$$\int_{0}^{r_0} 4\pi r^2 n_e(r) dr = Z \tag{1}$$

$$\begin{array}{l} n_e(r) = \frac{8\pi}{3h^3} [2m_e(e_F + eV(r))]^{\frac{3}{2}} \ (1.29 \ \mathrm{apuntes}) \\ x = \frac{r}{\mu a_0} \implies r = xa_0 (\frac{9\pi^2}{128Z})^{\frac{1}{3}} \\ \Phi(x) = \frac{e_F + eV(r)}{\frac{Ze^2}{4\pi\epsilon_0 r}} \implies \\ e_F + eV(r) = \Phi(x) \frac{Ze^2}{4\pi\epsilon_0 r} \implies \\ n_e(r) = \frac{8\pi}{3h^3} (2m_e \Phi(x) \frac{Ze^2}{4\pi\epsilon_0 r})^{\frac{3}{2}} \\ \mathrm{reemplazando en } (1): \\ \int_0^{r_0} 4\pi r^2 \frac{8\pi}{3h^3} (\Phi(x) \frac{m_e Ze^2}{2\pi\epsilon_0 r})^{\frac{3}{2}} dr = Z \implies \\ \frac{32\pi^2}{3h^3} (\frac{m_e Ze^2}{2\pi\epsilon_0 r})^{\frac{3}{2}} \int_0^{r_0} r^{\frac{1}{2}} \Phi(x)^{\frac{3}{2}} dr = Z \\ \mathrm{Cambio de variable } r \mathrm{por } x \\ dr = dxa_0 (\frac{9\pi^2}{128Z})^{\frac{1}{3}} \\ \frac{32\pi^2}{3h^3} (\frac{m_e Ze^2}{2\pi\epsilon_0 r})^{\frac{3}{2}} a_0^{\frac{3}{2}} (\frac{9\pi^2}{128Z})^{\frac{1}{2}} \int_0^{x_0} x^{\frac{1}{2}} \Phi(x)^{\frac{3}{2}} dx = Z \\ \Phi(x)^{\frac{3}{2}} = x^{\frac{1}{2}} \frac{d^2\Phi}{2a_0^2} \implies \\ \frac{32\pi^2}{3h^3} (\frac{m_e Ze^2}{2\pi\epsilon_0 r})^{\frac{3}{2}} a_0^{\frac{3}{2}} (\frac{9\pi^2}{128Z})^{\frac{1}{2}} \int_0^{x_0} x \frac{d^2\Phi}{dx^2} dx = Z \implies \\ \frac{32\pi^2}{3h^3} (\frac{m_e Ze^2}{2\pi\epsilon_0 r})^{\frac{3}{2}} a_0^{\frac{3}{2}} (\frac{9\pi^2}{128})^{\frac{1}{2}} \int_0^{x_0} x \frac{d^2\Phi}{dx^2} dx = Z \implies \\ \frac{32\pi^2}{3h^3} (\frac{m_e Ze^2}{2\pi\epsilon_0 r})^{\frac{3}{2}} a_0^{\frac{3}{2}} (\frac{9\pi^2}{128})^{\frac{1}{2}} \int_0^{x_0} x \frac{d^2\Phi}{dx^2} dx = 1 \implies \\ \mathrm{Notamos} \mathrm{C} = (\frac{32\pi^2}{3h^3} (\frac{m_e e^2}{128Z})^{\frac{3}{2}} a_0^{\frac{3}{2}} (\frac{9\pi^2}{2\pi\epsilon_0 r})^{\frac{3}{2}} a_0^{\frac{3}{2}} (\frac{9\pi^2}{2\pi\epsilon_0 r})^{\frac{3}{2}} a_0^{\frac{3}{2}} (\frac{9\pi^2}{128})^{\frac{3}{2}} -1 \\ \int_0^x x \frac{d^2\Phi}{dx^2} dx = C \\ \mathrm{Integrando por partes:} \\ \int_0^x x \frac{d^2\Phi}{dx^2} dx = (x\Phi'(x))|_0^{x_0} - \int_0^{x_0} \Phi'(x) dx = x_0\Phi'(x_0) - \Phi(x_0) + \Phi(0) \implies \\ x_0\Phi'(x_0) - \Phi(x_0) = C - 1 \\ C = 1 \ (\iff a_0 = \frac{4\pi\epsilon_0h^2}{m_e e^2}) \\ \implies \Phi'(x_0) = \frac{\Phi(x_0)}{x_0} \end{aligned}$$

H1 p11 all ionized
$$\Longrightarrow \frac{1}{\mu} = 2X + \frac{3}{4}Y + \frac{1}{2}Z$$

$$\frac{1}{\mu} = 1.3793 \text{ g/mol}$$
eq 1.40, $M = M_{\odot} \Longrightarrow$

$$C = 6.65 \cdot 10^4 \frac{\mu}{Z(1+X)} = 91.7241 \cdot 10^4 \text{ erg } s^{-1}K^{\frac{-7}{2}}$$

$$T_c = (\frac{L}{C})^{\frac{7}{7}}$$

$$L = 0.03L_{\odot} = 0.117 \cdot 10^{33} \text{ erg/s} \Longrightarrow$$

$$T_c = 2.0697 \cdot 10^6 \text{ K}$$

$$T_s = (\frac{C}{4\pi R^2 \sigma})^{\frac{1}{4}} T_c^{\frac{7}{8}}$$

$$\sigma = 5.67 \cdot 10^{-5} \text{ erg } cm^{-2}K^{-4}s^{-1}$$

$$R = R_{\odot} = 6.96 \cdot 10^{10} \text{ cm}$$

$$\Longrightarrow T_s = 2412.9238 \text{ K}$$

H2 p4 caso no relativista (baja densidad: $\rho << 6 \cdot 10^{15} \text{ g/cm}^3$)

$$\gamma = 5/3, K = \frac{3^{\frac{2}{3}}\pi^{\frac{4}{3}}\hbar^{2}}{5^{\frac{5}{m_{n}^{3}}}} = 5.38752 \cdot 10^{9}$$

en la ecuación Lane Emden n = 1.5 igual que en el caso de las enanas blancas de baja densidad \Longrightarrow tiene la misma resolución: $\xi_1 = 3.65375 \text{ y } |\theta'(\xi_1)| = 0.203302$

polítropos apuntes eq 1.18, 1.19:

R =

H3 p2 partícula de masa = 1 parte del reposo de una distancia R $\implies E = c^2 \sqrt{1 - \frac{r_s}{R}}$ (la energía total es la energía de su masa en reposo) eq 3.6 apuntes: $c^2(1 - \frac{r_s}{r})\frac{dt}{d\tau} = E \implies \frac{d\tau}{dt} = (1 - \frac{r_s}{r})^{-\frac{1}{2}}$ apuntes (parte de una distancia R): $\tau(r) = \frac{1}{c}(\frac{R^3}{r_s})^{\frac{1}{2}}[(\frac{r}{R} - \frac{r^2}{R^2})^{\frac{1}{2}} + arccos(\sqrt{\frac{r}{R}})]$

$$\begin{split} r &= R \frac{1 + \cos \eta}{2} \implies \tau(\eta) = \frac{1}{c} (\frac{R^3}{r_s})^{\frac{1}{2}} \big[(\frac{1 + \cos \eta}{2} - (\frac{1 + \cos \eta}{2})^2)^{\frac{1}{2}} + \arccos(\sqrt{\frac{1 + \cos \eta}{2}}) \big] \\ \frac{d\tau}{d\eta} &= \frac{1}{c} (\frac{R^3}{r_s})^{\frac{1}{2}} (\frac{\frac{1}{2} \sin(\eta)(\cos(\eta) + 1) - \frac{\sin(\eta)}{2}}{2\sqrt{\frac{1}{2}}(\cos(\eta) + 1) - \frac{1}{4}(\cos(\eta) + 1)^2}} + \frac{\sin(\eta)}{2\sqrt{2}\sqrt{\frac{1}{2}}(-\cos(\eta) - 1) + 1} \sqrt{\cos(\eta) + 1}} \big) \\ \frac{d\tau}{d\eta} \frac{d\eta}{dt} &= 1 - \frac{r_s}{r} \implies \frac{dt}{d\eta} = \frac{1}{c} (\frac{R^3}{r_s})^{\frac{1}{2}} (1 - \frac{2r_s}{R(1 + \cos \eta)})^{-1} (\frac{\frac{1}{2} \sin(\eta)(\cos(\eta) + 1) - \frac{\sin(\eta)}{2}}{2\sqrt{\frac{1}{2}}(\cos(\eta) + 1) - \frac{1}{4}(\cos(\eta) + 1)^2}} + \frac{\sin(\eta)}{2\sqrt{2}\sqrt{\frac{1}{2}}(-\cos(\eta) - 1) + 1} \sqrt{\cos(\eta) + 1}} \big) \\ \implies t(\eta) &= \frac{(\cos(\eta) + 1)^{3/2} \tan(\frac{\eta}{2}) \sec^2(\frac{\eta}{2}) \left(4r_s^{3/2} \tanh^{-1}\left(\frac{\sqrt{r_s}\tan(\frac{\eta}{2})}{\sqrt{R} - r_s}\right) + \sqrt{R} - r_s(\eta(R + 2r_s) + R\sin(\eta))}\right)}{4R\sqrt{R} - r_s} \sqrt{1 - \cos(\eta)}} \\ \text{singularidad en } r &= r_s \end{split}$$

H3 p4 partícula con masa m = 1 parte del reposo $(E = c^2)$ desde el infinito:

$$\frac{dr}{d\tau} = -c(\frac{r_s}{r})^{\frac{1}{2}} \Longrightarrow
r^{\frac{1}{2}}dr = -cr_s^{\frac{1}{2}}d\tau \Longrightarrow
\tau(r) = C - \frac{2}{3}c^{-1}r_s^{-\frac{1}{2}}r^{\frac{3}{2}}
\tau(R) = 0 \Longrightarrow C = \frac{2}{3}c^{-1}r_s^{-\frac{1}{2}}R^{\frac{3}{2}}
\tau(r) = \frac{2}{3}c^{-1}r_s^{-\frac{1}{2}}R^{\frac{3}{2}} - \frac{2}{3}c^{-1}r_s^{-\frac{1}{2}}r^{\frac{3}{2}}
\tau(r_s) = \frac{2}{3}c^{-1}r_s^{-\frac{1}{2}}R^{\frac{3}{2}} - \frac{2}{3}c^{-1}r_s^{-\frac{1}{2}}r^{\frac{3}{2}} = \frac{2}{3}c^{-1}r_s^{-\frac{1}{2}}R^{\frac{3}{2}} - \frac{2}{3}c^{-1}r_s$$