

# Meeting 31st March

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## Rossby waves

### Introduction

Rossby waves are inertial waves which occur in rotating fluids due to the conservation of vorticity. In the atmosphere, these take the form of large-scale meanders in the high level winds which circle the poles. In the ocean they propagate along the thermocline and are caused by wind stress anomalies at the surface.

### Physical basis

Rossby waves occur in an inviscid, barotropic fluid of constant depth, such as the atmosphere.

- Inviscid - viscous (frictional) forces are considered to be zero in the free atmosphere.
- Barotropic - surfaces of constant pressure and constant density coincide ( $\nabla\rho \times \nabla p = 0$ ). Pressure depends only on density.

In this case, the divergence of the horizontal velocity is zero, given by:

$$\nabla_h \cdot \tilde{\mathbf{u}} = \frac{\partial u}{\partial x} + \frac{\partial v}{\partial y} = 0 \quad (1)$$

The vorticity of the fluid is given by the curl of the velocity:

$$\nabla \times \tilde{\mathbf{u}}_a = \frac{\partial v}{\partial x} - \frac{\partial u}{\partial y} + f = \eta \quad (2)$$

The absolute vorticity is given by the curl of the absolute velocity ( $\tilde{\mathbf{u}}_a$ ), which is the air velocity relative to an inertial frame. For this we consider only the vertical component of the relative vorticity,  $\zeta = \frac{\partial v}{\partial x} - \frac{\partial u}{\partial y}$ , and the vorticity of the earth,  $f$  to give the absolute vorticity:  $\eta = \zeta + f$ .

For the barotropic, non-divergent (equation 1) fluid case, the absolute vorticity is conserved following the motion:

$$\frac{D}{Dt}(\zeta + f) = 0 \quad (3)$$

In the more general case, we consider a baroclinic atmosphere where density depends on both temperature and pressure, where the geostrophic wind varies with height. In this case, the Rossby wave is a potential vorticity-conserving motion which occurs due to gradients in potential vorticity. The potential vorticity is given by:

$$\mathcal{Q} = \alpha(2\mathbf{\Omega} + \nabla \times \mathbf{u}) \cdot \nabla \theta \quad (4)$$

Where  $\mathbf{\Omega}$  is the angular velocity of the earth,  $\alpha$  is the specific volume ( $1/\rho$ ),  $\mathbf{u}$  is the three dimensional vector velocity, and  $\theta$  is the potential temperature ( $\theta = T(p_0/p)^\kappa$ ). In the absence of friction and heating, the potential vorticity (ertel PV,  $\mathcal{Q}$ ) is conserved following the motion of the fluid.

### **Emergence of Rossby waves**

For an idealized case, we consider a barotropic fluid with a constant density and depth, where variation in the coriolis parameter is given by:

$$f = f_0 + \beta y \quad (5)$$

Where:

$f_0$  is the coriolis parameter at  $\phi_0$ , the latitude of the equator ( $y = 0$ ).

$\beta \equiv (df/dy)_{\phi_0} = 2\Omega \cos \phi_0/a$  is the Rossby parameter which accounts for the meridional variation of the Coriolis force caused by the spherical shape of the earth. Where  $\Omega$  is the angular velocity of the earth, and  $a$  is the radius of the earth.

$y$  is the distance from the equator.

### **Relevance**

### **Eddy mean flow**

### **Introduction**

### **Physical basis**

### **Relevance**

### **PhD update**