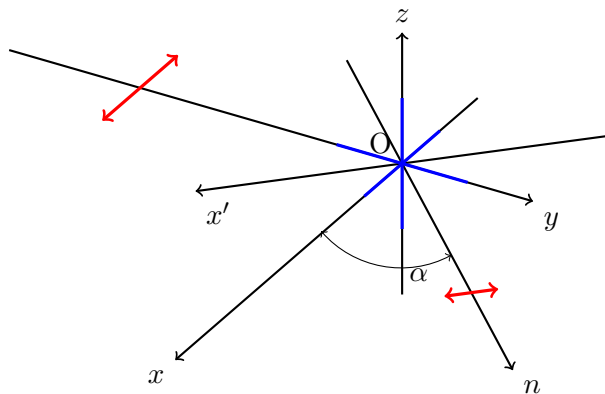


Radiation of a non relativistic charge accelerated free or bound  
(electron in the classical description with oscillators)

$$\vec{E}_{rad} \propto \vec{n} \times (\vec{n} \times \vec{a}) \quad (1)$$

# Geometric representation

Electron accelerated by a force due to the electric field of polarized radiation (Thomson scattering)

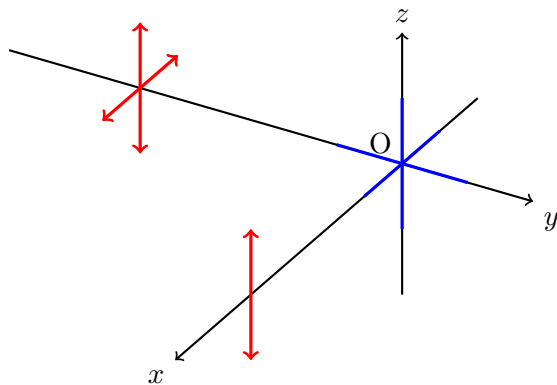


$\vec{x}, \vec{y}, \vec{n}, \vec{x}'$  same plane

- Anisotropic radiation creates a net dipole moment in scattering atoms or molecules (Rayleigh scattering)
- shorter wavelengths more scattered

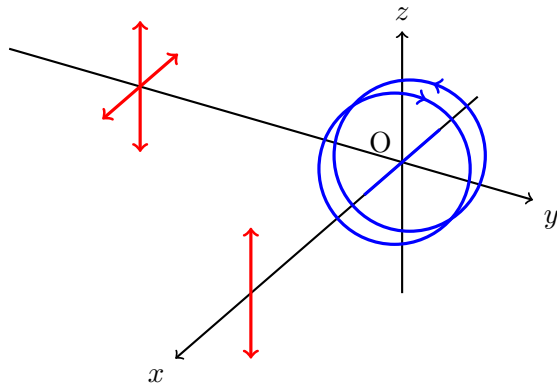
# Geometric representation

Incident unpolarized light, cartesian oscillators



# Geometric representation

Incident unpolarized light, circular oscillators

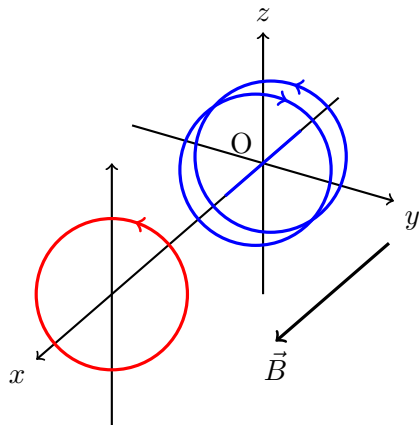


oscillators

coherence between circular

# Geometric representation

Magnetic field along l.o.s , circular oscillators



ciclotron radiation

# Geometric representation

Unpolarized incident light and magnetic field along l.o.s , circular oscillators

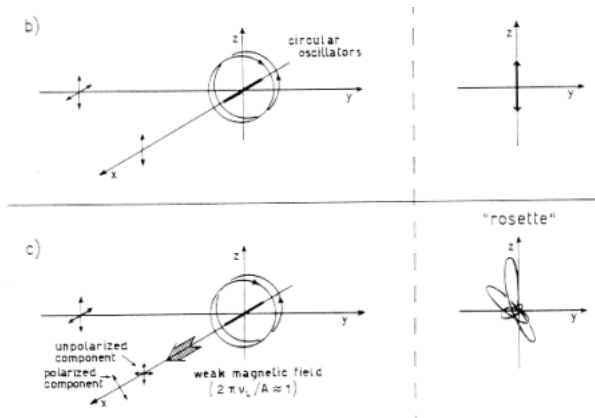


Figure 1: Incoherence in the circular oscillators - Hanle effect

# Quantic representation (bound electron)

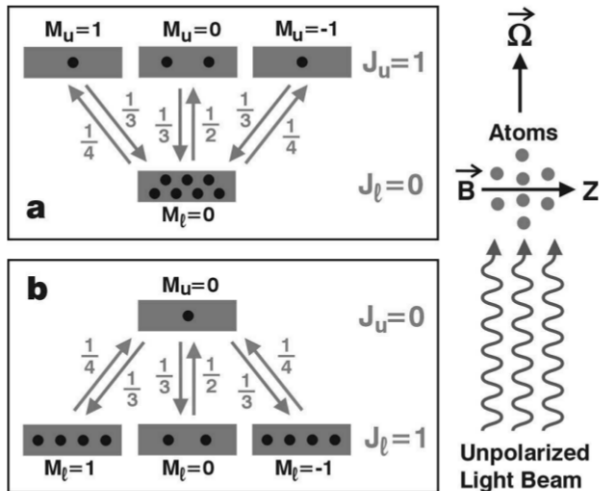


Figure 2: Emission/Absorption



# Quantic representation (bound electron)

- Correspondence between atomic transitions and radiation components
  - $\Delta M = 1, \sigma_+ (\nu + \nu_L)$ , clockwise oscillator
  - $\Delta M = 0, \pi (\nu)$ , x oscillator
  - $\Delta M = -1, \sigma_- (\nu - \nu_L)$ , anticlockwise oscillator
- similar mechanism(Fig 2 b) for absorption lines (experimentally done by observing same line in filaments: absorption and in prominences: emission)
- Anisotropic radiation populates/depopulates twice  $M=0$  sublevel
- Magnetic field:
  - splitting of energy level  $J$  into  $2J+1$  sublevels (Zeeman effect)
  - creates incoherence between population of levels corresponding to  $\sigma$  components which leads to incoherence in the 2 circular oscillators (Hanle effect)

## Second specter of the sun

With the same geometry: Q represents linear polarization parallel to the closest limb

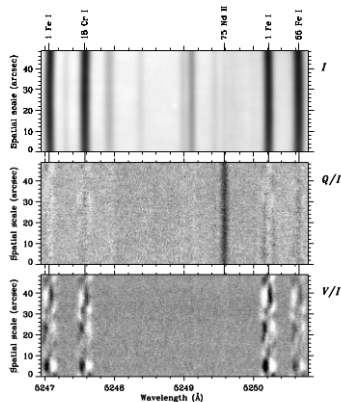


Fig. 1. Examples of the different characters of the ordinary intensity spectrum (Stokes  $I$ , top diagram), the linearly polarized spectrum (Stokes  $Q/I$ , here called the "second solar spectrum"), and the circularly polarized spectrum (Stokes  $V/I$ , bottom diagram). While the circular polarization shows Zeeman-effect signatures due to magnetic canopies at the supergranulation size scale, the linear polarization is dominated by the signature of coherent scattering, which here unexpectedly is due to a transition in ionized neodymium. The recording

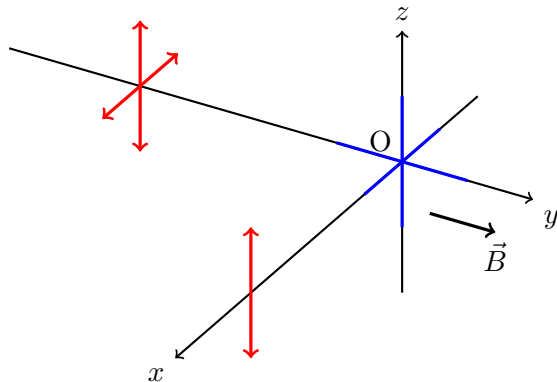
# Second specter of the sun

Unknown I 0 level

Eliminating Q signal generated by other mechanisms

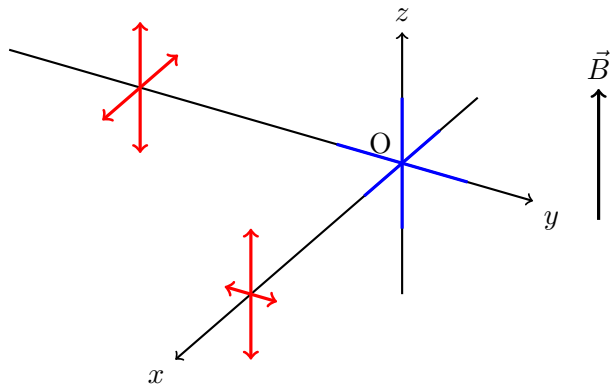
- atmospheric seeing and CCD noise: resolved by modulating the signal at high frequency: 2 modulators one at 42 kHz and 84 kHz
- position near the (geographic) north/south pole where influence of the magnetic field is smaller
- cross talk Stokes V generated by the magnetic field recognized (sign reversed with spatial period = the size of a granulation) which confirms the existence of horizontal magnetic field oriented along the line of sight
- vertical magnetic field
- horizontal magnetic field perpendicular to line of sight
- polarization by reflection

# Vertical magnetic field



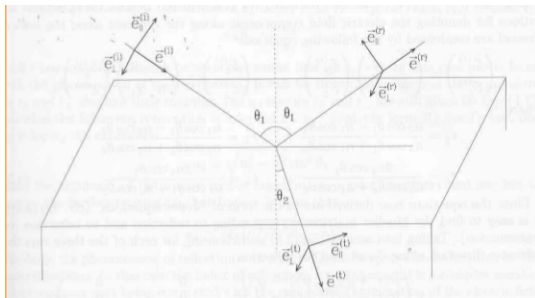
effect adding to that of the radiation: linear polarization (no V signal, no Hanle effect)

# Horizontal magnetic field, perpendicular to l.o.s.



horizontal component(Hanle effect), no V signal

# Polarization by reflection



# Polarization by reflection

$$\begin{pmatrix} E_{\parallel}^{(r)} \\ E_{\perp}^{(r)} \end{pmatrix} = \begin{pmatrix} r_{\parallel} & 0 \\ 0 & r_{\perp} \end{pmatrix} \begin{pmatrix} E_{\parallel}^{(i)} \\ E_{\perp}^{(i)} \end{pmatrix}$$

$$r_{\parallel} = \frac{n_2 \cos \theta_1 - n_1 \cos \theta_2}{n_2 \cos \theta_1 + n_1 \cos \theta_2}$$

$$r_{\perp} = \frac{n_1 \cos \theta_1 - n_2 \cos \theta_2}{n_1 \cos \theta_1 + n_2 \cos \theta_2}$$

$$\begin{pmatrix} E_{\parallel}^{(t)} \\ E_{\perp}^{(t)} \end{pmatrix} = \begin{pmatrix} t_{\parallel} & 0 \\ 0 & t_{\perp} \end{pmatrix} \begin{pmatrix} E_{\parallel}^{(i)} \\ E_{\perp}^{(i)} \end{pmatrix}$$

$$t_{\parallel} = \frac{2n_1 \cos \theta_1}{n_2 \cos \theta_1 + n_1 \cos \theta_2}$$

$$t_{\perp} = \frac{2n_1 \cos \theta_1}{n_1 \cos \theta_1 + n_2 \cos \theta_2}$$

$$n_1 = 1 < n_2, \cos \theta_1 < \cos \theta_2 \implies (n_2 \cos \theta_1 = \cos \theta_2 \implies r_{\parallel} = 0)$$

$$\sin \theta_1 = n_2 \sin \theta_2 \implies 1 - (\cos \theta_1)^2 = n_2^2 (1 - \cos^2 \theta_1) \implies 1 - \cos^2 \theta_1 = n_2^2 (1 - \cos^2 \theta_1)$$

$$\theta_1 = \arccos\left(\sqrt{\frac{n_2^2 - 1}{n_2^4 - 1}}\right) \implies \text{reflected light polarized along } r_{\perp}$$

- instrument: ZIMPOL, the Zurich Imaging Stokes Polarimeter
- telescope: McMath-Pierce(National Solar Observatory (Kitt Peak, USA).)
- accuracy:  $10^{-5}$
- April 1995, slit positioned 5 seconds of arc inside the north polar limb (where the cosine  $\mu$  of the heliocentric angle is 0.1)