

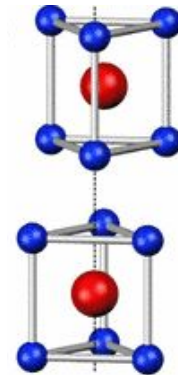
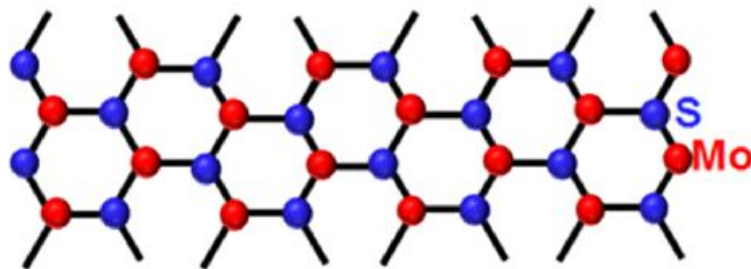
Control of Valley Polarization in Monolayer MoS₂ by Optical Helicity

Ben Safvati

APPPHYS 204 - Quantum Materials

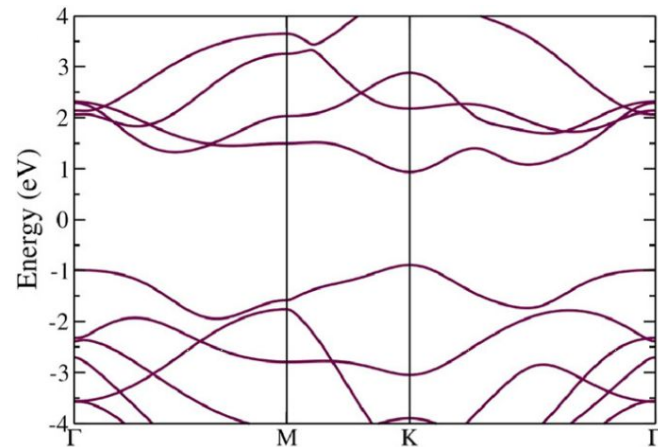
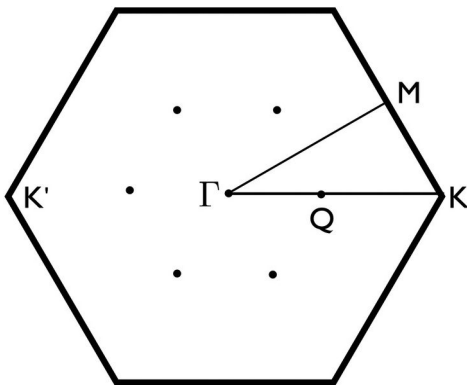
The Valley Degree of Freedom

- Band structures of certain materials contain energy extrema called “valleys.”



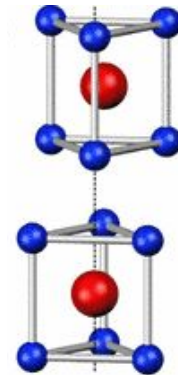
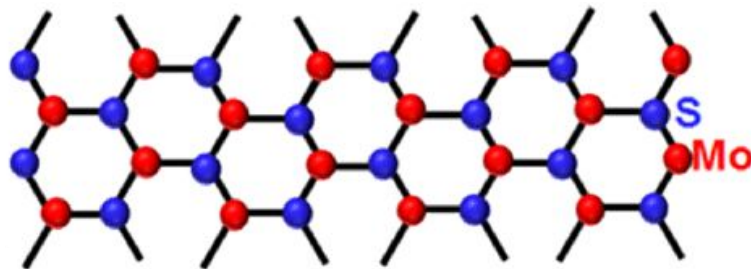
For MoS₂:

- 2 degenerate valleys at K, K' points, opposite in momentum space.
- Direct ~1.8 eV band gap.



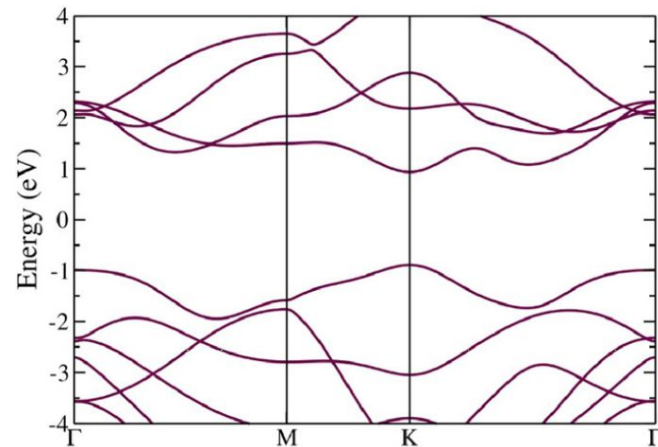
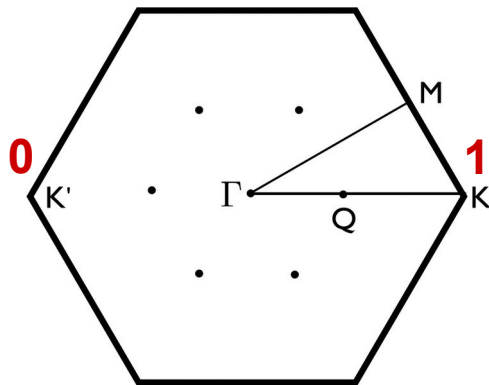
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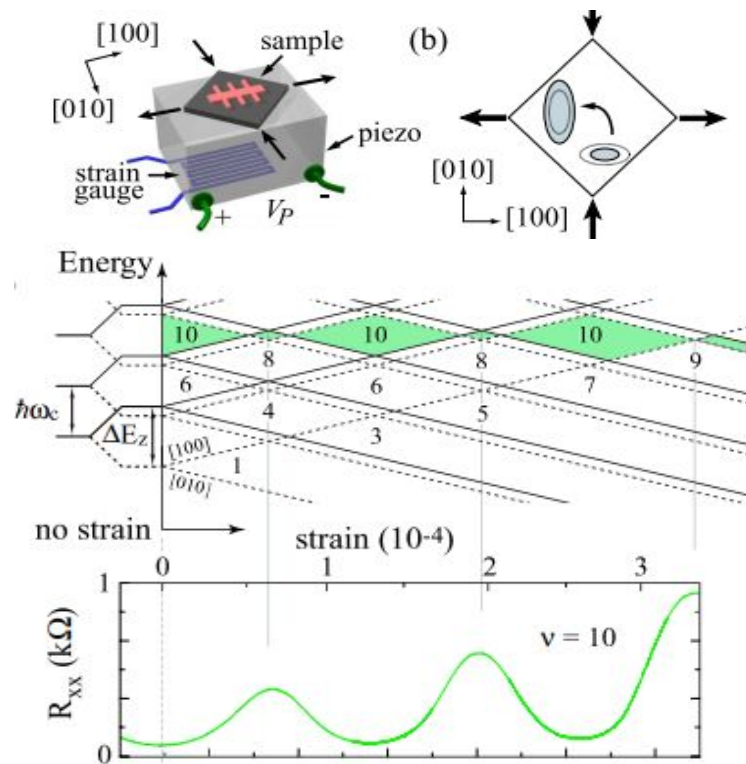
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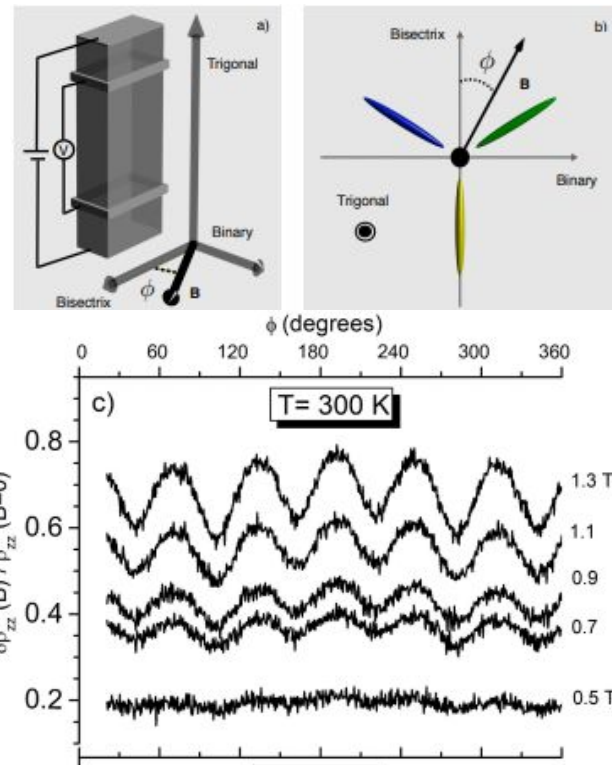
Prior Work on Valley Polarization

Strain-induced valley splitting of FQHE states in AIAs



Bishop, N. C. et al. *Phys. Rev. Lett.* 98, 266404 (2007).

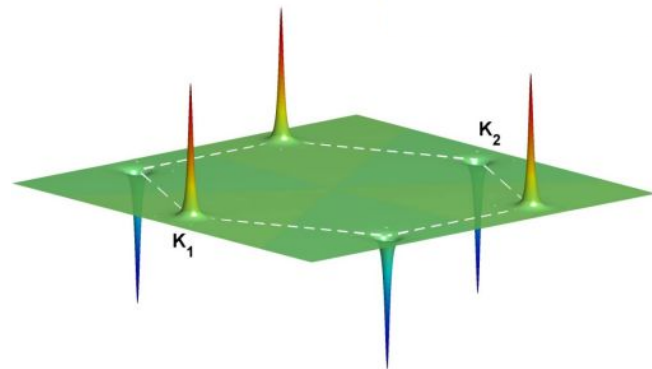
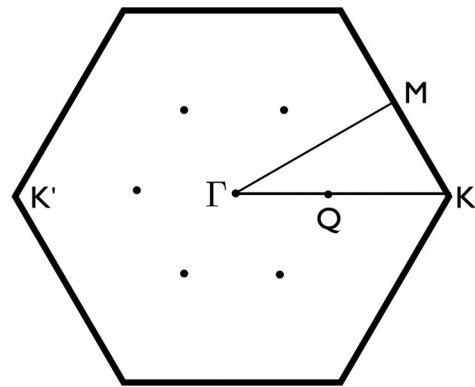
Valley valve in Bi with rotating B field



Zhu, Z., Collaudin, A., Fauqué, B., Kang, W. & Behnia, K. *Nature Phys.* 8, 89–94 (2012).

MoS₂ Symmetry Considerations

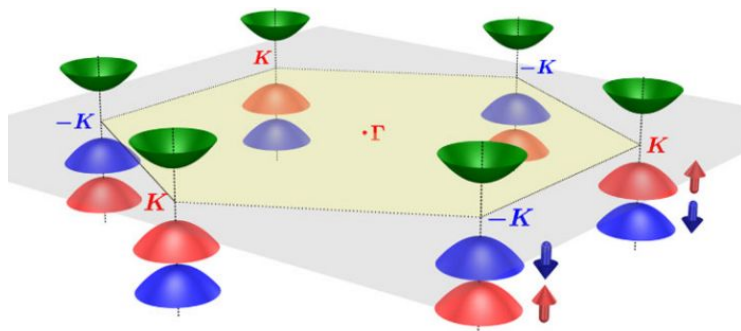
- Consider valley-differentiated orbital magnetic moment $m_v \propto$ valley index $\tau = \pm 1$.
 - τ odd under time reversal and spatial inversion
 - m_v odd under time reversal, but even under spatial inversion.
- Thus, m_v can only be nonzero for inversion symmetry breaking
 - \Rightarrow spatial inversion symmetry breaking is necessary to observe valley-contrasting magnetic moment.**



TMD Minimal Band Model

Symmetry-adapted basis states
between band-edge electrons

$$|\phi_c\rangle = |d_z^2\rangle, \quad |\phi_v^\tau\rangle = \frac{1}{\sqrt{2}}(|d_{x^2-y^2}\rangle + i\tau|d_{xy}\rangle).$$



Spin-Valley Locking

$\mathbf{k} \cdot \mathbf{p}$ Hamiltonian
expanded around K point

$$\hat{H}_0 = at(\tau k_x \hat{\sigma}_x + k_y \hat{\sigma}_y) + \frac{\Delta}{2} \hat{\sigma}_z.$$

Related by
time reversal

Spin-orbit coupling
expanded around K point

$$H_{soc} = \lambda\tau \frac{\hat{\sigma}_z - 1}{2} \hat{s}_z$$

Chiral Optical Selection Rules

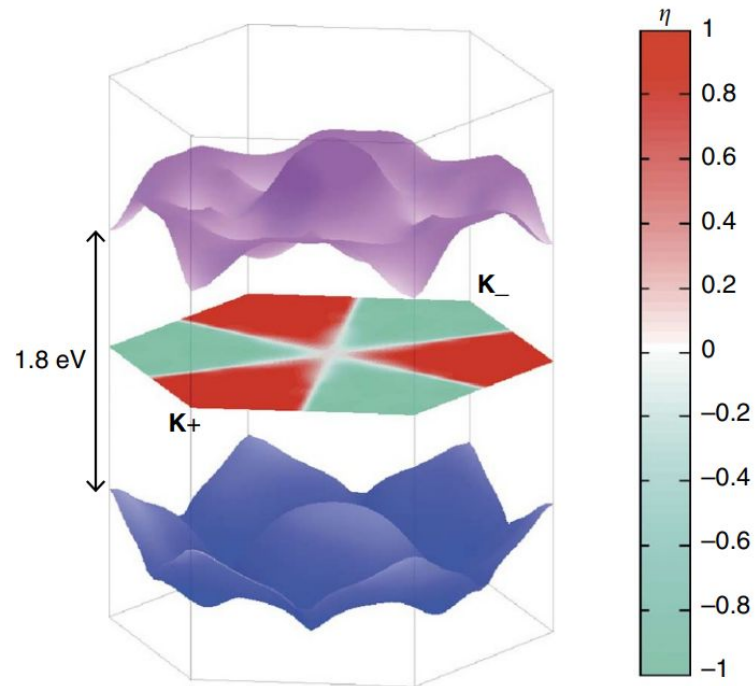
Optical interband transition elements

$$\mathcal{P}_{\pm}(k) \equiv \mathcal{P}_x(k) \pm i\mathcal{P}_y(k)$$

$$\mathcal{P}_{\alpha}(k) \equiv m_0 \langle u_c(k) | \frac{1}{\hbar} \frac{\partial \hat{H}}{\partial k_{\alpha}} | u_v(k) \rangle$$

$$|\mathcal{P}_{\pm}(k)|^2 = \frac{m_0^2 a^2 t^2}{\hbar^2} \left(1 \pm \tau \frac{\Delta'}{\sqrt{\Delta'^2 + 4a^2 t^2 k^2}} \right)^2.$$

$\Delta' \gg atk \Rightarrow$ **Nearly perfect valley selection rules!**



$$\eta(\mathbf{k}, \omega_{cv}) = \frac{|\mathcal{P}_+^{cv}(\mathbf{k})|^2 - |\mathcal{P}_-^{cv}(\mathbf{k})|^2}{|\mathcal{P}_+^{cv}(\mathbf{k})|^2 + |\mathcal{P}_-^{cv}(\mathbf{k})|^2}$$

Chiral Optical Selection Rules

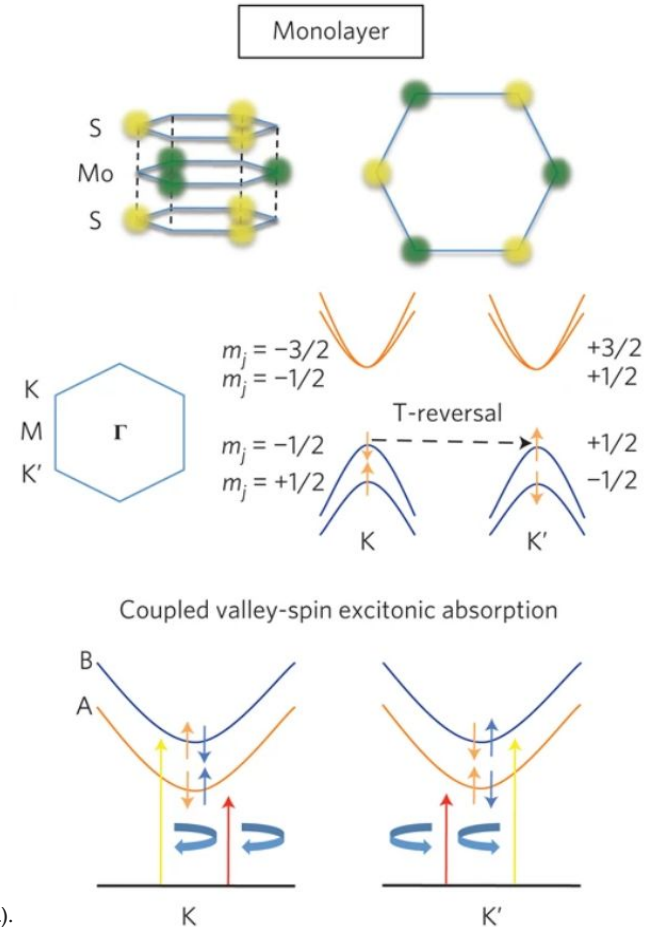
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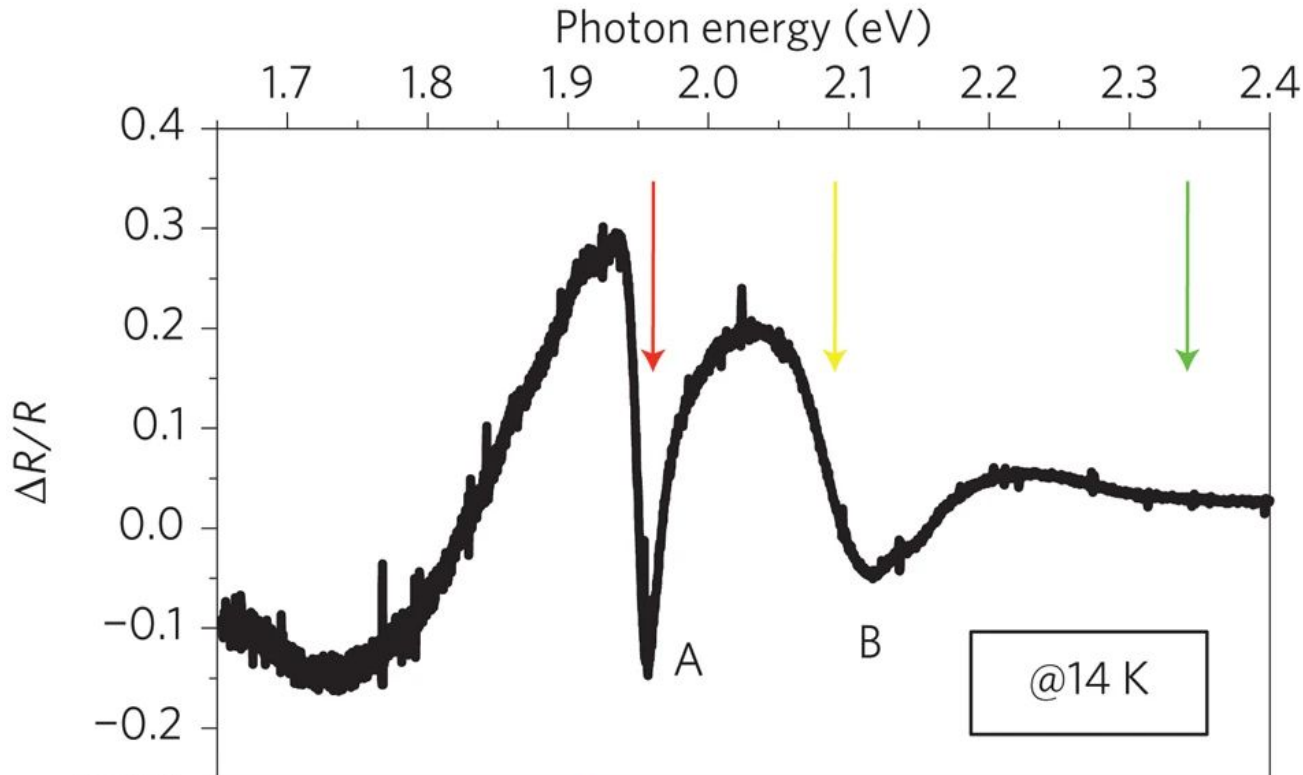
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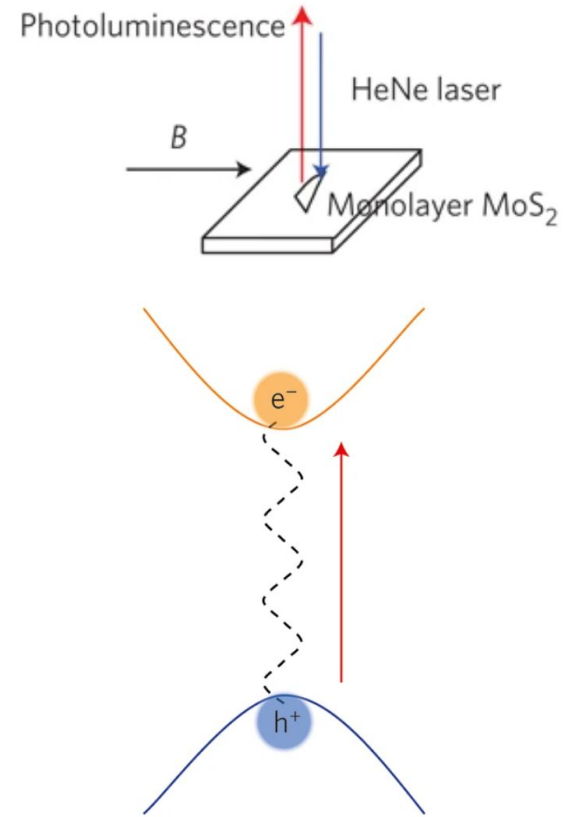
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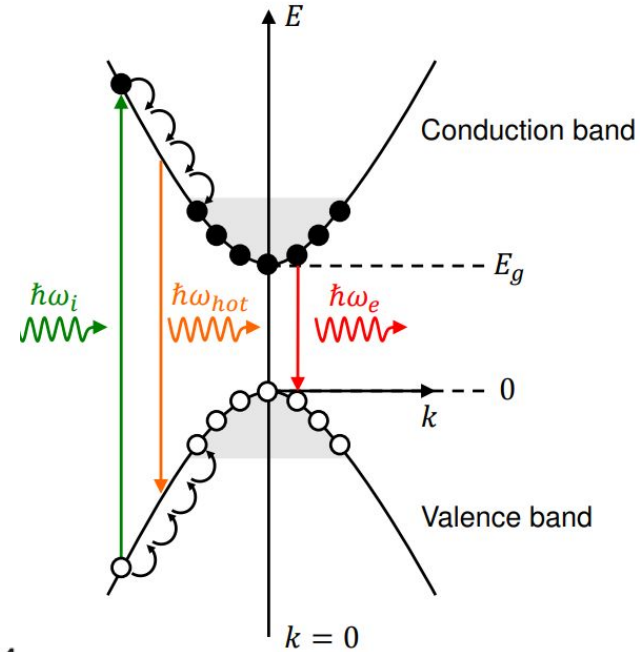
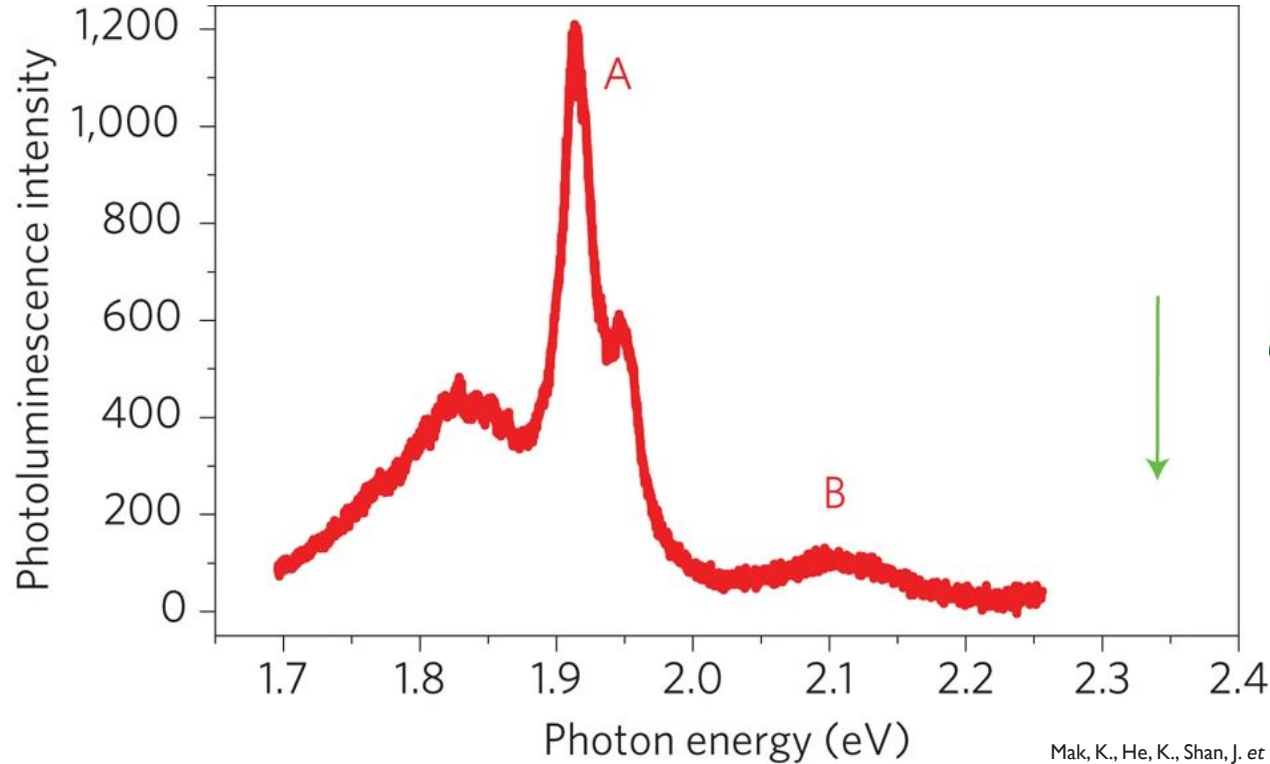
Excitonic Absorption Spectrum



Mak, K., He, K., Shan, J. et al. *Nature Nanotech* 7, 494–498 (2012).



Excitonic Photoluminescence Spectrum



A Exciton Resonance

Perfect helicity

Bilayer MoS₂

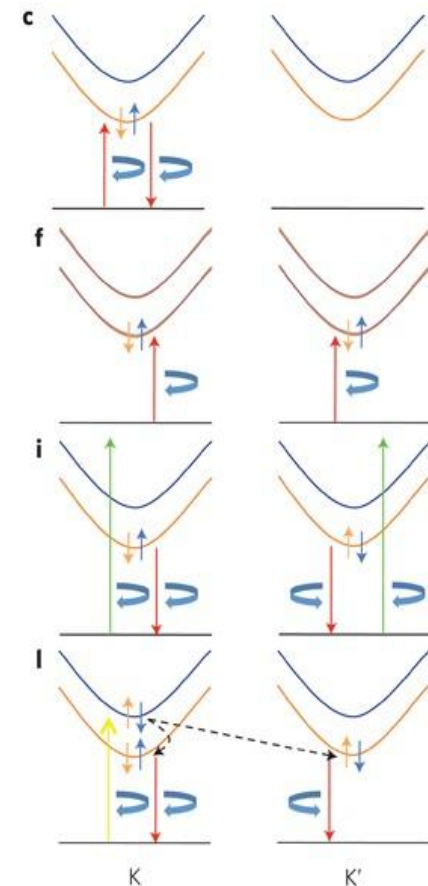
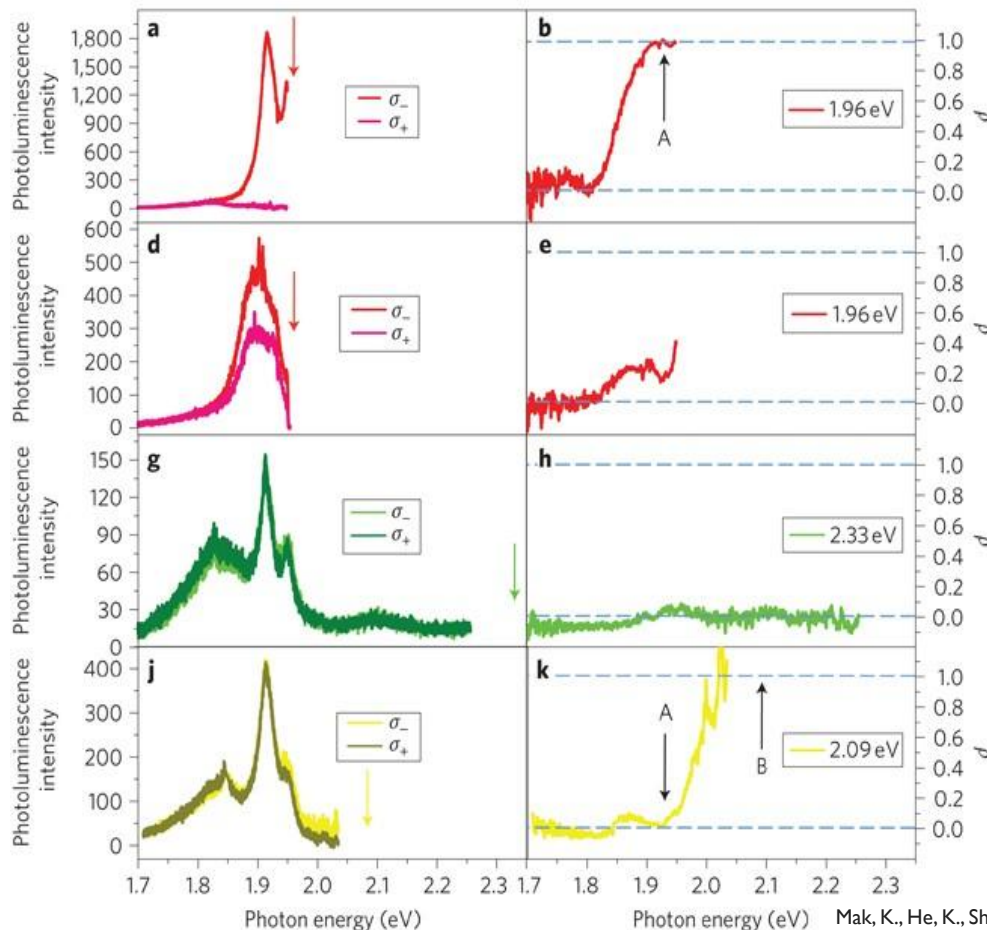
No valley polarization

Off-Resonance

Simultaneous population

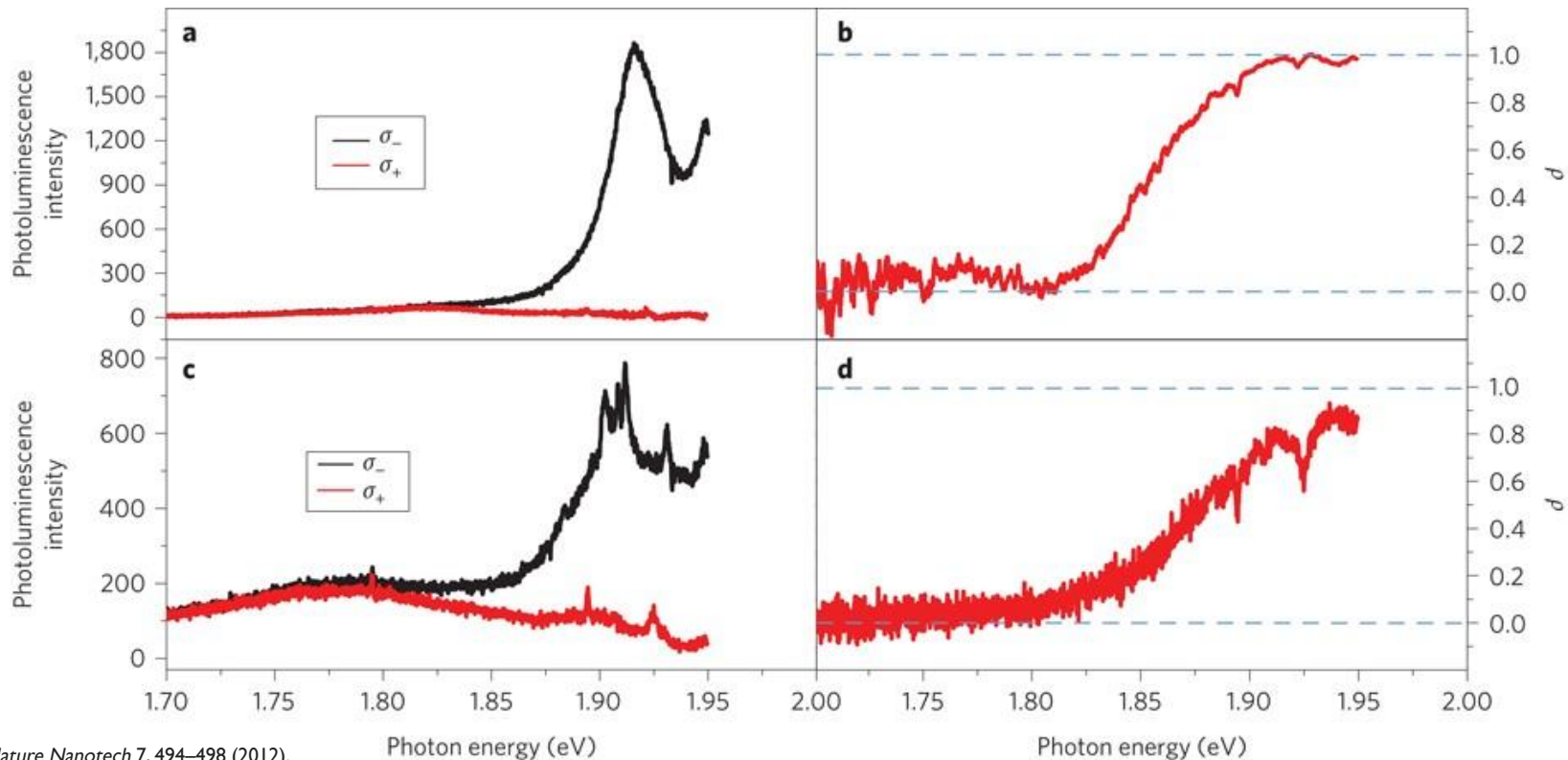
B exciton Resonance

Perfect helicity
(hot luminescence)

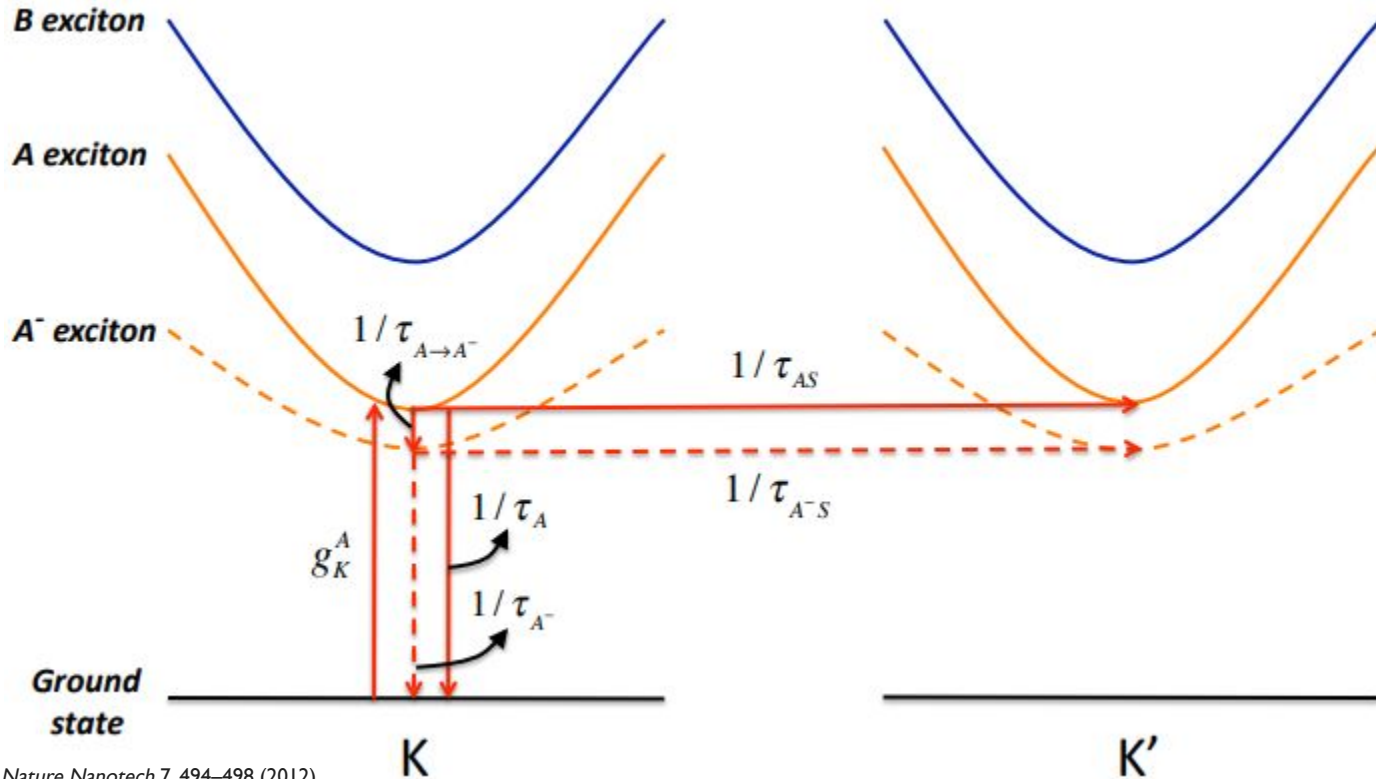


Substrate-Independent Optical Helicity

h-BN



Future Research: Relaxation Mechanisms



Thank you!

Chiral Optical Selection Rules

