

# Thermo-Mechanics

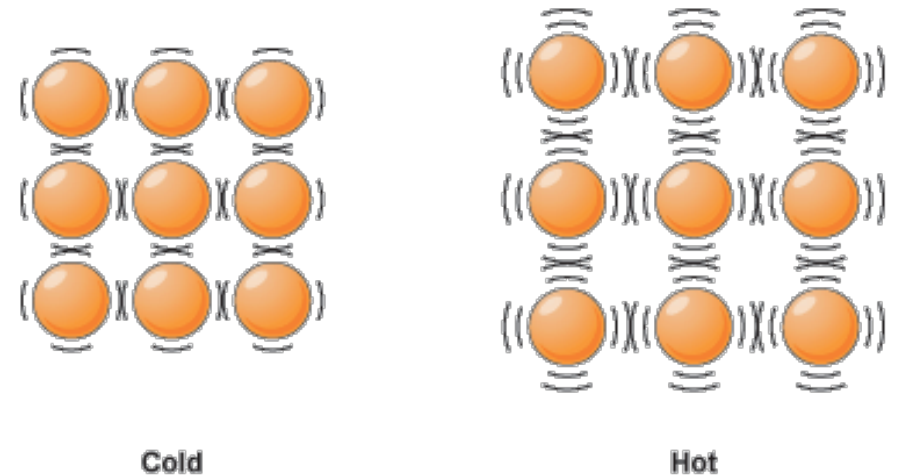
NE 591

## Last Time

- Solid mechanics predicts the deformation of a body from its applied load
  - The strain defines the deformation
  - The stress defines the material's internal response to the strain
- Materials can have recoverable and permanent deformation
  - Elastic deformation is recoverable and occurs due to the stretching of the atomic lattice
  - Plastic deformation is permanent and results from the breaking of bonds during dislocation (line defect) migration
- We derived two analytical solutions for the stress in a pressurized cylinder
  - Thin walled cylinder
  - Any size wall

# Thermal Expansion

- As the temperature increases, atoms have larger vibrations, causing the material to expand
- This expansion doesn't raise the energy of the material when unrestrained
- This expansion can be treated as a strain, but as one that doesn't cause stress



$$\sigma = C (\epsilon - \epsilon_0)$$

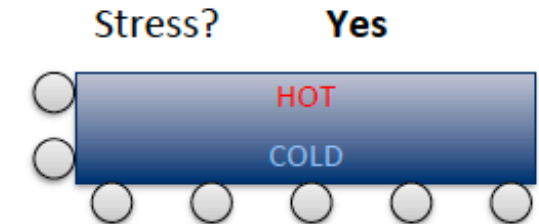
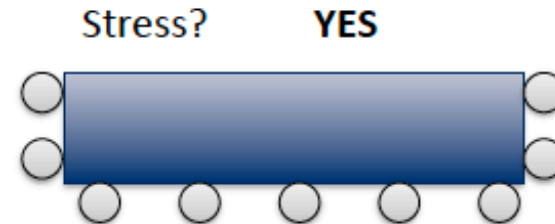
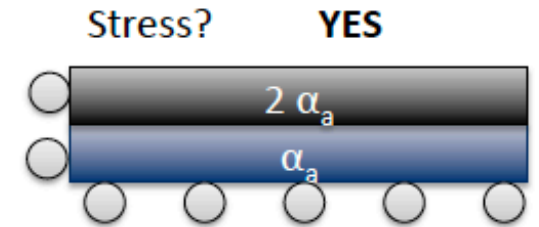
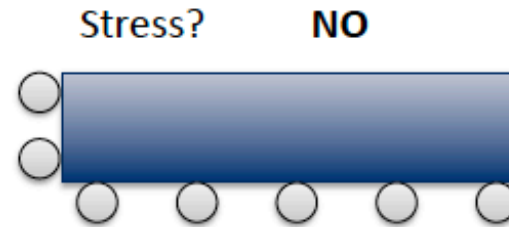
# Thermal Expansion

- In isotropic materials, thermal expansion happens equally in all directions and is linear with temperature
- $\epsilon_0 = (T - T_0)\alpha I$
- In this equation
  - T is the current temperature
  - $T_0$  is the temperature the original size was measured
  - $\alpha$  is the linear thermal expansion coefficient
  - I is the identity tensor

Material	$\alpha$ ( $\times 10^{-6}$ 1/K)
Aluminum	24
Copper	17
Steel	13
UO <sub>2</sub>	11
Zircaloy (Axial)	5.5
Zircaloy (radial)	7.1

# Thermal Expansion

- Though thermal expansion doesn't directly cause stress, it can still lead to thermal stress



## What is the stress in a thin constrained rod of length $L$ when it is heated to $\Delta T$ ?

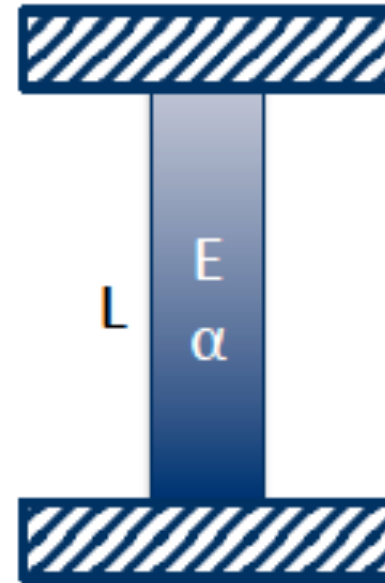
- The rod has a Young's modulus of  $E$  and an expansion coefficient of  $\alpha$

$$\epsilon_0 = (T - T_0)\alpha \quad \sigma = E (\epsilon - \epsilon_0)$$

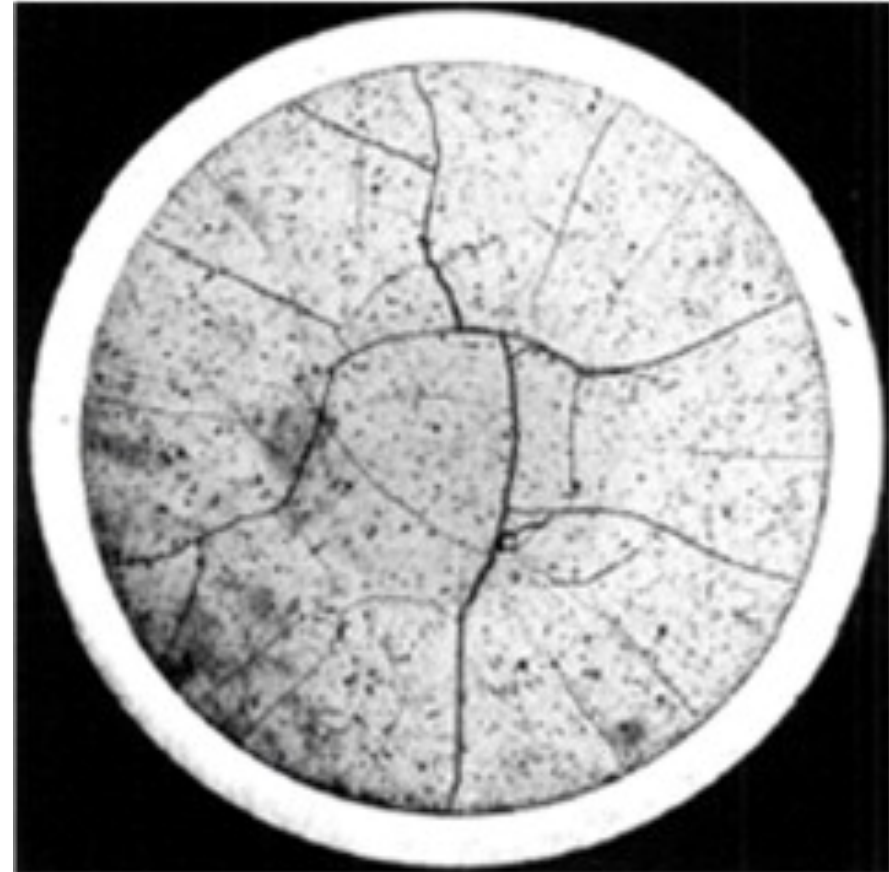
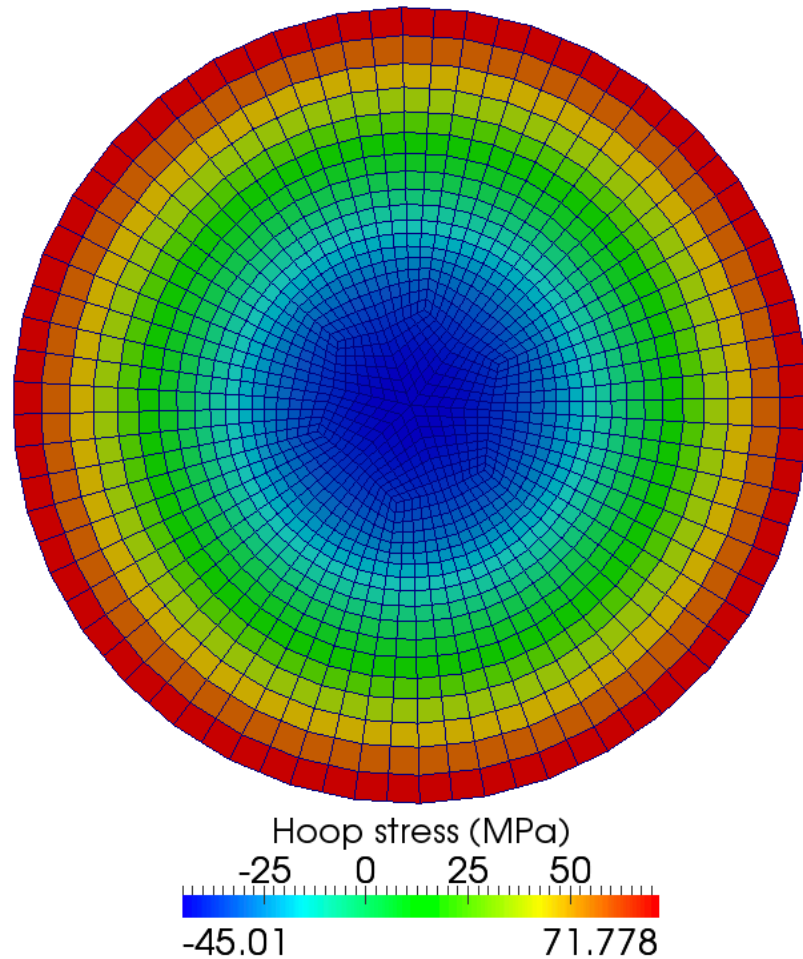
$$\epsilon_0 = (T - T_0)\alpha$$

$$\sigma = E (0 - \Delta T\alpha)$$

$$\sigma = -E\Delta T\alpha$$



**The large temperature gradient within a fuel pellet results in large thermal stresses**



# Consider the material response of the axisymmetric body

- We assume small strains, so the strain is defined as

$$\epsilon_{rr} = \frac{\partial u_r}{\partial r}, \quad \epsilon_{\theta\theta} = \frac{u_r}{r}, \quad \epsilon_{zz} = \frac{\partial u_z}{\partial z}, \quad \epsilon_{rz} = \frac{1}{2} \left( \frac{\partial u_r}{\partial z} + \frac{\partial u_z}{\partial r} \right)$$

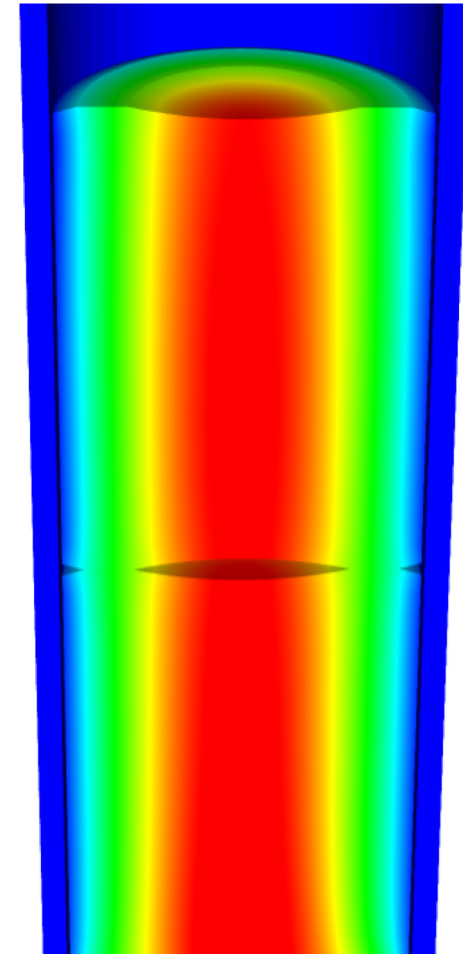
- We assume isotropic material response, so

$$\epsilon_{rr} = \frac{1}{E} (\sigma_{rr} - \nu(\sigma_{\theta\theta} + \sigma_{zz})) + \alpha\Delta T$$

$$\epsilon_{\theta\theta} = \frac{1}{E} (\sigma_{\theta\theta} - \nu(\sigma_{rr} + \sigma_{zz})) + \alpha\Delta T$$

$$\epsilon_{zz} = \frac{1}{E} (\sigma_{zz} - \nu(\sigma_{\theta\theta} + \sigma_{rr})) + \alpha\Delta T$$

$$\epsilon_{rz} = \frac{1}{2G} \sigma_{rz}$$





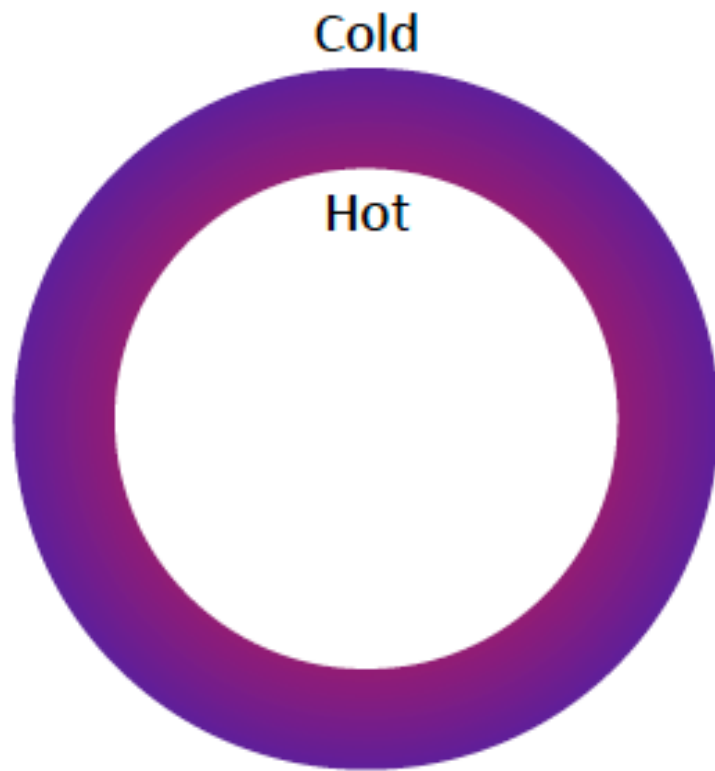
## Consider a cylinder with thermal expansion but not pressure

- $\sigma_{rr}(R_i) = \sigma_{rr}(R_o) = 0$
- Similar to the equations we worked through last time
- $\frac{1}{r^3} \frac{d}{dr} \left( r^3 \frac{d\sigma_{rr}}{dr} \right) = - \left( \frac{\alpha E}{1-\nu} \right) \frac{1}{r} \frac{dT}{dr}$
- Solving this ODE:

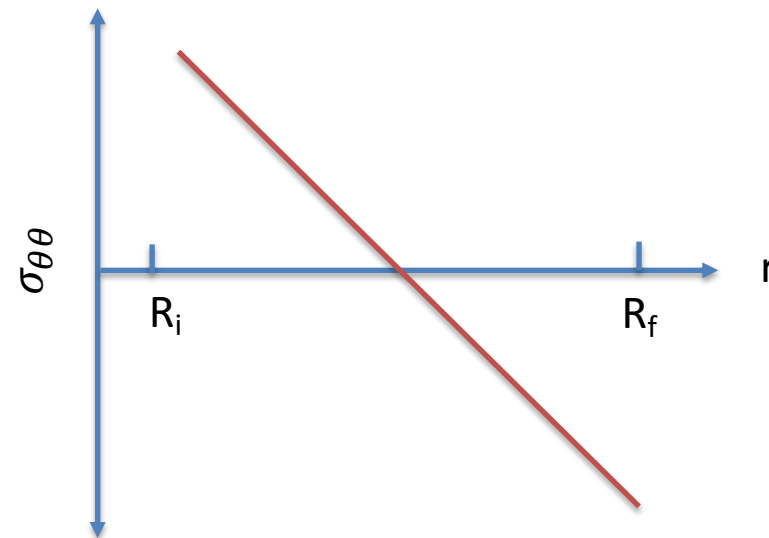
$$\begin{aligned}\sigma_{rr}(r) &= \frac{1}{2} \Delta T \frac{\alpha E}{1-\nu} \left( \frac{r}{R_i} - 1 \right) \left( 1 - \frac{R_i}{\delta} \left( \frac{r}{R_i} - 1 \right) \right) \\ \sigma_{\theta\theta}(r) &= \frac{1}{2} \Delta T \frac{\alpha E}{1-\nu} \left( 1 - 2 \frac{R_i}{\delta} \left( \frac{r}{R_i} - 1 \right) \right) \\ \sigma_{zz}(r) &= \frac{1}{2} \Delta T \frac{\alpha E}{1-\nu} \left( 1 - 2 \frac{R_i}{\delta} \left( \frac{r}{R_i} - 1 \right) \right)\end{aligned}$$



## What is the hoop stress in the cladding?




$$\sigma_{\theta\theta}(r) = \frac{1}{2} \Delta T \frac{\alpha E}{1 - \nu} \left( 1 - 2 \frac{R_i}{\delta} \left( \frac{r}{R_i} - 1 \right) \right)$$



## Where is hoop stress equal to zero?

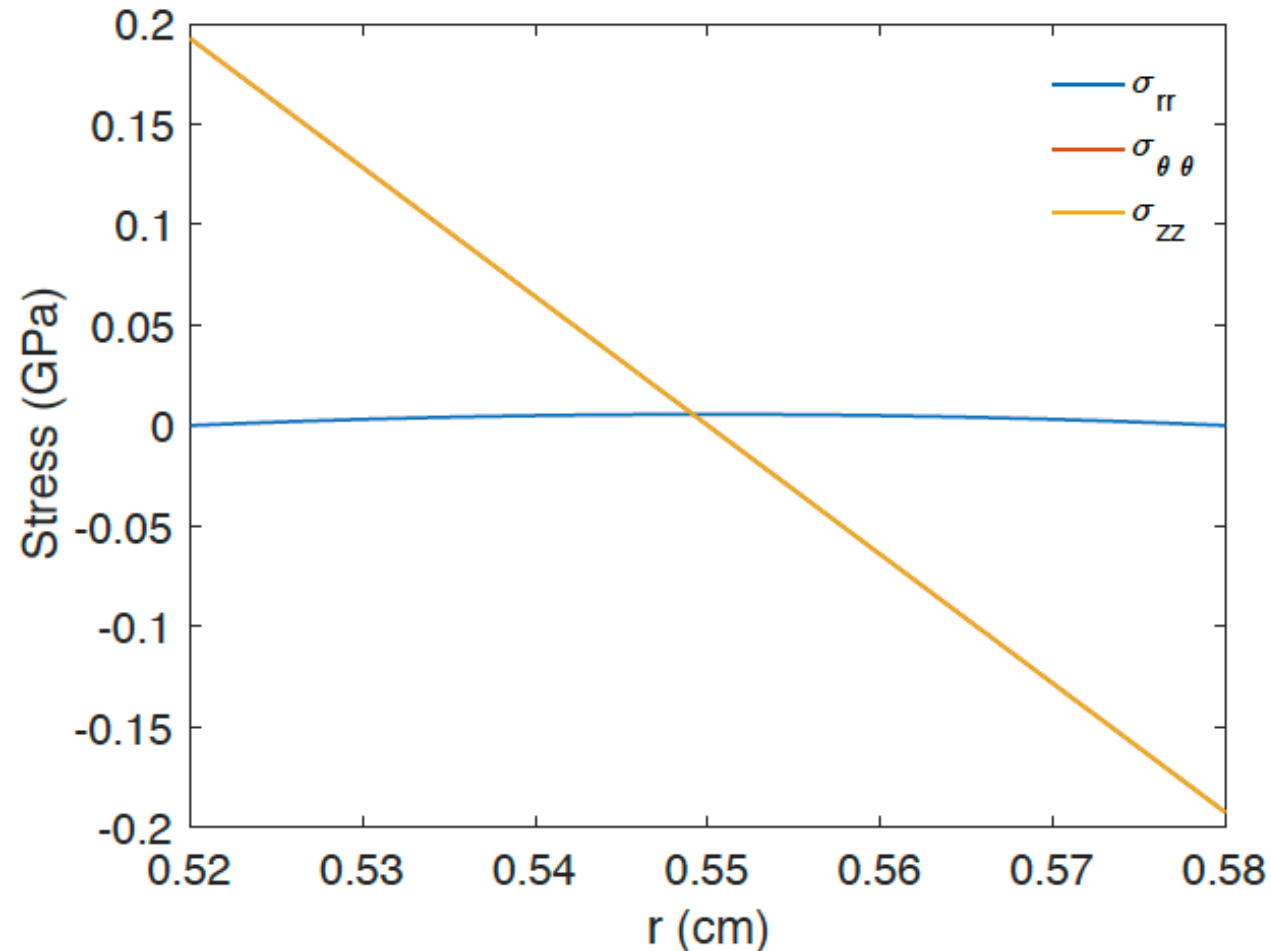
$$\sigma_{\theta\theta}(r) = \frac{1}{2}\Delta T \frac{\alpha E}{1-\nu} \left( 1 - 2\frac{R_i}{\delta} \left( \frac{r}{R_i} - 1 \right) \right)$$

$$\sigma_{\theta\theta}(r) = \frac{1}{2}\Delta T \frac{\alpha E}{1-\nu} \left( 1 - 2\frac{R_i}{\delta} \left( \frac{r}{R_i} - 1 \right) \right) = 0$$


$$\left( 1 - 2\frac{R_i}{\delta} \left( \frac{r}{R_i} - 1 \right) \right) = 0 \quad \longrightarrow \quad 2\frac{R_i}{\delta} \left( \frac{r}{R_i} - 1 \right) = 1 \quad \longrightarrow \quad \left( \frac{r}{R_i} - 1 \right) = \frac{\delta}{2R_i}$$

$$r = \frac{\delta}{2} + R_i$$

# The linear temperature gradient across the cladding causes axial thermal stresses



# Same approach to the thermal stress in a fuel pellet

- The thermal stress is due to the temperature gradient

$$T(r) = \frac{LHR}{4\pi k} \left( 1 - \frac{r^2}{R_f^2} \right) + T_s$$

$$T_0 - T_s = \frac{LHR}{4\pi k}$$

$$T - T_s = (T_0 - T_s) \left( 1 - \frac{r^2}{R_f^2} \right)$$

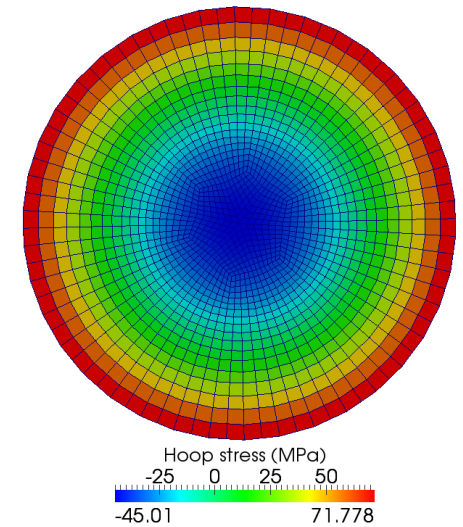
$$\frac{\partial T}{\partial r} = 2(T_0 - T_s) \left( \frac{r}{R_f^2} \right)$$

$$\frac{1}{r^3} \frac{d}{dr} \left( r^3 \frac{d\sigma_{rr}}{dr} \right) = - \left( \frac{\alpha E}{1 - \nu} \right) \frac{1}{r} \frac{dT}{dr}$$

$$\eta = \frac{r}{R_f}$$

$$\sigma^* = \frac{\alpha E (T_0 - T_s)}{4(1 - \nu)}$$

$$\frac{d}{d\eta} \left( \eta^3 \frac{d\sigma_{rr}}{d\eta} \right) = 8\sigma^* \eta^3$$



## Solve this stress ODE

- The boundary conditions are:

$$\frac{d\sigma_{rr}}{d\eta} = 0 \text{ at } \eta = 0$$

$$\sigma_{rr} = 0 \text{ at } \eta = 1$$

$$\frac{d}{d\eta} \left( \eta^3 \frac{d\sigma_{rr}}{d\eta} \right) = 8\sigma^* \eta^3 \quad \sigma^* = \frac{\alpha E (T_0 - T_s)}{4(1 - \nu)}$$

- Once we solve it, we obtain
- Then we can solve the hoop stress
- The axial stress is much more complicated to obtain, but you end up with

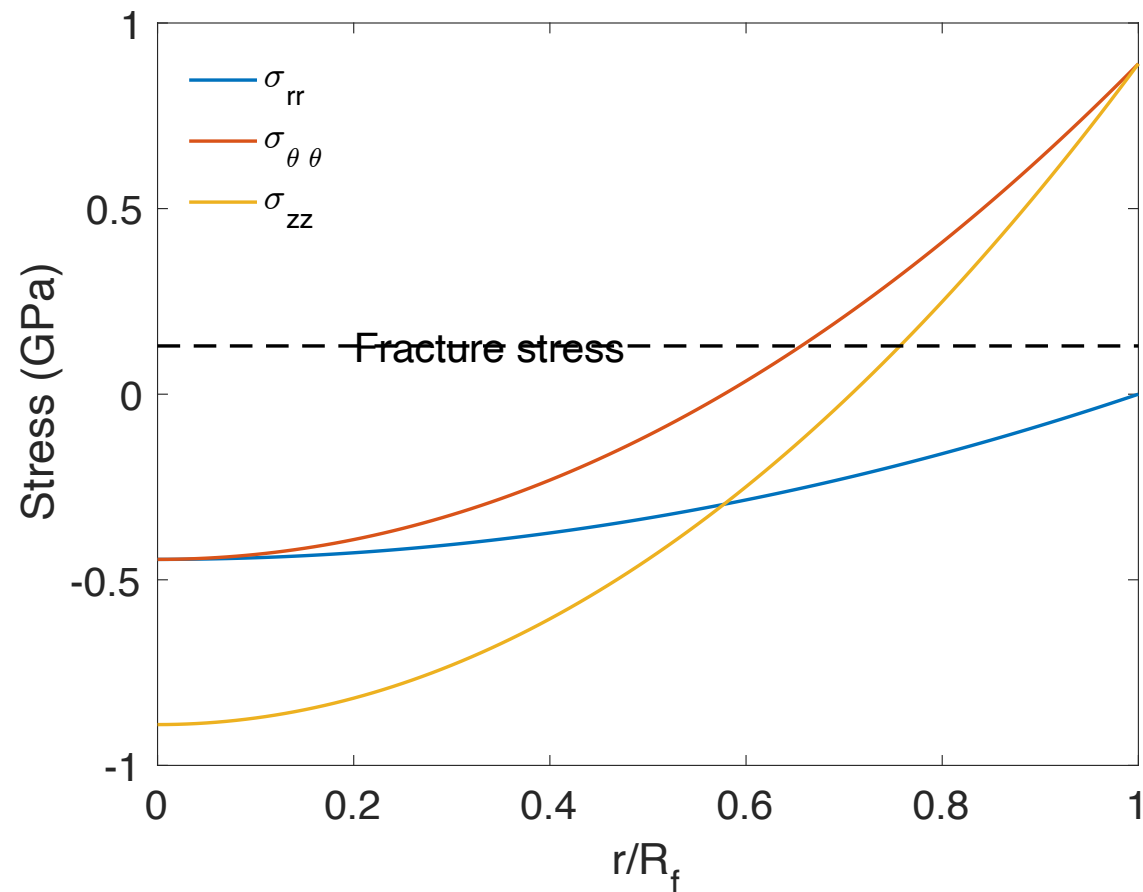
$$\sigma_{rr}(\eta) = -\sigma^*(1 - \eta^2)$$

$$\sigma_{\theta\theta} = \sigma_{rr} + r\sigma_{rr,r}$$

$$\sigma_{\theta\theta}(\eta) = -\sigma^*(1 - 3\eta^2)$$

$$\sigma_{zz}(\eta) = -2\sigma^*(1 - 2\eta^2)$$

# The fuel temperature gradient causes large thermal stresses



## How far do fuel cracks extend?

$$\sigma_{\theta\theta}(\eta) = -\sigma^*(1 - 3\eta^2) \quad \sigma^* = \frac{\alpha E(T_0 - T_s)}{4(1 - \nu)} \quad \eta = \frac{r}{R_f}$$

- $E = 200 \text{ GPa}$ ,  $\nu = 0.345$ ,  $\alpha = 11.0\text{e-}6 \text{ 1/K}$ ,  $\sigma_{\text{fr}} = 130 \text{ MPa}$ ,  $\Delta T = 550 \text{ K}$
- Solve for  $\eta$ 
  - $-\sigma_{\text{fr}} / \sigma^* = 1 - 3 \eta^2$
  - $3 \eta^2 = 1 + \sigma_{\text{fr}} / \sigma^*$
  - $\eta = ( (1 + \sigma_{\text{fr}} / \sigma^*) / 3 )^{1/2}$
- $\sigma^* = 11.0\text{e-}6 * 200 * 550 / (4 * (1 - 0.345)) = 461.8 \text{ MPa}$
- $\eta = \text{sqrt}( (1 + 130/461.8) / 3 ) = 0.65$



# Analytical Thermomechanics Summary

- Increases in temperature cause most materials to expand
- We call this thermal expansion, which is a strain that doesn't inherently cause stress
- Thermal expansion can cause stress when
  - Deformation is constrained
  - There are gradients in the expansion coefficient
  - There is a temperature gradient
- We have analytical equations for thermal stresses in the cladding and in the fuel

## The gap changes as a function of time

- Both the pellet and the cladding swell

$$\Delta\delta_{gap} = \delta_{gap} - \delta_{gap}^0$$

$$\Delta\delta_{gap} = \Delta\bar{R}_C - \Delta R_f$$

$$\frac{\Delta R_f}{\bar{R}_C} = \alpha_f (\bar{T}_f - T_{fab})$$

$$\frac{\Delta R_C}{\bar{R}_C} = \alpha_C (\bar{T}_C - T_{fab})$$

$$\Delta\delta_{gap} = \bar{R}_c \alpha_C (\bar{T}_C - T_{fab}) - \bar{R}_f \alpha_f (\bar{T}_f - T_{fab})$$

- But, as the gap decreases, the temperature changes, which again makes the gap change
- The solution using the analytical equations is iterative, due to the dependence of the gap size and temperature.

## Calculate the steady state temperature profile in the rod, including thermal expansion

- $LHR = 200 \text{ W/cm}$ ,  $\delta_{gap}^0 = 30 \text{ }\mu\text{m}$ ,  $R_f = 0.5$ ,  $T_{cool} = 580 \text{ K}$ ,  $T_{EXP,0} = 373 \text{ K}$ ,  
 $k_{gap} = 0.0026 \text{ W/cm-K}$ ,  $t_c = 0.06 \text{ cm}$ ,  $\alpha_f = 11.0\text{e-}6 \text{ 1/K}$ ,  $\alpha_c = 7.1\text{e-}6 \text{ 1/K}$   

$$\Delta\delta_{gap} = \bar{R}_c\alpha_c (\bar{T}_c - T_{fab}) - \bar{R}_f\alpha_f (\bar{T}_f - T_{fab}) = \Delta R_c - \Delta R_f \quad \Delta T_{gap} = \frac{LHR}{2\pi R_f k_{gap}/\delta_{gap}}$$
- $\Delta T_{cool} = 25.5 \text{ K}$ ,  $\Delta T_{clad} = 22.5 \text{ K}$ ,  $\Delta T_{fuel} = 530.5 \text{ K}$
- So,  $T_{IC} = 580 + 25.5 + 22.5 = 628.0 \text{ K}$ ,  $T_s = 701.5 \text{ K}$ ,  $T_0 = 1232.0 \text{ K}$
- First, we will deal with expansion in the cladding
  - $Av(R_c) = 0.5 + 30\text{e-}4 + 0.06/2 = 0.533 \text{ cm}$
  - $Av(T_c) = 580 + 25.5 + 22.5/2 = 616.75 \text{ K}$
  - $\Delta R_c = 0.533 * 7.1\text{e-}6 * (616.75 - 373) = 9.22\text{e-}4 \text{ cm}$

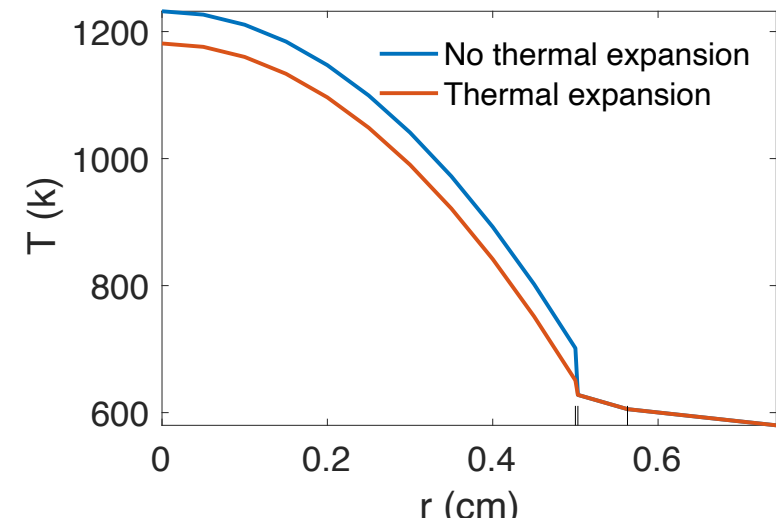
## Calculate the steady state temperature profile in the rod, including thermal expansion

- Second, we deal with the fuel
    - $A_v(T_f) = (1232 + 701.5)/2 = 966.7 \text{ K}$
    - $\Delta R_f = 0.5 \cdot 11 \text{e-}6 \cdot (966.7 - 373) = 0.0033 \text{ cm}$
- $$\Delta \delta_{gap} = \bar{R}_c \alpha_c (\bar{T}_c - T_{fab}) - \bar{R}_f \alpha_f (\bar{T}_f - T_{fab})$$
- The total change in the gap is  $9.22 \text{e-}4 - 0.0033 = -0.0023$
  - However, that means the gap is smaller and so our temperatures were wrong!

## This calculation is repeated until the gap width stops changing significantly

- The change in the gap does NOT effect the coolant or cladding temperatures, just the gap and fuel temperatures.
- We only need to repeat the calculation of the fuel and cladding temperatures and the change in the gap

Iteration	$\delta_{\text{gap}}$ (cm)	$T_s$ (K)	$T_o$ (K)
0	0.003	701	1232
1	0.00066	644	1174
2	0.00097	652	1182
3	0.00094	651	1181
4	0.00094	651	1181



## Next time

- Solve the temperature and the displacement vector for the full thermomechanical problem