

## Description of Numerical Approach

The governing equations for the strain probability density functions  $p_1(s, t)$ ,  $p_2(s, t)$ , and  $p_3(s, t)$  are

$$\begin{aligned}\frac{\partial p_1}{\partial t} + \frac{dL}{dt} \frac{\partial p_1}{\partial s} &= k_a \delta(s) U_{NR} P - \tilde{k}_d p_1 - \tilde{k}_1 e^{-\alpha_1 s} p_1 + k_{-1} e^{+\alpha_1 s} p_2 \\ \frac{\partial p_2}{\partial t} + \frac{dL}{dt} \frac{\partial p_2}{\partial s} &= \tilde{k}_1 e^{-\alpha_1 s} p_1 - k_{-1} e^{+\alpha_1 s} p_2 - k_2 e^{-\alpha_2 s} p_2 + \tilde{k}_{-2} p_3 \\ \frac{\partial p_3}{\partial t} + \frac{dL}{dt} \frac{\partial p_3}{\partial s} &= k_2 e^{-\alpha_2 s} p_2 - \tilde{k}_{-2} p_3 - \tilde{k}_3 e^{\alpha_3(s+s_3)^2} p_3.\end{aligned}\quad (1)$$

These equations are simulated by discretizing the probability density variables on a grid

$$p_1(s, t) \approx P_1(n, t)$$

$$p_2(s, t) \approx P_2(n, t)$$

$$p_3(s, t) \approx P_3(n, t)$$

where  $n \in (0, N)$  represents the index on the spatial grid and  $P_1(n, t)$ ,  $P_2(n, t)$ , and  $P_3(n, t)$  are discretized approximations of  $p_1(s, t)$ ,  $p_2(s, t)$ , and  $p_3(s, t)$ . Since the velocities simulated in this study are always negative (associated with muscle shortening), the strain variable  $s$  is discretized over a range of negative values:

$$s = -nh, \quad n = 0, 1, \dots, N$$

where  $h$  is the space step. Equation (1) is numerically integrated over a finite times step  $\Delta t$  by splitting the space and time operators into separate steps.

Specifically, the two operators are, an advection operator:

$$\begin{aligned}\frac{\partial P_1}{\partial t} + \frac{dL}{dt} \frac{\partial P_1}{\partial s} &= 0 \\ \frac{\partial P_2}{\partial t} + \frac{dL}{dt} \frac{\partial P_2}{\partial s} &= 0 \\ \frac{\partial P_3}{\partial t} + \frac{dL}{dt} \frac{\partial P_3}{\partial s} &= 0\end{aligned}\quad (2)$$

and an ordinary-differential equation operator:

$$\begin{aligned}\frac{dP_1(n, t)}{dt} &= k_a \delta(s) U_{NR} P - \tilde{k}_d P_1 - \tilde{k}_1 e^{+\alpha_1 nh} P_1 + k_{-1} e^{-\alpha_1 nh} P_2 \\ \frac{dP_2(n, t)}{dt} &= \tilde{k}_1 e^{+\alpha_1 nh} P_1 - k_{-1} e^{-\alpha_1 nh} P_2 - k_2 e^{+\alpha_2 nh} P_2 + \tilde{k}_{-2} P_3 \\ \frac{dP_3(n, t)}{dt} &= k_2 e^{+\alpha_2 nh} P_2 - \tilde{k}_{-2} P_3 - \tilde{k}_3 e^{\alpha_3(-nh+s_3)^2} P_3.\end{aligned}\quad (3)$$

For a finite time step  $\Delta t$ , intermediate estimates of strain densities are first calculated by integrating Equation (2) using an upwind differencing approach, with the time step

$$P_1^*(n, t) = P_1(n + 1, t)$$

$$P_2^*(n, t) = P_2(n + 1, t)$$

$$P_3^*(n, t) = P_3(n + 1, t)$$

where the time step is set to match the space step based on the constant velocity:  $\Delta t = -h/(dL/dt)$ . Next the ordinary differential equations of Equation (3) are integrated using a differential equation integrator package over the time interval  $(0, \Delta t)$  with initial condition  $P_1^*(n, t)$ ,  $P_2^*(n, t)$ ,  $P_3^*(n, t)$  to yield  $P_1(n, t + \Delta t)$ ,  $P_2(n, t + \Delta t)$ ,  $P_3(n, t + \Delta t)$ .