

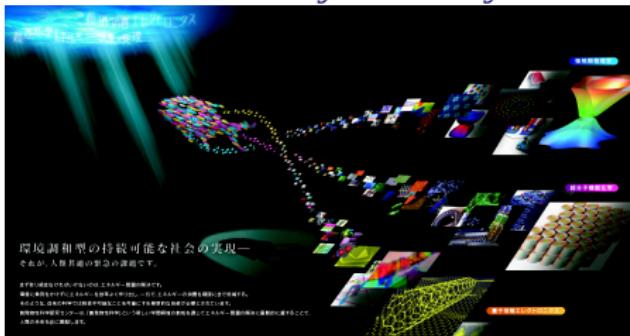
Spin motive force in the presence of spin-orbit interaction

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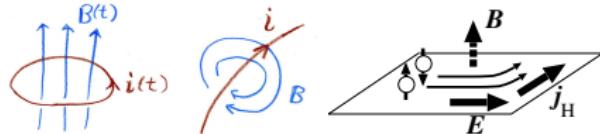
Electromagnetism for charge and spin

• Electromagnetism

$$\nabla \times E + \dot{B} = 0$$

$$\nabla \times B = \mu j + \epsilon \mu \dot{E}$$

$$j_{\text{Hall}} = \sigma_H (E \times B)$$



• Spin electromagnetism Ferromagnetic metals

Volovik'87

- Drives conduction electron spin
- *sd* exchange interaction

$$\nabla \times E_s + \dot{B}_s = 0$$

$$\nabla \times B_s = \mu_s j + \epsilon_s \mu_s \dot{E}_s$$

$$j_{\text{Hall}} = \sigma_{\text{SH}} (E \times B_s)$$



- E_s : Spin motive force, B_s : Spin Berry's phase

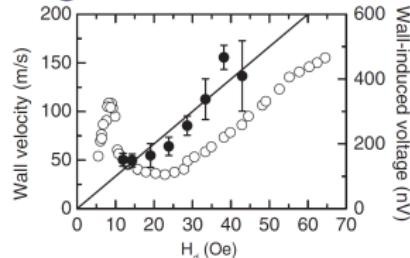
Berger'86, Stern'92, Barnes & Maekawa'07

- Real field Detectable by electric measurements

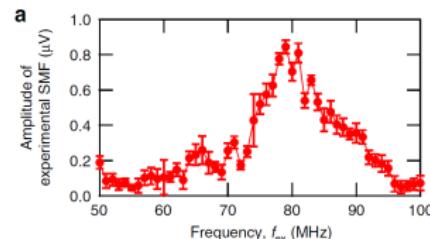
$j_s = P j$ in ferromagnetic metals $P \sim O(1)$

Spin electromagnetic field

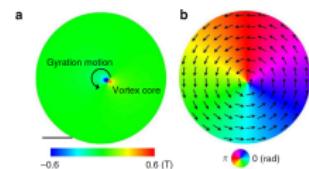
- E_s from motion of domain wall, vortex $V \sim \mu V \propto E_s \propto v_{dw}$



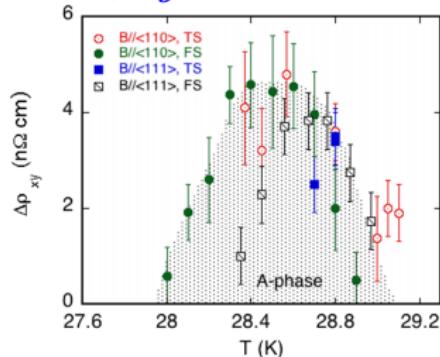
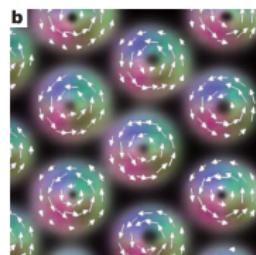
Domain wall Yang'09



Vortex Tanabe'12



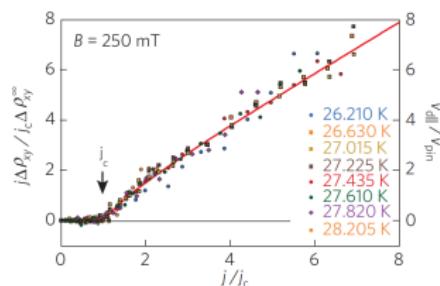
- Skyrmiion lattice $\rho_{xy} \sim 4n\Omega cm \propto B_s$



Topological Hall effect B_s

Yu'10

Neubauer'09



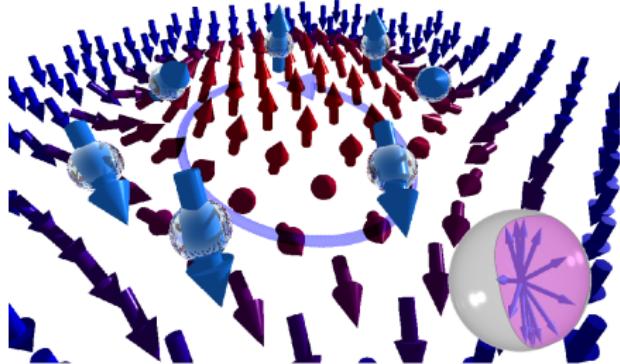
Voltage , $E_s \propto v_{sky}$

Schulz'12

Spin electromagnetic field

Volovik'87, Stern'92, Barnes&Maekawa'07

Adiabatic limit



- Electron spin rotation
⇒ Phase $e^{i\varphi}$

$$\varphi = \int_C \mathbf{dr} \cdot \mathbf{A}_s$$

$$A_s = \frac{1}{2}(1 - \cos \theta) \partial \phi$$

- Spin magnetic field

$$\varphi = \int_S dS \cdot \mathbf{B}_s$$

- Spin electric field (dynamics)

$$\dot{\varphi} = - \int_C \mathbf{dr} \cdot \mathbf{E}_s$$

- Faraday's law is satisfied

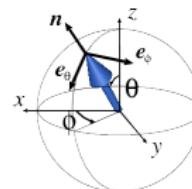
$$\nabla \times \mathbf{E}_s = - \frac{\partial \mathbf{B}_s}{\partial t}$$

Electromagnetic field coupled to spin

Topological spin electromagnetic field (Adiabatic limit)

- Effective U(1) gauge field Adiabatic limit

$$A_{\mathbf{s},\mu}^z = \frac{1}{2}(1 - \cos \theta)\partial_\mu \phi$$



- Spin electromagnetic fields

$$E_{\mathbf{s},i} = -\nabla_i A_{\mathbf{s},0}^z + \partial_t A_{\mathbf{s},i}^z = -\frac{1}{2}\mathbf{n} \cdot (\dot{\mathbf{n}} \times \nabla_i \mathbf{n})$$

$$B_{\mathbf{s},i} = (\nabla \times A_{\mathbf{s}}^z)_i = \frac{1}{4} \sum_{jk} \epsilon_{ijk} \mathbf{n} \cdot (\nabla_j \mathbf{n} \times \nabla_k \mathbf{n})$$



$\mathbf{B}_{\mathbf{s}}$: Spin Berry's phase

$\mathbf{E}_{\mathbf{s}}$: Spin motive force

$\mathbf{E}_{\mathbf{s}}$ and $\mathbf{B}_{\mathbf{s}}$ couple spin structure and electron transport

Effects of spin-orbit interaction

- Inverse (and direct) spin Hall effects
 - Spin-orbit \Rightarrow No longer in the adiabatic limit
 - Spin relaxation effects
- Spin relaxation effects on E_s
 - Spin relaxation β *Duine PRB'08*
$$\mathbf{n} \cdot (\dot{\mathbf{n}} \times \nabla_i \mathbf{n}) \Rightarrow \beta \mathbf{n} \cdot (\dot{\mathbf{n}} \times (\mathbf{n} \times \nabla_i \mathbf{n})) = \beta (\dot{\mathbf{n}} \nabla_i \mathbf{n})$$
 - Onsager relation *Saslow PRB'07, Tserkovnyak PRB'08*
- Rashba effects
 - Weak sd Transport (Diagrams) *Takeuchi> J.Phys.Soc.Jpn'12*
 - Strong sd "Chiral derivative" *Kim PRL'12*
 - Strong sd Transport (Diagrams) *GT PRB'13, Nakabayashi>'13*

Rashba-induced spin electromagnetic field

- Strong sd coupling Δ_{sd} + Rashba interaction $\alpha_{\mathbf{R}}$

$$H = \int d^3r c^\dagger \left[\left(\frac{\hbar^2}{2m} \nabla^2 + \epsilon_F \right) + \Delta_{\text{sd}} (\mathbf{n} \cdot \boldsymbol{\sigma}) - \frac{i}{2} \alpha_{\mathbf{R}} \cdot (\overleftrightarrow{\nabla} \times \boldsymbol{\sigma}) \right] c$$

- Diagram calculation

Nakabayashi & GT, arXiv:1308.0152

- Result

- Electromagnetic field description works

Linear order of $\alpha_{\mathbf{R}}$

$$\boxed{\begin{aligned} \nabla \times \mathbf{E}_{\mathbf{R}} + \dot{\mathbf{B}}_{\mathbf{R}} &= 0 \\ \nabla \times \mathbf{B}_{\mathbf{R}} - \epsilon_{\mathbf{R}} \mu_{\mathbf{R}} \dot{\mathbf{E}}_{\mathbf{R}} &= \mu_{\mathbf{R}} \mathbf{j} \\ \mathbf{j}_{\text{Hall}} &= \sigma_{\mathbf{R}} (\mathbf{E} \times \mathbf{B}_{\mathbf{R}}) \end{aligned}}$$

$$\begin{aligned} \mathbf{E}_{\mathbf{R}} &= -\frac{m}{e\hbar} (\alpha_{\mathbf{R}} \times \dot{\mathbf{n}}) \\ \mathbf{B}_{\mathbf{R}} &= \frac{m}{e\hbar} [\nabla \times (\alpha_{\mathbf{R}} \times \mathbf{n})] \end{aligned}$$

- Spin vector potential Linear order

$$\mathbf{A}_{\mathbf{R}} = -\frac{m}{e\hbar} (\alpha_{\mathbf{R}} \times \mathbf{n})$$

Kim'12, Nakabayashi & GT, arXiv

Total effective electromagnetic field

- Ferromagnetic metals
- Charge + Spin electromagnetic fields

$$E_{\text{eff}} = E + E_s,$$

$$E_s = E_{s,\text{top}} + E_R$$

$$B_{\text{eff}} = B + B_s,$$

$$B_s = B_{s,\text{top}} + B_R$$

- Ampère's law

$$\bullet \nabla \times (B + B_s) - \frac{\partial}{\partial t} (\epsilon \mu E + \epsilon_s \mu_s E_s) = \mu j$$

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$$\bullet \nabla \times (B + B_s) - \frac{\partial}{\partial t} (\epsilon \mu E + \epsilon_s \mu_s E_s) = \mu j$$

$$\boxed{\bullet \nabla \times B - \epsilon \mu \frac{\partial E}{\partial t} = \mu (j + j^{(\text{SEMF})})}$$

$$j^{(\text{SEMF})} \equiv -\frac{1}{\mu} \left(\nabla \times B_s - \epsilon_s \mu_s \frac{\partial E_s}{\partial t} \right)$$

$$= -\frac{1}{\mu} (\nabla \times B_s) + \sigma_s E_s \quad (\omega \tau \ll 1)$$

Electric current induced by spin electromagnetic field

Rashba-induced spin electromagnetic field

- Rashba-induced current

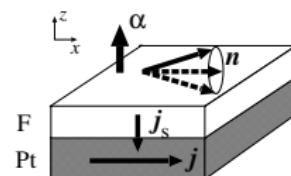
$$\begin{aligned} j^{(\text{SEMF})} &= -\frac{1}{\mu}(\nabla \times B_{\mathbf{R}}) + \sigma_s E_{\mathbf{R}} \\ &= -\frac{1}{\mu}(\nabla \times (\nabla \times (\alpha_{\mathbf{R}} \times n))) + \frac{m}{e\hbar} \sigma_s (\alpha_{\mathbf{R}} \times \dot{n}) \end{aligned}$$

- Dynamic part

- Spin pumping + Inverse spin Hall by Rashba (?)

$$j_x = \alpha_{\mathbf{R}} j_{s,z}^y, j_{s,z} \propto \dot{n}$$

- Present model is not directly applicable



- Rashba interaction
- Uniform system (Rashba and *sd* coexist)

Junction is argued by including electron diffusion

Rashba-induced spin electromagnetic field

- Where is essential spin pumping contribution ?
 - Include spin relaxation

$$\mathbf{E}_{\mathbf{R}} = \frac{m}{e\hbar} [(\boldsymbol{\alpha}_{\mathbf{R}} \times \dot{\mathbf{n}}) + \beta_{\mathbf{R}} [\boldsymbol{\alpha}_{\mathbf{R}} \times (\mathbf{n} \times \dot{\mathbf{n}})]]$$

$\beta_{\mathbf{R}}$ Spin relaxation rate

$$j^{(\text{SEMF})} = -\frac{m}{\mu e \hbar} \sigma_s \boldsymbol{\alpha}_{\mathbf{R}} \times [\dot{\mathbf{n}} + \beta_{\mathbf{R}} (\mathbf{n} \times \dot{\mathbf{n}})]$$

- Looks more like Spin pumping + Inverse spin Hall

$$j_{s,z} = g_{\uparrow\downarrow} (\mathbf{n} \times \dot{\mathbf{n}}) + g'_{\uparrow\downarrow} \dot{\mathbf{n}} \quad g_{\uparrow\downarrow} \simeq \frac{m}{\mu e \hbar} \sigma_s \beta_{\mathbf{R}}$$

- No undefined quantity

Rashba-induced spin electromagnetic field

- Where is essential spin pumping contribution ?
 - Include spin relaxation

$$E_{\mathbf{R}} = \frac{m}{e\hbar} [(\alpha_{\mathbf{R}} \times \dot{\mathbf{n}}) + \beta_{\mathbf{R}} [\alpha_{\mathbf{R}} \times (\mathbf{n} \times \dot{\mathbf{n}})]]$$

$\beta_{\mathbf{R}}$ Spin relaxation rate

$$\Downarrow$$

$$j^{(\text{SEMF})} = -\frac{m}{\mu e \hbar} \sigma_s \alpha_{\mathbf{R}} \times [\dot{\mathbf{n}} + \beta_{\mathbf{R}} (\mathbf{n} \times \dot{\mathbf{n}})]$$

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Rashba-induced spin electromagnetic field

- Rashba-induced spin electric field

$$\mathbf{E}_{\mathbf{R}} = \frac{m}{e\hbar} [(\boldsymbol{\alpha}_{\mathbf{R}} \times \dot{\mathbf{n}}) + \beta_{\mathbf{R}} [\boldsymbol{\alpha}_{\mathbf{R}} \times (\mathbf{n} \times \dot{\mathbf{n}})]]$$



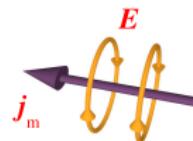
$$\nabla \times \mathbf{E}_{\mathbf{s}} + \dot{\mathbf{B}}_{\mathbf{s}} = \mathbf{j}_{\mathbf{m}} \neq 0$$

- Emergence of spin damping monopole from spin relaxation

Takeuchi, J. Phys.Soc.Japan '12

$$\mathbf{j}_{\mathbf{m}} = \beta_{\mathbf{R}} \nabla \times (\boldsymbol{\alpha}_{\mathbf{R}} \times \mathbf{N})$$

$\mathbf{N} \equiv \mathbf{n} \times \dot{\mathbf{n}}$: Spin damping vector



- Non-conservation of spin (dissipation) \Rightarrow Monopole (?)
- No topological meaning (?)
- But is physical

$$\boxed{\mathbf{n} \times \dot{\mathbf{n}}} \Rightarrow \mathbf{j}_{\mathbf{m}} \Rightarrow \nabla \times \mathbf{E}_{\mathbf{R}} = \mathbf{j}_{\mathbf{m}} \Rightarrow \boxed{\mathbf{E}}$$

Monopole current converts spin damping to electric voltage

Summary

- Strong sd + Rashba spin-orbit system
- Effective electromagnetic field description
 - Beautiful Physicists' job is to seek beauty ..?
 - No undefined parameters
 - More works necessary (junction, diffusion)
Post docs, good students

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RIKEN positions

- Postdoc/researcher positions available in Spin Physics Theory Group
gen.tatara@riken.jp
- Positions offered by RIKEN <http://www.riken.jp/en/careers/programs/>

- Postdocs, Graduate school students (PhD candidates)

For Doctoral Candidates International Program Associate	For Postdoctoral Researchers Foreign Postdoctoral Researcher Program
<p>Qualifications  IPA applicants must be (or soon to be) holders of master's degrees and be enrolled as doctoral candidates at universities that have signed (or are willing to sign) an agreement with RIKEN to participate in the Joint Graduate School Program. They should not have more than 6 years of research experience after earning their master's degree. The RIKEN researchers in charge of supervising successful applicants must also hold concurrent positions as visiting faculty members at the collaborating Japanese or overseas universities where the applicants are enrolled.</p> <p>Duration  In principle, IPAs can participate in the program for a maximum of 3 years.</p> <p>Support from RIKEN  A daily living allowance of 5,200 yen will be paid. Residence in RIKEN campus housing for the duration of the IPA's stay at RIKEN is free of charge. When campus housing is unavailable, RIKEN will pay a housing allowance of up to 70,000 yen per month for off-campus housing. (The IPA must pay his or her own utility costs.) The roundtrip transportation between RIKEN and the IPA's home country will be paid by RIKEN.</p>	<p>Qualifications  In principle, applicants should have no more than 5 years of post doctoral research experience.</p> <p>Duration  In principle, the contract period is for a maximum of 3 years.</p> <p>Remuneration  The base salary is 487,000 yen per month before taxes and the social insurance premium deduction. Commuting, housing, and relocation allowances will also be provided based on individual circumstances. An annual research budget of approximately 1 million yen is allocated to the laboratories hosting the FPRs.</p> <p>Application and selection process  Applications are publicly solicited each year and reviewed by a screening committee of scientists in and outside RIKEN working in fields relevant to applications.</p>