

Lecture 13: Forced Motion and Elastic Vehicle Mean-Axes

Textbook Sections 10.1 & 10.2

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Introduction

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 - Not purely structural/aeroelastic phenomena such as flutter or divergence

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 - Not purely structural/aeroelastic phenomena such as flutter or divergence
- Only lower frequency modes typically of interest for addition to rigid-body modes in 6-DOF EOMs
 - All real vehicles elastic & in many cases: vibration analysis only necessary to check if rigid-body assumption can be made for modeling flight vehicle
 - Larger flight vehicles typically require some elastic modeling as natural frequencies will be lower and more likely to interact with rigid-body or flight controller modes of vehicle

Introduction (continued)

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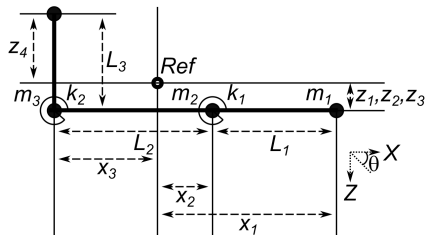
- Course treatment for elastic vehicles: Lagrangian mechanics and energy concepts stated without proof
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- Previous lectures: derived generalized coordinates to model vibration problems
 - Vibration modes orthogonal to rigid-body modes and demonstrated for some simple lumped-mass models to assist in visualizing construction of elastic-body dynamic models
- Types of coordinate frames satisfy mean-axis constraints: used for mean-axes derived in EOMs describing elastic-body flight dynamics in subsequent lectures
 - Before developing concept of mean-axes, first assume that navigation frame N inertial, i.e., “flat-Earth” model

Bi-Directional Beam Example

- Extend to multi-directional motion \rightarrow more generalized vectors and matrices
 - 1 element of mode shape/eigenvector \rightarrow 1 direction of motion

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- Extend to multi-directional motion → more generalized vectors and matrices
 - 1 element of mode shape/eigenvector → 1 direction of motion
- Demonstration: consider bi-directional example of unforced 2D truss



- Directions of motion: X & Z
- Reference point, Ref : center of mass of truss

Energies

- Kinetic energy of truss:

$$\begin{aligned}
 T &= \frac{1}{2} \left[m_1(\dot{X}_1^2 + \dot{Z}_1^2) + m_2(\dot{X}_2^2 + \dot{Z}_2^2) + m_3(\dot{X}_3^2 + \dot{Z}_3^2) + m_4(\dot{X}_4^2 + \dot{Z}_4^2) \right] \\
 &= \frac{1}{2} \begin{bmatrix} \dot{\vec{X}}^T & \dot{\vec{Z}}^T \end{bmatrix} \begin{bmatrix} M & 0 \\ 0 & M \end{bmatrix} \begin{bmatrix} \dot{\vec{X}} \\ \dot{\vec{Z}} \end{bmatrix}
 \end{aligned} \tag{1}$$

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 \end{aligned} \tag{1}$$

- Potential/strain energy of truss:

$$U = \frac{1}{2} k_1 \theta_1^2 + \frac{1}{2} k_2 \theta_2^2 = \frac{1}{2} \begin{bmatrix} \theta_1 & \theta_2 \end{bmatrix} \begin{bmatrix} k_1 & 0 \\ 0 & k_2 \end{bmatrix} \begin{bmatrix} \theta_1 \\ \theta_2 \end{bmatrix} \tag{2}$$

- 2 deflections, θ_1 & θ_2 : relative angular displacements between rods 1-2 & 2-3

Geometric Constraints

- Geometry of truss and small angles \rightarrow constraints for relative angular displacements

$$\begin{aligned}
 \theta_1 &= \frac{Z_1 - Z_2}{x_1 - x_2} - \frac{Z_2 - Z_3}{x_2 + x_3} \\
 &= \begin{bmatrix} \frac{1}{x_1 - x_2} & \left(\frac{-1}{x_1 - x_2} - \frac{1}{x_2 + x_3} \right) & \frac{1}{x_2 + x_3} & 0 \end{bmatrix} \begin{bmatrix} Z_1 \\ Z_2 \\ Z_3 \\ Z_4 \end{bmatrix} \\
 &= C_1 \vec{Z}
 \end{aligned} \tag{3}$$

Geometric Constraints (continued)

$$\begin{aligned}
 \theta_2 &= \frac{Z_3 - Z_2}{x_2 + x_3} - \frac{X_3 - X_4}{z_3 + z_4} \\
 &= \begin{bmatrix} 0 & 0 & \frac{-1}{z_3 + z_4} & \frac{1}{z_3 + z_4} \end{bmatrix} \begin{bmatrix} X_1 \\ X_2 \\ X_3 \\ X_4 \end{bmatrix} + \begin{bmatrix} 0 & \left(\frac{-1}{x_2 + x_3} - \frac{1}{x_2 + x_3} \right) & \frac{1}{x_2 + x_3} & 0 \end{bmatrix} \begin{bmatrix} Z_1 \\ Z_2 \\ Z_3 \\ Z_4 \end{bmatrix} \quad (4) \\
 &= \begin{bmatrix} C_2 & C_3 \end{bmatrix} \begin{bmatrix} \vec{X} \\ \vec{Z} \end{bmatrix}
 \end{aligned}$$

Equal Vibration Displacements

- Assume equal vibration displacements for colinear masses:

$$X_1 = X_2 = X_3 \quad \begin{bmatrix} X_1 \\ X_2 \\ X_3 \\ X_4 \end{bmatrix} = \begin{bmatrix} 1 & 0 \\ 1 & 0 \\ 1 & 0 \\ 0 & 1 \end{bmatrix} \begin{bmatrix} X_3 \\ X_4 \end{bmatrix} \quad \vec{X} = C'_x \vec{X}' \quad (5)$$

$$Z_3 = Z_4 \quad \begin{bmatrix} Z_1 \\ Z_2 \\ Z_3 \\ Z_4 \end{bmatrix} = \begin{bmatrix} 1 & 0 & 0 \\ 0 & 1 & 0 \\ 0 & 0 & 1 \\ 0 & 0 & 1 \end{bmatrix} \begin{bmatrix} Z_1 \\ Z_2 \\ Z_3 \end{bmatrix} \quad \vec{Z} = C'_z \vec{Z}' \quad (6)$$

- Degrees of freedom of truss reduced using constraints \rightarrow modified coordinate vectors:
 \vec{X}', \vec{Z}'

Kinetic Energy of Truss

$$\begin{aligned}
 T &= \frac{1}{2} \begin{bmatrix} \dot{\vec{X}}'^T & \dot{\vec{Z}}'^T \end{bmatrix} \begin{bmatrix} C'_x & 0 \\ 0 & C'_z \end{bmatrix}^T \begin{bmatrix} M & 0 \\ 0 & M \end{bmatrix} \begin{bmatrix} C'_x & 0 \\ 0 & C'_z \end{bmatrix} \begin{bmatrix} \dot{\vec{X}}' \\ \dot{\vec{Z}}' \end{bmatrix} \\
 &= \frac{1}{2} \begin{bmatrix} \dot{\vec{X}}'^T & \dot{\vec{Z}}'^T \end{bmatrix} [MM]' \begin{bmatrix} \dot{\vec{X}}' \\ \dot{\vec{Z}}' \end{bmatrix}
 \end{aligned} \tag{7}$$

Potential/Strain Energy of Truss

$$\begin{aligned}
 U &= \frac{1}{2} \begin{bmatrix} \vec{X}^T & \vec{Z}^T \end{bmatrix} \begin{bmatrix} 0 & C_1 \\ C_2 & C_3 \end{bmatrix}^T \begin{bmatrix} k_1 & 0 \\ 0 & k_2 \end{bmatrix} \begin{bmatrix} 0 & C_1 \\ C_2 & C_3 \end{bmatrix} \begin{bmatrix} \vec{X} \\ \vec{Z} \end{bmatrix} \\
 &= \frac{1}{2} \begin{bmatrix} \vec{X}^T & \vec{Z}^T \end{bmatrix} K_c \begin{bmatrix} \vec{X} \\ \vec{Z} \end{bmatrix} \\
 &= \frac{1}{2} \begin{bmatrix} \vec{X}'^T & \vec{Z}'^T \end{bmatrix} \begin{bmatrix} C'_x & 0 \\ 0 & C'_z \end{bmatrix}^T K_c \begin{bmatrix} C'_x & 0 \\ 0 & C'_z \end{bmatrix} \begin{bmatrix} \vec{X}' \\ \vec{Z}' \end{bmatrix} \\
 &= \frac{1}{2} \begin{bmatrix} \vec{X}'^T & \vec{Z}'^T \end{bmatrix} K'_c \begin{bmatrix} \vec{X}' \\ \vec{Z}' \end{bmatrix}
 \end{aligned} \tag{8}$$

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$$\frac{d}{dt} \left(\frac{\partial T}{\partial \dot{\vec{q}}} \right) - \frac{\partial T}{\partial \vec{q}} + \frac{\partial U}{\partial \vec{q}} = \vec{Q} \quad (9)$$

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- Q : generalized force

$$\vec{Q}^T = \frac{\partial \delta W}{\partial \delta \vec{q}} \quad (10)$$

- $\delta \vec{q}$: virtual displacement of generalized coordinates \vec{q}

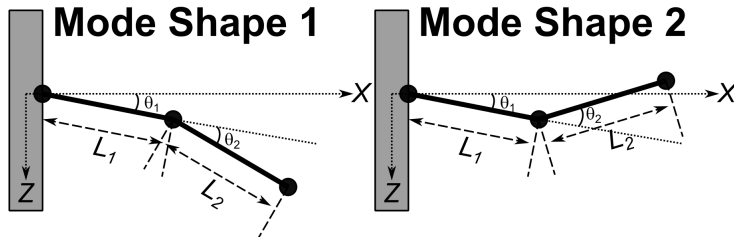
Virtual Work

- δW : **Virtual work**

$$\delta W = \sum_{i=1}^m \vec{F}_i \cdot \delta \vec{d}_i$$
$$\delta W = \begin{bmatrix} \vec{F}_1^T & \dots & \vec{F}_m^T \end{bmatrix}^T \begin{bmatrix} \delta \vec{d}_1 \\ \vdots \\ \delta \vec{d}_m \end{bmatrix} \quad (11)$$

- $\delta \vec{d}_i$: virtual physical displacement of point of application of force F_i

2-Lumped-Mass Example

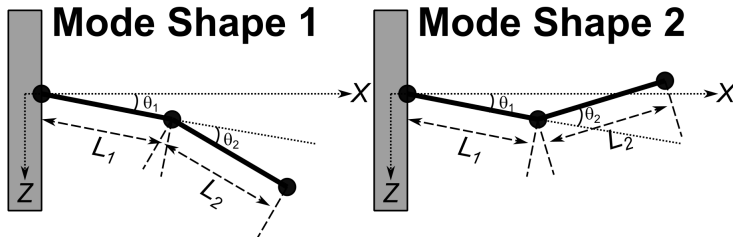


- Define vertical forces:

$$\vec{F}_1 = F_1 \hat{k} \quad \vec{F}_2 = F_1 \hat{k} \quad (12)$$

- \hat{k} : unit vector for Z direction

2-Lumped-Mass Example



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$$\vec{F}_1 = F_1 \hat{k} \quad \vec{F}_2 = F_1 \hat{k} \quad (12)$$

- \hat{k} : unit vector for Z direction
- Virtual physical displacements at points of application:

$$\begin{aligned} \delta \vec{d}_1 &= \delta Z_1 \hat{k} \\ \delta \vec{d}_2 &= \delta Z_2 \hat{k} \end{aligned} \quad (13)$$

2-Lumped-Mass Example (continued)

$$\delta W = F_1 \delta Z_1 + F_2 \delta Z_2$$

$$\delta W = \begin{bmatrix} F_1 & F_2 \end{bmatrix} \begin{bmatrix} \delta Z_1 \\ \delta Z_2 \end{bmatrix} \quad (14)$$

2-Lumped-Mass Example (continued)

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$$\delta W = \begin{bmatrix} F_1 & F_2 \end{bmatrix} \begin{bmatrix} \delta Z_1 \\ \delta Z_2 \end{bmatrix} \quad (14)$$

$$\begin{bmatrix} \delta Z_1 \\ \delta Z_2 \end{bmatrix} = \begin{bmatrix} L_1 & 0 \\ L_1 + L_2 & L_2 \end{bmatrix} \begin{bmatrix} \delta \theta_1 \\ \delta \theta_2 \end{bmatrix} \quad (15)$$

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- Assuming unforced (i.e. “free”) vibration problem solved in terms of θ_1 and θ_2
- Elements of free-vibration mode shapes correspond to angular displacements

2-Lumped-Mass Example (continued)

- Virtual displacements $\delta\theta_1$ and $\delta\theta_2$ expressed in terms of mode shapes and two vibration modal coordinates, η_1 and η_2 :

$$\begin{bmatrix} \delta\theta_1 \\ \delta\theta_2 \end{bmatrix} = [\vec{\nu}_1 \quad \vec{\nu}_2] \begin{bmatrix} \delta\eta_1 \\ \delta\eta_2 \end{bmatrix} \quad (16)$$
$$\delta\vec{\theta} = \Psi \delta\vec{\eta}$$

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- Substitution: virtual work

$$\delta W = [F_1 \quad F_2] \begin{bmatrix} L_1 & 0 \\ L_1 + L_2 & L_2 \end{bmatrix} \Psi \delta\vec{\eta}$$
$$\delta W = \vec{\mathcal{F}}^T \Psi \delta\vec{\eta} \quad (17)$$

- $\vec{\mathcal{F}}$: vector of applied forces relative to virtual displacements

2-Lumped-Mass Example (continued)

- Kinetic energy (modal coordinates):

$$T = \frac{1}{2} \dot{\theta}^T M \dot{\theta} = \frac{1}{2} \dot{\eta}^T \Psi^T M \Psi \dot{\eta} = \frac{1}{2} \dot{\eta}^T \mathcal{M} \dot{\eta} \quad (18)$$

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- Lagrange's equation with external forces applied:

$$\mathcal{M} \ddot{\vec{\eta}} + \mathcal{K} \vec{\eta} = Q = \Psi^T \vec{\mathcal{F}} = \vec{\nu}_{vib}^T \vec{\mathcal{F}} \quad (20)$$

Unrestrained 3-Lumped-Mass Extension

- Similar expression in form:

$$M_{tot} \ddot{\vec{Z}}_{Ref} = F_1 + F_2 + F_3 = \vec{1}^T \begin{bmatrix} F_1 \\ F_2 \\ F_3 \end{bmatrix}$$

$$I_G \ddot{\theta}_{Ref} = F_1 x_1 + F_2 x_2 - F_3 x_3 = \vec{x}^T \begin{bmatrix} F_1 \\ F_2 \\ F_3 \end{bmatrix} \quad (21)$$

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- For forced response of unrestrained beam's EOMs
 - Fundamental for elastic-body flight dynamics
- *Ref*: center of mass of lumped-mass system
 - Easily extended to n -lumped-mass systems

Kinetic Energy of Vehicle

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- \vec{v}_N : inertial velocity of mass element of vehicle

$$\vec{v}_N = \vec{v}_{B/N} + \vec{v}_B + \omega_{B/N} \times \vec{x}_B \quad (23)$$

- \vec{v}_B : velocity of mass element in body-fixed frame
- $\vec{v}_{B/N}$: velocity of body-fixed frame relative to navigation frame
- $\vec{\omega}_{B/N}$: angular velocity of body-fixed frame relative to navigation frame
 - Note: velocities can be represented in body-fixed or navigation frames
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$$\vec{x}_N = \vec{x}_B + \vec{x}_{B/N} \quad (24)$$

- $\vec{x}_{B/N}$: position of origin of body-fixed frame relative to navigation frame

Kinetic Energy of Vehicle (continued)

- By substitution:

$$\begin{aligned}
 T = \frac{1}{2} \int_{Vol} [& \vec{v}_{B/N,N} \cdot \vec{v}_{B/N,N} \\
 & + 2 \vec{v}_{B/N,N} \cdot \vec{v}_B \\
 & + 2 (\vec{v}_{B/N,N} + \vec{v}_B) \cdot (\vec{\omega}_{B/N} \times \vec{x}_B) \\
 & + \vec{v}_B \cdot \vec{v}_B \\
 & + (\vec{\omega}_{B/N} \times \vec{x}_B) \cdot (\vec{\omega}_{B/N} \times \vec{x}_B)] \rho_V dV
 \end{aligned} \tag{25}$$

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- Gravitational potential energy:

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$$U_g = - \int_{Vol} \vec{g} \cdot (\vec{x}_B + \vec{x}_{B/N}) \rho_V dV$$
(26)

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- \vec{g} : acceleration due to gravity
- Elastic strain energy: energy stored in elastic structure due to deformation resulting from applied force
 - Negative of work done on structure by applied force
 - Work: force acting over displacement \vec{d}_e

Potential Energy (continued)

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- D'Alembert's principle:
 - Express force on mass element in terms of mass of element and acceleration

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- D'Alembert's principle:
 - Express force on mass element in terms of mass of element and acceleration
- Elastic strain energy:

$$U_e = -\frac{1}{2} \int_{Vol} \ddot{\vec{d}}_{e,B} \cdot \vec{d}_{e,B} \rho_V dV \quad (29)$$

Mean-Axes Body Frame

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 - \rightarrow *additional* requirements for body-fixed frame to exhibit decoupled dynamic modes
- **Mean-axes constraints:** define coordinate axes about which relative translational and angular momenta (about center of mass) due to elastic vibrations are zero

$$\int_{Vol} \dot{\vec{d}}_{e,B} \rho_V dV = \int_{Vol} \vec{x}_B \times \dot{\vec{d}}_{e,B} \rho_V dV = \vec{0} \quad (30)$$

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- “Special” **mean-axes body-fixed frame:** always exist for elastic-body

Practical Mean-Axes

- By substitution in previous section:

$$\int_{Vol} \frac{d}{dt} \left(\vec{x}_{r-b} + \vec{d}_{e,B} \right) \rho_V dV = \vec{0} \quad (31)$$

$$\int_{Vol} \vec{x}_{r-b} \times \dot{\vec{d}}_{e,B} \rho_V dV + \int_{Vol} \vec{d}_{e,B} \times \vec{v}_B \rho_V dV = \vec{0} \quad (32)$$

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- If elastic displacement, $\vec{d}_{e,B}$, sufficiently small, i.e., only linear effects considered

Practical Mean-Axes

- By substitution in previous section:

$$\int_{Vol} \frac{d}{dt} \left(\vec{x}_{r-b} + \vec{d}_{e,B} \right) \rho_V dV = \vec{0} \quad (31)$$

$$\int_{Vol} \vec{x}_{r-b} \times \dot{\vec{d}}_{e,B} \rho_V dV + \int_{Vol} \vec{d}_{e,B} \times \vec{v}_B \rho_V dV = \vec{0} \quad (32)$$

- If elastic displacement, $\vec{d}_{e,B}$, sufficiently small, i.e., only linear effects considered
- neglect moment from elastic displacements: **practical mean-axes constraints**

$$\int_{Vol} \dot{\vec{d}}_{e,B} \rho_V dV = \int_{Vol} \vec{x}_{r-b} \times \dot{\vec{d}}_{e,B} \rho_V dV = \vec{0} \quad (33)$$

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- Analogous to modal orthogonality constraints for lumped-mass systems

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- Elastic vibrations:

$$\vec{d}_{e,B} = \sum_{i=1}^n \vec{v}_i(\vec{x}) \eta_i(t) \quad (34)$$

- $\eta_i(t)$: generalized coordinate associated with i -th vibration mode
- $\vec{v}_i(\vec{x})$: vector with components defined in body-fixed frame
- Each component function of \vec{x} location on *undeformed* structure

Practical Mean-Axes Constraints

$$\int_{Vol} \dot{\vec{d}}_{e,B} \rho_V dV = \sum_{i=1}^n \dot{\eta}_i(t) \left(\int_{Vol} \vec{v}_i(\vec{x}) \rho_V dV \right) = 0 \quad (35)$$

$$\int_{Vol} \vec{x}_{r-b} \times \dot{\vec{d}}_{e,B} \rho_V dV = \sum_{i=1}^n \left(\int_{Vol} \vec{x}_{r-b} \times \vec{v}_i(\vec{x}) \rho_V dV \right) = 0 \quad (36)$$

- Satisfied as integrals inside parentheses correspond to momenta conservation requirements
- Selected vibration modes, by design, orthogonal to rigid-body body translation and rotation modes (w.r.t. mass distribution)

Constraints to Energies

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- m : total mass of vehicle

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- m : total mass of vehicle
- 2nd term of kinetic energy = 0:

$$\int_{Vol} \vec{v}_{B/N,N} \cdot \dot{\vec{d}}_{e,B} \rho_V dV = \vec{v}_{B/N,N} \cdot \int_{Vol} \dot{\vec{d}}_{e,B} \rho_V dV = 0\tag{38}$$

3rd Term of Kinetic Energy

- Recall: requirement for origin of mean-axis body-fixed frame = center of mass:

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$$\begin{aligned} & \int_{Vol} \left(\vec{v}_{B/N,N} + \dot{\vec{d}}_{e,B} \right) \cdot (\vec{\omega}_{B/N} \times \vec{x}_B) \rho_V dV \\ &= \vec{v}_{B/N,N} \cdot \left(\vec{\omega}_{B/N} \times \int_{Vol} \vec{x}_B \rho_V dV \right) = 0 \end{aligned} \quad (41)$$

4th Term of Kinetic Energy

- Use mode shapes and generalized coordinate summations for displacement *velocity*:

$$\begin{aligned}\int_{Vol} \vec{v}_B \cdot \vec{v}_B \rho_V dV &= \frac{1}{2} \int_{Vol} \left(\sum_{i=1}^n \vec{v}_i(\vec{x}) \dot{\eta}_i \cdot \sum_{i=1}^n \vec{v}_i(\vec{x}) \dot{\eta}_i \right) \rho_V dV \\ &= \int_{Vol} \left(\sum_{i=1}^n \vec{v}_i(\vec{x}) \cdot \vec{v}_i(\vec{x}) \dot{\eta}_i^2 \right) \rho_V dV \\ &= \sum_{i=1}^n \mathcal{M}_i \dot{\eta}_i^2\end{aligned}\tag{42}$$

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 &= \int_{Vol} \left(\sum_{i=1}^n \vec{v}_i(\vec{x}) \cdot \vec{v}_i(\vec{x}) \dot{\eta}_i^2 \right) \rho_V dV \\
 &= \sum_{i=1}^n \mathcal{M}_i \dot{\eta}_i^2
 \end{aligned} \tag{42}$$

- Simplification due to mutual orthogonality of vibration modes:

$$\int_{Vol} \vec{v}_i \cdot \vec{v}_j \rho_V dV = \begin{cases} 0 & i \neq j \\ \mathcal{M}_i & i = j \end{cases} \tag{43}$$

5th Term and Final of Kinetic Energy

- In terms of inertia matrix:

$$\int_{Vol} (\vec{\omega}_{B/N} \times \vec{x}_B) \cdot (\vec{\omega}_{B/N} \times \vec{x}_B) \rho_V dV = \vec{\omega}_{B/N}^T I_G \vec{\omega}_{B/N} \quad (44)$$

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- Kinetic energy:

$$T = \frac{1}{2} m \vec{v}_{B/N,N}^T \vec{v}_{B/N,N} + \frac{1}{2} \vec{\omega}_{B/N}^T I_G \vec{\omega}_{B/N} + \frac{1}{2} \sum_{i=1}^n \mathcal{M}_i \dot{\eta}_i^2 \quad (45)$$

Gravitational Potential Energy

- Applying *mean-axis constraints*:

$$\begin{aligned}
 U_g &= - \int_{Vol} \vec{g} \cdot (\vec{x}_B + \vec{x}_{B/N}) \rho_V dV \\
 U_g &= - \vec{g} \cdot \int_{Vol} \vec{x}_B \rho_V dV - \vec{g} \cdot \vec{x}_{B/N} \int_{Vol} \rho_V dV \\
 U_g &= -m \vec{g}^T \vec{x}_{B/N}
 \end{aligned} \tag{46}$$

Elastic Strain Energy

- Applying *mean-axis constraints*:

$$U_e = -\frac{1}{2} \int_{Vol} \sum_{i=1}^n \vec{v}_i(\vec{X}) \ddot{\eta}_i(t) \cdot \sum_{i=1}^n \vec{v}_i(\vec{X}) \eta_i(t) \rho_V dV \quad (47)$$

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- Recall sinusoidal solution for modal coordinates:

$$\eta_i(t) = A_i \cos(\omega_i t + \Gamma_i) \quad (49)$$

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 - Appropriate for rigid-body bodies: relative position of point on body represented with location of center of mass *and* attitude
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 - Appropriate for rigid-body bodies: relative position of point on body represented with location of center of mass *and* attitude
 - Not true for elastic vehicles
- B/N subscript dropped to easily compare elastic vehicle dynamics with rigid-body vehicle dynamics: selected generalized coordinates

$$\vec{x}_{B/N,N} = [x_N \quad y_N \quad z_N]^T \quad (52)$$

- Recall $z_N = -h$ for flat-Earth

Generalized Coordinates (continued)

- Velocity of body-fixed frame's origin (in navigation frame coordinates)

$$\vec{V}_{B/N,N} = [\dot{x}_N \quad \dot{y}_N \quad \dot{z}_N]^T \quad (53)$$

Generalized Coordinates (continued)

- Velocity of body-fixed frame's origin (in navigation frame coordinates)

$$\vec{v}_{B/N,N} = [\dot{x}_N \quad \dot{y}_N \quad \dot{z}_N]^T \quad (53)$$

- (in body-fixed frame coordinates)

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- Select angular velocity of body-fixed frame relative to navigation frame:

$$\vec{\omega}_{B/N} = [p \quad q \quad r]^T \quad (55)$$

Generalized Coordinates (continued)

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- Either body or navigation frame coordinates
- Generalized coordinates use body-fixed frame coordinates

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- Select the 3-2-1 Euler angles ψ , θ , ϕ : attitude/orientation of mean-axes body-fixed frame w.r.t. navigation frame

Generalized Coordinates (continued)

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- Recall: Euler angles related to angular velocity

$$\begin{bmatrix} \dot{p} \\ \dot{q} \\ \dot{r} \end{bmatrix} = \begin{bmatrix} 1 & 0 & -\sin \theta \\ 0 & \cos \phi & -\sin \phi \cos \theta \\ 0 & -\sin \phi & \cos \phi \cos \theta \end{bmatrix} \begin{bmatrix} \dot{\phi} \\ \dot{\theta} \\ \dot{\psi} \end{bmatrix} \quad (56)$$

$$\begin{bmatrix} \dot{\phi} \\ \dot{\theta} \\ \dot{\psi} \end{bmatrix} = \begin{bmatrix} 1 & \sin \phi \tan \theta & \cos \phi \tan \theta \\ 0 & \cos \phi & -\sin \phi \\ 0 & \sin \phi \sec \theta & \cos \phi \sec \theta \end{bmatrix} \begin{bmatrix} \dot{p} \\ \dot{q} \\ \dot{r} \end{bmatrix} \quad (57)$$

Generalized Coordinates (continued)

- Selected generalized coordinates:

$$\vec{q} = [x_N \quad y_N \quad z_N \quad \phi \quad \theta \quad \psi \quad \eta_i, i = 1, 2, \dots]^T \quad (58)$$

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- Selected generalized coordinates:

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- Same for rigid-body vehicle dynamics *with* addition of vibration modal coordinates corresponding to mutually orthogonal vibration mode shapes

Rewritten Energies

- Kinetic energy of elastic vehicle:

$$T = \frac{1}{2}m \begin{bmatrix} \dot{x}_N & \dot{y}_N & \dot{z}_N \end{bmatrix} \begin{bmatrix} \dot{x}_N \\ \dot{y}_N \\ \dot{z}_N \end{bmatrix} + \frac{1}{2} \begin{bmatrix} p & q & r \end{bmatrix} I_G \begin{bmatrix} p \\ q \\ r \end{bmatrix} + \frac{1}{2} \sum_{i=1}^n \mathcal{M}_i \dot{\eta}_i^2 \quad (59)$$

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- Gravitational potential energy:

$$U_g = -mgz_N = mgh \quad (60)$$

- Elastic strain energy (same):

$$U_e = -\frac{1}{2} \sum_{i=1}^n \omega_i^2 \eta_i^2(t) \mathcal{M}_i$$

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- Elastic flight vehicle dynamics: decoupling translation, rotation, and vibration modes possible through
 - Body-fixed frame origin: center of mass
 - Body-fixed frame axes: mean-axes
 - Separate terms for kinetic energy
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 - Position of center of mass
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