvalue

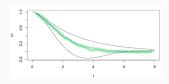
- You only get one linearly independent eigenvector.
- So, arbitrarily choose the second column of your V matrix (making sure it is linearly independent from the eigenvector v<sub>1</sub>)
- Compute  $A_v$  using the change of basis diagram  $A_v = V^{-1}AV$
- This gives you an upper-triangular matrix
- The bottom row of the matrix gives you an equation in one variable (only depends on the second component of  $x_v$ ), which you can solve to get something in the form of  $x_{v,2} = ae^{\lambda t}$

# Critical Damping: repeated eigenvalue

- The top row of the matrix depends on both components of  $x_{\nu}$ , but you can plug in the value you got for  $x_{\nu,2}$  and then solve for  $x_{\nu,1}$
- If you solve this differential equation (which is a diff eq with a non-constant input\*), you'll get an equation in the form:

$$x(t) = c_1 e^{-\lambda t} + c_2 t e^{-\lambda t}$$

The critically damped case gives us the fastest-decaying exponential that doesn't oscillate



\*You can solve this using the formula  $x_p(t) = x_0 e^{\lambda t} + \int_0^t u(\tau) e^{\lambda(t-\tau)} d\tau$ 

2 min break!

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## **Phasors**

#### **Phasors**

- Phasors express the response of a circuit to a sinusoidal input
- Any real, periodic signal can be expresses as the sum of sinusoids - we can apply this technique very broadly!
- A phasor encodes information about amplitude and phase, but not frequency

$$v_s = A\cos(\omega t + \phi) \rightarrow \tilde{V}_s = Ae^{j\phi}$$

### **Impedance**

- Impedance (*Z*) is a generalized form of resistance expresses a component's response to an alternating current
- Each common passive circuit element (resistor, capacitor, and inductor) has a characteristic impedance

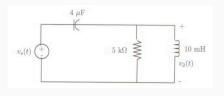
$$Z_R = R$$

$$Z_C = \frac{1}{j\omega C}$$

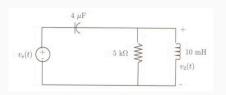
$$Z_L = j\omega L$$

### **Phasor Analysis Procedure**

- 1. Express all time domain signals as cosines
- 2. Convert voltages, currents, and impedances to their phasor equivalents
- 3. Set up phasor equations and solve for unknowns
- 4. Transform back to time domain

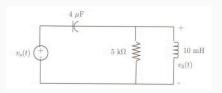


Let 
$$v_s(t) = 5\cos(5000t + \frac{\pi}{4})$$
. Find  $v_2(t)$ .



#### 1. Translate to phasor domain

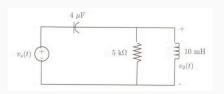
- $\omega = 5000$
- $v_s(t) = 5\cos(5000t + \frac{\pi}{4}) \rightarrow \tilde{V}_s = 5e^{j\frac{\pi}{4}}$
- $C = 4\mu F \rightarrow Z_C = \frac{1}{j(4*10^{-6})(5000)} = -j50$
- $R = 5k\Omega \rightarrow Z_R = 5*10^3$
- $L = 10mH \rightarrow Z_L = j(5000)(10 * 10^{-3}) = j50$



## 2. Solve for $\tilde{V}_2$

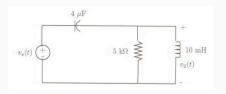
• Set up as impedance divider:

$$\begin{split} \frac{\tilde{V}_{s}}{Z_{C} + (Z_{R} \| Z_{L})} &= \frac{\tilde{V}_{2}}{Z_{R} \| Z_{L}} \\ \tilde{V}_{2} &= \tilde{V}_{s} \frac{Z_{R} \| Z_{L}}{Z_{C} + (Z_{R} \| Z_{L})} &= \tilde{V}_{s} \frac{\frac{Z_{R} Z_{L}}{Z_{R} + Z_{L}}}{\frac{Z_{R} Z_{L} + Z_{C} Z_{C} + Z_{L} Z_{C}}{Z_{R} + Z_{L}}} \end{split}$$



# 2. Solve for $\tilde{V}_2$

$$\tilde{V}_2 = \tilde{V}_s \frac{Z_R Z_L}{Z_R Z_L + Z_R Z_C + Z_L Z_C} = \tilde{V}_s \frac{j2.5 * 10^5}{j2.5 * 10^5 - j2.5 * 10^5 + 2500}$$
$$\tilde{V}_2 = j100 \tilde{V}_s = 100 \tilde{V}_s e^{j\frac{\pi}{2}}$$



#### 3. Find time-domain solution

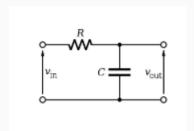
• 
$$\tilde{V}_2 = 100e^{j\frac{\pi}{2}}\tilde{V}_s \rightarrow v_2(t) = 100v_s(t + \frac{\pi}{2}) = 500\cos(5000t + \frac{3\pi}{4})$$

#### **Transfer Function Review**

A transfer function of a circuit or system describes the output response to an input excitation as a function of the angular frequency  $\omega$ .

$$H(\omega) = rac{V_{out}(\omega)}{V_{in}(\omega)}$$

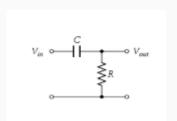
#### **RC Circuits**



$$H(\omega) = \frac{1}{1+j\omega RC}$$

#### Low pass:

$$\lim_{\omega \to \infty} H(\omega) = 0, H(0) = 1$$



$$H(\omega) = \frac{j\omega RC}{1+j\omega RC}$$

# High pass:

$$\lim_{\omega\to\infty}H(\omega)=1, H(0)=0$$

#### **Bode Plot Review**

$$V[dB] = 20\log_{10}(\frac{V}{V_0})$$

- Bode plots provide a way for us to easily visualize the output response of our system, depending on the input frequency  $\omega$ .
- Because Bode plots are in log scale for  $\omega$ , we are able to take advantage of the properties of logs.
  - $G = XY \implies G[dB] = X[dB] + Y[dB]$
  - $G = \frac{X}{Y} \implies G[dB] = X[dB] Y[dB]$

Hence, we can break our transfer function  $H(\omega)$  into a product of familiar transfer functions (simple poles, quadratic zeros, etc. - "functional forms"), graph them out individually, and then add them together on the graph.

#### **Bode Plot Review**

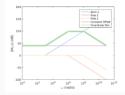
- We want to be able to plot the transfer function as a function of the frequency  $\omega$ .
- However, it's a complex valued function so it's easier to have a plot for the magnitude and one for the phase.
- Because the magnitude plot is in log scale, and the phase plot is in semi-log scale we are able to take advantage of the properties of logs.
  - $\log(XY) = \log(X) + \log(Y)$
- Hence, we can break our transfer function  $H(\omega)$  into a product of familiar transfer functions (simple poles, quadratic zeros, etc. "functional forms"), graph them out individually, and then add them together on the graph.

#### **Bode Plot Steps**

1. Break transfer function into product of functional forms

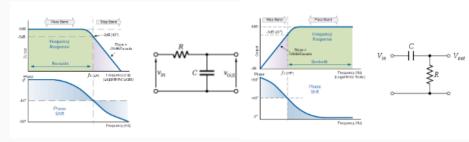
1.1 
$$H(\omega) = \frac{n(\omega)}{d(\omega)} = \frac{(j\omega)^n \alpha_n + (j\omega)^{n-1} \alpha_{n-1} + \dots + j\omega \alpha_1 + \alpha_0}{(j\omega)^n \beta_n + (j\omega)^{n-1} \beta_{n-1} + \dots + j\omega \beta_1 + \beta_0} = \kappa \frac{(j\frac{\omega}{\omega_{21}} + 1)(j\frac{\omega}{\omega_{22}} + 1)\dots}{(j\frac{\omega}{\omega_{\rho_1}} + 1)(j\frac{\omega}{\omega_{\rho_2}} + 1)\dots} = \kappa \frac{(j\frac{\omega}{\omega_{\rho_1}} + 1)(j\frac{\omega}{\omega_{\rho_2}} + 1)\dots}{(j\frac{\omega}{\omega_{\rho_1}} + 1)(j\frac{\omega}{\omega_{\rho_2}} + 1)\dots}$$

- 2. Graph each transfer function individually
- 3. Add them up on the graph (thanks to log scale)



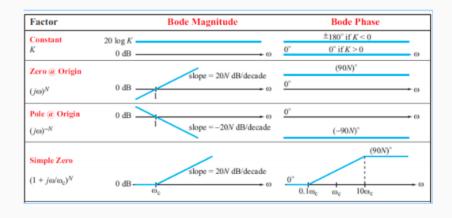
3.1

#### RC Circuits, revisited

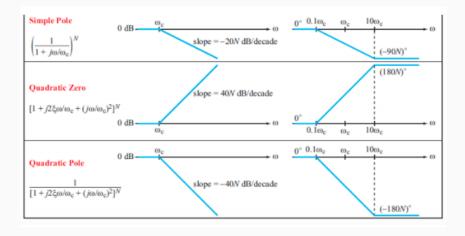


$$H(\omega) = \frac{1}{1+j\omega RC}$$
  $H(\omega) = \frac{j\omega RC}{1+j\omega RC}$ 

#### **Important Functional Forms**



# **Important Functional Forms**



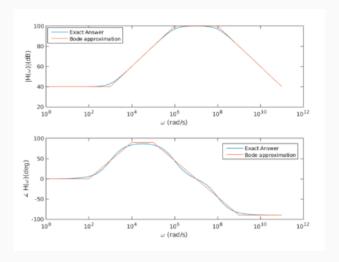
#### **Practice Problem**

$$H(\omega) = \frac{\frac{j\omega}{10} + 100}{1 + \frac{(j\omega)^2}{10^14} + \frac{j\omega}{10^8} + \frac{j\omega}{10^6}}$$

Draw the Bode plot for this transfer function.

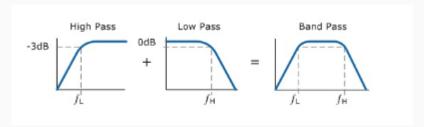
$$H(\omega) = \frac{\frac{j\omega}{10} + 100}{1 + \frac{(j\omega)^2}{10^14} + \frac{j\omega}{10^8} + \frac{j\omega}{10^6}}$$
$$= \frac{100(\frac{j\omega}{10^3} + 1)}{(\frac{j\omega}{10^6} + 1)(\frac{j\omega}{10^8} + 1)}$$

Main idea: break our transfer function into the product of standard forms- a constant, one zero, and two poles



Final Result

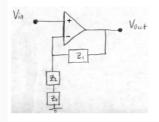
#### Common Filters and their Bode Plots



In the high pass filter, the frequencies greater than the corner frequency have a gain of OdB (so they aren't changed), while frequencies less than the corner frequency have a gain ¡ OdB (so they are multiplied by something less than 1). In the lowpass filter, the opposite is true.

# **Design Problem**

Design an active high pass filter with cutoff frequency of 1KHz and a gain of 1000.



- Hint 1: Find the transfer function of the following circuit.
- Hint 2: Which transfer function is a high pass filter?
  - $H_1(\omega) = 1 + \frac{j\omega R_1 C}{1 + j\omega R_2 C}$   $H_2(\omega) = 1 + \frac{R_1 R_2}{1 + i\omega R_1 C}$
- Hint 3: What is the impedance of a resistor in series with a capacitor?
- Hint 4: Replace the Z's with resistors/capacitors/wires/open

## Design Problem: Hint 1

- Step 1: Apply the golden rules of Op. Amps. This means that  $V^+ = V^-$ .
- Step 2: By ohm's law, we can calculate the current from V-to ground.
  - $\bullet \ \ I = \frac{V_{in}}{Z_2 + Z_3}$
- Step 3: By Ohm's law, we can calculate the output voltage.
  - $V_{out} = V_{in}(1 + \frac{Z_1}{Z_2 + Z_3})$

#### Design Problem: Hint 2 + Hint 3

- The high pass filter transfer function is:  $H_1(\omega)=1+rac{j\omega R_1 C}{1+j\omega R_2 C}$
- To more easily separate the gain from the frequency filtering,
   we can approximate it as (at high frequencies):

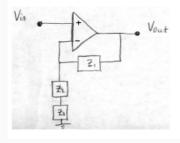
$$H_1(\omega) = \frac{R_1}{R_2} \left( \frac{j\omega R_2 C}{1 + j\omega R_2 C} \right)$$

• A resistor in series with a capacitor has impedances of:

$$R + \frac{1}{j\omega C} = \frac{j\omega RC + 1}{j\omega C}$$

## **Design Problem: Summary of Hints**

Now, use all of the earlier hints to get a high pass filter! (Cutoff Frequency is 1KHz, gain of 1000)



$$V_{out} = V_{in} \left(1 + \frac{Z_1}{Z_2 + Z_3}\right)$$

$$H_1(\omega) = 1 + \frac{j\omega R_1 C}{1 + j\omega R_2 C}$$

$$R + \frac{1}{j\omega C} = \frac{j\omega RC + 1}{j\omega C}$$

# **Design Problem Solution**

We implement the circuit such that it has the transfer function of :

$$H_1(\omega) = 1 + \frac{j\omega R_1 C}{1 + j\omega R_2 C}$$

To meet our specifications:

$$\frac{R_1}{R_2} = 1000$$

$$R_2C = \frac{1}{2\pi(1000)}$$

One Possible Solution:

