

# Lecture 30: Mutual Inductance, Magnetic Energy

ECE221: Electric and Magnetic Fields

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# Outline

1 Mutual Inductance

2 Magnetic Energy

3 Faraday's Law

## Mutual Inductance

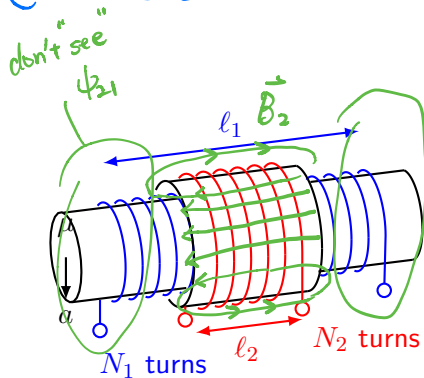
### Mutual flux linkage

$$\Lambda_{12} = N_2 \Psi_{12} \quad \Lambda_{21} = N_1 \Psi_{21}$$

### Mutual inductance

$$L_{12} = \frac{\Lambda_{12}}{I_1} \quad L_{21} = \frac{\Lambda_{21}}{I_2}$$

# Reciprocity



$$n_1 = \frac{N_1}{\ell_1} \text{ turns/m}$$

$$n_2 = \frac{N_2}{\ell_2} \text{ turns/m}$$

When coil 1 is the source,

$$\Lambda_{12} = N_2 \Psi_{12} = N_2 \frac{\mu N_1 \pi a^2 I_1}{\ell_1}$$

$$L_{12} = \frac{\Lambda_{12}}{I_1} = \mu n_1 n_2 \ell_2 \pi a^2$$

When coil 2 is the source, coil 1 only sees a flux linkage of

$$\Lambda_{21} = N_1 \Psi_{21} = \boxed{n_1 \ell_2} \frac{\mu N_2 \pi a^2 I_2}{\ell_2} \psi_{21}$$

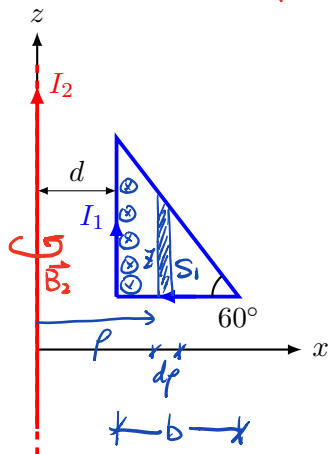
$$L_{21} = \frac{\Lambda_{21}}{I_2} = \mu n_1 n_2 \ell_2 \pi a^2 = L_{12}!$$

↑  
part of coil ① that sees  $\psi_{21}$

$$\ell_{21} = \ell_{12} = \ell$$

## Example: Mutual Inductance Between Wire and Loop

Determine the mutual inductance between an infinitely long wire along the  $z$ -axis and the triangular loop as shown.



Find the flux  $\psi_{21}$  produced by  $I_2$  over coil 1

$$\vec{B}_2 = \frac{\mu_0 I_2}{2\pi \rho} \hat{\phi} \quad \psi_{21} = \int_{S_1} \vec{B}_2 \cdot d\vec{S}_1$$

$$d\vec{S}_1 = z \, dp \, \hat{\phi} \quad z = [p - (d+b)] \tan 60^\circ$$

$$\psi_{21} = \Lambda_{21} = \frac{\sqrt{3} \mu_0 I_2}{2\pi} \int_d^{d+b} \frac{1}{\rho} [p - (d+b)] dp$$

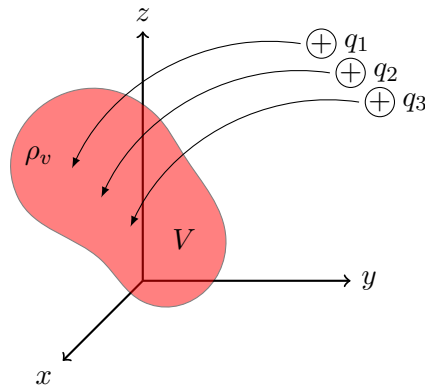
$$= \frac{-\sqrt{3} \mu_0 I_2}{2\pi} [(d+b) \ln(1 + b/d) - b]$$

$$L_{21} = \frac{\Lambda_{21}}{I_2}$$

## Reminder: Electric Energy

Recall what we did to find electric energy stored in a charge distribution:

- 1 Starting in empty space we brought in a charge from infinity;
- 2 Then we brought in another charge and calculated the work needed to position the charge;
- 3 Then another...
- 4 Repeat until we have a volume charge density  $\rho_v$



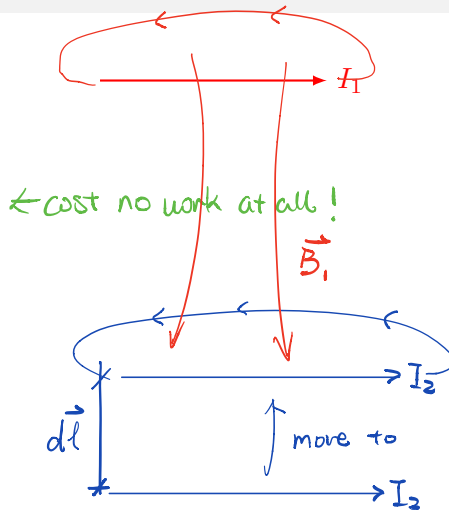
$$W_E = \frac{1}{2} \sum_{m=1}^N q_m V_m = \frac{1}{2} \iiint_V \rho_v V d\nu$$

# Magnetic Energy

Can we apply the same process to compute magnetic energy? Let's try with some line currents.

- ① Starting in empty space we bring in a line current  $I_1$  from infinity;
- ② Then we bring in another line current  $I_2$  from infinity:
  - There is a force between the two conductors, so work is involved in moving the current  $I_2$
  - How would we compute the incremental work done?

$$dW_m = ?$$



# Magnetic Energy

For now, let's leverage the dual relationship between electric and magnetic fields:

## Electrostatics

$$W_e = \frac{1}{2} \iiint_V \mathbf{D} \cdot \mathbf{E} d\nu \text{ [J]}$$

$$\frac{\partial W_e}{\partial \nu} = \frac{1}{2} \mathbf{D} \cdot \mathbf{E} \text{ [J/m}^3\text{]}$$

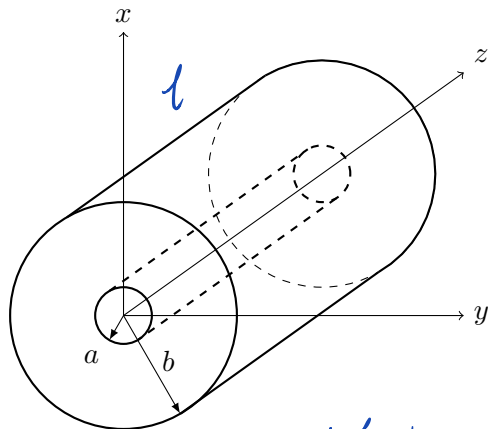
## Magnetostatics

$$W_m = \frac{1}{2} \iiint_V \mathbf{B} \cdot \mathbf{H} d\nu \text{ [J]}$$

$$\frac{\partial W_m}{\partial \nu} = \frac{1}{2} \mathbf{B} \cdot \mathbf{H} \text{ [J/m}^3\text{]}$$



# Example: Magnetic Energy Stored in a Coaxial Cable



$$\text{but } L = \frac{\mu_0 l}{2\pi} \ln \frac{b}{a}$$

$$\vec{H} = \frac{I}{2\pi\rho} \hat{\phi}$$

$$\vec{B} = \frac{\mu_0 I}{2\pi\rho} \hat{\phi}$$

$$\frac{1}{2} \vec{B} \cdot \vec{H} = \frac{\mu_0 I^2}{(2\pi)^2 \rho^2} \cdot \frac{1}{2}$$

$$W_m = \int_0^l \int_0^{2\pi} \int_a^b \frac{1}{2} \frac{\mu_0 I^2}{(2\pi)^2 \rho^2} \cdot \rho \, d\rho \, d\phi \, dz$$

$$= 2\pi l \cdot \frac{\mu_0 I^2}{2(2\pi)^2} \int_a^b \frac{1}{\rho} \, d\rho$$

$$= \frac{\mu_0 I^2 l}{4\pi} \ln \left( \frac{b}{a} \right)$$

$$W_m = \frac{1}{2} L I^2 //$$

# Summary of Electrostatics and Magnetostatics

Maxwell's equations so far are:

Law	Point Form	Differential Form
Gauss' Law	$\nabla \cdot \mathbf{D} = \rho_v$	$\oint_S \mathbf{D} \cdot d\mathbf{s} = Q$
Conservative E-field	$\nabla \times \mathbf{E} = 0$	$\oint_C \mathbf{E} \cdot d\boldsymbol{\ell} = 0$
No magnetic monopoles	$\nabla \cdot \mathbf{B} = 0$	$\oint_S \mathbf{B} \cdot d\mathbf{s} = 0$
Ampère's Law	$\nabla \times \mathbf{H} = \mathbf{J}$	$\oint_C \mathbf{H} \cdot d\boldsymbol{\ell} = I$

The top two equations are **completely decoupled** from the bottom two.

# Faraday's Law

- We have seen that a current can produce a magnetic field
- Michael Faraday wondered: Is the reverse possible?

