



Universidad de  
**los Andes**



**FACULTAD  
DE INGENIERÍA  
Y CIENCIAS  
APLICADAS**

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# Finite Elements

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# 1. Chapter 1: Direct stiffness method

## 1.1. Truss element 2D

A truss element is a structural element that can only carry axial loads. It is assumed that the truss element is made of a linear elastic material and that the cross-sectional area is constant along its length. Each truss element is defined by two nodes and has two degrees of freedom (DOF's) at each node: Vertical and horizontal displacements. Also, it is important to note, that for each DOF, there is a corresponding force.

$$\mathbf{f} = \begin{bmatrix} f_{x0} \\ f_{y0} \\ f_{x1} \\ f_{y1} \\ \vdots \\ f_{x_{n-1}} \\ f_{y_{n-1}} \end{bmatrix}, \quad \mathbf{K} = \begin{bmatrix} k_{x0x0} & k_{x0y0} & k_{x0x1} & k_{x0y1} & \cdots \\ k_{y0x0} & k_{y0y0} & k_{y0x1} & k_{y0y1} & \cdots \\ k_{x1x0} & k_{x1y0} & k_{x1y1} & k_{x1y1} & \cdots \\ k_{y1x0} & k_{y1y0} & k_{y1y1} & k_{y1y1} & \cdots \\ \vdots & \vdots & \vdots & \vdots & \ddots \end{bmatrix}, \quad \mathbf{u} = \begin{bmatrix} u_{x0} \\ u_{y0} \\ u_{x1} \\ u_{y1} \\ \vdots \\ u_{x_{n-1}} \\ u_{y_{n-1}} \end{bmatrix} \quad (1)$$

So, if we have the forces acting on the nodes of the truss element and the stiffness matrix of the element, we can calculate the displacements of the nodes. Which is the main goal of the finite element method.

$$\mathbf{f} = \mathbf{K} \cdot \mathbf{u} \quad (2)$$

## 2. Finite element method for elastostatic problems

Now we pass to the finite element method for elastostatic problems. The main goal of this method is to find the displacements of the nodes of the structure. To understand this method, we can use the problem of plane stress as an example.

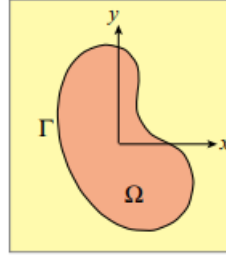


Figura 1: Plane stress problem

As we can see in the figure 1, we have a the domain geometry  $\Gamma$  and Specified interior Forces, which are known forces that act in the interior  $\Omega$  of the plate.

Also one of the most important things are the Specified Surface Forces, these are known forces that act on the boundary  $\Gamma$  and the displacement boundary conditions, these specify how the plate is supported.

### 2.1. Problem Unknowns

The unknown fields are the displacements, strains and stresses.

$$\mathbf{u}(\mathbf{x}, \mathbf{y}) = \begin{bmatrix} u_x(x, y) \\ u_y(x, y) \end{bmatrix}, \quad \boldsymbol{\varepsilon}(\mathbf{x}, \mathbf{y}) = \begin{bmatrix} e_{xx}(x, y) \\ e_{yy}(x, y) \\ 2e_{xy}(x, y) \end{bmatrix}, \quad \boldsymbol{\sigma}(\mathbf{x}, \mathbf{y}) = \begin{bmatrix} \sigma_{xx}(x, y) \\ \sigma_{yy}(x, y) \\ \sigma_{xy}(x, y) \end{bmatrix} \quad (3)$$

### 2.2. Governing equations

The governing equations are the equilibrium equations, the strain-displacement equations and the stress-strain equations.

$$\boldsymbol{\varepsilon} = D\mathbf{u} \quad (4)$$

$$\boldsymbol{\sigma} = E\boldsymbol{\varepsilon} \quad (5)$$

$$D^T \boldsymbol{\sigma} + \mathbf{b} = 0 \quad (6)$$

Is important to note, that  $\mathbf{b}$  is the body force vector,  $E$  is the 3x3 stress-strain matrix of plane stress elastic moduli  $D$  is the 3x2 symmetric-gradient operator and its transpose the  $2 \times 3$  tensor-divergence operator

$$\mathbf{D} = \begin{bmatrix} \frac{\partial}{\partial x} & 0 \\ 0 & \frac{\partial}{\partial y} \\ \frac{\partial}{\partial y} & \frac{\partial}{\partial x} \end{bmatrix}, \quad \mathbf{E} = \begin{bmatrix} E_{11} & E_{12} & E_{13} \\ E_{21} & E_{22} & E_{23} \\ E_{31} & E_{32} & E_{33} \end{bmatrix}, \quad \mathbf{b} = \begin{bmatrix} b_x \\ b_y \end{bmatrix} \quad (7)$$

### 2.3. boundary conditions

There are two boundary conditions prescribed on  $\Gamma$ :

- Displacement boundary conditions ( $\Gamma_u$ ): These are the displacements that are prescribed in the form of  $u = \hat{u}$ .
- Force boundary conditions ( $\Gamma_t$ ): These are the forces that are prescribed in the form  $\sigma_n = \hat{t}$ .

To see it better, we can see the following figure:

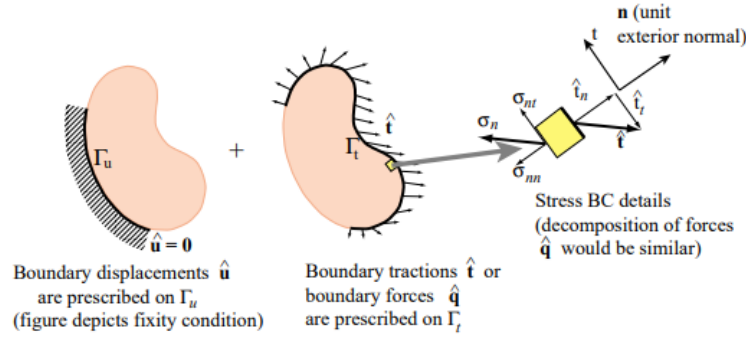


Figura 2: Boundary conditions

### 2.4. displacement interpolation

The displacement field  $u^e(x, y)$  over the element is interpolated from the node displacements. We shall assume that the same interpolation functions are used for both displacement components

$$u_x(x, y) = \sum_{i=1}^n N_i^e(x, y) u_{xi} \quad u_y(x, y) = \sum_{i=1}^n N_i^e(x, y) u_{yi} \quad (8)$$

where  $N_i^e(x, y)$  are the element shape functions. This  $N$  (with superscript  $e$  omitted to reduce clutter) is called the shape function matrix. It has dimensions  $2 \times 2n$

$$\mathbf{N} = \begin{bmatrix} N_1^e & 0 & N_2^e & 0 & \cdots & N_n^e & 0 \\ 0 & N_1^e & 0 & N_2^e & \cdots & 0 & N_n^e \end{bmatrix} \quad (9)$$

Differentiating the finite element displacement field yields the strain-displacement relations:

$$\boldsymbol{\varepsilon}(x, y) = \mathbf{D} \mathbf{N} \cdot \mathbf{u}^e(x, y) = \mathbf{B} \cdot \mathbf{u}^e(x, y) \quad (10)$$

This  $\mathbf{B} = \mathbf{D} \mathbf{N}$  is called the strain-displacement matrix. It is dimensioned  $3 \times 2n$

$$\mathbf{B} = \begin{bmatrix} \frac{\partial N_1^e}{\partial x} & 0 & \frac{\partial N_2^e}{\partial x} & 0 & \cdots & \frac{\partial N_n^e}{\partial x} & 0 \\ 0 & \frac{\partial N_1^e}{\partial y} & 0 & \frac{\partial N_2^e}{\partial y} & \cdots & 0 & \frac{\partial N_n^e}{\partial y} \\ \frac{\partial N_1^e}{\partial y} & \frac{\partial N_1^e}{\partial x} & \frac{\partial N_2^e}{\partial y} & \frac{\partial N_2^e}{\partial x} & \cdots & \frac{\partial N_n^e}{\partial y} & \frac{\partial N_n^e}{\partial x} \end{bmatrix} \quad (11)$$

## 2.5. Element Stiffness Equations

With the following relations, we can calculate the element stiffness matrix and the element force vector.

$$u = Nu^e \quad (12)$$

$$\varepsilon = Bu^e \quad (13)$$

$$\sigma = E\varepsilon \quad (14)$$

The element stiffness matrix is given by the following equation:

$$\mathbf{K}^e = \int_{\Omega^e} hB^T E B d\Omega^e \quad (15)$$

$$(16)$$

and the consistent element nodal force vector is:

$$\mathbf{f}^e = \int_{\Omega^e} hN^T b d\Omega^e + \int_{\Gamma^e} hN^T \hat{t} d\Gamma^e \quad (17)$$

$$(18)$$

where  $h$  is the thickness of the element,  $b$  is the body force vector and  $\Omega^e$  is the volume of the element.

Finally, we can describe the element stiffness matrix and the element force vector in a more compact form:

$$\int_{\Omega^e} hB^T E B d\Omega^e \cdot u = \int_{\Omega^e} hN^T b d\Omega^e + \int_{\Gamma^e} hN^T \hat{t} d\Gamma^e \rightarrow \mathbf{K}^e \cdot u = \mathbf{f}^e \quad (19)$$

## 2.6. 2D Elasticity

### 2.6.1. Plane stress

Plane stress is occupied when the stress in the  $z$  direction is negligible. In other words when the plane has a very small thickness. The stress-strain relations are given by the following equations:

$$\sigma = \frac{E}{1-\nu^2} \cdot \begin{bmatrix} 1 & \nu & 0 \\ \nu & 1 & 0 \\ 0 & 0 & \frac{1-\nu}{2} \end{bmatrix} \cdot \varepsilon \quad (20)$$

### 2.6.2. Plane strain

Plane strain is occupied when the strain in the  $z$  direction is infinite. In other words when the plane has a very big thickness. The stress-strain relations are given by the following equations:

$$\sigma = \frac{E}{(1+\nu)(1+2\nu)} \cdot \begin{bmatrix} 1 & \nu & 0 \\ \nu & 1 & 0 \\ 0 & 0 & \frac{1-\nu}{2} \end{bmatrix} \cdot \varepsilon \quad (21)$$

### 3. The constant strain/stress triangle

The triangle is a 2D element with three nodes. Each node has its own coordinates  $(x_i, y_i)$  for  $i = 1, 2, 3$ . But we can also use the local coordinates  $(\zeta_i)$  for  $i = 1, 2, 3$ . The local coordinates are used to simplify the calculations. The only restriction is that:

$$\zeta_1 + \zeta_2 + \zeta_3 = 1$$

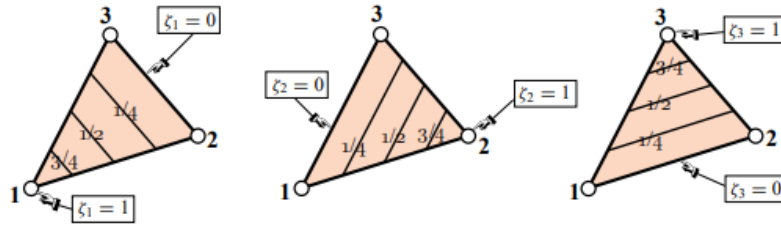


Figura 3: Constant strain triangle

To transform the global coordinates to the local coordinates, we can use the following equations:

$$\begin{bmatrix} x \\ y \\ 1 \end{bmatrix} = \begin{bmatrix} x_1 & x_2 & x_3 \\ y_1 & y_2 & y_3 \\ 1 & 1 & 1 \end{bmatrix} \cdot \begin{bmatrix} \zeta_1 \\ \zeta_2 \\ \zeta_3 \end{bmatrix} \quad (22)$$

And now we need to specify the interpolation matrix  $N(x, y)$ , because:

$$B(x, y) = D \cdot N(x, y)$$

So, we can use the following equations to calculate the interpolation matrix:

$$D(x, y) = \begin{bmatrix} N_1(x, y) & 0 & N_2(x, y) & 0 & N_3(x, y) & 0 \\ 0 & N_1(x, y) & 0 & N_2(x, y) & 0 & N_3(x, y) \end{bmatrix} = \begin{bmatrix} \zeta_1 & 0 & \zeta_2 & 0 & \zeta_3 & 0 \\ 0 & \zeta_1 & 0 & \zeta_2 & 0 & \zeta_3 \end{bmatrix} \quad (23)$$

$$(24)$$

With the interpolation matrix, we can calculate the strain-displacement matrix  $B(x, y)$  using the following equations:

$$B(x, y) = \frac{1}{2 \cdot A^e} \begin{bmatrix} \frac{\partial N_1}{\partial x} & 0 & \frac{\partial N_2}{\partial x} & 0 & \frac{\partial N_3}{\partial x} & 0 \\ 0 & \frac{\partial N_1}{\partial y} & 0 & \frac{\partial N_2}{\partial y} & 0 & \frac{\partial N_3}{\partial y} \\ \frac{\partial N_1}{\partial y} & \frac{\partial N_1}{\partial x} & \frac{\partial N_2}{\partial y} & \frac{\partial N_2}{\partial x} & \frac{\partial N_3}{\partial y} & \frac{\partial N_3}{\partial x} \end{bmatrix} = \begin{bmatrix} \frac{y_2 - y_3}{2A^e} & 0 & \frac{y_3 - y_1}{2A^e} & 0 & \frac{y_1 - y_2}{2A^e} & 0 \\ 0 & \frac{x_3 - x_2}{2A^e} & 0 & \frac{x_1 - x_3}{2A^e} & 0 & \frac{x_2 - x_1}{2A^e} \\ \frac{x_2 - x_3}{2A^e} & \frac{y_2 - y_3}{2A^e} & \frac{x_3 - x_1}{2A^e} & \frac{y_3 - y_1}{2A^e} & \frac{x_1 - x_2}{2A^e} & \frac{y_1 - y_2}{2A^e} \end{bmatrix} \quad (25)$$

This changes the form to calculate the Stiffnes matrix and the force vector. The stiffness matrix is given by the following equation:

$$\mathbf{K}^e = \int_{\Omega^e} B^T E B d\Omega^e \rightarrow B^T(x, y) \cdot E_x \cdot B(x, y) \cdot A^e \cdot t^e \quad (26)$$

While the force vector is given by the following equation:

$$\mathbf{f}^e = \int_{\Omega^e} N^T b d\Omega^e \rightarrow \frac{A^e \cdot t^e}{3} \cdot \begin{bmatrix} b_x \\ b_y \\ b_x \\ b_y \\ b_x \\ b_y \end{bmatrix} \quad (27)$$

## 4. Iso-parametric finite elements

This term is used to describe the finite elements that **have the same shape functions and the same interpolation functions**. This means that the shape functions are used to interpolate the displacements and the shape functions are used to interpolate the coordinates of the nodes. The iso-parametric elements are used to solve problems with complex geometries.

There are two types of mesh refinement:

- p-refinement: This is when we have the same number of elements in each direction, but we increase the order of the polynomial.
- h-refinement: This is when we have the same order of the polynomial, but we increase the number of elements in each direction.

The p-refinement is used when we want to increase the accuracy of the solution, but leads to a Stiffness matrix with larger bandwidth.

### 4.1. Iso-parametric concept

A Finite element is iso-parametric if its shape and interpolation functions are the same. This means, that:

- Convergence: The iso-parametric elements converge to the exact solution as the number of elements increases.
- Consistency: The iso-parametric elements are consistent with the exact solution.
- Stability: The iso-parametric elements are stable, which means that the solution does not oscillate or diverge.

In other words, needs to satisfy the following conditions:

$$1 = \sum N_i^e \quad \forall i = 1, 2, \dots, n \quad (28)$$

$$x = \sum N_i^e \cdot x_i \quad \forall i = 1, 2, \dots, n \quad (29)$$

$$y = \sum N_i^e \cdot y_i \quad \forall i = 1, 2, \dots, n \quad (30)$$

$$u_x = \sum N_i^e \cdot u_{xi} \quad \forall i = 1, 2, \dots, n \quad (31)$$

$$u_y = \sum N_i^e \cdot u_{yi} \quad \forall i = 1, 2, \dots, n \quad (32)$$

$$(33)$$

## 5. Quad 4 element

The natural coordinates for a quadrilateral element are  $\xi$  and  $\eta$ . As shown in the figure below:

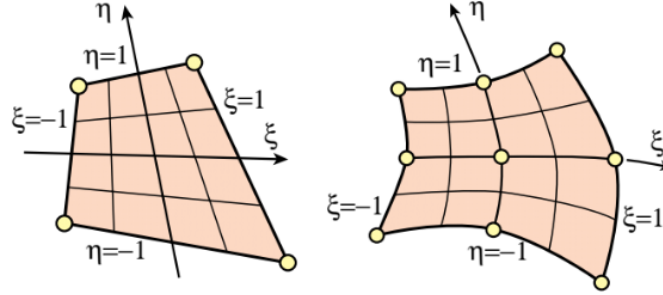


Figura 4: Quad 4 element

Due to the fact that the quad 4 element is represented by a bilinear polynomial, we can use the following equations to calculate the shape functions:

$$N_1(\xi, \eta) = \frac{1}{4}(1 - \xi)(1 - \eta) \quad (34)$$

$$N_2(\xi, \eta) = \frac{1}{4}(1 + \xi)(1 - \eta) \quad (35)$$

$$N_3(\xi, \eta) = \frac{1}{4}(1 + \xi)(1 + \eta) \quad (36)$$

$$N_4(\xi, \eta) = \frac{1}{4}(1 - \xi)(1 + \eta) \quad (37)$$

### 5.1. Jacobian matrix

In quadrilateral element derivations we will need the Jacobian of two-dimensional transformations that connect the differentials of  $x, y$  to those of  $\xi, \eta$  and vice-versa. The Jacobian matrix is given by the following equation:

$$\begin{bmatrix} dx \\ dy \end{bmatrix} = \begin{bmatrix} \frac{\partial x}{\partial \xi} & \frac{\partial x}{\partial \eta} \\ \frac{\partial y}{\partial \xi} & \frac{\partial y}{\partial \eta} \end{bmatrix} \cdot \begin{bmatrix} d\xi \\ d\eta \end{bmatrix} = \mathbf{J}^T \cdot \begin{bmatrix} d\xi \\ d\eta \end{bmatrix} \quad (38)$$

The scalar symbol  $J$  is the Jacobian determinant.

### 5.2. Stiffness matrix and force vector

The stiffness matrix and the force vector are given by the following equations:

$$\mathbf{K}^e = \int_{\Omega^e} \mathbf{B}^T \mathbf{E} \mathbf{B} d\Omega^e \rightarrow \int_{-1}^1 \int_{-1}^1 \mathbf{B}^T \mathbf{E} d\det(I) d\xi d\eta \quad (39)$$

$$\mathbf{f}^e = \int_{\Omega^e} \mathbf{N}^T \mathbf{b} d\Omega^e \rightarrow \int_{-1}^1 \int_{-1}^1 \mathbf{N}^T \mathbf{b} d\det(J) d\xi d\eta \quad (40)$$

$$(41)$$



## 6. Numerical integration by Gauss Rules

The standard practice has been to use Gauss integration because such rules use a minimal number of sample points to achieve a desired level of accuracy

### 6.1. 1D Gauss quadrature

The 1D Gauss quadrature is given by the following equation:

$$\int_{-1}^1 f(x) dx = \sum_{i=1}^n w_i f(x_i) \quad (42)$$

$$(43)$$

where  $w_i$  are the weights and  $x_i$  are the sample points. The weights and sample points are given by the following table:

Points	Rule
1	$\int_{-1}^1 F(\xi) d\xi \approx 2F(0)$
2	$\int_{-1}^1 F(\xi) d\xi \approx F(-1/\sqrt{3}) + F(1/\sqrt{3})$
3	$\int_{-1}^1 F(\xi) d\xi \approx \frac{5}{9}F(-\sqrt{3/5}) + \frac{8}{9}F(0) + \frac{5}{9}F(\sqrt{3/5})$
4	$\int_{-1}^1 F(\xi) d\xi \approx w_{14}F(\xi_{14}) + w_{24}F(\xi_{24}) + w_{34}F(\xi_{34}) + w_{44}F(\xi_{44})$
5	$\int_{-1}^1 F(\xi) d\xi \approx w_{15}F(\xi_{15}) + w_{25}F(\xi_{25}) + w_{35}F(\xi_{35}) + w_{45}F(\xi_{45}) + w_{55}F(\xi_{55})$

For the 4-point rule,  $\xi_{34} = -\xi_{24} = \sqrt{(3 - 2\sqrt{6/5})/7}$ ,  $\xi_{44} = -\xi_{14} = \sqrt{(3 + 2\sqrt{6/5})/7}$ ,  $w_{14} = w_{44} = \frac{1}{2} - \frac{1}{6}\sqrt{5/6}$ , and  $w_{24} = w_{34} = \frac{1}{2} + \frac{1}{6}\sqrt{5/6}$ .

For the 5-point rule,  $\xi_{55} = -\xi_{15} = \frac{1}{3}\sqrt{5 + 2\sqrt{10/7}}$ ,  $\xi_{45} = -\xi_{35} = \frac{1}{3}\sqrt{5 - 2\sqrt{10/7}}$ ,  $\xi_{35} = 0$ ,  $w_{15} = w_{55} = (322 - 13\sqrt{70})/900$ ,  $w_{25} = w_{45} = (322 + 13\sqrt{70})/900$  and  $w_{35} = 512/900$ .

Figura 5: One-Dimensional Gauss Rules with 1 through 5 Sample Points

#### 6.1.1. Example: Case 2 points

The 2 points Gauss quadrature is given by the following equation:

$$\int_{-1}^1 f(x) dx = w_1 f(x_1) + w_2 f(x_2) \quad (44)$$

Now we can calculate the weights and the sample points. In order to do that, there are as many degrees of x as there are sample points. So, we can use the following Equations:

Grade 0:

$$w_1 + w_2 = 2 \quad (45)$$

Grade 1:

$$w_1 \cdot x_1 + w_2 \cdot x_2 = 0 \quad (46)$$

Grade 2:

$$w_1 \cdot x_1^2 + w_2 \cdot x_2^2 = \frac{2}{3} \quad (47)$$

Solving the equations, we get the following results:

$$w_1 = 1 \quad (48)$$

$$w_2 = 1 \quad (49)$$

$$x_1 = -\frac{1}{\sqrt{3}} \quad (50)$$

$$x_2 = \frac{1}{\sqrt{3}} \quad (51)$$

That is the same as in the table 5.

## 6.2. 2D Gauss quadrature

The 2D Gauss quadrature is given by the following equation:

$$\int_{-1}^1 \int_{-1}^1 f(x, y) dx dy = \sum_{i=1}^n \sum_{j=1}^n w_i w_j f(x_i, y_j) \quad (52)$$

$$(53)$$