

• Stiffness and compliance matrices

Within the context of force-based approaches, the PCVW (and Menabrea's theorem) can be used for evaluating:

1. Unknown reaction forces (structurally indeterminate systems)

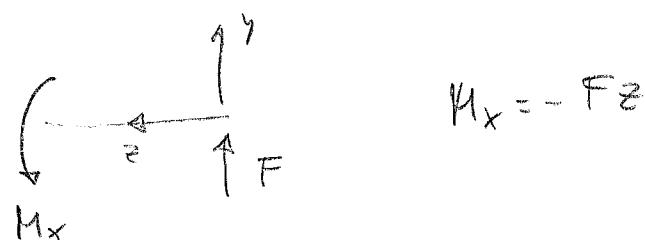
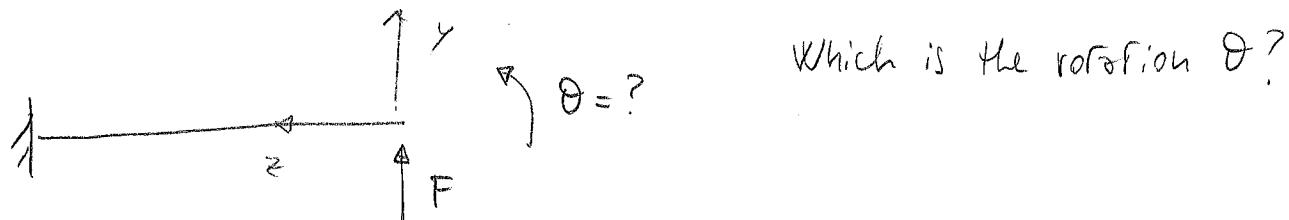
2. Displacements in given points for given set of loads

Consider now this second case. The PCVW for the DSV beam was found to be

$$\int_e \left(\delta T_z \frac{T_z}{EA} + \delta M_x \frac{M_x}{EJ_{xx}} + \delta M_y \frac{M_y}{EJ_{yy}} + \delta T_x \frac{T_x}{GA_x^*} + \delta T_y \frac{T_y}{GA_y^*} + \delta M_z \frac{M_z}{GJ} \right) dz = \delta F_r \vec{u}_r + \delta M_r \vec{\theta}_r + \delta M_s \vec{\theta}_s$$

→ the components due to T_x, T_y, M_z have not been dismissed yet, and are reported here for generality purposes.

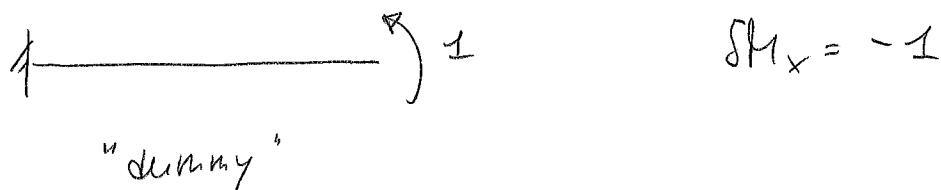
To understand how to evaluate displacement/rotations, consider this simple example



The real generalized deformation is then $\frac{\delta M_x}{EI_{xx}} = -\frac{F_z}{EI_{xx}}$

Consider how a properly selected "dummy" system for evaluating a suitable set of equilibrated variation of forces..

The dummy system is chosen with regard to the displacement component to be evaluated, in such a manner that the unknown, in this case θ , appears in the external work.



The application of the PCVW is:

$$\int_0^l \frac{F_z}{EI_{xx}} dz = \theta \Rightarrow \theta = \frac{F_z l^2}{2EI_{xx}}$$

- In this example no displacements were imposed, so

$$\delta W_e^* = \left[\delta F_i u_i + \delta M_j \dot{\theta}_j \right] + \delta F_r u_r + \delta M_s \dot{\theta}_s$$

The dummy system was taken by considering a variation of the applied moment $\delta M = 1$, while $\delta F = 0$

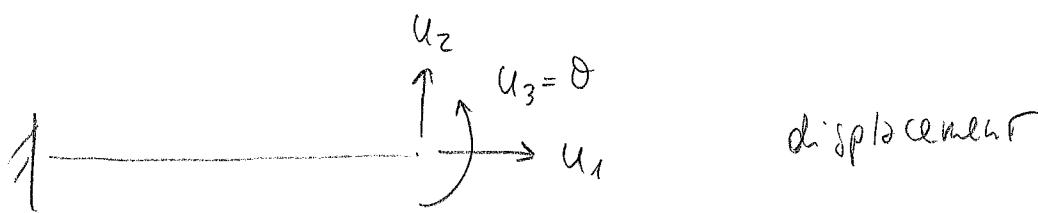
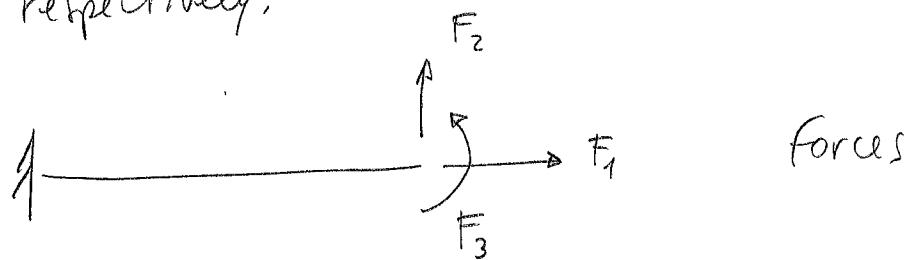
so

$$\delta W_e^* = \delta M \theta = \theta$$

The quantity δM_x in the expression of δW_e^* is then the equilibrated variation of bending moments due to the application of δM .

Stiffness matrix for a cantilever beam

Consider a cantilever beam and denote with F_i and u_i the generalized forces and displacements, respectively.



The convention for F_i and u_i is completely arbitrary.

Obtaining the compliance and the stiffness matrices means obtaining a relation in the form:

$\underline{u} = \underline{C} \underline{F}$	Compliance
$\underline{F} = \underline{K} \underline{u}$	stiffness with $\underline{K} = \underline{C}^{-1}$

Remarks

1. The procedure relies upon the linearity assumption of the problem. It follows that the principle of superposition holds

(the displacement u_i due to the application of F_A and F_B is $u_i = u_A + u_B$ where u_A is the displacement due to F_A and u_B is the displacement due to F_B)

2. The compliance/stiffness matrix will be derived in the context of a force-based approach. The result will be the compliance matrix while the stiffness matrix will be derived indirectly, after inverting $\underline{\underline{C}}$.

- Force-based approaches $\rightarrow \underline{\underline{C}} \quad (\underline{\underline{k}} = \underline{\underline{C}}^{-1})$
- Displacement-based approaches $\rightarrow \underline{\underline{k}} \quad (\underline{\underline{C}} = \underline{\underline{k}}^{-1})$

Evaluation of the coefficients

$$\begin{Bmatrix} u_1 \\ u_2 \\ u_3 \end{Bmatrix} = \begin{bmatrix} C_{11} & C_{12} & C_{13} \\ C_{21} & C_{22} & C_{23} \\ C_{31} & C_{32} & C_{33} \end{bmatrix} \begin{Bmatrix} F_1 \\ F_2 \\ F_3 \end{Bmatrix}$$

The generic term C_{ij} denotes the i -th displacement due to the j -th unitary force

① C_{11}

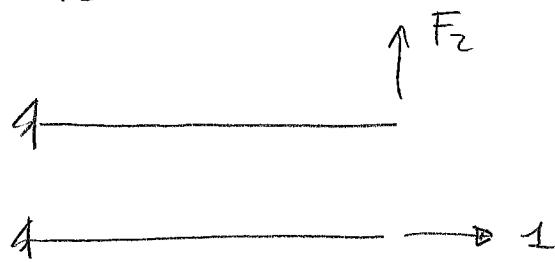
$$1 \longrightarrow F_1 \quad T_2 = F_1$$

$$1 \longrightarrow 1 \quad \delta T_2 = 1$$

$$\int_0^l f T_2 \frac{T_2}{EA} = 1 \cdot u_1$$

$$F_1 \frac{l}{EA} = u_1 \Rightarrow \boxed{C_{11} = \frac{l}{EA}}$$

② C_{12}

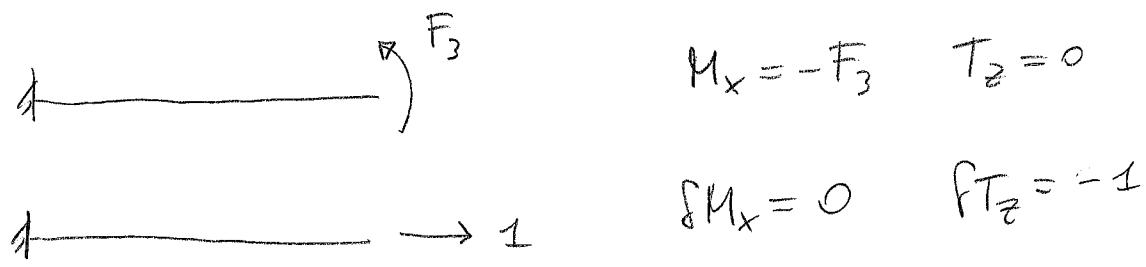


$$M_x = 0 \quad T_z = 0$$

$$\delta M_x = 0 \quad \delta T_z = 1$$

$$\boxed{C_{12} = 0}$$

③ C_{13}

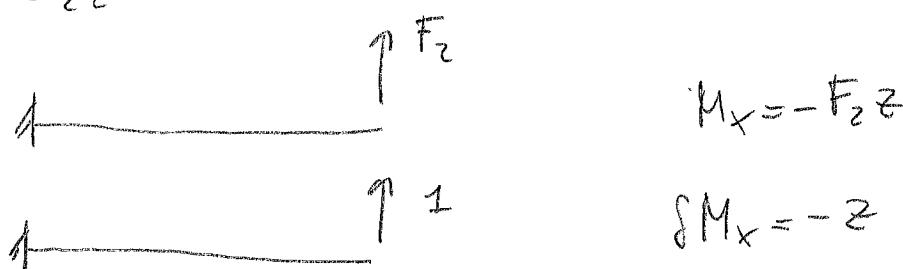


$$M_x = -F_3 \quad T_z = 0$$

$$\delta M_x = 0 \quad \delta T_z = -1$$

$$\boxed{C_{13} = 0}$$

④ C_{22}



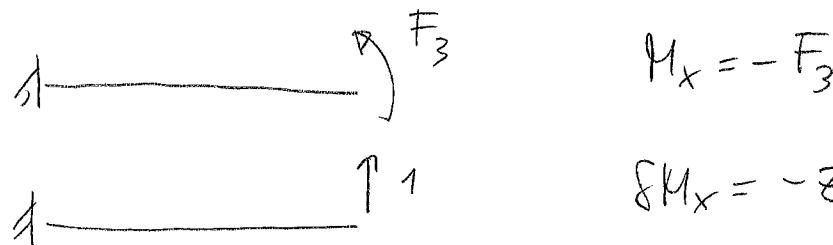
$$M_x = -F_2 z$$

$$\delta M_x = -z$$

$$U_2 = \int_0^l \frac{F_2 z^2}{EJ} dz = \frac{F_2 l^3}{3EJ}$$

$$\boxed{C_{22} = \frac{l^3}{3EJ}}$$

⑤ C_{23}



$$M_x = -F_3$$

$$\delta M_x = -z$$

$$U_2 = \int_0^l \frac{F_2 z}{EJ} dz = \frac{F_2 l^2}{2EJ}$$

$$\boxed{C_{23} = \frac{l^2}{2EJ}}$$

⑥ C_{33}

$$\begin{array}{c} \text{F}_3 \\ \text{F}_3 \end{array} \quad \mu_x = -\text{F}_3$$

$$\begin{array}{c} \text{F}_3 \\ \text{F}_3 \end{array} \quad \delta \mu_x = -1$$

$$u_3 = \int_0^l \frac{F}{EI} dz = \frac{Fl}{EI} \quad \boxed{C_{33} = \frac{l}{EI}}$$

The compliance matrix is then obtained as:

$$\begin{Bmatrix} u_1 \\ u_2 \\ u_3 \end{Bmatrix} = \begin{bmatrix} \frac{l}{EA} & 0 & 0 \\ 0 & \frac{l^3}{3EI} & \frac{l^2}{2EI} \\ 0 & \frac{l^2}{2EI} & \frac{l}{EI} \end{Bmatrix} \begin{Bmatrix} F_1 \\ F_2 \\ F_3 \end{Bmatrix} \quad \underline{\text{Compliance}}$$

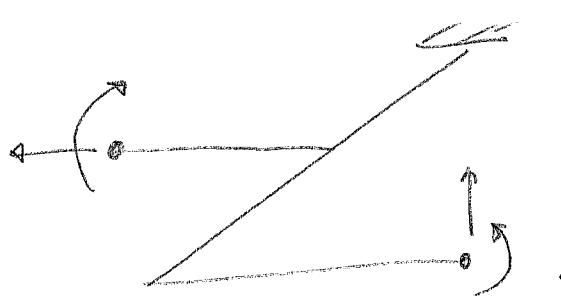
The stiffness matrix is available after inverting C :

$$\begin{Bmatrix} F_1 \\ F_2 \\ F_3 \end{Bmatrix} = \begin{bmatrix} EA/l \\ 12EI/l^3 & -6EI/l^2 \\ -6EI/l^2 & 4EI/l \end{bmatrix} \begin{Bmatrix} u_1 \\ u_2 \\ u_3 \end{Bmatrix} \quad \underline{\text{Stiffness}}$$

Extension to a generic beam system

The same approach can be adopted when dealing with a more complex system of beams

The derivation of the compliance matrix represents a sort of condensation of the structure's response to a few points of interest



After defining the forces and displacements of interest a relation can be found in the form $\underline{U} = \underline{C} \underline{F}$

Where the generic contribution C_{ij} represents the i -th displacement due to a unitary force F_j .

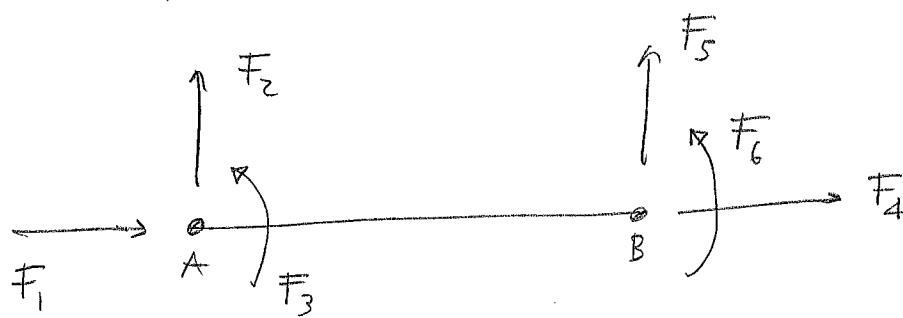
$\Rightarrow C_{ij}$ is obtained by considering a real system where the force F_j is applied, and a dummy system with unitary force F_i

C_{ij} is obtained with i-th dummy system j-th real system

Free-Free beam

The most general case is given by a beam which is free at the two ends.

- From this case any set of boundary conditions can be derived as a special case.
- The resulting stiffness matrix will be useful for the computer implementation of a simple code for analyzing relatively complex systems of beams.



The goal is to derive a relation in the form

$$\begin{Bmatrix} \underline{F}_A \\ \underline{F}_B \end{Bmatrix} = \begin{bmatrix} k_{AA} & k_{AB} \\ k_{BA} & k_{BB} \end{bmatrix} \begin{Bmatrix} \underline{u}_A \\ \underline{u}_B \end{Bmatrix}$$

where

$$\underline{F}_A = \{F_1 \ F_2 \ F_3\}^T \quad \underline{u}_A = \{u_1 \ u_2 \ u_3\}^T$$

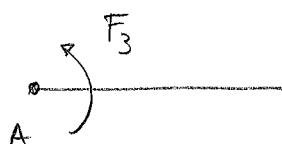
$$\underline{F}_B = \{F_4 \ F_5 \ F_6\}^T \quad \underline{u}_B = \{u_4 \ u_5 \ u_6\}^T$$

① k_{BB} is available from the previous results of the cantilever beam

② k_{AA} can be obtained with a similar procedure (by repeating all the steps)

A smarter approach is to consider that:

1. The direct terms $C_{\alpha\alpha}$ do not change
(the displacement u_1 due to F_1 is equal to
the displacement u_4 due to F_4 , and similarly
for the other terms)
2. The mixed contributions $C_{\alpha\beta}$ may change
the sign. (C_{23} and C_{32} , as $C_{13} = C_{23} = 0$)



F_3 determines
a displacement
 u_2 in the downward
direction



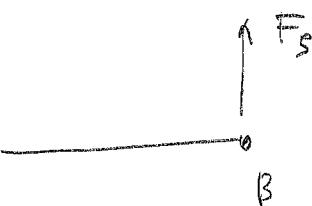
F_6 determines
a displacement
 u_5 in the Upward
direction

C_{23} changes sign!

Similarly



F_2 determines
a rotation in
the clockwise direction
(negative in the
conventions here
used)



F_5 determines
a rotation in
the councclockwise
direction
(positive in the
conventions here
used)

C₃₂ changes sign

Thus:

$$C_{AA} = \begin{bmatrix} C_{11} & 0 & 0 \\ 0 & C_{22} - C_{23} \\ 0 & C_{-23} & C_{33} \end{bmatrix}$$

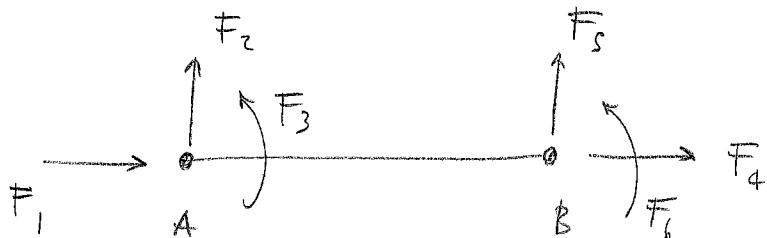
and so:

$$K_{AA} = \begin{bmatrix} EA/\epsilon & & \\ & 12EI/\epsilon^3 & 6EI/\epsilon^2 \\ & 6EI/\epsilon^2 & 4EI/\epsilon \end{bmatrix}$$

③ Evaluation of K_{AB}

Even in this case, it is possible to obtain K_{AB} term by term.

A smarter strategy is to proceed as follows.



From equilibrium:

$$F_1 = -F_4$$

$$F_2 = -F_5$$

$$F_3 = -F_6 - F_5 \ell$$

$$\text{or } F_A = T F_B \text{ with } T =$$

$$T = \begin{bmatrix} -1 & 0 & 0 \\ 0 & -1 & 0 \\ 0 & -\ell & -1 \end{bmatrix}$$

Recall now that:

$$\begin{Bmatrix} F_A \\ F_B \end{Bmatrix} = \begin{bmatrix} k_{AA} & k_{AB} \\ k_{AB}^T & k_{BB} \end{bmatrix} \begin{Bmatrix} u_A \\ u_B \end{Bmatrix}$$

- Assume the set $u_A = 0$, then:

$$F_A = k_{AB} u_B$$

$$F_B = k_{BB} u_B$$

but $F_A = T F_B$ so:
$$\boxed{k_{AB} = T k_{BB}}$$

It is then obtained that:

$$k_{AB} = \begin{bmatrix} -EA/e & & \\ & -12EI/e^3 & 6EI/e^2 \\ & -6EI/e^2 & 2EI/e \end{bmatrix}$$

- With a similar approach it is found $k_{BA} = k_{AB}^T$

Summarizing, the stiffness matrix reads

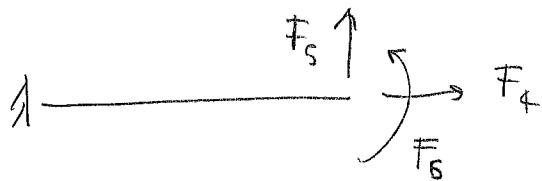
$$k = \begin{bmatrix} EA/e & -EA/e & & \\ 12EI/e^3 & 6EI/e^2 & -12EI/e^3 & 6EI/e^2 \\ 6EI/e^2 & 4EI/e & -6EI/e^2 & 2EI/e \\ \hline EA/e & & & \\ & 12EI/e & -6EI/e^2 & \\ & -6EI/e^2 & 4EI/e & \end{bmatrix}$$

SYM

Axial and bending stiffness

Evaluate how the role played by the axial and the bending stiffness and the corresponding load paths.

Consider the cantilever beam here below



The stiffness matrix was found to be:

$$K_{BB} = \begin{bmatrix} EA/l & & \\ & 12EI/l^3 & -6EI/l^2 \\ & -6EI/l^2 & 4EI/l \end{bmatrix}$$

For comparing the role played by the various contributions it is useful to highlight some nondimensional parameters.

1. EA is, dimensionally, a force so

- Forces can be normalized with respect to EA
- Moments " to $EA \cdot l$

2. The displacements can be normalized with respect to l

$$\bar{F} = \frac{F}{EA} \quad \bar{M} = \frac{M}{EA \cdot l} \quad \bar{u} = \frac{u}{l}$$

$$\left\{ \begin{array}{l} F_4 = \frac{EA}{\ell} u_4 \\ F_5 = \frac{12EI}{\ell^3} u_5 - \frac{6EI}{\ell^2} u_6 \\ F_6 = -\frac{6EI}{\ell^2} u_5 + \frac{4EI}{\ell} u_6 \end{array} \right. \quad \begin{array}{l} \text{Multiply by } \frac{1}{EA} \\ \text{Multiply by } \frac{1}{EA} \\ \text{Multiply by } \frac{1}{EA\ell} \end{array}$$

$$\left\{ \begin{array}{l} \frac{F_4}{EA} = \frac{u_4}{\ell} \\ \frac{F_5}{EA} = \frac{12EI}{EA\ell^2} \frac{u_5}{\ell} - \frac{6EI}{EA\ell^2} u_6 \\ \frac{F_6}{EA\ell} = -\frac{6EI}{EA\ell^2} \frac{u_5}{\ell} + \frac{4EI}{EA\ell^2} u_6 \end{array} \right.$$

Defining now $\lambda^2 = \frac{EA\ell^2}{EI}$ slenderless ratio , it

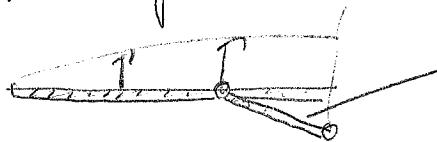
is possible to write the relation between nondimensional forces and non dimensional displacements

$$\left\{ \begin{array}{l} F_4 \\ F_5 \\ F_6 \end{array} \right\} = \begin{bmatrix} 1 & & \\ 12/\lambda^2 & -6/\lambda^2 & \\ -6/\lambda^2 & 4/\lambda^2 & \end{bmatrix} \left\{ \begin{array}{l} \bar{u}_4 \\ \bar{u}_5 \\ \bar{u}_6 \end{array} \right\}$$

Remarks

- When λ^2 is high, the bending-related contributions tend to zero
 - λ^2 is high when $EAl^2 > EI$. This condition is verified whenever:
 1. EA is high
 2. l is high (for fixed sectional properties EA and EI , the value of λ increases as l is increased)
- so: the major contributions to the load carrying capability are due to the axial behavior whenever λ^2 is high.
- This is true for beams characterized high values of axial stiffness EA (high has to be intended with respect to EI) and/or slender beams with high l .
- A structure which is designed for sustaining the loads via axial behaviour is an extremely efficient structure. On the contrary, the efficiency reduces when the loads are carried via bending response.
 - Systems of trusses are highly efficient structures and for this reason they are widely employed in the space field.
 - Even the early aeronautical structures were based on the axial load carrying capabilities of

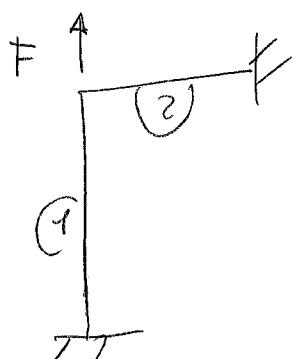
the Wihs



truss

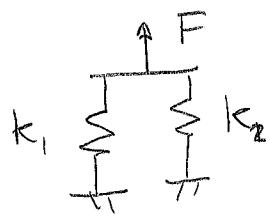
This kind of configuration was changed for aerodynamic reasons. The structure is efficient!

- According to what has been discussed, a statically indeterminate structure is below



will be characterized by a load path such that most of the load is carried via axial response by beam 1.

In fact this structure can be seen as

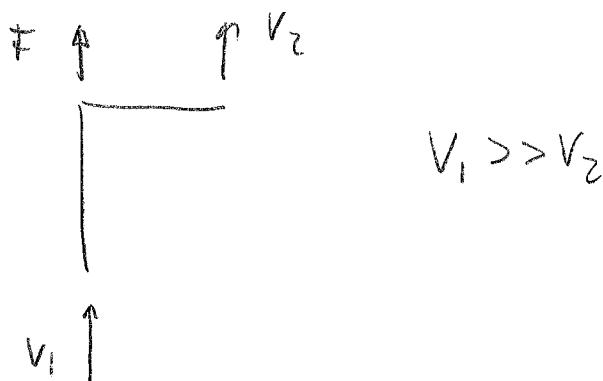


$$\text{where } k_1 = \frac{EA}{l}$$

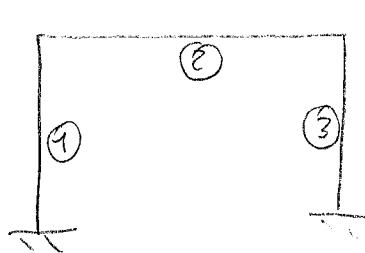
$$k_2 = \frac{3EI}{l^3}$$

$$\begin{aligned} \text{Thus } F_1 &= k_1 u \\ F_2 &= k_2 u \Rightarrow \frac{F_1}{F_2} = \frac{\frac{EA}{l}}{\frac{3EI}{l^3}} = \frac{EA l^2}{3EI} = \frac{l^2}{3} \end{aligned}$$

Clearly the reaction forces behave accordingly

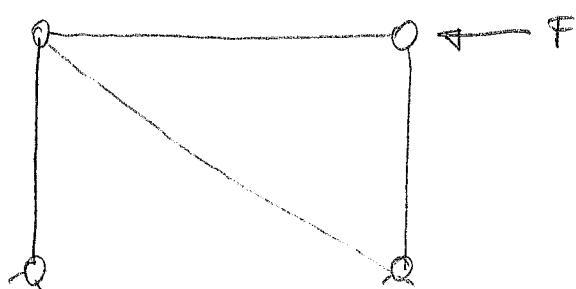


- For the same reasons a beam system as in the figure



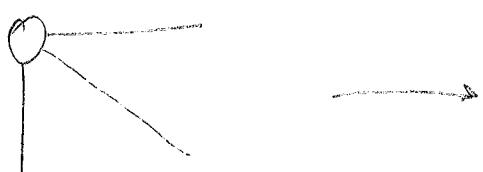
is characterized by a bending response for beams ④ and ③, while beam ② carry the load via axial behaviour.

A much more efficient design would be the following:

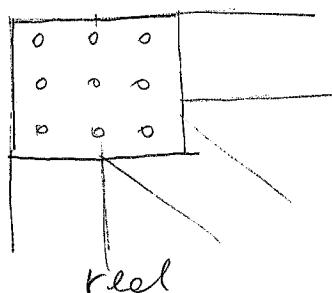


where all of the beams are working in axial mode
(no bending energy is involved)

- Note that the diagonal element is loaded in traction. If the diagonal were placed in the other direction, a buckling requirement would enter the design.
- In the previous sketch the elements were connected with hinges, thus no bending moments could be introduced. This is clearly an idealization as, in a real structure, the constraint will not be a perfect hinge and the introduction of some sort of bending moments would be permitted. Consider for instance this case



ideal

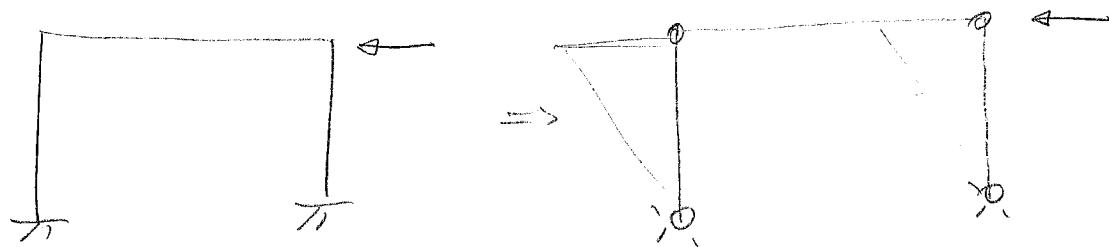


real

If the real structure, after idealizing all the junctions as hinges, is statically determined (thus no kinematisms are possible) then the behaviour will resemble the ideal one, and the loads will be carried via axial behaviour.

In other words, if a purely axial load path is possible, then the structure will use it for carrying loads.
(it is stiffer with respect to bending).

Clearly, if the axial path is not possible due to offset of a mechanism, the response will require the contribution of bending.



(not possible to have a
purely axial load path)

• Contribution of shear to the internal complementary virtual work

It was mentioned that in many cases (more specifically, for slender beams) the energy contribution associated with the shear can be neglected.

To justify this assumption consider the example here below



$$\delta W_{i, \text{bending}}^* = \int_0^l \frac{Fz^2}{EI} dz = \frac{Fl^3}{3EI}$$

$$\delta W_{i, \text{shear}}^* = \int_0^l \delta F \frac{I}{GA^*} dz = \frac{\delta Fl}{GA^*}$$

The ratio between the two contribution is then:

$$\frac{\delta W_{i, \text{bend}}^*}{\delta W_{i, \text{shear}}^*} = \frac{Fl^3}{3EI} \frac{GA^*}{\delta Fl} = \frac{GA^* l^2}{3EI} \propto \frac{Al^2}{J}$$

(A^* is of the order of A ; G is of the order of E)

Thus:

$\frac{\delta W_{i, \text{bend}}^*}{\delta W_{i, \text{shear}}^*} \propto \lambda^2$
--

with λ slenderness ratio

For slender beams the strain energy due to bending is much higher in comparison to the shear contribution.

For thin-walled beams this will (not) be, in general, true

As a matter of fact

$$\frac{6A^*l^2}{3EI} \neq \frac{Al^2}{J} \quad \text{if} \quad A^* \neq A$$

This is the case of a rectangular beam, where

$$A^*/A = 5/6 = \chi \quad (\text{shear factor})$$

For thin walled beams $\chi << 1$ then, in many cases, the shear deformability will not be negligible.