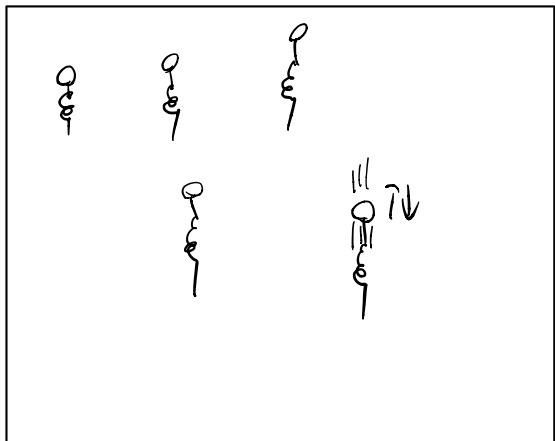


Quantum mechanics

Origins of QM:

Blackbody radiation

accelerated charged particles radiate light



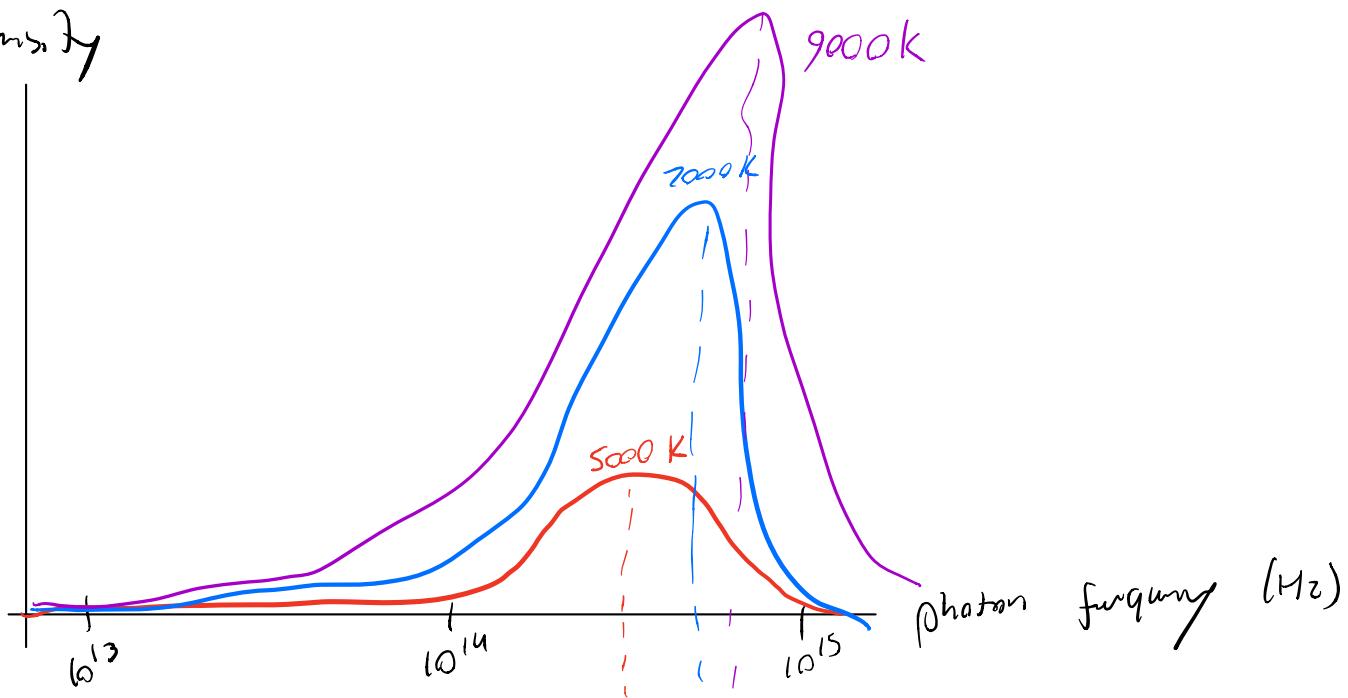
shorter jiggling = higher frequency and more compton light

Sun emits visible spectrum because $T \approx 6000K$

you glow in IR $\odot T \approx 300K$

Blackbody spectrum carefully measured by end of 1800s

Intensity



The law of shape of curve was still unknown

Rayleigh + Jeans: equilibrium theory

→ heat jiggles particles in all ways they can be jiggled

→ energy is equally partitioned over all possible energy states

Rayleigh - Jeans Law:



The problem:

assumes vibrations can be infinitely subdivided



can subdivide
into infinitesimally
small high-w states

Max Planck: (out of desperation) particles can only vibrate with $E = hf \cdot n$

↑
quantized energy states!

- h had yet to be measured, just a small number
- would become Planck's constant!
- quantizing energies allowed Planck to quantize radiation in a way that solved the ultraviolet catastrophe

~ why?

→ limited how much energy the high-freq vibrations could hold!

Planck's law

$$B = \frac{2hf^3}{c^2} \frac{1}{e^{hf/k_B T} - 1}$$

described the shape of a BB spectrum perfectly!

Planck didn't think quantization was real

↳ expected $n=0 \rightarrow$ no interval between any states,
no quantization,
no min energy

Einstein: photoelectric effect (1921 nobel prize)

→ realized that actually light is quantized!

↳ vibration quantized because can only emit/absorb
1 photon at a time

Photoelectric effect

Cs + Zn emit electrons when exposed to light

If light is just a wave, what happens
if we turn up brightness?

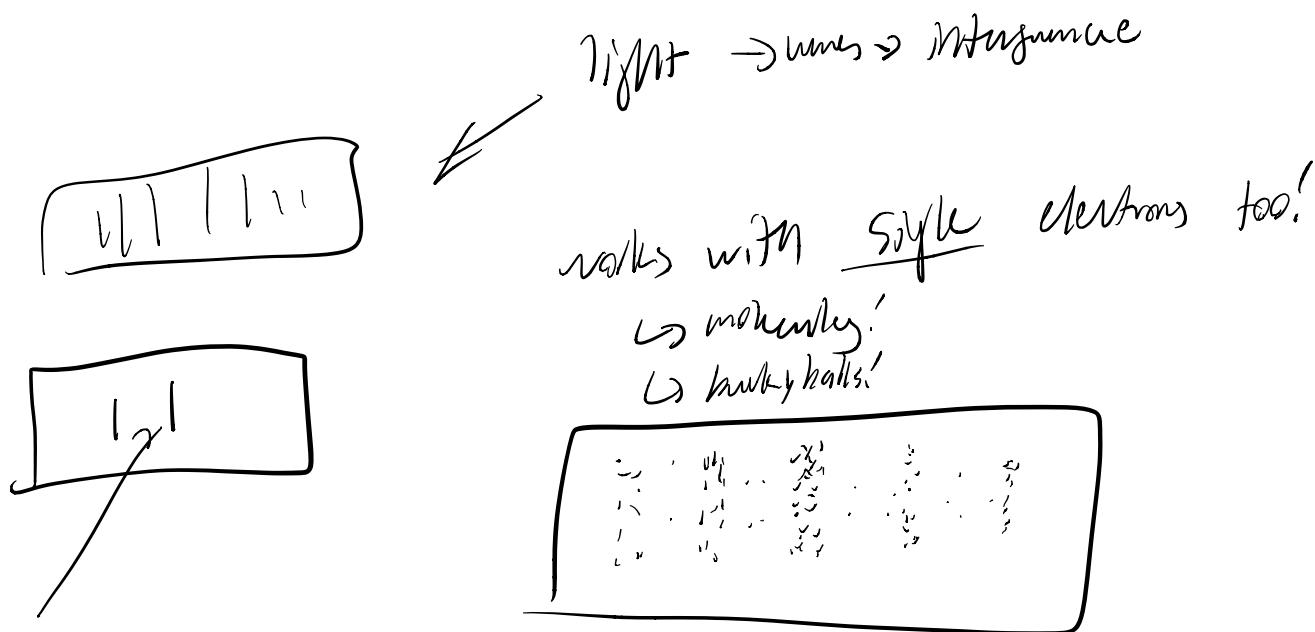
↪ electrons have more energy

Not the case!

Energy of emitted electrons directly
proportional to frequency of photons,
independent of brightness

1905: Einstein proposes particles of light
called photons with E = hf

Doubt sft experiment



electron chooses "landig" spots randomly, but distribution follows interference pattern!

↳ each electron travels through both slits as well! never interacts with itself

↳ somehow, we reduce it to a single position when you name it

interference peaks \rightarrow more probable regions to find particle
"probability wave"

other properties (momentum, pos) follow similar wave behaviors
distribution of properties: wavefunction Ψ

Deterministic wavefunction is the heart of QM!

What is their worth of? Why does it collapse to a point when you measure position?

↳ Open question, open to interpretation

Bohr + Heisenberg: Copenhagen interpretation

wavefunctions don't have physical location

wave function has all possible configurations of particles

movement \rightarrow wavefunction collapse

unlike Schrödinger all possible realities at once!

↳ realities can interfere with each other!

outcome is fundamentally random

- QM + Copenhagen producing stunningly accurate predictions of reality!

- But, not only interpretation:

- Many worlds

- De Broglie - Bohm \rightarrow particles in definite positions guided by pilot wave eqn

- Objective collapse (Ghirardi - Rimini - Weber)

- Consciousness causes collapse (von Neumann - Wigner)

Dynamics of the wavefunction

containing!

$\psi(x, t)$ — wavefunction

$\int_a^b |\psi(x, t)|^2 dx = \text{prob. of finding particle between } a \text{ and } b \text{ at time } t$

$$\hookrightarrow |\psi|^2 = \psi^* \psi$$

(conserving)

$$\int_{-\infty}^{\infty} |\psi(x, t)|^2 dx = 1 \quad \leftarrow \text{normalization}$$

wavefunction evolution determined by Schrödinger eqn:

$$\frac{\hbar^2}{2m} \frac{\partial^2 \psi}{\partial x^2} = -i\hbar \frac{\partial \psi}{\partial t} + V\psi(x, t) \quad \text{time dependent}$$

Wave equation!

$$\frac{\hbar^2}{2m} \frac{\partial^2 \psi}{\partial x^2} = E\psi(x) - V(x)\psi(x) \quad \text{time independent}$$

Expectation of an operator

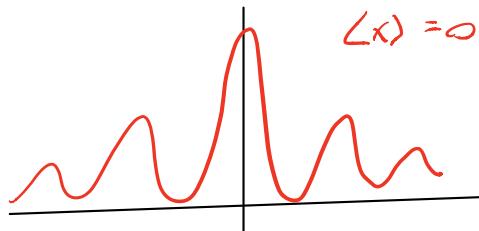
properties got promoted to "operators"

$\langle \hat{Q} \rangle$ = expected value of \hat{Q}

$$\langle \hat{Q} \rangle = \int_{-\infty}^{\infty} \psi^*(x,t) \hat{Q} \psi(x,t) dx$$

$$\hat{x} = x$$

$$\langle x \rangle = \int_{-\infty}^{\infty} \psi^*(x,t) x \psi(x,t) dx = \int_{-\infty}^{\infty} x |\psi(x,t)|^2 dx$$

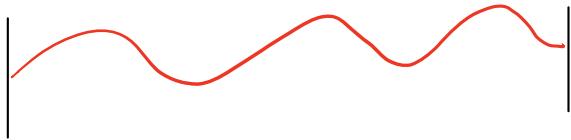


$$\hat{p} = -i\hbar \frac{\partial}{\partial x}$$

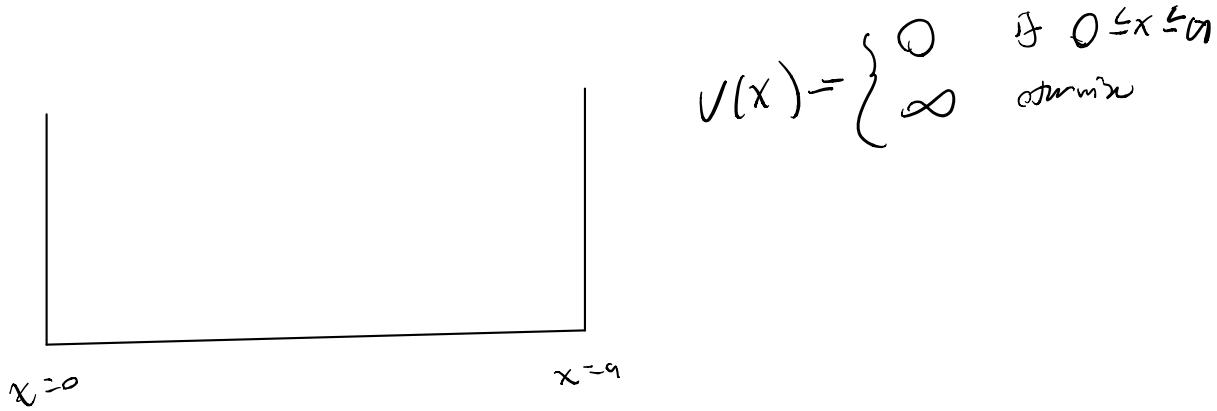
$$\langle p \rangle = \int_{-\infty}^{\infty} \psi^*(x,t) \frac{\partial}{\partial x} \psi(x,t) dx$$

Partikel in a box

Wur munition:



Potential "Box"



At all points:

$$\frac{\hbar^2}{2m} \frac{\partial^2}{\partial x^2} \psi(x) = (E - V) \psi(x)$$

Outside well: $\psi(x) = 0$ is only sensible solution

Inside well: $\psi(x) = A \sin kx + B \cos kx$

$$\psi(0) = 0 : B = 0$$

$$\psi(a) = 0 : k_a = \pm n\pi \rightarrow k = \frac{n\pi}{a}$$

1

$$\psi(x) = A \sin\left(\frac{n\pi}{a}x\right)$$

find A : $\int_{-\infty}^{\infty} \psi^2(x) dx = \int_0^a A^2 \sin^2\left(\frac{n\pi}{a}x\right) dx = A^2 \frac{a}{2} = 1$

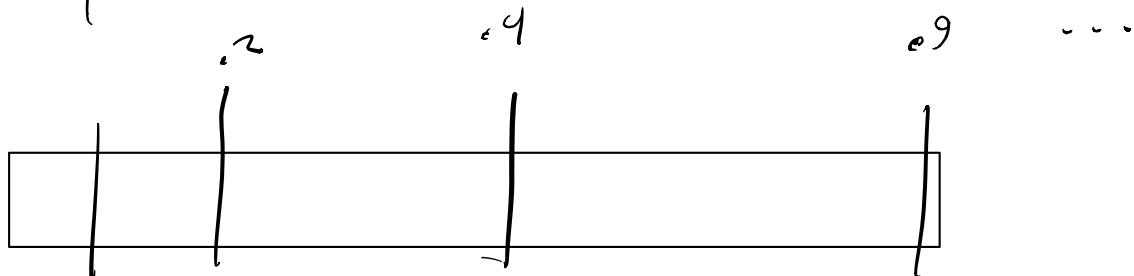
\downarrow
 $A = \sqrt{\frac{2}{a}}$

$$\psi(x) = \sqrt{\frac{2}{a}} \sin\left(\frac{n\pi}{a}x\right)$$

possible energies: $E_n = \frac{\hbar^2}{2m} k^2$

$$= \frac{\hbar^2}{2m} \left(\frac{n\pi}{a}\right)^2$$

Spectrum:



$$E = \frac{\hbar^2 \pi^2}{2ma^2} \cdot 1$$

Eigen states: \leftarrow ignoring normalization

$$\psi_1(x) = \sin\left(\frac{\pi x}{a}\right) \text{ is a valid waveform... so is}$$

$$\psi_2(x) = \sin\left(2\frac{\pi x}{a}\right).$$

Any linear combination $A\psi_1(x) + B\psi_2(x)$ is also a valid solution!

$$\psi_1 \text{ has energy } E_1 = \frac{\hbar^2}{2m} \cdot 1$$

$$\psi_2 \text{ has energy } E_2 = \frac{\hbar^2}{2m} \cdot 4$$

\rightarrow If we have $\psi = \psi_1 + \psi_2$ and we measure energy of particle, we get 50/50 chance of

$$E_1 \text{ or } E_2$$

\rightarrow If we have $\psi = 1\psi_1 + 2\psi_2 + 3\psi_3$, we get

$$\frac{1}{1^2+2^2+3^2} = \frac{1}{14} \text{ chance } E_1, \frac{2^2}{14} \text{ chance } E_2, \frac{3^2}{14} \text{ chance } E_3$$

\rightarrow States with definite property values are called eigenstates of the observable operator. The value of the property is the eigenvalue.

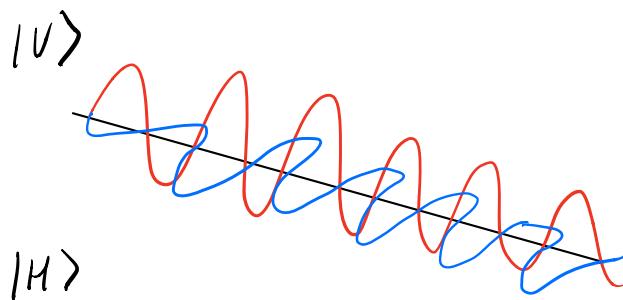
Linear algebra: eigenvectors: $A\vec{x} = \lambda\vec{x}$



Measurement examples: Linear polarisation

Light has polarization, which denotes orientation of E-field

Linear, vertical:

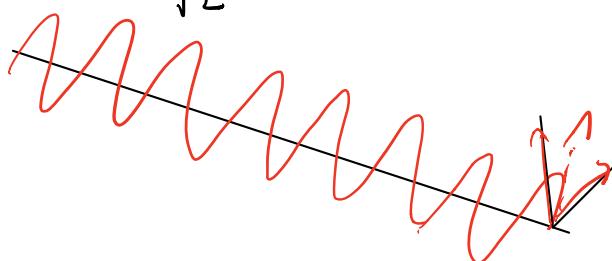


Linear, horizontal:



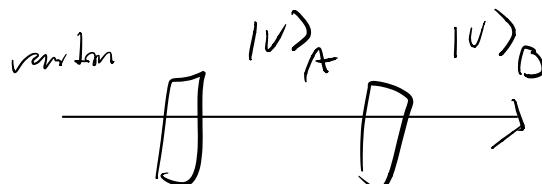
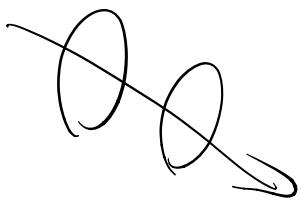
$$|D\rangle = |V\rangle + |H\rangle$$

Linear, diagonal:



Polariser means polarization & only allows one polarization to pass

2 polarisers:



rotate by θ : ant. transmitted B
 $\cos^2 \theta$

Bras & Kets:

$\psi(x)$ is position of state vector $|\psi\rangle$, which encodes everything about the particle's state

other properties: polarization

$$|\psi\rangle = A|V\rangle + B|H\rangle$$

Probability of observing particle in state $|\phi\rangle$ is $\langle \phi | \psi \rangle = \int_{-\infty}^{\infty} \phi(x) \psi(x) dx$

↑ inner product, just like vectors but complex

$$\langle A | B \rangle = 0 \text{ if } A, B \text{ orthogonal}$$

2 polarizers @ 90° :

Pr-A: ^{up-down polarization}

A: $|V\rangle$, 50% transmitted

B: $|H\rangle$, 0% transmitted since $\langle H | V \rangle = 0$

Add a 3^{rd} polarizer:

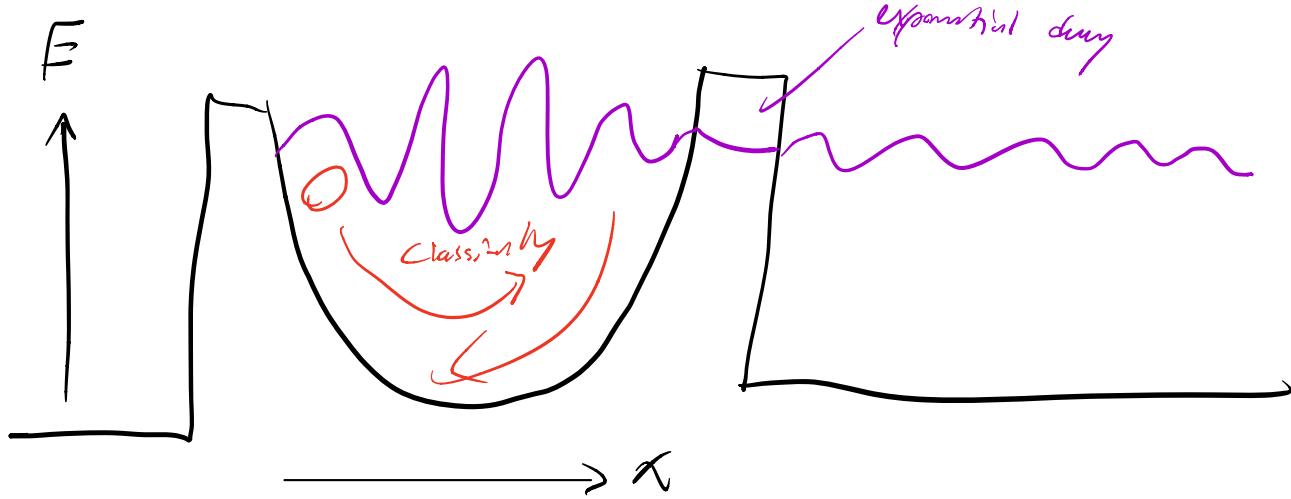


$A: |V\rangle, 50\%$ transmitted

$X: \langle X | V \rangle, \cos^2 \theta \text{ transmitted}$
 $\hookrightarrow = \frac{1}{2} \text{ if } \theta = 45^\circ$

$B: \langle B | X \rangle, \cos^2(90 - \theta) \text{ transmity}$
 $= \frac{1}{2} \text{ if } \theta = 45^\circ$

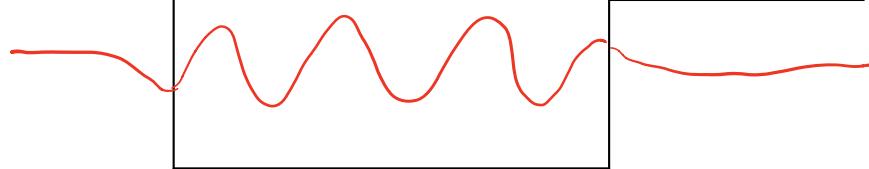
Quantum tunneling



protons, neutrons, & protons \rightarrow fusion!
no sunlight without tunneling!

Example: finite square well

$$\psi(x) = \begin{cases} \sim V_0 & -a \leq x \leq a \\ 0 & \text{otherwise} \end{cases}$$



\leftarrow particle has some energy $E < 0$

If $-a \leq x \leq a$:

$$-\frac{\hbar^2}{2m} \frac{\partial^2 \psi}{\partial x^2} = (E - V_0) \psi$$

$$\frac{\partial^2 \psi}{\partial x^2} = -\frac{\psi}{l^2}, \quad l \equiv \sqrt{\frac{2m(E + V_0)}{\hbar^2}}$$

If $|x| > a$

$$-\hbar^2 \frac{\partial^2 \psi}{\partial x^2} = E \psi$$

$$\frac{\partial^2 \psi}{\partial x^2} = -K^2 \psi$$

$$K = \frac{\sqrt{2mE}}{\hbar}$$

$$\psi \propto e^{ikx}$$

$$i: \frac{\sqrt{2m(E+V_0)}}{\hbar} > 0$$

$$K : \frac{\sqrt{2mE}}{\hbar}$$

$$\psi \propto e^{ikx} \quad \downarrow$$

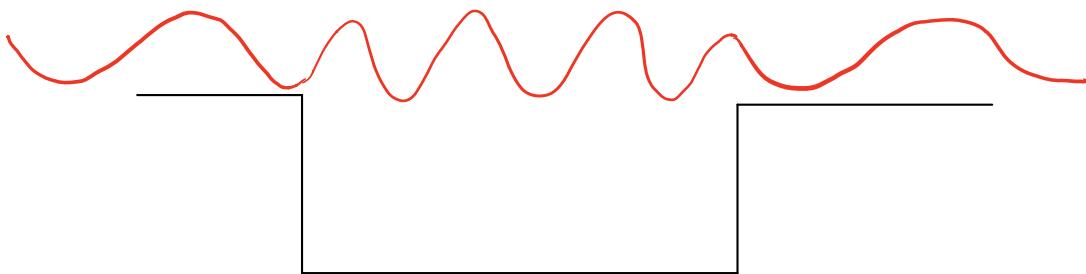
$$A \cos(kx) + B \sin(kx)$$

$$so \quad t \sim e^{ikx}$$

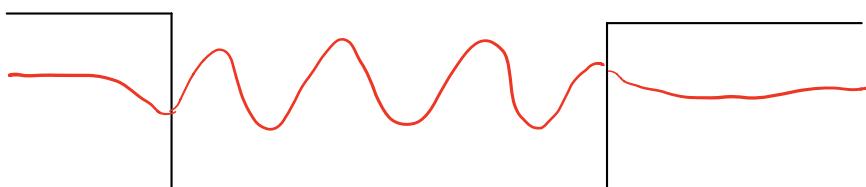
if $E > 0$: K is real



$$\psi = A \cos(Kx) + B \sin(Kx)$$



if $E < 0$: K is imaginary so $e^{ikx} \rightarrow e^{-\frac{\sqrt{-2mE}}{\hbar} x}$



The Heisenberg uncertainty principle

uncertainty principle: fundamental limit to the precision with which certain complementary variables, (e.g. $x \& P$) can be known

$$\sigma_x \sigma_p \geq \frac{\hbar}{2}$$

NOT a statement about:

- limited measuring device precision
- observation altering state of system

Consider sound waves:

sound of project tor: sine wave

sound of orchestra: complicated

BUT

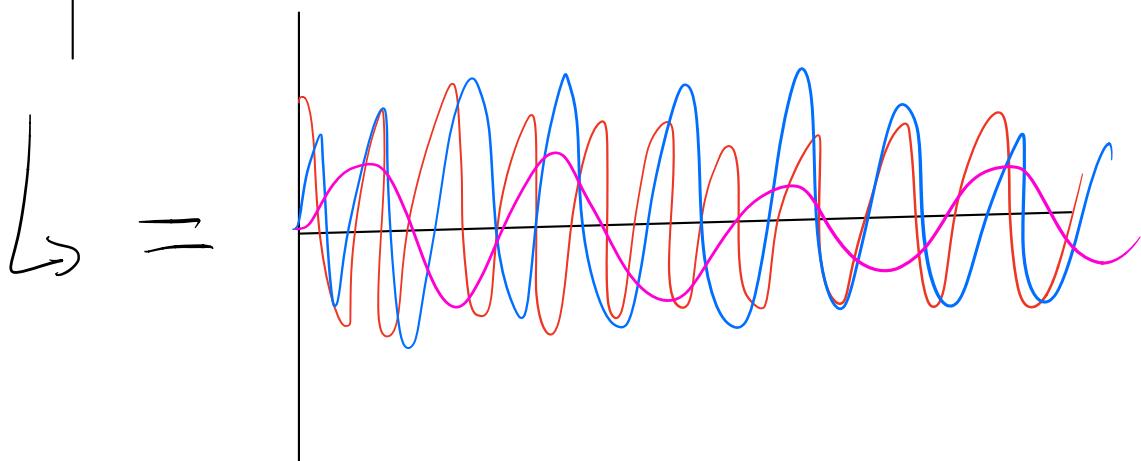
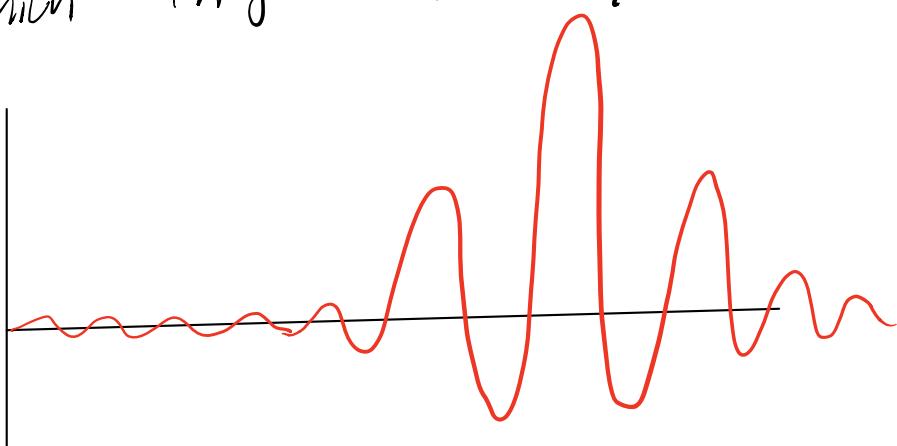
Any wave can be expressed as sum of pure sine waves

$$\text{Fourier Series: } f(x) = \frac{1}{2}a_0 + \sum_{n=1}^{\infty} a_n \cos(nx) + \sum_{n=1}^{\infty} b_n \sin(nx)$$

Sound wave can be represented in amplitude or frequency basis

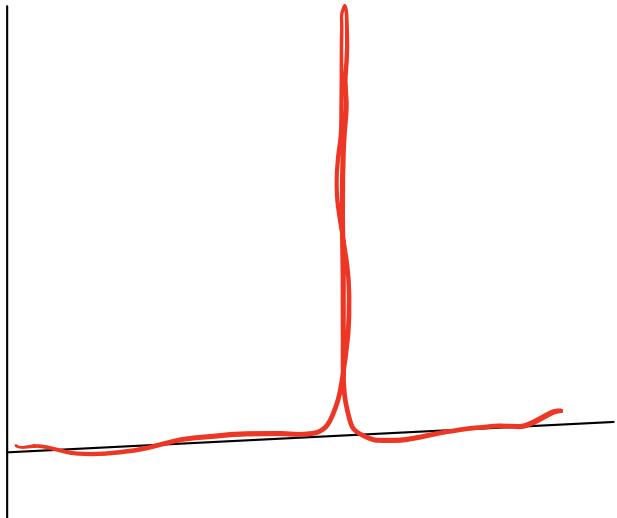
amplitude, frequency are Fairin pairs,
conjugate variables

Can make a new packet with my frequency
which is non distantly except @ location



To get steep edges around peak, need higher & higher
frequencies because high freq. always give more edges
in intensity.

To compress wave to an arbitrarily high spike,
red $\propto \rightarrow$ very short waves



How's this relate to Quantum Jiffy?

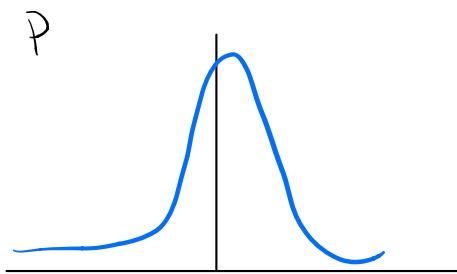
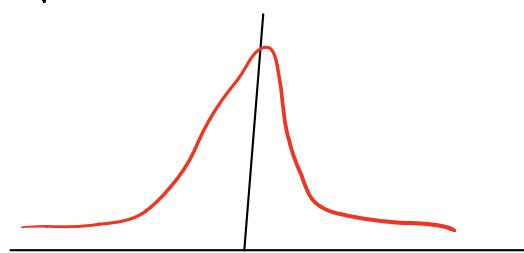
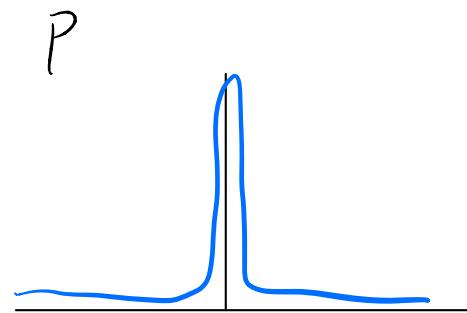
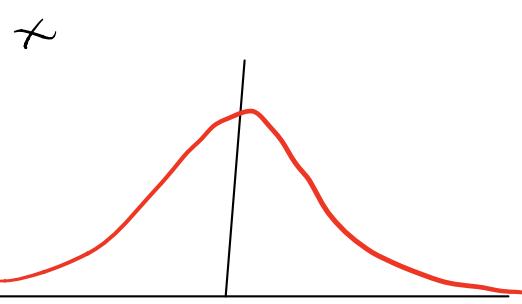
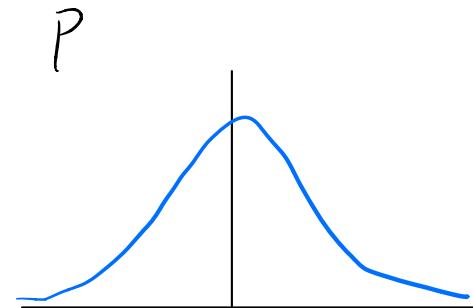
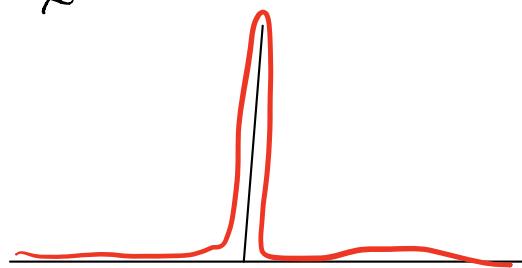
position: amplitude of wave at position x
momentum: ^{wavelength} of wavefunction $P = \frac{h}{\lambda}$

$|\psi(x)|^2 \sim$ probability of particle at position x

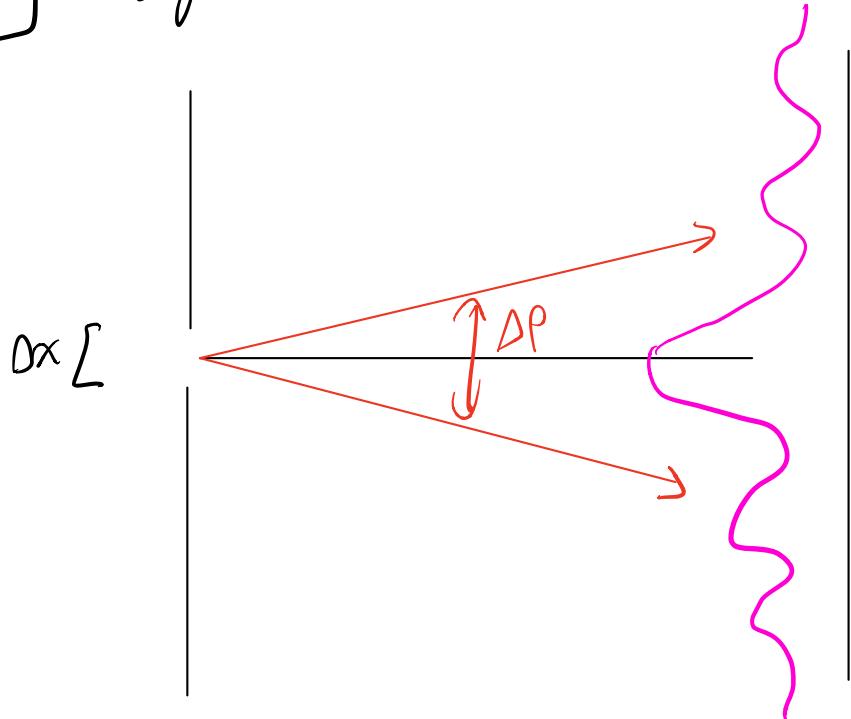
$|\psi(p)|^2 \sim$ probability of particle at momentum P

Mean position \rightarrow spike at $x \rightarrow$ wide range of momenta

Ex)



Ex] Single Slit Diffraction



Quantum entanglement

Einstein hated Copenhagen interpretation

↳ interpretation is incomplete!

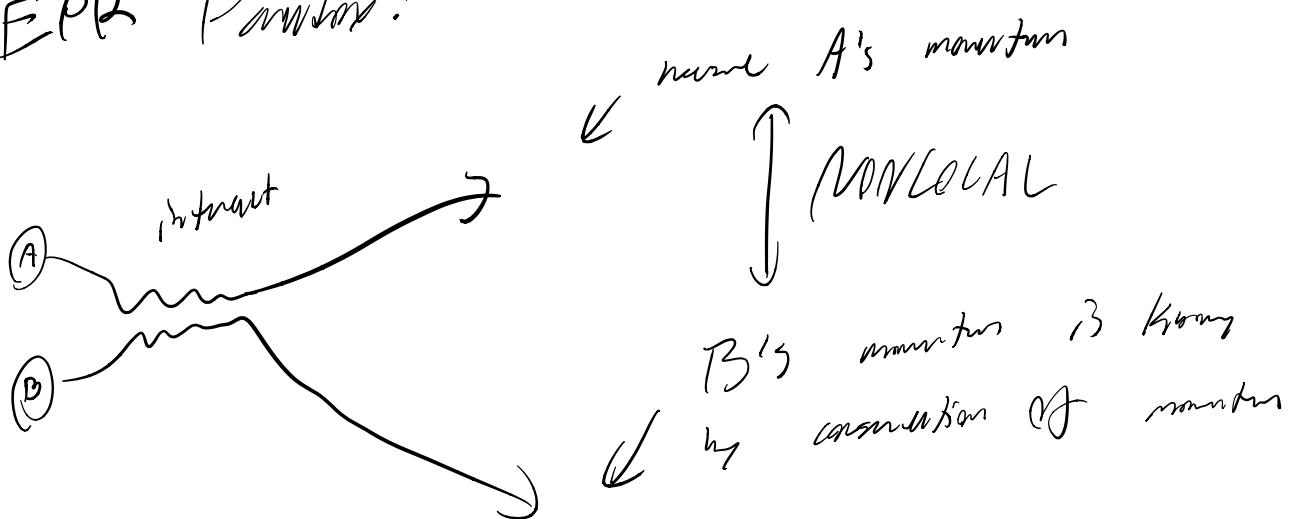
↳ nature should have "hidden variables"

To demonstrate validity of WF + Copenhagen

Einstein + Podolsky + Rosen showed that to abandon assumption of realism, must abandon assumption of locality

↳ each bit of info only acts on immediate surroundings

EPR Paradox:



ψ can only be described by a joint measurement

$$|\psi\rangle \neq |\psi_A\rangle \otimes |\psi_B\rangle$$

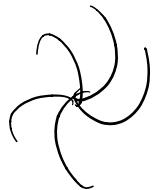
entanglement can violate locality & possibly causality!
↳ rink + morty problem

Bell's theorem:

compte $e^- + e^+ \rightarrow$ spin is connly so

$$|\psi\rangle = \frac{|\uparrow\downarrow\rangle + |\downarrow\uparrow\rangle}{\sqrt{2}}$$

(A)



(B)

no physical theory of local hidden variables
can ever reproduce all predictions of QM

non-locality allows particles to affect
each other instantaneously

BUT

Contextualism doesn't allow you to send
real information FTL

↳ correlations require (local) comparison of results

ONE by exception

Supervenientism:

There is a way to escape the inference of superluminal speeds and spooky action at a distance. But it involves absolute determinism in the universe, the complete absence of free will. Suppose the world is super-deterministic, with not just inanimate nature running on behind-the-scenes clockwork, but with our behavior, including our belief that we are free to choose to do one experiment rather than another, absolutely predetermined, including the 'decision' by the experimenter to carry out one set of measurements rather than another, the difficulty disappears. There is no need for a faster-than-light signal to tell particle A what measurement has been carried out on particle B, because the universe, including particle A, already 'knows' what that measurement, and its outcome, will be. [5]

Do we have free will?