

Figure 1: *Temperature vs z plot. logarithmic y scale*

1a) In order to identify the layers I put conditions on temperature:

[http://www.nasa.gov/mission\\_pages/iris/multimedia/layerzoo.html](http://www.nasa.gov/mission_pages/iris/multimedia/layerzoo.html)

Algorithm for getting the layers: start with values at the top (the values from the file 'atmosphere.dat' are ordered downwards by height) the corona is while temperature  $\geq 500000$  K (T is decreasing), transition region until  $T = 8000$  K, the chromosphere until T reaches the (only) minimum, (afterwards the temperature starts to raise) the photosphere is before T reaches 6500 K and the solar interior afterwards until the end

The exact values matching these conditions are:

corona between [39.802200, 2.535930] Mm temperatures: [1.080180e+06, 5.025160e+05] K

transition region between [2.516350, 0.991115] Mm temperatures: [4.991350e+05, 8.067640e+03] K

chromosphere between [0.971556, 0.305708] Mm temperatures: [7.306160e+03, 2.843670e+03] K

photosphere between [0.286093, -0.303487] Mm temperatures: [2.848470e+03, 6.297540e+03] K

solar interior between [-0.323184, -2.592960] Mm temperatures: [6.837750e+03, 2.068340e+04] K

1b) 
$$\mu = \frac{n_H + 4n_{He}}{n_e + n_H + n_{He}}$$
  

$$n_H = 10n_{He} \implies \mu = \frac{1.4n_H}{n_e + 1.1n_H}$$

- totally ionized H and He  $\implies n_e = n_H + 2n_{He} = 1.2n_H \implies \frac{n_H}{n_e} = \frac{5}{6}$  and  $\mu = 0.6087$
- neutral H and He  $\implies n_e = 0 \implies \mu = 1.2727$

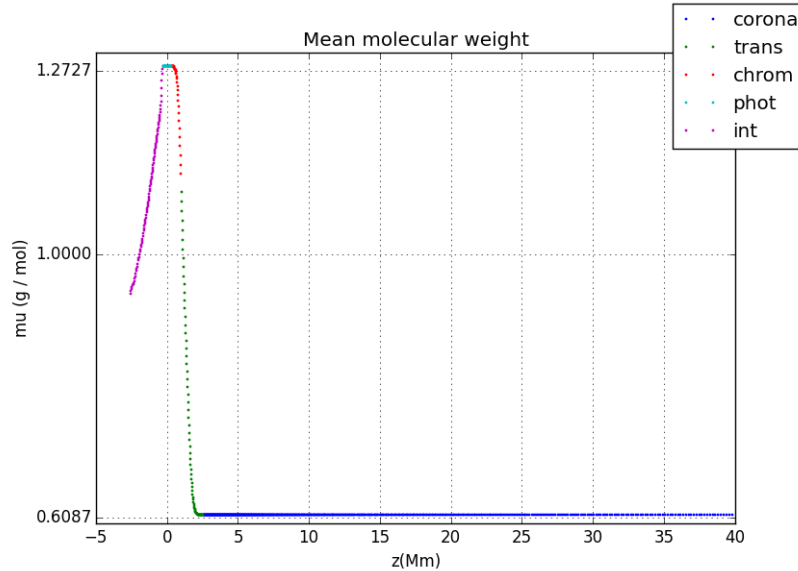


Figure 2: *Mean molecular weight(g/mol) vs z plot* Maximum close to  $1.2727 = \mu$  in the case of neutral H and He and minimum close to  $0.6087 = \mu$  calculated in the case of completely ionized H and He

$$\frac{n_H}{n_e} = \frac{\mu}{1.4 - 1.1\mu}$$

In the case of neutral H and He  $n_e \rightarrow 0 \implies \frac{n_H}{n_e} \rightarrow \infty$

We expect to have big values of this variable in the photosphere

and as we can see in the graphic of  $\mu$  there are some values of  $z$  for which

$$\mu > 1.2727 \implies 1.4 - 1.1\mu < 0 \implies \frac{n_H}{n_e} < 0$$

I will limit oy axis values to  $[0, 4]$  in order to avoid these negative values and the big ones

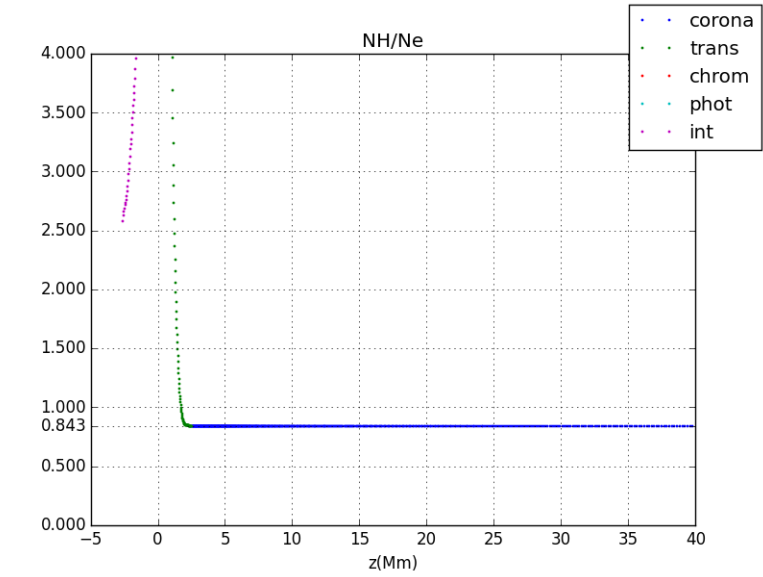


Figura 3:  $nH / ne$

We can see a constant value  $\frac{n_H}{n_e} = 0.843 \approx \frac{5}{6}$  which is the value we calculate in the case of totally ionized H and He (we expect to have totally ionized H and He because of the high values of the temperature in the corona)

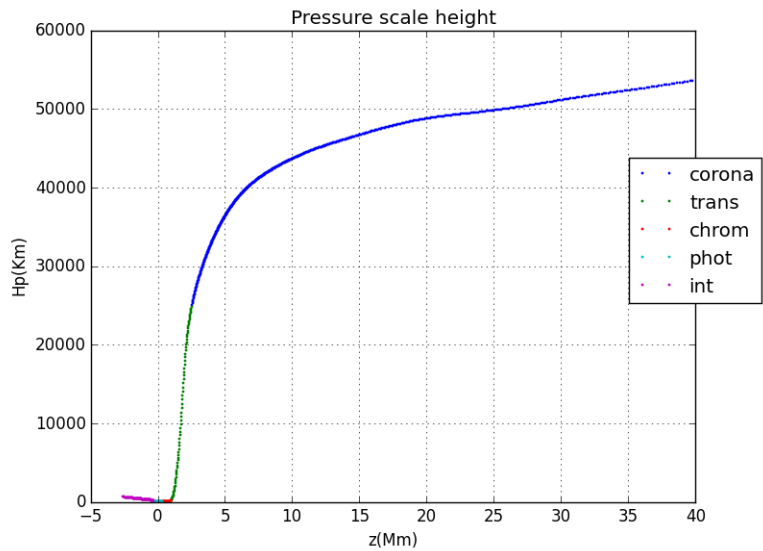


Figura 4: Pressure scale height

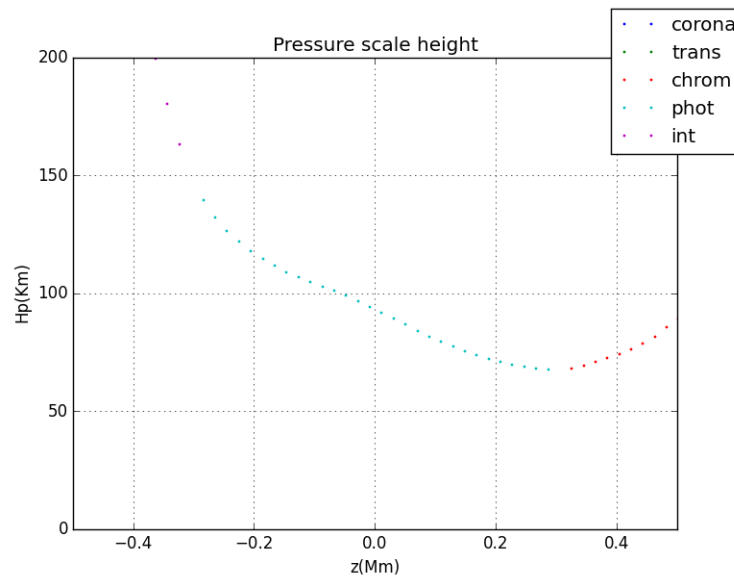


Figure 5: Checking  $H_p$  in the photosphere (between approx 90 - 200 km)

$H_p$  has the minimum at the bottom of the chromosphere (is where  $T$  has the minimum and  $\mu$  the maximum and  $H_p \propto \frac{T}{\mu}$ )

$H_p > 0 \implies$  pressure is a decreasing function.  $H_p$  is the distance in which pressure will decrease by a factor  $e$  so a small value like in the photosphere and chromosphere means that pressure will decrease fast in this portion

$$2) \quad \frac{d \ln p}{dz} = -\frac{1}{H_p}, H_p \text{ const} \implies \ln p(z) - \ln p(z_0) = -\frac{1}{H_p}(z - z_0) \implies p(z) = p(z_0) \exp\left(-\frac{z - z_0}{H_p}\right)$$

$$\rho(z) = \frac{1}{g H_p} p(z) = \frac{p(z_0)}{g H_p} \exp\left(-\frac{z - z_0}{H_p}\right) = \rho(z_0) \exp\left(-\frac{z - z_0}{H_p}\right)$$

Analytic test for  $H_p$  constant (with values 1 and 1e10) with  $\rho(z_f)$  taking values:  $1e-10, 1e-5, 1e-2, 1, 1e2, 1e3, 1e7, 1e10$   
Integrating downward or forward in height makes no difference (using  $\ln p$ )

We see that analytic solution matches exactly numerical solution (we plot  $\ln p(z) - \ln p(z_i)$  vs  $z$ ) and

that the graphic is a line with slope  $\frac{\ln p(z_f) - \ln p(z_i)}{z_f - z_i} = -\frac{1}{H_p}$

where  $z_f = z_{max}$  (z at the top of the atmosphere) and  $z_i = z_{min}$  (z at the bottom of the atmosphere)

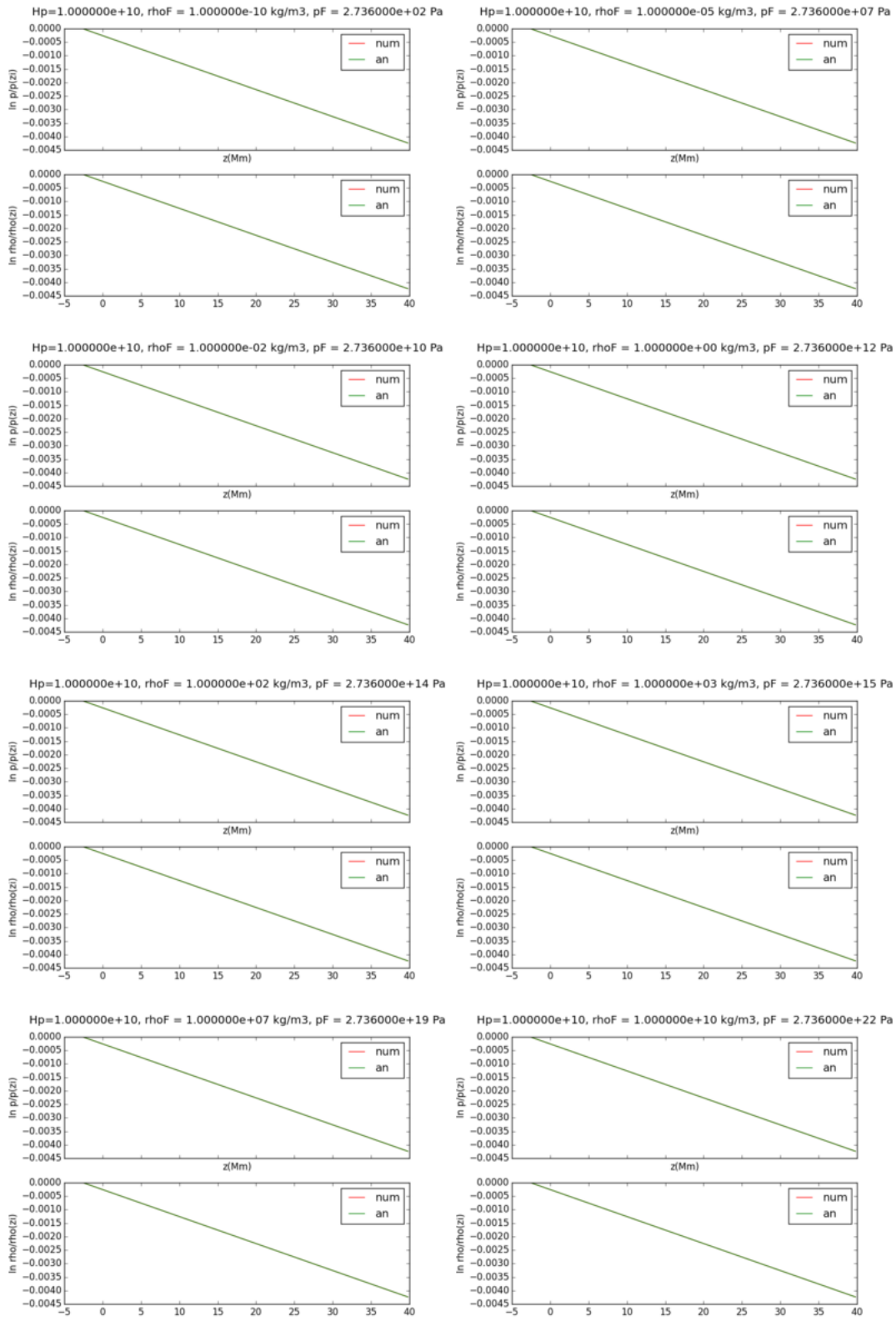


Figura 6: Analytic test  $H_p=1e10$

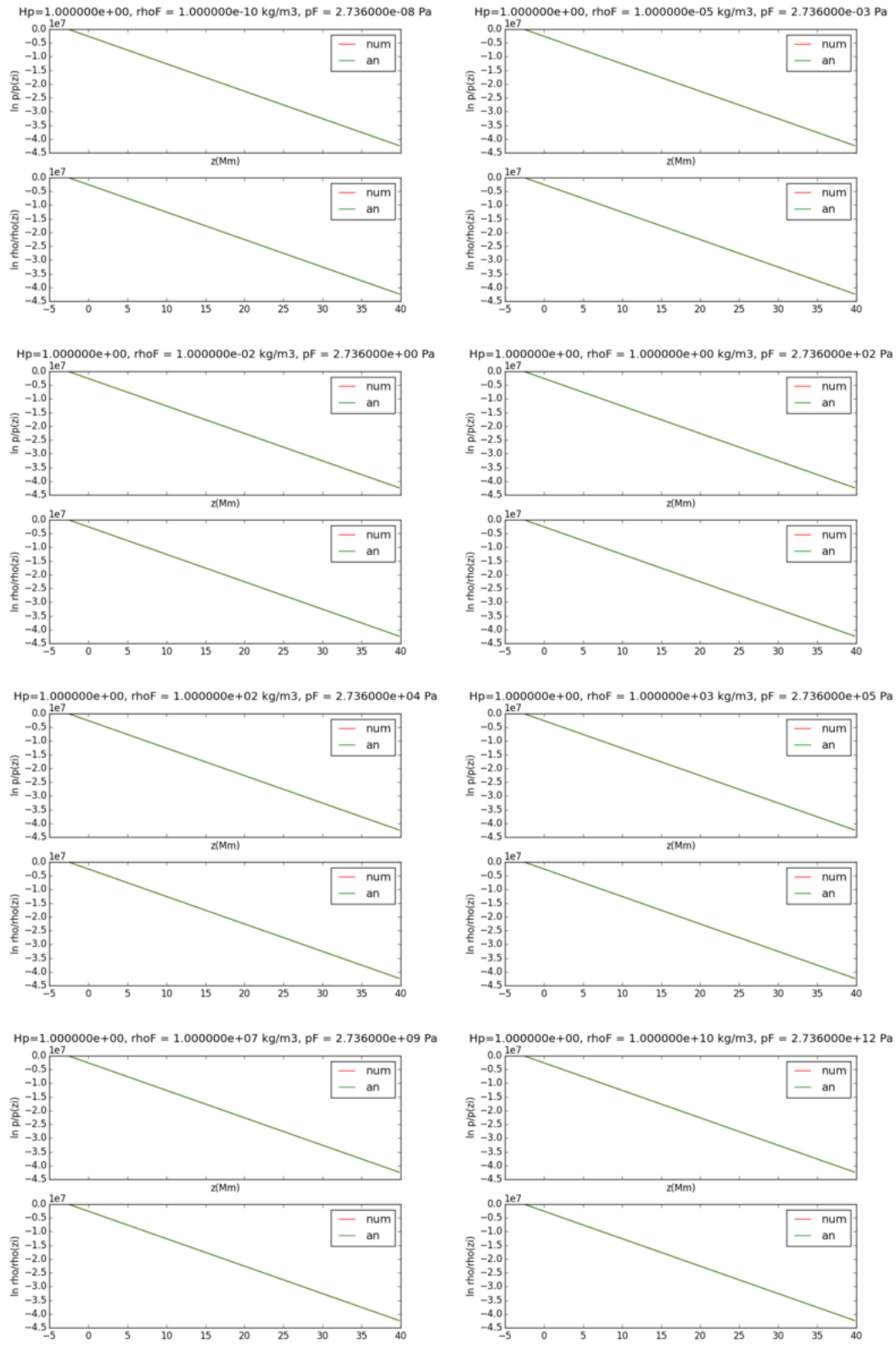


Figura 7: Analytic test  $H_p=1$

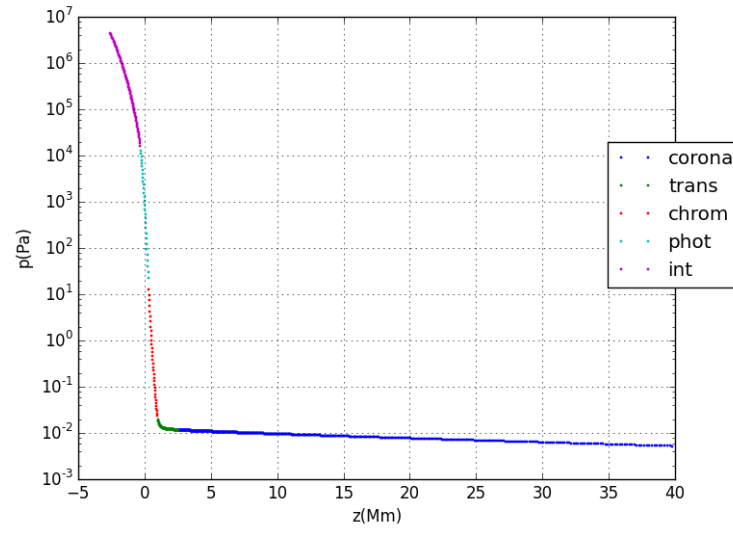


Figura 8: pres log10 oy scale

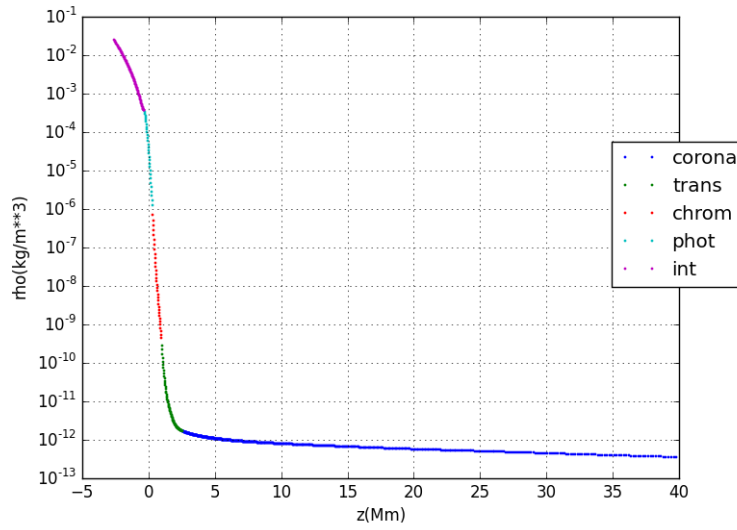


Figura 9: rho log10 oy scale

Pressure will decrease very fast (a few orders of magnitude in a short distance) in the photosphere and chromosphere ( $H_p$  is very small in this portion )

In the transition zone and corona pressure will decrease slowly because  $H_p$  has now bigger values

Density ( $\rho \propto \frac{p}{T}$ ) will decrease fast in the transition zone as well because temperature raises very fast in this portion

**2b)** Notation:  $\mu_0$  = magnetic permeability

$$\beta = \frac{p}{p_{mag}} \text{ where } p_{mag} = \frac{B^2}{2\mu_0}$$

$$v_A^2 = \frac{B^2}{\mu_0 \rho}$$

$$c_s = \sqrt{\frac{\gamma p}{\rho}}$$

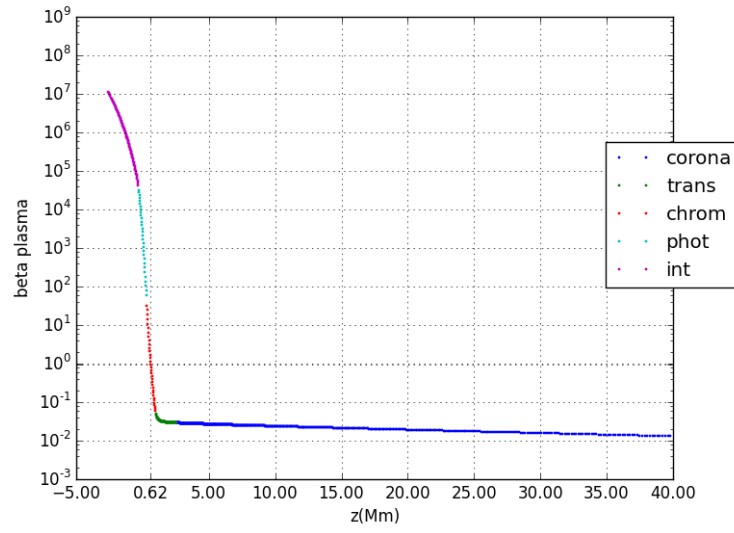


Figura 10: plasma beta log10 oy scale

Plasma beta is a decreasing function and has value 1 at  $z = 0.62$  Mm (in the chromosphere)

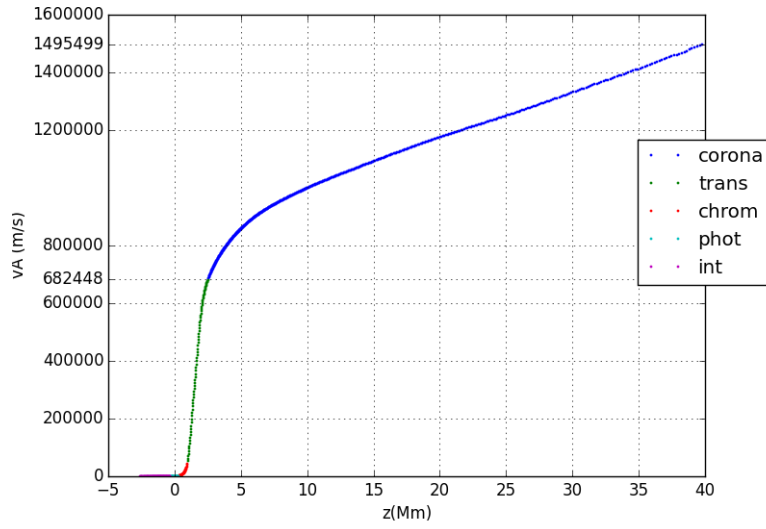


Figura 11: vA

In the corona we observe big values of vA (between approx. 700 - 1500 km /s) ( $vA \propto \rho^{-0.5}$ )

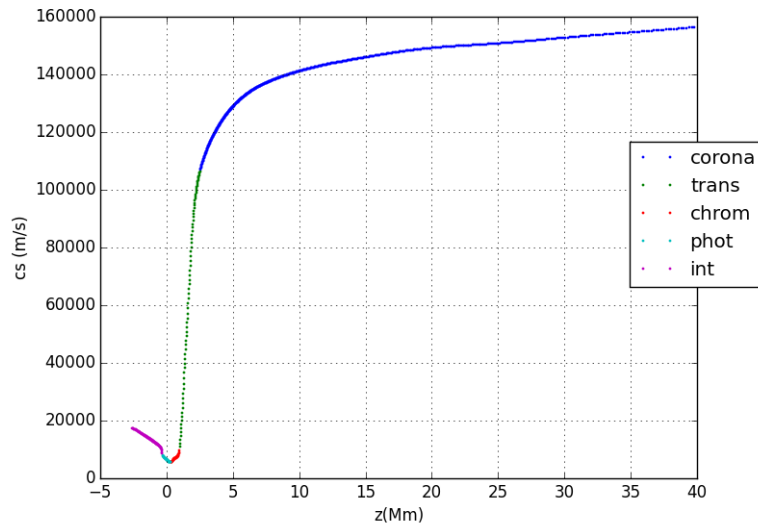


Figura 12: cs

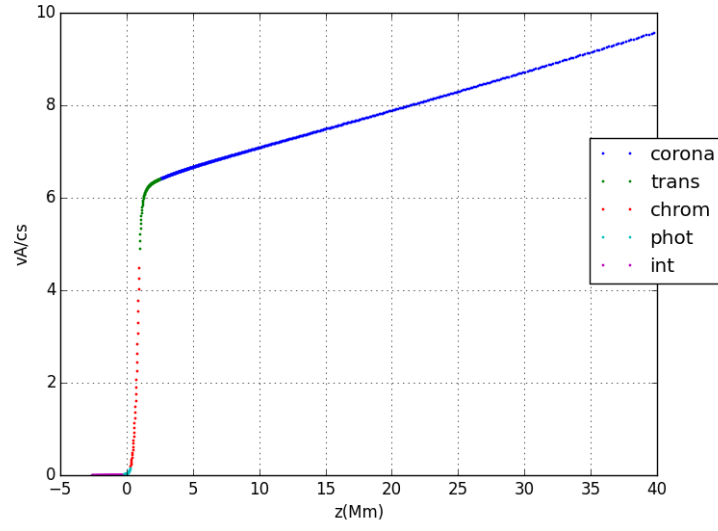


Figura 13:  $v_A/c_s$

In the corona  $v_A > c_s$

$\beta = \frac{2p\mu_0}{B^2} = \frac{2p}{\rho v_A^2} = \frac{2}{\gamma} \left( \frac{c_s}{v_A} \right)^2 \implies \beta \left( \frac{v_A}{c_s} \right)^2 \frac{\gamma}{2} = 1$  We call this function  $\text{func}(\beta, \frac{v_A}{c_s})$  in the graphic below and expect it to be 1

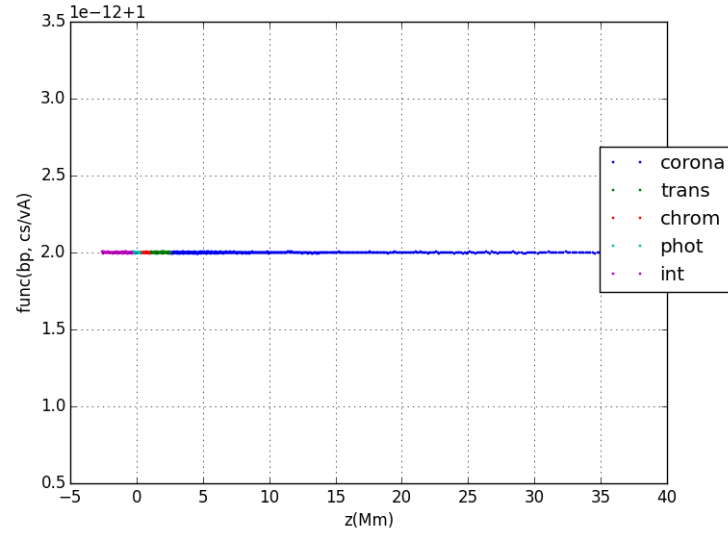


Figura 14:  $\text{func}(\text{bp}, v_A/c_s) \approx 1$

**3a)**  $L_r = \Lambda n_e n_H \implies$

in cgs:  $\frac{\text{erg}}{\text{cm}^3 \text{s}} = [\Lambda] \frac{1}{\text{cm}^6}$

units of  $\Lambda$  in c.g.s are  $\frac{\text{erg cm}^3}{\text{s}}$



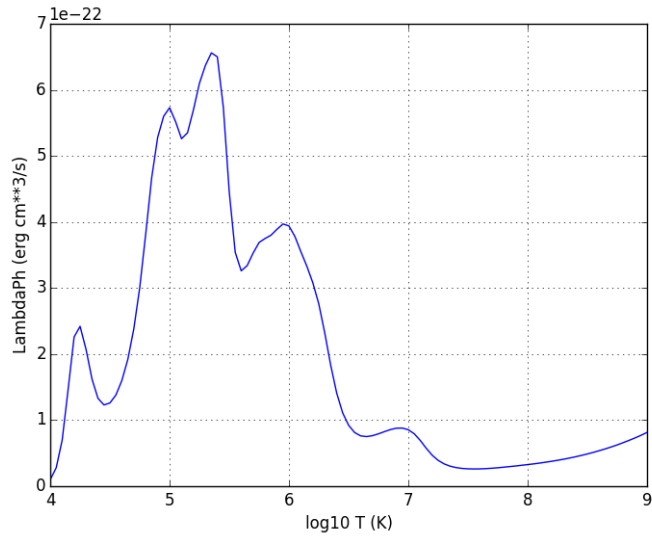


Figura 15: Lambda phot

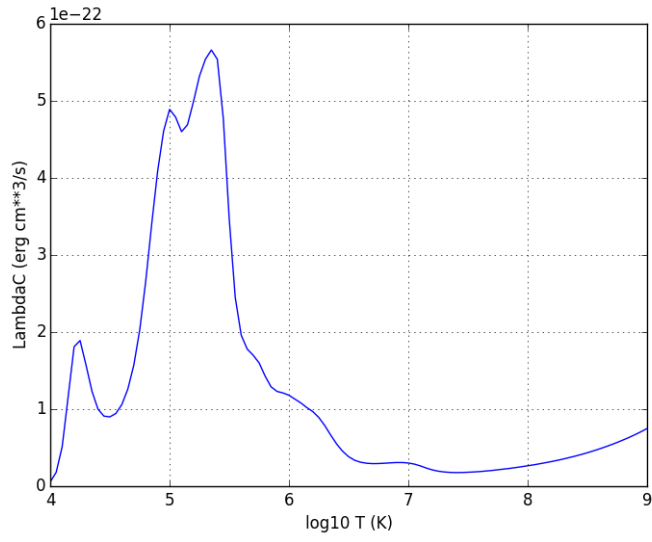


Figura 16: Lambda corona

Both functions have the maximum for  $T = 2.238721e+05$  K

**3b)**  $\rho = \sum_i n_i a_i m_H = (n_H + 4n_{He})m_H$   
 $n_H = 10n_{He} \implies \rho = 1.4n_H m_H$

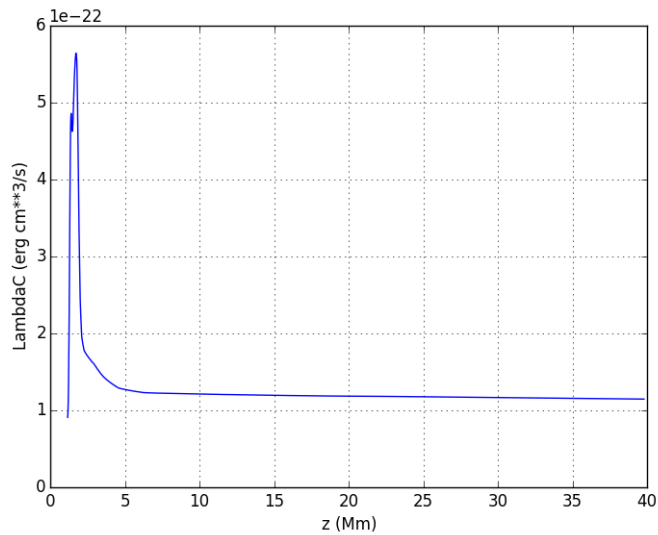


Figure 17: Lambda corona interpolated for atm. temperatures  $> 3 \times 10^4$  K in 'atmosphere.dat' plotted vs z

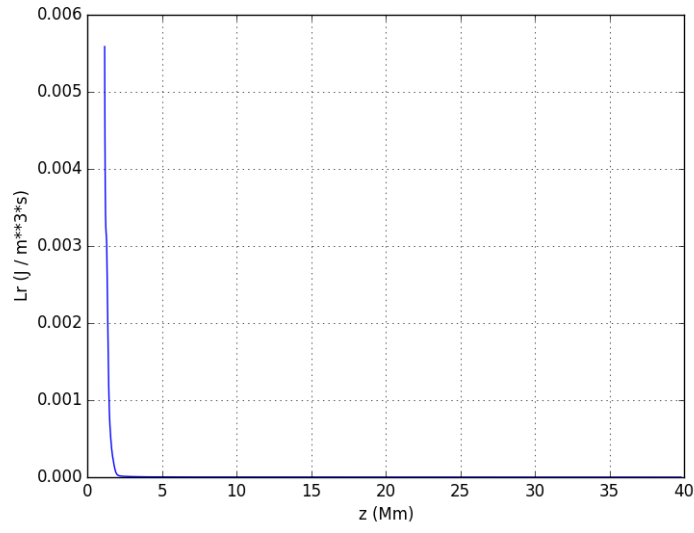


Figura 18:  $L_r$

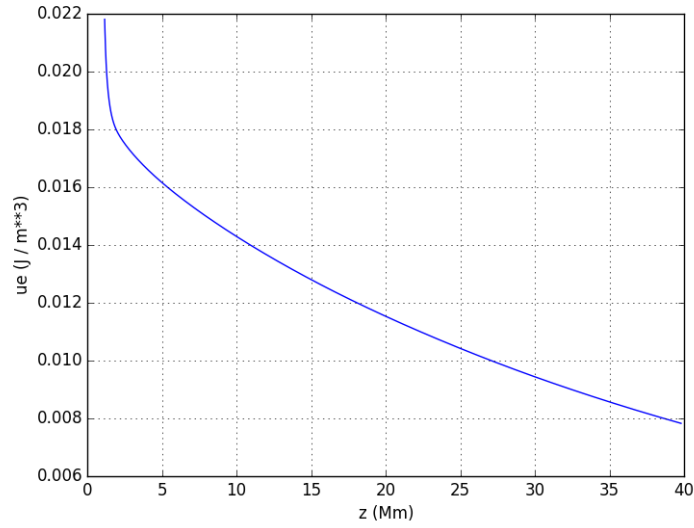


Figure 19: Internal energy calculated as  $1.5 * p$

**3c)**

**3d)** Equation of energy when  $\vec{q} = 0, \vec{v} = 0, \vec{j} = 0$ :

$$\frac{\partial u_e}{\partial t} = -L_r$$

if we consider  $L_r$  constant in time (in fact  $L_r$  will be decreasing in time because of its dependence on  $n_e$  and  $T$ ):

$u_e(t) = u_e(t=0) - L_r t \implies \frac{u_e(t=0)}{L_r}$  is the (minimum, if we think that  $L_r$  will decrease in time) time needed to convert all internal energy into radiation energy

units of  $\frac{u_e(t=0)}{L_r}$  are units of time: s

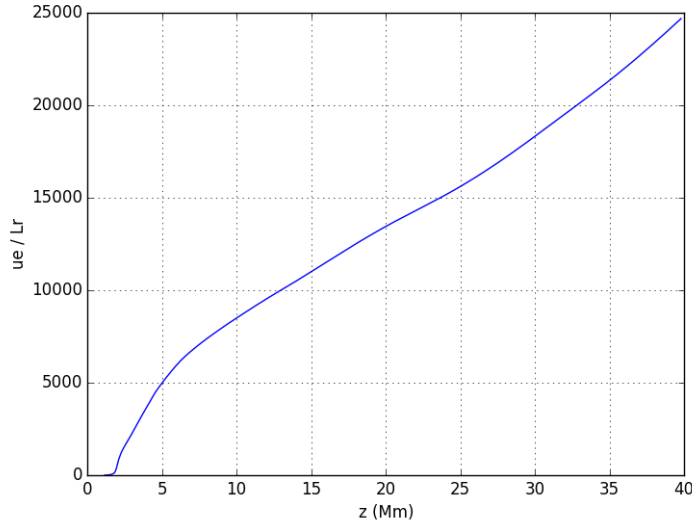


Figura 20: Internal energy / Lr

The maximum value of  $\frac{u_e(t=0)}{L_r}$  is about 25000 s at the top the corona, it needs more than 7 hours to cool completely

**Testing** Derivate  $\ln p$  (obtained after integrating  $-1/H_p$  calculated with data taken from the file) by the following scheme:

for an array of n elements  $f[i]$ ,  $i$  in  $[0..n-1]$  we calculate  $df[i] = (f[i+1] - f[i-1]) / (dz[i-1] + dz[i])$ , for  $i$  in  $[1..n-2]$ ;  $df[0] = df[1]$ ;  $df[n-1] = df[n-2]$

using  $dz[i] = z[i+1] - z[i]$  for  $i$  in  $[0..n-2]$  ( $z$  taken from the file, we have both  $z$  and  $\ln p$  reversed)

and for the array  $d \ln p$  obtained this way plot  $-1 / d \ln p$  and compare with  $H_p$

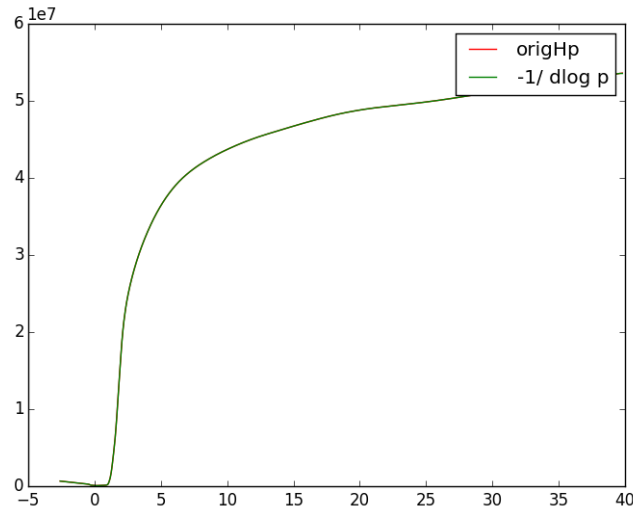


Figura 21:  $-1/d \ln p = H_p$

**Resolution** decrease resolution by taking off points: in one step we keep only points with even index in the array (first index is 0)

increase resolution artificially by introducing points : in one step we add points in the middle between each 2 consecutive points and calculate the functions from the table:  $\mu$  and  $T$  by lineal or cubic interpolation

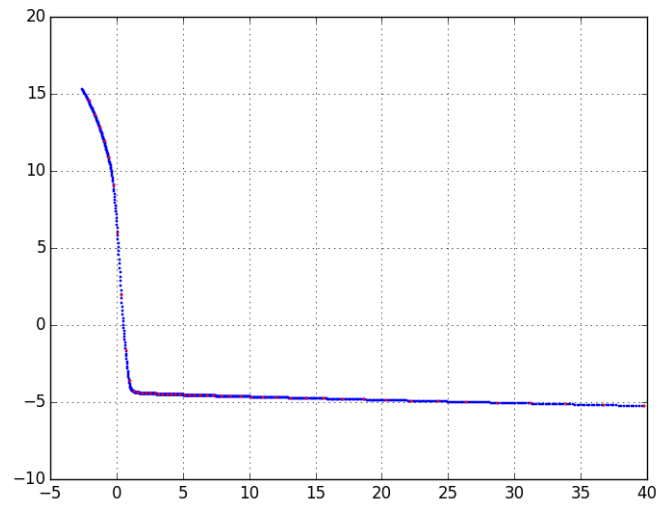


Figura 22:  $\ln p$  decreased resolution (4 steps): 48 points with red points plotted on top, original resolution 768 points

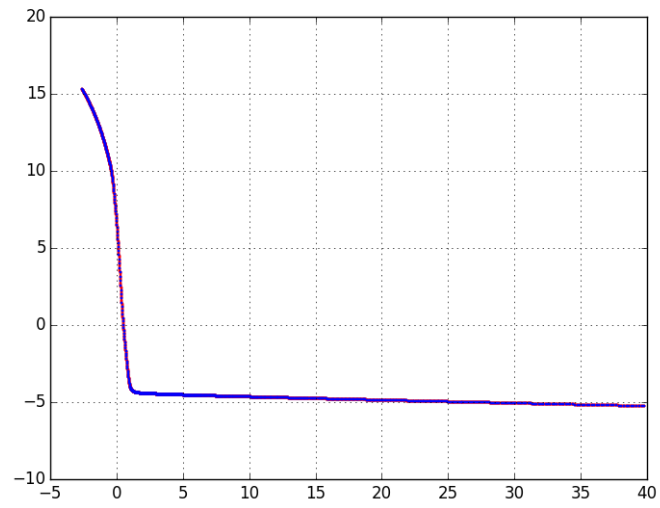


Figura 23:  $\ln p$  increased resolution (2 steps, cubic interpolation): 3069 points with red points, original resolution 768 points plotted on top