Two-fluid simulations of waves and reconnection with Mancha code

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Sun atmosphere layers

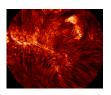
Photosphere

- collisions dominated: LTE, MHD
- relatively easy observations
- diagnostics techniques well developed Chromosphere
 - not fully collisionally coupled: NLTE, No MHD (frequently not taken into account)
 - very few spectral lines
 - complicated radiative diagostics

Corona

- magnetically dominated
- very low density
- all ionized, MHD can be applied







2 fluids model

- 13 variables: 10 variables p , ρ and v of the 2 fluids: charges(_c) and neutrals(_n) + magnetic field
- hydrostatic equilibrium(variables: $p_{c0}, p_{n0}, \rho_{c0}, \rho_{n0}, \vec{B_0}$), $\vec{v_{c0}} = \vec{v_{n0}} = 0$
 - charges:

$$\rho_{c0}\vec{g} - \vec{\nabla}p_{c0} + \frac{1}{\mu_0}(\nabla \times \vec{B_0}) \times \vec{B_0} = 0 \tag{1}$$

• neutrals:

$$\rho_{n0}\vec{g} - \vec{\nabla}p_{n0} = 0 \tag{2}$$

• 13 partial differential equations for the evolution of the perturbation (variables: $p_{c1}, p_{n1}, \vec{v_{c1}}, \vec{v_{n1}}, \rho_{c1}, \rho_{n1}, \vec{B1}$)

Boltzmann equation

$$\frac{\partial f_{\alpha}}{\partial t} + \vec{v} \cdot \vec{\nabla} f_{\alpha} - \vec{a} \cdot \vec{\nabla}_{v} f_{\alpha} = (\frac{\partial f_{\alpha}}{\partial t})_{coll}$$
(3)

with $\alpha \in i, e, n$ (charges = e+i)

radiation not taken into account

single ionized H plasma $(n_i = n_e)$

collision terms:
$$C_{\alpha} \stackrel{\text{not}}{=} (\frac{\partial f_{\alpha}}{\partial t})_{coll} = C_{\alpha}^{elastic} + C_{\alpha}^{inelastic}$$

In order to calculate the 0^{th} , first and second moment of Boltzmann equation we need to calculate

$$S_{\alpha} \stackrel{\text{not}}{=} \int_{V} C_{\alpha} d^{3} \vec{v}$$

$$\vec{R}_{\alpha} \stackrel{\text{not}}{=} \int_{V} \vec{v} C_{\alpha} d^{3} \vec{v}$$

$$M_{\alpha} \stackrel{\text{not}}{=} \int_{V} v^{2} C_{\alpha} d^{3} \vec{v}$$

Collision terms(inelastic)

$$C_{\alpha}^{inelastic} = \sum_{\alpha'} \left(n_{\alpha'} C_{\alpha'\alpha}^{inelastic} - n_{\alpha} C_{\alpha\alpha'}^{inelastic} \right)$$

$$C_{\alpha\alpha'}^{inelastic} = \sum_{\beta} C_{\alpha\alpha',\beta}^{inelastic}$$

$$C_{\alpha\alpha',\beta}^{inelastic} = \sigma_{\alpha\alpha'} f_{\beta} v_{\beta}$$
where $\sigma_{\alpha'} = \sigma_{\alpha'} f_{\beta} v_{\beta}$

where $\sigma_{\alpha\alpha'} = \sigma_{\alpha\alpha'}(v_{\beta})$ is the collisional cross section of α and β considering ionization and recombination processes for inelastic collisions:

neutrals:

$$\begin{split} &C_{n}^{inelastic} = n_{i}C^{rec} - n_{n}C^{ion} \\ &\text{where } C^{ion} \stackrel{\text{not}}{=} C^{inelastic}_{ni,e} = \sigma_{ion}f_{e}v_{e} \\ &C^{rec} \stackrel{\text{not}}{=} C^{inelastic}_{in,e} = \sigma_{rec}f_{e}v_{e} \\ &\sigma_{ion} \stackrel{\text{not}}{=} \sigma_{ni} = \sigma_{ni}(v_{e}) \\ &\sigma_{rec} \stackrel{\text{not}}{=} \sigma_{in} = \sigma_{in}(v_{e}) \end{split}$$

charges:

$$C_c^{inelastic} = -C_n^{inelastic}$$

Collision terms(inelastic)

0^{th} moment

$$S_{\alpha\alpha',\beta}^{inelastic} = \int_{V} \sigma_{\alpha\alpha'} f_{\beta} v_{\beta} d^{3} \vec{v} = n_{\beta} < \sigma_{\alpha\alpha'} v_{\beta} > S^{ion} = n_{e} < \sigma_{ion} v_{e} > S^{rec} = n_{e} < \sigma_{rec} v_{e} >$$

$$x^{2} + 2xy + y^{2}$$

expressions for the collision sections in Leake article:

$$<\sigma_{ion}v_{e}> = \frac{1}{\sqrt{T_{e}^{*}}} 2.6 \cdot 10^{-19} m^{3}/s$$

$$<\sigma_{rec}v_{e}> = A \frac{1}{X + \frac{\phi_{ion}}{T_{e}^{*}}} (\frac{\phi_{ion}}{T_{e}^{*}})^{K} e^{-\frac{\phi_{ion}}{T_{e}^{*}}} m^{3}/s$$

where $\phi_{ion} = 13.6 \text{eV}$, T_e^* is electron temperature in eV and they define $A = 2.91 \cdot 10^{-14}$, K = 0.39, X = 0.232

$$S_n^{inelastic} = n_i S^{rec} - n_n S^{ion}$$

Collision terms(inelastic)

First moment

$$\begin{split} R_{\alpha\alpha',\beta}^{inelastic} &= \int_{V} \vec{v_{\alpha}} \sigma_{\alpha\alpha'} f_{\beta} v_{\beta} d^{3} \vec{v} \\ \vec{v_{\alpha}} &= \vec{u_{\alpha}} + \vec{w_{\alpha}} \\ \int_{V} \vec{w_{\alpha}} \sigma_{\alpha\alpha'} f_{\beta} v_{\beta} d^{3} \vec{v} &= 0 \\ \int_{V} \vec{u_{\alpha}} \sigma_{\alpha\alpha'} f_{\beta} v_{\beta} d^{3} \vec{v} &= \vec{u_{\alpha}} S_{\alpha\alpha',\beta}^{inelastic} \\ R_{n}^{inelastic} &= n_{i} \vec{u_{i}} S^{rec} - n_{n} \vec{u_{n}} S^{ion} \end{split}$$

Second moment

$$\begin{split} M_{\alpha\alpha',\beta}^{inelastic} &= \int_{V} v_{\alpha}^{2} \sigma_{\alpha\alpha'} f_{\beta} v_{\beta} d^{3} \vec{v} \\ \int_{V} u_{\alpha}^{2} \sigma_{\alpha\alpha'} f_{\beta} v_{\beta} d^{3} \vec{v} &= u_{\alpha}^{2} S_{\alpha\alpha',\beta}^{inelastic} \\ \int_{V} 2 \vec{w_{\alpha}} \vec{u_{\alpha}} \sigma_{\alpha\alpha'} f_{\beta} v_{\beta} d^{3} \vec{v} &= 0 \\ \int_{V} w_{\alpha}^{2} \sigma_{\alpha\alpha'} f_{\beta} v_{\beta} d^{3} \vec{v} &= \frac{3k_{B} T_{\alpha}}{m_{\alpha}} S_{\alpha\alpha',\beta}^{inelastic} \\ M_{n}^{inelastic} &= n_{i} u_{i}^{2} S^{rec} - n_{n} u_{n}^{2} S^{ion} + 3k_{B} (\frac{n_{i} T_{i}}{m_{i}} S_{rec} - \frac{n_{n} T_{n}}{m_{n}} S_{ion}) \end{split}$$

Collision terms(elastic)

$$C_{lpha}^{elastic} = \sum_{eta} C_{lphaeta}^{elastic}$$
 neutrals: $C_{n}^{elastic} = C_{ni}^{elastic} + C_{ne}^{elastic}$ charges: $C_{n}^{elastic} = -C_{n}^{elastic}$

$$C_c^{elastic} = -C_n^{elasti}$$

 0^{th} moment

$$S_{\alpha}^{elastic} = \int_{V} C_{\alpha}^{elastic} d^{3}\vec{v} = 0$$

First moment

$$R_{\alpha}^{elastic} = \int_{V} \vec{v_{\alpha}} C_{\alpha}^{elastic} d^{3} \vec{v} = \int_{V} \vec{w_{\alpha}} C_{\alpha}^{elastic} d^{3} \vec{v}$$

 $\int_{V} \vec{w_{\alpha}} C_{\alpha\beta}^{elastic} d^{3} \vec{v} = n_{\alpha} \nu_{\alpha\beta} (\vec{u_{\beta}} - \vec{u_{\alpha}}) \text{ where } \nu_{\alpha\beta} \text{ is the collision frequency}$

and for
$$\alpha \in i, e$$
 and $\beta = n$ it is expressed as: $n_{\beta} \sqrt{\frac{8k_B T_{\alpha\beta}}{\pi m_{\alpha\beta}}} \Sigma_{\alpha\beta}$

where $m_{\alpha\beta} = \frac{m_{\alpha} + m_{\beta}}{2}$, $T_{\alpha\beta} = \frac{T_{\alpha} + T_{\beta}}{2}$ and $\Sigma_{\alpha\beta}$ is the cross section for

$$(n_{\alpha}\nu_{\alpha\beta} = n_{\beta}\nu_{\beta\alpha})$$

$$R_n^{elastic} = n_i(\nu_{in} + \nu_{en})(\vec{u_c} - \vec{u_n})$$

Collision terms(elastic)

$$u_{in} = n_n \sqrt{\frac{8k_B T_{ni}}{\pi m_{ni}}} \Sigma_{ni}, \ \nu_{en} = n_n \sqrt{\frac{8k_B T_{ne}}{\pi m_{ne}}} \Sigma_{ne}$$
and we use the values: $\Sigma_{ne} = 10^{-19} m^2, \Sigma_{ni} = 5 \cdot 10^{-19} m^2$
Second moment

 $M_{\alpha}^{elastic} = \int_{V} v_{\alpha}^{2} C_{\alpha}^{elastic} d^{3} \vec{v} = 2 \vec{u_{\alpha}} \int_{V} \vec{w_{\alpha}} C_{\alpha}^{elastic} d^{3} \vec{v} + \int_{V} w_{\alpha}^{2} C_{\alpha}^{elastic} d^{3} \vec{v}$ We neglect the second term for the moment (braginskii 'heat generation') then:

$$M_{\alpha}^{elastic} = 2\vec{u_{\alpha}}R_{\alpha}^{e\vec{lastic}}$$

From Boltzmann equation to the evolution of the perturbations

transport equation:

$$\frac{\partial (n_{\alpha} < \chi >_{\alpha})}{\partial t} + \nabla \cdot (n_{\alpha} < \chi \vec{v} >_{\alpha}) - n_{\alpha} < \vec{a} \cdot \vec{\nabla_{v}} \chi >_{\alpha} = \int_{V} \chi (\frac{\partial f_{\alpha}}{\partial t})_{coll} d^{3} \vec{v}$$

$$\tag{4}$$

with $\alpha \in n, c$ and $\chi = m_{\alpha}, m_{\alpha} \vec{v_{\alpha}}, \frac{1}{2} m_{\alpha} v_{\alpha}^2$ in the Boltzmann transport equation and using the 0^{th} , first and second moment of Boltzmann equation will result in equations for $u = \rho_{\alpha}, \rho_{\alpha} \vec{v_{\alpha}}, \epsilon_{\alpha} + \frac{1}{2} \rho_{\alpha} v_{\alpha}^2$ The equations from the code:

$$\frac{\partial u}{\partial t} + \nabla \cdot \vec{F_u} = S_u \tag{5}$$

$$\vec{F_u} = \vec{F_u}^{id} - \vec{F_u}^{idff}, S_u = S_u^{id} + S_u^{coll} + S_u^{diff}$$

Diffusivity(artificial)

Artificial diffusivity coefficients (as in mancha 1 fluid): $\nu_u^{diff}_{x_i} = \nu_u^{diff_shock}_{x_i} + \nu_u^{diff_const}_{x_i} + \nu_u^{diff_var}_{x_i}$ for $i \in 1, 2, 3$ $\nu_u^{diff_shock} = \nu_u^{shock} \cdot max(|\nabla \cdot \vec{v}_u^{shock}|, 0.5) \cdot dx_i^2 \text{ for } \nabla \cdot \vec{v}_u^{shock} < 0 \text{ and }$ 0 otherwise where we define $\vec{v}_{\rho_{c1}}^{shock} = \vec{v}_{\epsilon_{c1}}^{shock} = \vec{v}_{\vec{v}_{c1}}^{shock} = \vec{v}_{c1}^{,i}, \ \vec{v}_{\rho_{n1}}^{shock} = \vec{v}_{\epsilon_{n1}}^{shock} = \vec{v}_{\vec{v}_{n1}}^{shock} = \vec{v}_{n1}^{,i},$ $\vec{v}_{\vec{B}_1}^{shock} = \vec{v}_{c\perp_{\vec{B}}} \text{ or } \vec{v}_{c\perp_{\vec{B}_0}}$ $\nu_u^{diff_var}_{r_i} = \nu_u^{var}_{r_i} \cdot vflow_u \cdot dx_i \cdot hyper_u^{x_i}$ **hyper** is defined (for example for i = 1): for each $k \in 2$...number of points discretized in dimension x_1 - 2: $hyper_{u}^{x_{i}}(k,:,:) = \frac{max(\Delta 3_{u}(k-1,:,:),\Delta 3_{u}(k,:,:),\Delta 3_{u}(k+1,:,:))}{max(\Delta 1_{u}(k-1,::),\Delta 1_{u}(k,:,:),\Delta 1_{u}(k+1,:,:))} \text{ where}$ $\Delta 3_u(k,:,:) = |3(u(k+1,:,:) - u(k,:,:)) - (u(k+2,:,:) - u(k-1,:,:))|$ and $\Delta 1_u(k,:,:) = |u(k+1,:,:) - u(k,:,:)|$ (we use T_{α} for $u = \epsilon_{\alpha}$ when calculating hyper)

 $vflow_u = |\vec{v_{c1}}| + c_{sc} + v_A$ for u related to charges and magnetic field

and $v f low_u = |\vec{v_{n1}}| + c_{sn}$ for u related to neutrals

Diffusivity(artificial)

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\nu_u^{const\_var}_{x_i} = \nu_u^{const}_{x_i} \cdot vel_u \cdot dx_i \cdot const

(const is a matrix introduced as a h5 file)
vel_u = c_{sc} + v_A for u related to charges and magnetic field and vel_u = c_{sn} for u related to neutrals
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Fluxes due to the diffusivity:

$$\begin{aligned} F_{\rho_{\alpha}}^{diff}{}_{x_{i}} &= \nu_{\rho_{\alpha}}^{diff}{}_{x_{i}} \cdot \frac{\partial \rho_{\alpha 1}}{\partial x_{i}} \\ F_{\epsilon_{\alpha}}^{diff}{}_{x_{i}} &= \rho_{\alpha} \cdot \nu_{\epsilon_{\alpha}}^{diff}{}_{x_{i}} \cdot \frac{\partial T_{\alpha 1}}{\partial x_{i}} \end{aligned}$$

symmetric diffusivity matrix for the velocities (artificial viscosity):

$$F_{\rho_{\alpha}v_{\alpha_{x_{j}}x_{i}}}^{diff} = \frac{1}{2}\rho_{\alpha} \cdot \left(\nu_{\rho_{\alpha}v_{\alpha_{x_{j}}x_{i}}}^{diff} \cdot \frac{\partial v_{\alpha_{x_{j}}}}{\partial x_{i}} + \nu_{\rho_{\alpha}v_{\alpha_{x_{i}}x_{j}}}^{diff} \cdot \frac{\partial v_{\alpha_{x_{i}}}}{\partial x_{j}}\right)$$

magnetic artificial diffusivity:

$$\begin{split} E_{x_{1}}^{artif_diff} &= F_{B_{x_{3}}}^{diff} - F_{B_{x_{2}}}^{diff}, \ E_{x_{2}}^{artif_diff} = F_{B_{x_{1}}}^{diff} - F_{B_{x_{3}}}^{diff}, \\ E_{x_{3}}^{artif_diff} &= F_{B_{x_{2}}}^{diff} - F_{B_{x_{1}}}^{diff}, \\ \text{where } F_{B_{x_{j}}}^{diff} &= \nu_{B_{x_{j}}}^{diff}, \\ &= \nu_{B_{x_{j}}}^{diff} \cdot \frac{\partial B_{x_{j}}}{\partial x_{i}}. \end{split}$$

Continuity equations

$$\begin{split} \frac{\partial \rho_{n1}}{\partial t} + \nabla \cdot (\rho_n \vec{v_n} - \vec{F}_{\rho_n}^d) &= S_n \\ \frac{\partial \rho_{c1}}{\partial t} + \nabla \cdot (\rho_c \vec{v_c} - \vec{F}_{\rho_c}^d) &= -S_n \\ S_n &= \rho_c S^{rec} - \rho_n S^{ion} \end{split}$$

where

Continuity equations2

$$F_{\rho_{\alpha}}^{id} = \rho_{\alpha} \vec{v_{\alpha}}$$

$$S_{n}^{'} \stackrel{\text{not}}{=} S_{\rho_{n}}^{coll} = m_{H}(n_{i}S^{rec} - n_{n}S^{ion}) = \rho_{c}S^{rec} - \rho_{n}S^{ion}$$

$$S_{\rho_{c}}^{coll} = -S_{n}^{'}$$

Momentum equations

$$\frac{\partial(\rho_{n1}\vec{v_n})}{\partial t} + \nabla \cdot (\begin{bmatrix} p_{n1} + \rho_n v_{nx}^2 & \rho_n v_{nx} v_{ny} & \rho_n v_{nx} v_{nz} \\ \rho_n v_{nx} v_{ny} & p_{n1} + \rho_n v_{ny}^2 & \rho_n v_{ny} v_{nz} \\ \rho_n v_{nx} v_{nz} & \rho_n v_{nx} v_{nz} & p_{n1} + \rho_n v_{nz}^2 \end{bmatrix} - \bar{\vec{F}}_{\vec{v_n}}^{\vec{d}}) = -\rho_n \vec{g} + \vec{K}_n$$

$$\frac{\partial(\rho_{c1}\vec{v_c})}{\partial t} + \nabla \cdot (\bar{\bar{A}} - \bar{\bar{F}}_{\vec{v_c}}^d) = -\rho_c \vec{g} - \vec{R_n}$$

where
$$A_{ij} = \rho_n v_{cx_j} v_{cx_j} - \frac{1}{\mu_0} (B_{x_i0} B_{x_j1} + B_{x_j0} B_{x_i1} + B_{x_i1} B_{x_j1})$$

$$A_{ii} = \rho_n v_{cx_i}^2 + p_{c1} - \frac{1}{2\mu_0} (B_{x_11} (B_{x_11} + 2B_{x_10}) + B_{x_21} (B_{x_21} + 2B_{x_20}) + B_{x_31} (B_{x_31} + 2B_{x_30}))$$

$$\vec{R_n} = \rho_c \vec{v_c} S^{rec} - \rho_n \vec{v_n} S^{ion} + \rho_c \rho_n \alpha^{elastic} (\vec{v_c} - \vec{v_n})$$

$$\alpha^{elastic} \stackrel{\text{not}}{=} \frac{1}{m_n^2} (\sqrt{\frac{8k_B T_{nc}}{\pi m_{in}}} m_{in} \Sigma_{in} + \sqrt{\frac{8k_B T_{nc}}{\pi m_{en}}} m_{en} \Sigma_{en})$$

Momentum equations2

symmetric flux matrices:

neutrals

$$F_{\rho_n v_{n_{x_i} x_j}}^{id} = \rho_n v_{n_{x_j}} v_{n_{x_j}}$$

$$F_{\rho_n v_{n_{x_i} x_i}}^{id} = \rho_n v_{n_{x_j}}^2 + p_{n_1}$$

charges

$$\begin{split} F^{id}_{\rho_c v_{cx_i} x_j} &= \rho_n v_{cx_j} v_{cx_j} - \frac{1}{\mu_0} (B_{x_i 0} B_{x_j 1} + B_{x_j 0} B_{x_i 1} + B_{x_i 1} B_{x_j 1}) \\ F^{id}_{\rho_c v_{cx_i} x_i} &= \rho_n v_{cx_i}^2 + p_{c1} - \frac{1}{2\mu_0} (B_{x_1 1} (B_{x_1 1} + 2 B_{x_1 0}) + B_{x_2 1} (B_{x_2 1} + 2 B_{x_2 0}) + B_{x_3 1} (B_{x_3 1} + 2 B_{x_3 0})) \\ \text{with gravity: } S^{id}_{\rho_\alpha \vec{v_\alpha}} &= -\rho_\alpha \vec{g} \\ \text{collision sources:} \\ \vec{R}^{\prime}_n &\stackrel{\text{not}}{=} S^{coll}_{\rho_n \vec{v_\alpha}} = m_H (n_i \vec{v_c} S^{rec} - n_n \vec{v_n} S^{ion} + \vec{R}^{elastic}_n) \end{split}$$

$$\vec{R'_n} \stackrel{\text{not}}{=} S_{\rho_n \vec{v_n}}^{coll} = m_H (n_i \vec{v_c} S^{rec} - n_n \vec{v_n} S^{ion} + \vec{R_n}^{elastic})$$

$$= \rho_c \vec{v_c} S^{rec} - \rho_n \vec{v_n} S^{ion} + \rho_c \rho_n \alpha^{elastic} (\vec{v_c} - \vec{v_n})$$
where $\alpha^{elastic} \stackrel{\text{not}}{=} \frac{1}{m_n^2} (\sqrt{\frac{8k_B T_{nc}}{\pi m_{in}}} m_{in} \Sigma_{in} + \sqrt{\frac{8k_B T_{nc}}{\pi m_{en}}} m_{en} \Sigma_{en})$

$$S_{\alpha, \vec{v_c}}^{coll} = -\vec{R'_n}$$

Momentum equations

Stiff terms
$$\rho_c \rho_n \alpha^{elastic} (\vec{v_c} - \vec{v_n})$$
 calculated as: $\frac{\rho_c \rho_n \alpha^{elastic} dt}{1 + \alpha^{elastic} (\rho_c + \rho_n) dt} (\vec{v_c} - \vec{v_n})$

Total energy equations

$$\begin{split} E_c &= \epsilon_c + \frac{1}{2} \rho_c v_c^2 + \frac{1}{2\mu_0} B^2 \\ E_n &= \epsilon_n + \frac{1}{2} \rho_n v_n^2 \\ \vec{F}_{E_c}^{id} &= (E_c + p_c + \frac{B^2}{2\mu_0}) \vec{v_c} - \frac{1}{\mu_0} (\vec{v_c} \cdot \vec{B}) \vec{B} \\ \vec{F}_{E_n}^{id} &= (E_n + p_n) \vec{v_n} \\ \text{with gravity: } S_{E_\alpha}^{id} &= \rho_\alpha \vec{v_\alpha} \vec{g} \\ S_{E_n}^{coll} &= M_n'^{inelastic} + m_H \vec{v_n} \vec{R_n}^{elastic} \\ \text{where} \\ M_n'^{inelastic} &\stackrel{\text{not}}{=} m_H (\frac{1}{2} n_i v_c^2 S^{rec} - \frac{1}{2} n_n v_n^2 S^{ion} + \frac{3}{2} k_B (\frac{n_i T_c}{m_i} S_{rec} - \frac{n_n T_n}{m_n} S_{ion})) \\ &= \frac{1}{2} \rho_c v_c^2 S^{rec} - \frac{1}{2} \rho_n v_n^2 S^{ion} + \frac{3k_B}{2m_H} (\rho_c T_c S_{rec} - \rho_n T_n S_{ion}) \\ S_{E_c}^{coll} &= -M_n'^{inelastic} + m_H \vec{v_c} \vec{R_n}^{elastic} \\ \vec{F}_{E_\alpha}^{diff} &= \vec{F}_{\epsilon\alpha}^{diff} - \frac{1}{\mu_0} \vec{E}^{diff} \times \vec{B} + \bar{\bar{F}}_{\rho_n \vec{v_\alpha}}^{diff} \cdot \vec{v_\alpha} \end{split}$$

Internal energy equations

 $S_{\epsilon}^{diff} = \vec{i} \cdot \vec{E}^{diff}$

internal energy $\epsilon_{\alpha} = \frac{p_{\alpha}}{\gamma - 1}$ (If we take into account the ionization energy: $\epsilon_{c} = \frac{p_{c}}{\gamma - 1} + n_{e}\phi_{ion}$) $\vec{F}_{\epsilon_{\alpha}}^{id} = \epsilon_{\alpha}\vec{v_{\alpha}}$ $S_{\epsilon_{\alpha}}^{id} = p_{\alpha}\nabla \cdot \vec{v_{\alpha}}$ $S_{\epsilon_{n}}^{coll} = S_{E_{n}}^{coll} - \vec{R_{n}'} \cdot \vec{v_{n}} + \frac{1}{2}v_{n}^{2}S_{n}'$ $S_{\epsilon_{c}}^{coll} = S_{E_{c}}^{coll} + \vec{R_{n}'} \cdot \vec{v_{c}} - \frac{1}{2}v_{c}^{2}S_{n}'$

Ohm law and induction equation

$$\begin{split} \vec{E} &= -\vec{v} \times \vec{B} - \vec{E}^{artif_diff} + \vec{E}^{plasma_diff} \\ \vec{E}^{plasma_diff} &= \nu_c \vec{j} + c_{jb} \vec{j} \times \vec{B} - c_{jb} \vec{\nabla} p_e + \nu_A (\vec{v_n} - \vec{v_c}) \\ \text{where } \nu_C &= \frac{m_e (\nu_{ei} + \nu_{en})}{n_e q_e^2}, \ \nu_A = \frac{m_e (\nu_{en} - \nu_{in})}{q_e}, \ c_{jb} = \frac{1}{n_e q_e} \\ \nu_{ei} &= n_e \Lambda_C T_{nc}^{-\frac{3}{2}} 3.7 \cdot 10^{-6}, \ \Lambda_C = 23.4 - 1.15 log_{10}(n_e) + 3.45 log_{10}(\frac{T_{cn} k_B}{q_e}) \\ \text{evolution of magnetic field:} \\ \frac{\partial B_1}{\partial t} &= -\nabla \times \vec{E} \end{split}$$

Orszag test

extended from mancha 1 fluid test no collision terms, variable artificial diffusivity and filtering Initial conditions:

extended from mancha 1 fluid test **Initial conditions**:

hydrostatic equilibrium in an isothermal gravity stratified atmosphere:

$$\vec{B_0} = (0,0,B_{z0}), \ B_{z0} = 50 \cdot 10^{-4} \text{ T}, \ \nabla \times \vec{B_0} = 0$$
 $\frac{\partial p_{\alpha}}{\partial z} = -\rho_{\alpha}g$ we define total pressure at the base $p_{00} = p_0(z=0) = 1.17 \cdot 10^4 \text{ Pa}$ and uniform temperature equal for neutrals and charges: $T_0 = 10000 \text{ K}$ assuming hydrogen plasma we calculate from Saha equation the pressure of neutrals and charges at the base: p_{n00}, p_{c00} we have different pressure scale heights for charges and neutrals: $H_{\alpha} = \frac{RT_0}{\mu_{\alpha}g}$ beacuse of different $\mu_c = \frac{1}{2}\mu_n$ and $\mu_n = 1g/mol$ (only H) we calculate then equilibrium pressure of charges and neutrals: $p_{\alpha 0}(z) = p_{\alpha 00} exp(-\frac{z}{H_0})$

$$p_{\alpha 0}(z) - p_{\alpha 00} exp(-\frac{1}{H_{\alpha}})$$

and density from ideal gas law: $\rho_{\alpha} = \frac{p_{\alpha}\mu_{\alpha}}{RT_0}$

perturbation - a gaussian shaped(in the xy plane) sound wave generated permanently at the base of the gravity stratified atmosphere: we specify the amplitude A=100, the period P=50 ($\omega = \frac{2\pi}{P}$) of the wave, x_0 , y_0 , σ_x , σ_y the center and the standard deviation of the gaussian and z_f the end of the perturbed region the cutoff frequency: $\omega_{c\alpha} = \frac{\gamma g}{2c_{s\alpha}}$ the pressure scale height: $H_{\alpha} = \frac{c_{s\alpha}^2}{2a}$

$$k = \begin{cases} -\frac{\sqrt{\omega^2 - \omega_c^2}}{c_{s_{\alpha}}} - \frac{i}{2H_{\alpha}} & \omega \ge \omega_c \\ i(\frac{\sqrt{\omega^2 - \omega_c^2}}{c_{s_{\alpha}}} - \frac{1}{2H_{\alpha}}) & \omega < \omega_c \end{cases}$$

$$g = exp(-\frac{(x-x_0)^2}{2\sigma_x^2} - \frac{(y-y_0)^2}{2\sigma_y^2})$$

$$rr = \frac{1}{\omega}(-k - \frac{i}{H})$$

$$pp = \frac{1}{\omega}(-k\gamma - \frac{i}{H})$$

$$\begin{split} p_{\alpha 1}(x,y,z) &= A \cdot g(x,y) \cdot pe_{\alpha 0} |pp| exp(Im(k)(z_f-z)) sin(Re(k)(z_f-z) + \omega t + atan(\frac{Im(pp)}{Re(pp)})) \\ \rho_{\alpha 1}(x,y,z) &= \\ A \cdot g(x,y) \cdot \rho_{\alpha 0} |rr| exp(Im(k)(z_f-z)) sin(Re(k)(z_f-z) + \omega t + atan(\frac{Im(rr)}{Re(rr)})) \\ v_{\alpha 1}(x,y,z) &= A \cdot g(x,y) exp(Im(k)(z_f-z)) sin(Re(k)(z_f-z) + \omega t)) \\ \text{in this test we used a Perfectly Matched Layer to avoid reflection at the upper boundary (described in previous version of Mancha) for S^{ion} and S^{rec} instead of expression in Leake the derivation of Saha equation in time (only derivating T and not n_{tot}):
$$\frac{\partial n_e}{\partial t} = ((\frac{2\pi m_e k_B}{h^2})^{1.5} \frac{\partial T_{cn}}{\partial t} exp(-\frac{\phi_{ion}}{k_B T_{cn}}) (1.5\sqrt{T_{cn}} + \frac{\phi_{ion}}{k_B T_{cn}^2}) (\frac{2A + n_{tot}}{2\sqrt{A^2 + An_{tot}}} - 1) \\ \text{min}(\frac{dt}{te \cdot relaxation \cdot timescale}, 1) \\ \text{where } A = (\frac{2\pi m_e k_B T_{cn}}{h^2})^{1.5} exp(-\frac{\phi_{ion}}{(k_B T_{cn})}) \\ \text{and te relaxation timescale is the relaxation time for saha: parameter set to 10 in this test} \end{split}$$$$

$$\begin{split} &\text{where} \left(\frac{\partial n_e}{\partial t} < 0 \right) \, S^{rec} = - \frac{1}{\rho_c} \frac{\partial n_e}{\partial t}, \, S^{ion} = 0 \\ &\text{elsewhere} \, \, S^{ion} = \frac{1}{\rho_n} \frac{\partial n_e}{\partial t}, \, S^{rec} = 0 \end{split}$$

Reconnection

3d