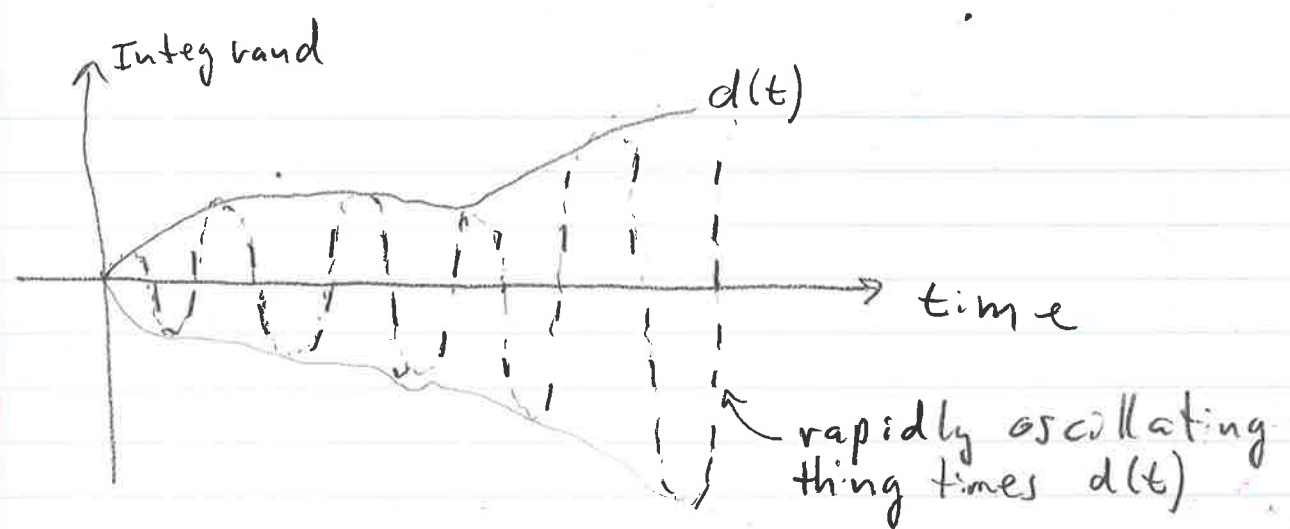


- Q1 The torque is still into the page. So we still need to gain L into the page. Since L is zero initially, it just needs to start rotating about the pivot to develop the necessary L .
- Q2 It precesses (rotates) around z , first coming out of the page.
- Q3 It can't do that because it would violate angular momentum conservation. There is no torque in that direction. We would need a torque, i.e. change its angular momentum.
- Q4 The $|+z\rangle$ state is a stationary state for this operator. So the probability does not change, because the $|z\rangle$ component of $|\psi(0)\rangle$ is constant.
- Q5 1) Initially in the $|+z\rangle$ state means $a=1, b=0$ so $|\langle +x | \psi(t) \rangle|^2 = \frac{1}{2}$, again $|+z\rangle$ is stationary so no time dependence, and $|+z\rangle$ has 50% chance to yield $|+x\rangle$
- 2) This is $a=b=\frac{1}{\sqrt{2}}$ so $|\langle +x | \psi(t) \rangle|^2 = \frac{1}{2} + \frac{1}{2} \cos \omega_0 t$. So we see that the state is "swinging around the z axis"
- 3) $|\langle +x | \psi(t) \rangle|^2 = \frac{1}{2} + \frac{3}{10} \cos \omega_0 t$. Since it's mostly aligned with z , I get a small "swinging component" but mostly a 50% mix.
- Q6 They should be. See comments above.

Q7



Since the rapidly oscillating thing $\times d(t)$ is negative as much as positive, it integrates to 0. On the other hand, my sample $d(t)$ is always positive.

Q8

$d(t)$ is the new stuff associated with B_x . B_x is supposed to be a small perturbation, that slowly rotates $|t\rangle$ to $|-\rangle$. Since B_z is a stronger field ω_0 will be bigger.

Q9

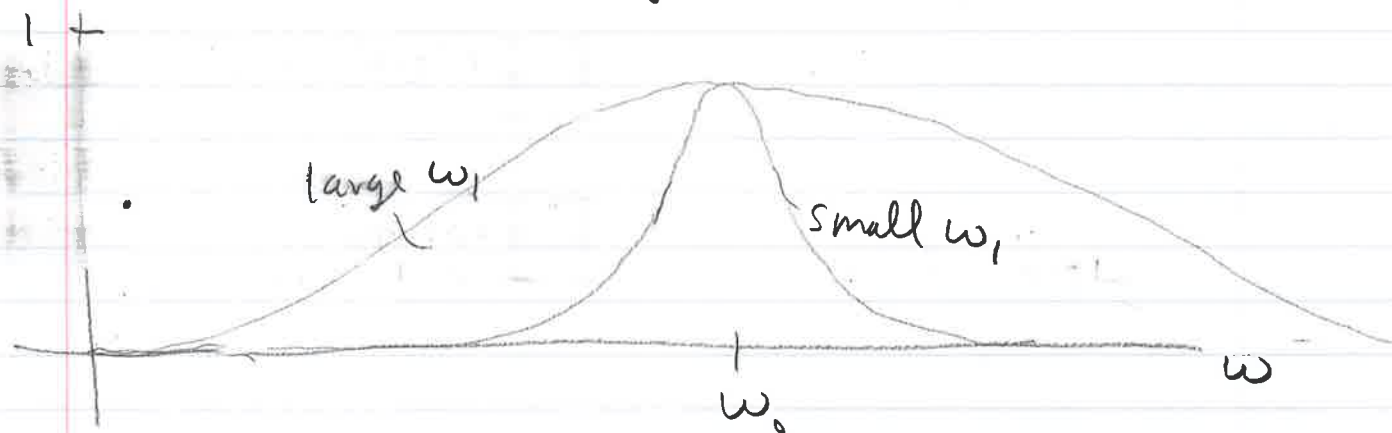
$$\begin{aligned} \ddot{c} &= -i \frac{\omega_1}{4} e^{i(\omega_0 - \omega)t} d + \frac{\omega_1}{4} (\omega_0 - \omega) e^{i(\omega_0 - \omega)t} d \\ &= -i \frac{\omega_1}{4} e^{i(\omega_0 - \omega)t} d - i \frac{\omega_1}{4} e^{-i(\omega_0 - \omega)t} c \\ &\quad + \frac{\omega_1}{4} (\omega_0 - \omega) e^{i(\omega_0 - \omega)t} d - \frac{i \omega_1}{4} e^{-i(\omega_0 - \omega)t} \dot{c} \\ &= -\frac{\omega_1^2}{4} c + i(\omega_0 - \omega) \dot{c} \\ \ddot{c} - i(\omega_0 - \omega) \dot{c} + \left(\frac{\omega_1}{4}\right)^2 c &= 0 \end{aligned}$$

Q10 We get $|\langle -z | \psi(t) \rangle|^2 = 0$, so if we start in $|+z\rangle$ we'll never be found in $|-z\rangle$ without the B_x field that makes $\omega_1 \neq 0$

Q11 $\omega = \omega_0$, like driving on resonance.

Q12 Less time, the thing in front of t is bigger

Q13 We can flip spins faster, and the resonance curve gets broader



We have a higher probability to flip off resonance.