#### **Master Thesis Defense**

# Thermophoresis in Liquids and its Connection to Equilibrium Quantities

Benjamin F. Maier

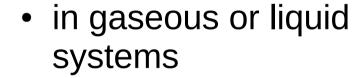
23 June 2014





# Thermophoresis

- directed movement of particles
- Induced by heat



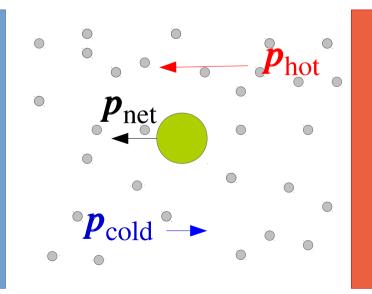
here: gas

solute moves to cold



http://bit.ly/1qy7Gs7, 06/20/14

 $T_c$ 



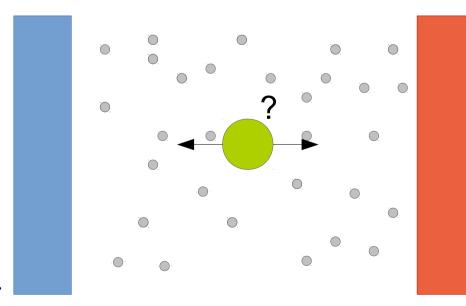
 $T_{l}$ 

# Thermophoresis - Liquids

- Different picture for liquids
- Hot region sometimes preferred
- No coherent prediction possible



http://i.imgur.com/FKLjmUV.jpg, 06/20/14

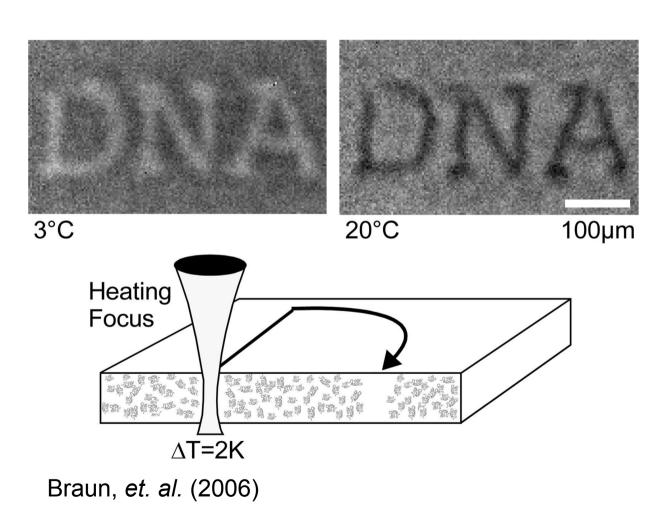


 $T_c$ 

 $T_c < T_h$ 

# Thermophoresis – A Powerful Tool!

#### Powerful addition to electrophoresis



**DNA** migration

Problem: Unpredictable!

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# Thermophoresis in Formulae

#### phenomenological flux

$$\mathbf{j}(x) = -D\nabla\rho(x) - \rho D_{\mathrm{T}}\nabla T(x) = 0$$
 diffusion thermal flux

#### Soret equilibrium

$$0 = -\nabla \rho - \rho S_{\mathrm{T}} \nabla T$$

T Temperature

ρ Density

D Diffusion coefficient

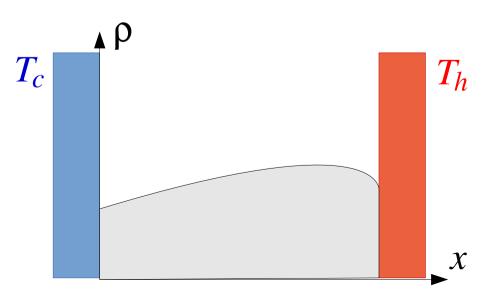
D<sub>T</sub> Thermal diffusive mobility

Soret coefficient

$$S_{\mathrm{T}} = \frac{D_{\mathrm{T}}}{D}$$

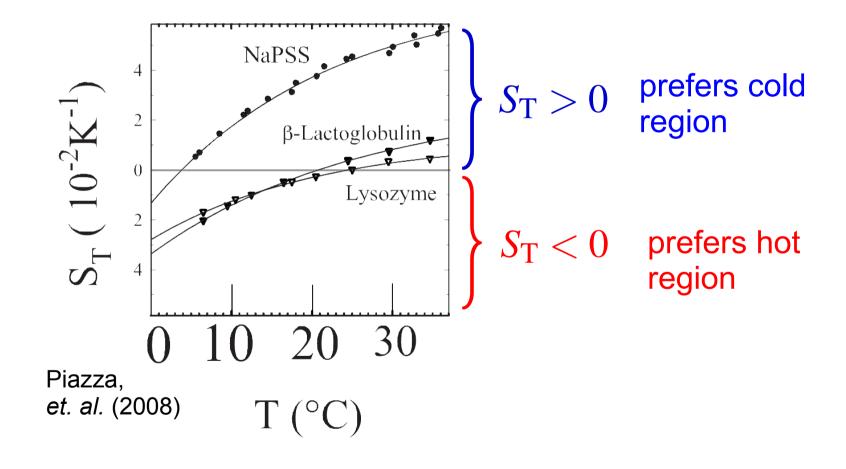
measured from Soret equilibrium density

$$S_{\rm T} = -\frac{\nabla \rho}{\rho \, \nabla T}$$



#### **Soret Coefficient**

• sign of  $S_{\rm T}$  dictates direction of movement



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# **Local Equilibrium**

- assumptions
  - length scale of solute-solvent interaction
  - small heat flux  $m{j}_O \simeq 0$
- solvation enthalpy Würger (2014) solvation free energy force on solute  $\frac{F}{T} \propto \nabla \left(\frac{\Delta G}{T}\right) + \text{id.gas.} \Rightarrow S_T^{(H)} = -\frac{\beta \Delta H}{T} + \frac{1}{T}$ 
  - Braun/Dhont (2007)

# Which one is true?

$$\frac{F}{T} \propto \frac{\nabla (\Delta G)}{T} + \text{id.gas.} \implies S_{\text{T}}^{(S)} = -\beta \Delta S + \frac{1}{T}$$
solvation entropy

solvation entropy

# **Open Problems**

- local equilibrium applicable?
- if yes, is driving force the entropy or the enthalpy?
- connection of T(r) and D(r)
- sign change of  $S_{\mathrm{T}}$  reproducible?

# Theory – Brownian Motion

Overdamped Langevin equation (an SDE)

$$\mathrm{d} \boldsymbol{r}_i = -\begin{bmatrix} \frac{\nabla_i V_{\mathrm{ext}}(\boldsymbol{r}_i)}{\gamma} + \frac{1}{\gamma} \nabla_i \sum_{j=1}^N V(\boldsymbol{r}_i - \boldsymbol{r}_j) \end{bmatrix} \mathrm{d} t + \sqrt{2D(\boldsymbol{r}_i)} \; \mathrm{d} \boldsymbol{B}_{t,i}$$
 position change 
$$\underbrace{ \begin{cases} \mathrm{external\ force} \\ \gamma - \mathrm{friction} \end{cases}}_{\text{external\ force}} \underbrace{ \begin{cases} \mathrm{ocal\ diffusion\ coefficient\ (not\ necessarily\ Einstein's\ relation) \end{cases}}_{\text{external\ force}}$$

Gaussian

interpretation: Ito or Stratonovich?

$$\begin{split} \mathrm{d} \boldsymbol{r}_i &= -\left[\frac{\nabla_i V_{\mathrm{ext}}(\boldsymbol{r}_i)}{\gamma} + \frac{1}{\gamma} \nabla_i \sum_{j=1}^N V(\boldsymbol{r}_i - \boldsymbol{r}_j)\right] \; \mathrm{d}t + \sqrt{2D(\boldsymbol{r}_i)} \; \mathrm{d}\boldsymbol{B}_{t,i} + \\ & \frac{+\alpha \nabla_i D(\boldsymbol{r}_i) \; \mathrm{d}t}{\mathbf{E}_{t,i}} \\ & \frac{1}{2} \end{split} \qquad \qquad \text{Ito} \quad \alpha = 0 \\ \text{Stratonovich} \quad \alpha = \frac{1}{2} \end{split}$$

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# Ideal Gas Soret Equilibrium Density

ideal gas without external potential

$$\rho(\mathbf{r}) = \frac{\mathscr{N}}{[D(\mathbf{r})]^{1-\alpha}} = \frac{\mathscr{N}}{[T(\mathbf{r})]^{1-\alpha}}$$

ideal gas equation of state

$$\rho k_{\rm B}T = {\rm const}$$

- interpretation of  $D(\mathbf{r})$ 
  - Einstein's relation follows for *spatially invariant* diffusion coefficient

$$\nabla D = \nabla V_{\rm ext} \left( D\beta - \gamma^{-1} \right) = 0 \qquad \Rightarrow \qquad D = k_{\rm B} T / \gamma$$

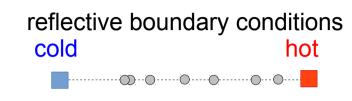
- may be wrong (Astumian, 2008)
- Will assume

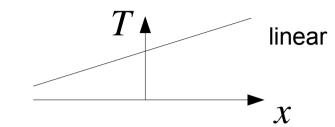
$$D(\mathbf{r}) = k_{\rm B}T(\mathbf{r})/\gamma$$

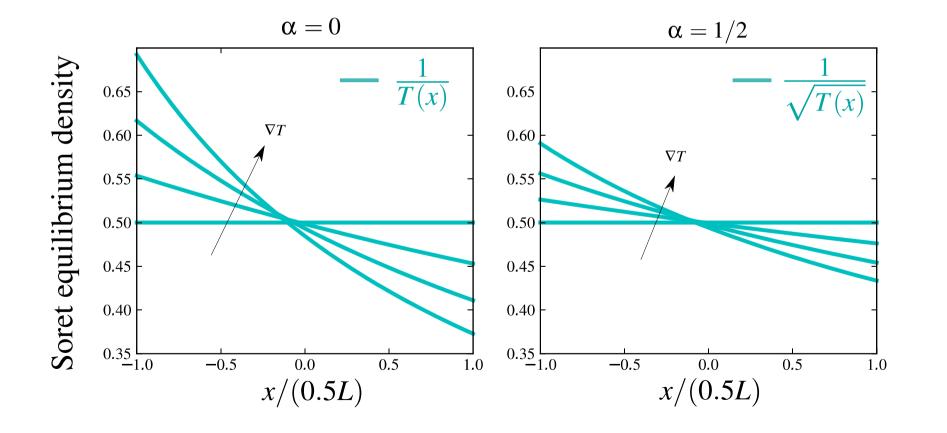
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#### **BD Simulation for 1D Ideal Gas**

$$\rho(x) = \frac{\mathscr{N}}{[T(x)]^{1-\alpha}}$$

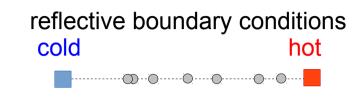


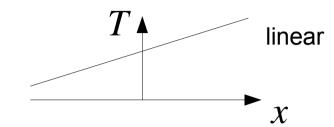


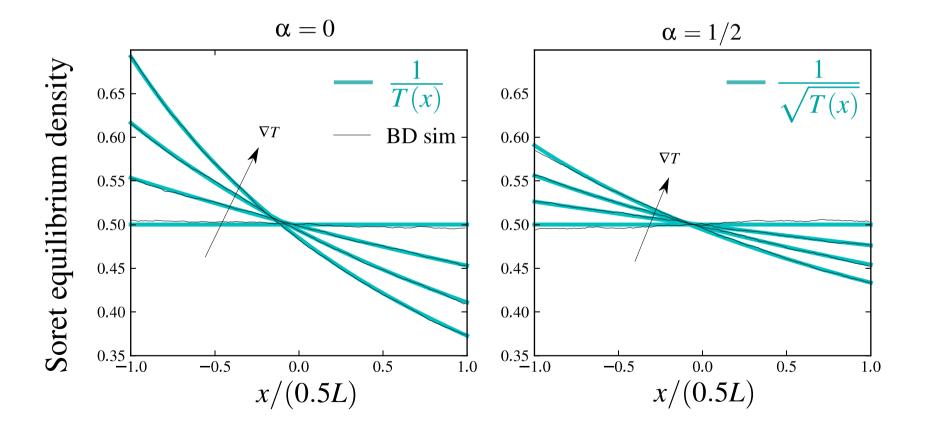


## **BD Simulation for 1D Ideal Gas**

$$\rho(x) = \frac{\mathscr{N}}{[T(x)]^{1-\alpha}}$$



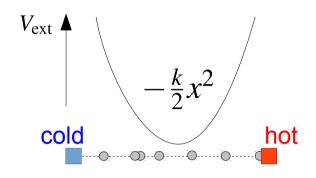




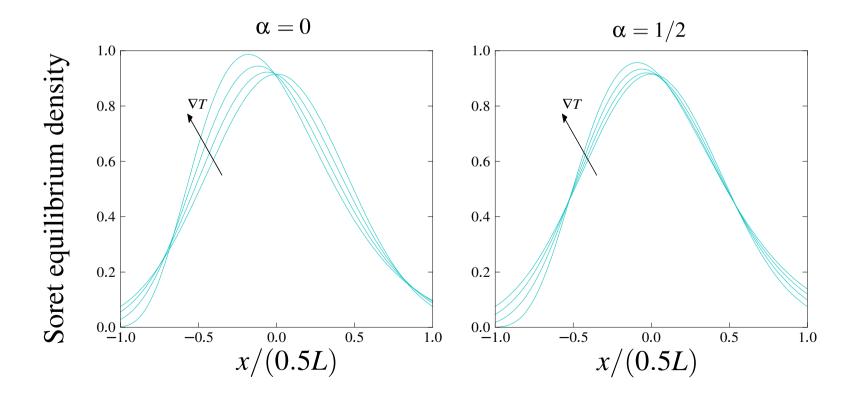
$$k_{\rm B} = 1; \ \gamma = 1$$

#### 1D Ideal Gas with External Potential

$$\rho(\mathbf{r}) = \frac{\mathcal{N}}{[T(\mathbf{r})]^{1-\alpha}} \exp\left(-\int_{\mathbf{r}_0}^{\mathbf{r}} d\tilde{\mathbf{r}} \, \frac{\nabla_{\tilde{\mathbf{r}}} V_{\text{ext}}(\tilde{\mathbf{r}})}{k_{\text{B}} T(\tilde{\mathbf{r}})}\right)$$

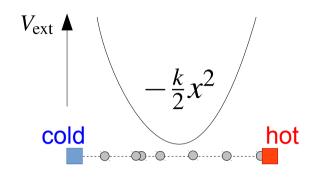


naïve Boltzmann expectation  $ho\left(m{r}
ight)=\mathscr{N}e^{-\beta\left(m{r}
ight)V_{\mathrm{ext}}\left(m{r}
ight)}$ 

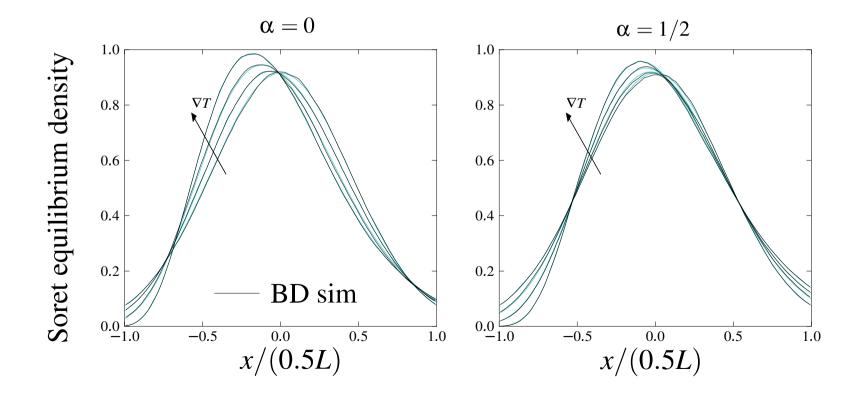


#### 1D Ideal Gas with External Potential

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naïve Boltzmann expectation  $ho\left(m{r}
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ight)V_{\mathrm{ext}}\left(m{r}
ight)}$ 



#### Dynamical Density Functional Theory (DDFT)

- local equilibrium assumption: use equilibrium functional
- traditional functional, dimensionless

$$\beta \mathscr{F}_{\text{id,ext}}[\rho] = \int d^3 \boldsymbol{r} \, \rho(\boldsymbol{r}) \left\{ \ln \left( \rho(\boldsymbol{r}) \Lambda^3 \right) - 1 \right\} + \beta \int d^3 \boldsymbol{r} \, \rho(\boldsymbol{r}) V_{\text{ext}}(\boldsymbol{r})$$

new scaling, adapted ideal gas functional

$$\overline{\mathscr{F}}_{\mathrm{id,ext}}[\rho] = \int \mathrm{d}^{3} \boldsymbol{r} \, \rho(\boldsymbol{r}) \left\{ \ln \left( \rho(\boldsymbol{r}) T^{1-\alpha}(\boldsymbol{r}) \right) - 1 \right\} + \int \mathrm{d}^{3} \boldsymbol{r} \, \rho(\boldsymbol{r}) \int \mathrm{d} \boldsymbol{\tilde{r}} \frac{\nabla V_{\mathrm{ext}}(\boldsymbol{\tilde{r}})}{T(\boldsymbol{\tilde{r}})}$$

interacting particles: how to scale excess functional?

# Interacting Particles – Solvent and Dilute Solute

two possibilities to scale the excess functional

$$\overline{\mathscr{F}}_{\rm exc}^{(H)} = \frac{\mathscr{F}_{\rm exc}}{k_{\rm B}T(r)}$$

$$\overline{\mathscr{F}}_{\rm exc}^{(S)} = \int d\tilde{\boldsymbol{r}} \; \frac{\nabla \mathscr{F}_{\rm exc}[\rho(\tilde{\boldsymbol{r}})]}{k_{\rm B}T(\tilde{\boldsymbol{r}})}$$

solute densities  $\rho^{(\#)}(r)$ 

$$\rho^{(H)}(\mathbf{r}) = \frac{\mathscr{N}}{T^{1-\alpha}} \exp\left(-\frac{\Delta G[T(\mathbf{r})]}{k_{\mathrm{B}}T(\mathbf{r})}\right)$$

$$\rho^{(H)}(\mathbf{r}) = \frac{\mathcal{N}}{T^{1-\alpha}} \exp\left(-\frac{\Delta G[T(\mathbf{r})]}{k_{\rm B}T(\mathbf{r})}\right) \qquad \rho^{(S)}(\mathbf{r}) = \frac{\mathcal{N}}{T^{1-\alpha}} \exp\left(-\int_{-\infty}^{T(\mathbf{r})} \frac{\mathrm{d}\tilde{T}}{k_{\rm B}\tilde{T}} \frac{\partial \Delta G(\tilde{T})}{\partial \tilde{T}}\right)$$

Soret coefficient  $S_{\mathbf{T}}^{(\#)}$ 

$$S_{\mathrm{T}}^{(H)} = \frac{1 - \alpha}{T} - \frac{\beta \Delta H}{T}$$

$$S_{\mathrm{T}}^{(S)} = \frac{1-\alpha}{T} - \beta \Delta S$$

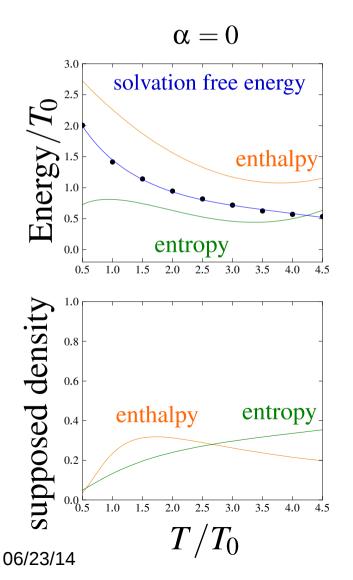
Würger, '14

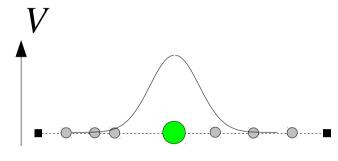
Braun/Dhont, '07

# Gaussian Solute in Ideal Gas Solvent (1D)

**equilibrium**  $\Delta G(T)$  from thermodynamic integration

fit function 
$$\Delta G(T) = a + bT^2 + cT^2 + dT \log(T/T_0)$$

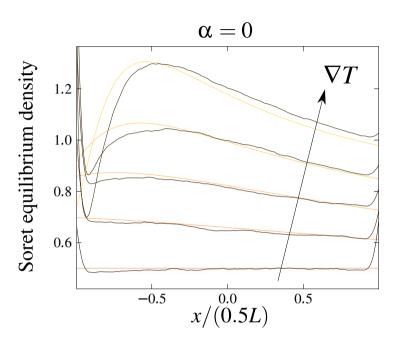


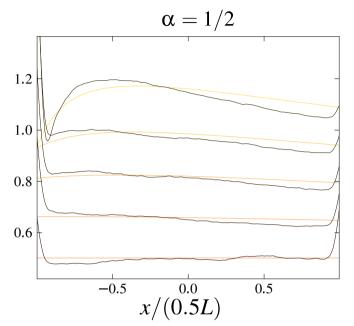


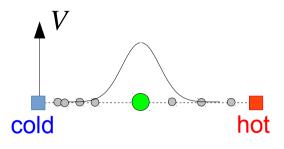
periodic boundary conditions

Significant difference in predictions from enthalpy and entropy approaches

# Thermophoretic BD Sim. – Gaussian Solute



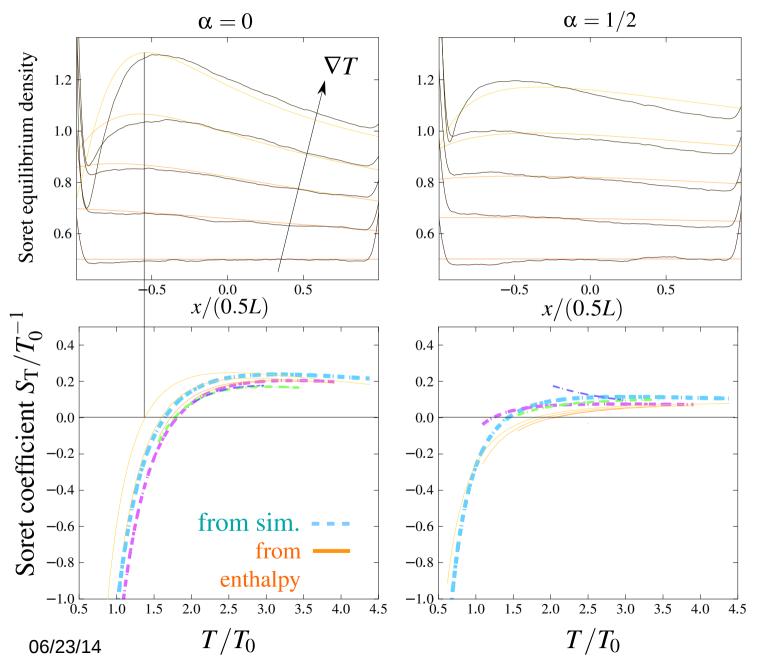


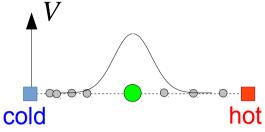


looks like the prediction from the enthalpy

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## Thermophoretic BD Sim. – Gaussian Solute





looks like the prediction from the enthalpy

Soret coefficient

$$S_{\mathrm{T}} = -\frac{\nabla \rho}{\rho \, \nabla T}$$

Sign change!

## Homogeneous 1D System of Gaussian Particles

Soret equilibrium density from enthalpy connected to equation of state

second order virial eq. of state

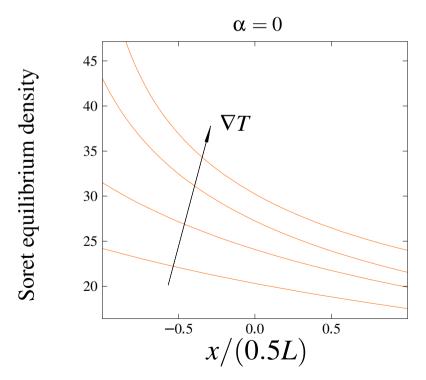
$$\rho = -\frac{1}{2B_2} + \sqrt{\frac{1}{4B_2^2} + \frac{P_{\alpha}}{k_{\rm B}T^{1-\alpha}B_2}}$$

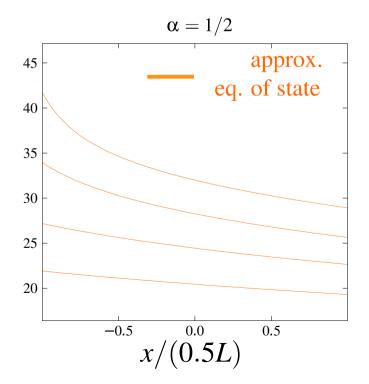
pressure

hot

second virial coefficient

$$B_2 = \frac{1}{2} \int_{-\infty}^{+\infty} \mathrm{d}r \left( 1 - e^{-\beta V(r)} \right)$$





cold

const.

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# Homogeneous 1D System of Gaussian Particles

Soret equilibrium density from enthalpy connected to equation of state

const.

pressure

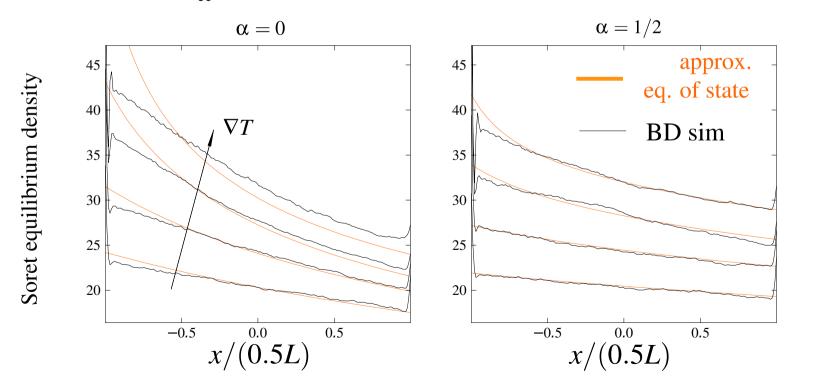
second order virial eq. of state

$$\rho = -\frac{1}{2B_2} + \sqrt{\frac{1}{4B_2^2} + \frac{P_{\alpha}}{k_{\rm B}T^{1-\alpha}B_2}}$$

cold

second virial coefficient

$$B_2 = \frac{1}{2} \int_{-\infty}^{+\infty} \mathrm{d}r \left( 1 - e^{-\beta V(r)} \right)$$



Thermophoretic system seems to follow the equation of state!

# Summary So Far

#### achieved

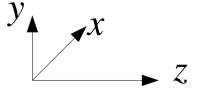
- local equilibrium assumption seems to be appropriate
- Soret coefficient connected to enthalpy
- sign change of  $S_{
  m T}$  reproducible
- crucial assumption:  $D({m r})=k_{
  m B}T({m r})/\gamma$
- open
  - does it describe real systems?
  - choice of  $\alpha$  and its origin

⇒ MD simulations

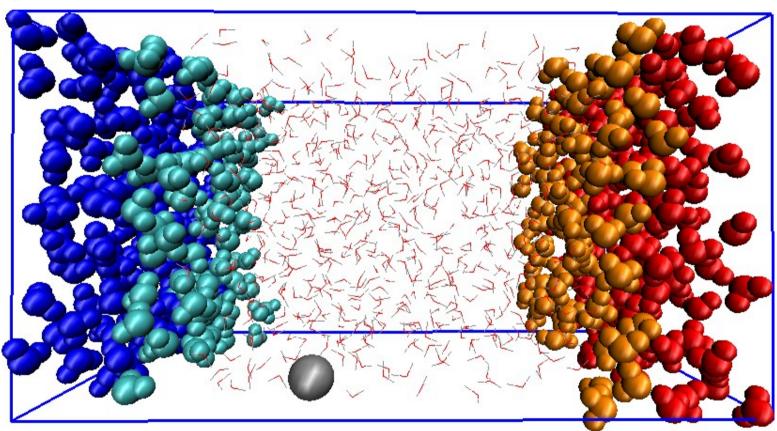
$$S_{\mathrm{T}}^{(H)} = \frac{1 - \alpha}{T} - \frac{\beta \Delta H}{T}$$

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# MD Setup



- periodic boundary conditions in z
- reflective boundary conditions in xy
- SPC/E: water model (extended simple point charge)



 $T_c$ 

- modified SPC/E thermostats (Nosé-Hoover)
- SPC/E solvent
- Lennard-Jones noble gas solute

NPT equilibration

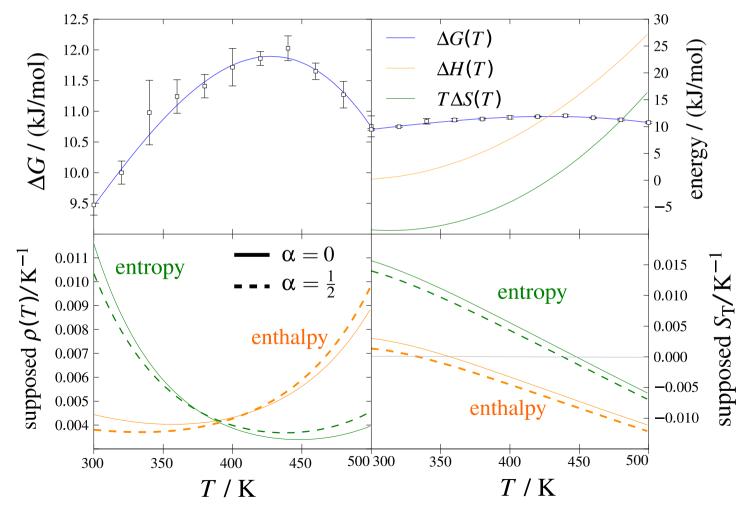
- NVT runs
- measuring the Soret equilibrium density

 $T_h$ 

# First: Solvation Free Energy

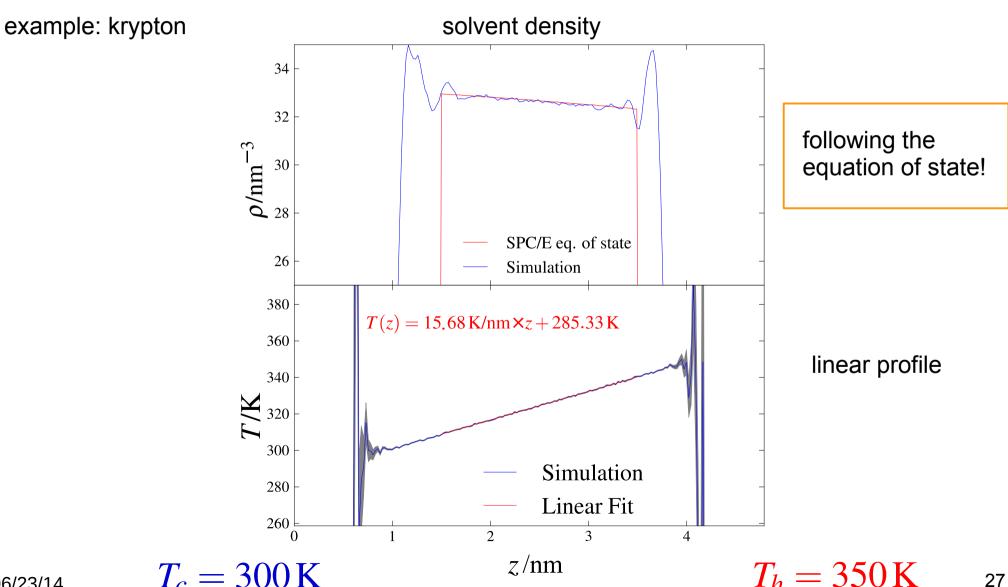
Method: Widom Insertion in thermodynamic equilibrium bulk SPC/E for Ar, Kr and Xe fit function  $\Delta G(T) = a + bT^2 + cT^2 + dT \log(T/1 \, \mathrm{K})$ 

example: xenon



## Thermophoretic simulations - Solvent

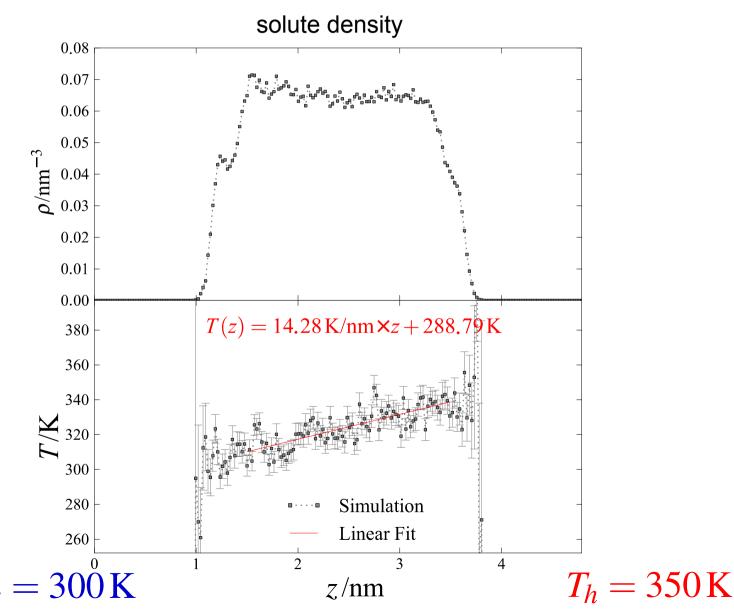
- 7 simulations of Ar, Xe, Kr with run times t > 200 ns in temperature range 300 K < T < 450 K
- measured the Soret equilibrium density and temperature profile of solvent and solute



## Thermophoretic simulations - Solute

- 7 simulations of Ar, Xe, Kr with run times t > 200 ns in temperature range 300 K < T < 450 K
- measured the Soret equilibrium density and temperature profile of solvent and solute

example: krypton



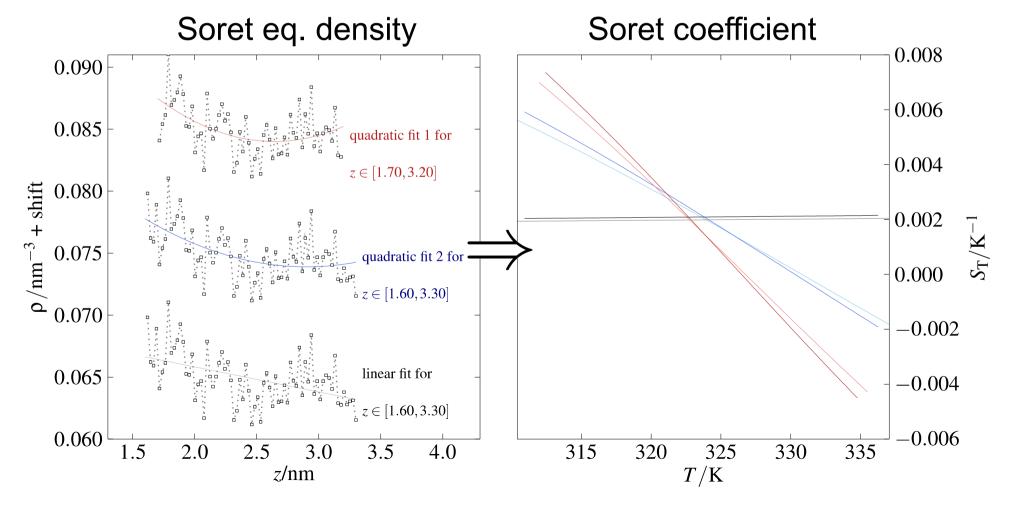
# Soret Coefficient from Density – Krypton

Calculate Soret coefficient as

Kr, 300 K < *T* < 350 K

$$S_{\mathrm{T}} = -rac{
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ho \, 
abla T}$$

derivative – how to model  $\rho$ 



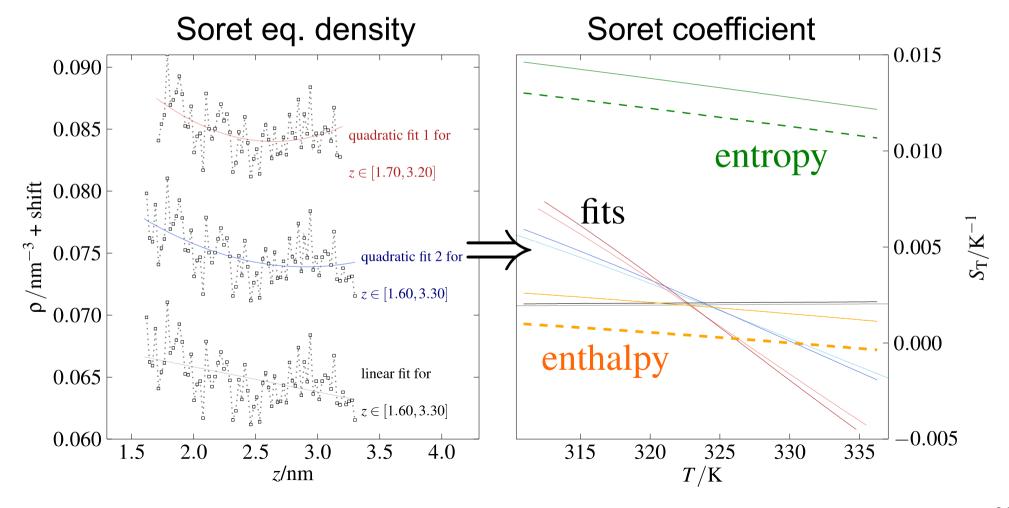
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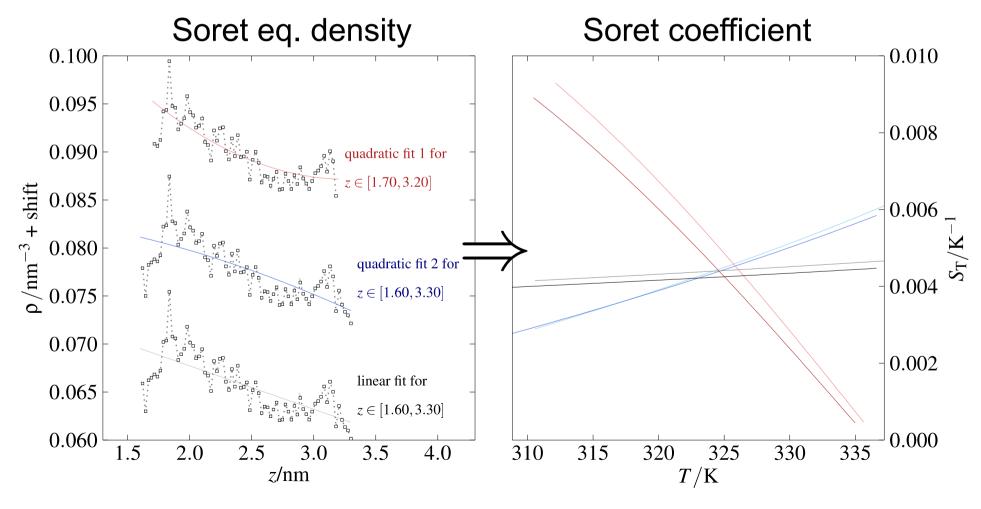


# Soret Coefficient from Density – Argon

#### Calculate Soret coefficient as

Ar, 300 K < T < 350 K

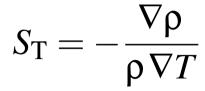
$$S_{\mathrm{T}} = -rac{
abla
ho}{
ho\,
abla T}$$



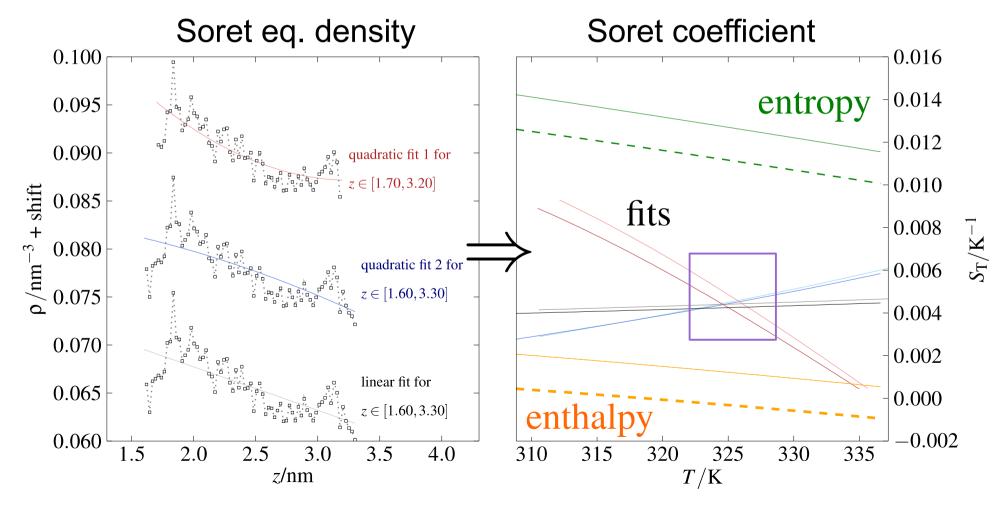
# Soret Coefficient from Density – Argon

#### Calculate Soret coefficient as

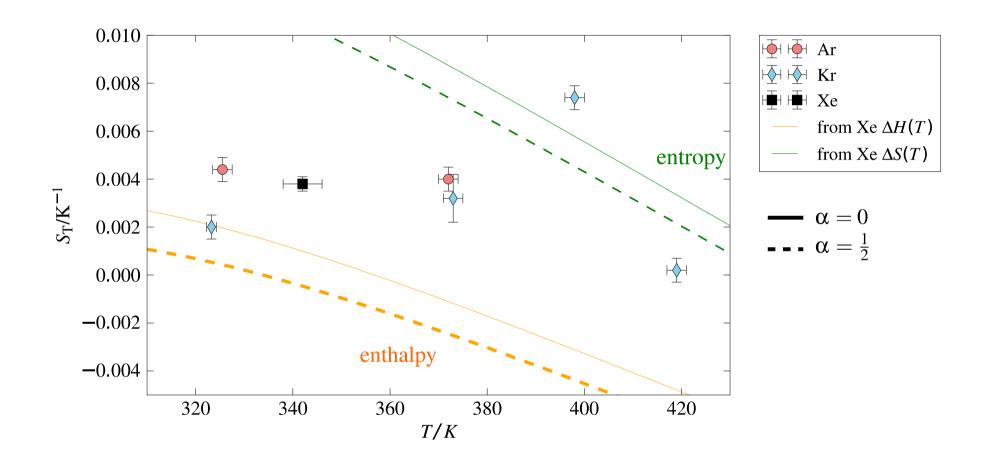
Ar, 300 K < T < 350 K



For every simulation: Soret coefficient is region where all fits coincide



#### Soret Coefficient from all Simulations



no significant agreement/results in MD simulations no sign change

# Summary & Outlook

- theoretical derivation of Soret equilibrium density for various systems via DDFT
- connection between Soret coefficient and solvation enthalpy/equation of state found in BD simulations
- no significant agreement between Soret coefficient and solvation enthalpy in MD simulations
- agreement between Soret eq. density and eq. of state in MD simulations

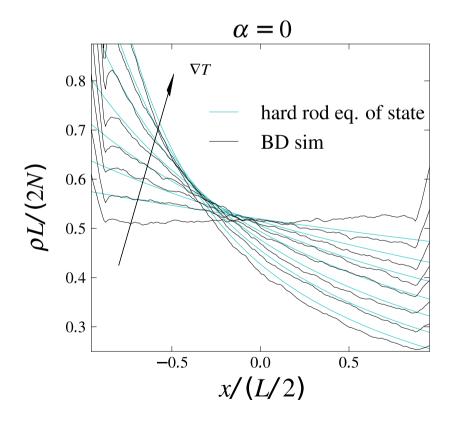
- Ito/Stratonovich
- more realistic BD simulations / difference between BD and MD, extract meaningful data from MD
- connection between local diffusion coefficient and local temperature
- consideration of electrostatics and hydrodynamics

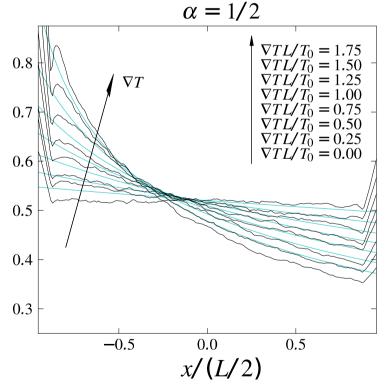
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## 1D Hard Rods

equation of state

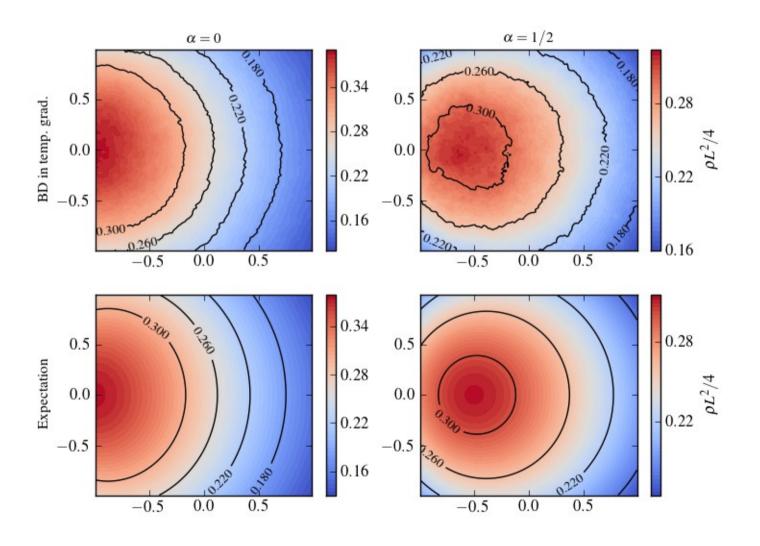
$$P = \frac{Nk_{\rm B}T}{L - (N - 1)\sigma}$$
  $\Rightarrow$   $\rho_{\alpha} = \frac{P_{\alpha}(1 + \sigma/L)}{T^{1 - \alpha} + P_{\alpha}\sigma}$ 





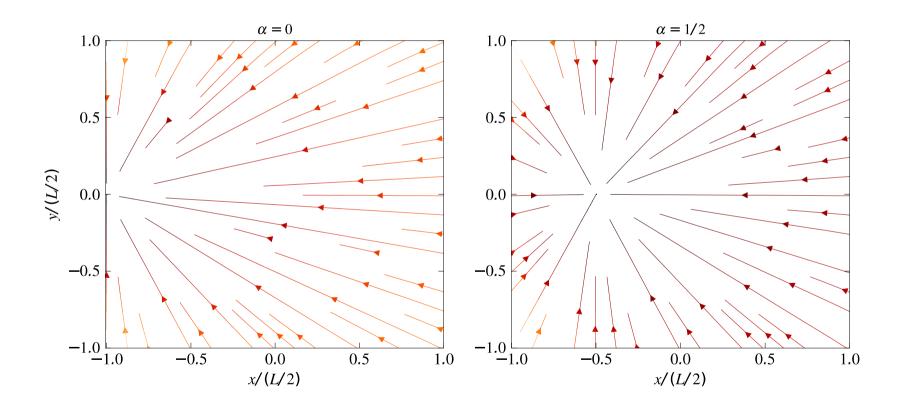
#### 2D Ideal Gas in Harmonic Potential

$$\rho(\mathbf{r}) = \frac{\mathscr{N}}{[T(\mathbf{r})]^{1-\alpha}} \exp\left(-\int_{\mathbf{r}_0}^{\mathbf{r}} d\tilde{\mathbf{r}} \, \frac{\nabla_{\tilde{\mathbf{r}}} V_{\text{ext}}(\tilde{\mathbf{r}})}{k_{\text{B}} T(\tilde{\mathbf{r}})}\right)$$



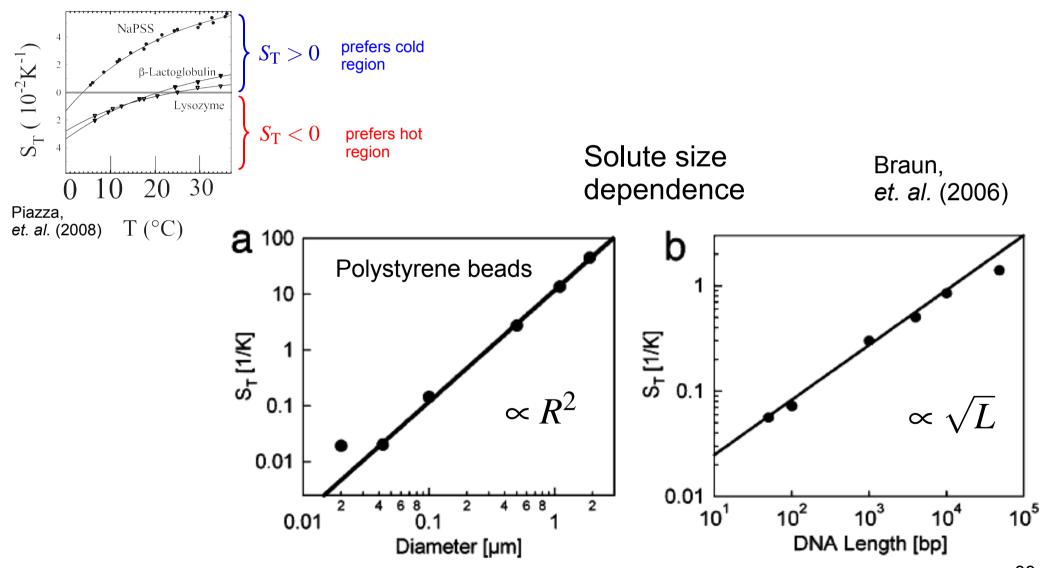
## 2D Ideal Gas in Harmonic Potential

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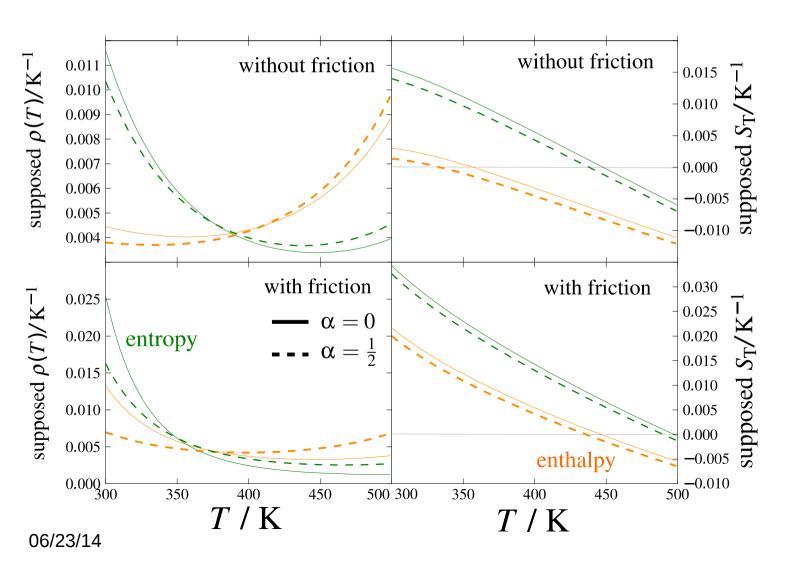
#### **Soret Coefficient**

• sign of  $S_{\rm T}$  dictates direction of movement



# Hypothetical Influence of Friction

Stokes' law  $\gamma(\mathbf{r}) = 6\pi R \eta\left(T(\mathbf{r})\right)$   $\Rightarrow$  consider SPC/E viscosity  $\Rightarrow$  additional factor in density  $\Rightarrow$  additional term in Soret coeff.



- 8 hypotheses
- sign change in  $300 \,\mathrm{K} < T < 500 \,\mathrm{K}$