# Contents

Ι	OI	otical	Coherence Tomography	6
1	Time Domain OCT			7
	1.1	Theor	y	7
		1.1.1	Michelson Interferometer	7
		1.1.2	Optical interferometry with coherent light	9
		1.1.3	Low-coherence interferometry	9
		1.1.4	Low coherence interferometry with multilayer structures	10
		1.1.5	Detector signal	12
	1.2	Reflec	tivity Sensitivity	12
		1.2.1	Effective reflectivity vs interface reflectivity	13
		1.2.2	Signal to noise ratio	13
		1.2.3	Noise sources	13
			1.2.3.1 Shot noise	13
			1.2.3.2 Relative intensity noise	13
			1.2.3.3 Thermal noise	14
		1.2.4	Sensitivity optimization	14
2	Fou	rier D	omain OCT	15
	2.1	Overv	iew of fourier domain OCT	15
	2.2	Freque	ency domain OCT	17
		2.2.1	FD OCT image reconstruction	19
			2.2.1.1 DC level removal	19
			2.2.1.2 Image in clean area	19
			2.2.1.3 Autointerference supression double spectrum	19

CONTENTS 2

		2.2.1.4 Differential Fourier domain method	19
		2.2.1.5 Heterodyne complex FDOCT	19
		2.2.1.6 Quadrature Fourier domain method	20
		2.2.1.7 Triple spectrum complete reconstruction	20
		2.2.1.8 Hibert transform	20
		2.2.2 Spectral resolution effects	20
	2.3	Processing	20
3	Sne	ectral Radar	23
J	3.1	Spectrometer	23
	0.1	Specifolicies	20
4	Swe	ept Source	25
	4.1	Light Source	25
	4.2	Detection and Digitalization	25
	4.3	Signal Processing	25
	4.4	Probe	25
5	Sign	nal Processing	26
6	App	plications	27
U			41
U	6.1	Materials and Methods	
U	6.1 6.2		28
	6.2	Materials and Methods	28
	6.2	Materials and Methods	28 29 <b>31</b>
	6.2 <b>Der</b>	Materials and Methods	28 29 <b>31</b>
	6.2 <b>Der</b> 7.1	Materials and Methods	28 29 <b>31</b> 31 31
	6.2 <b>Der</b> 7.1 7.2	Materials and Methods	28 29 <b>31</b> 31 31 32
	6.2 Den 7.1 7.2 7.3	Materials and Methods  Results  Atal fiber-reinforced composite analisis  abstract  Intro  Experimental details	28 29 <b>31</b> 31 31 32 34
	6.2  Den 7.1 7.2 7.3 7.4 7.5	Materials and Methods  Results  Atal fiber-reinforced composite analisis  abstract  Intro  Experimental details  Fiber reinforced composite studies	28 29 <b>31</b> 31 31 32 34
7	6.2  Den 7.1 7.2 7.3 7.4 7.5	Materials and Methods  Results  Intal fiber-reinforced composite analisis  abstract  Intro  Experimental details  Fiber reinforced composite studies  Conclusions	28 29 31 31 31 32 34 36 restorations
7	6.2  Den 7.1 7.2 7.3 7.4 7.5  Clin	Materials and Methods Results  Intal fiber-reinforced composite analisis abstract Intro Experimental details Fiber reinforced composite studies Conclusions  Intal use of optical coherence tomography to evaluate the integrity of dental	28 29 31 31 31 32 34 36 restorations

CONTENTS 3

	8.4	Results	42
	8.5	Discussion	46
	8.6	Conclusion	49
II	O	ptics on chip	51
	8.7	Broad picture	52
	8.8	Slab waveguide	52
	8.9	Dispersion in Silicon	55
	8.10	Rectangular waveguide	55
		8.10.1 fem	55
	8.11	Bending radius	55
	8.12	Ring resonators	56
	8.13	Coupling light to the chip	56
	8.14	Active components	57
		8.14.1 Thermo optic effect	57
		8.14.2 Eletrooptic effect	57
		8.14.3 Effect $\chi^3$	57
		8.14.4 Plasma dispersion	57
	8.15	Light sources	57
	8.16	Detectors	57
	8.17	Fabrication Tecniques	58
		8.17.1 Fabrication difficulties discussion	58
9	Diff	raction Grating Spectrometers	59
	9.1	<u> </u>	59
	9.2	Kirchhoff's diffraction theory	
	9.3	Diffraction grating	
	0.0	9.3.1 Free spectral range	
	9.4	Planar devices	
	9.5	Spectrometer Architectures	
	5.0	9.5.1 Czerny-Turner	
		9.5.2 Rowland	
		DIGIE INCHIMITE I I I I I I I I I I I I I I I I I I	$\sim$

CONTENTS 4

		9.5.3	Arrayed Waveguide Grating (AWG)	69
	9.6	Grating	g illuminated by a Gaussian Beam	70
	9.7	Size .		72
	9.8	Spectra	al defects	72
		9.8.1	Rowland Ghosts	72
		9.8.2	Lyman Ghosts	72
	9.9	Aberra	tions	72
		9.9.1	Aberration free conditions	73
		9.9.2	One point stigmatic correction	74
		9.9.3	Two points stigmatic correction	75
	9.10	Blazing	g and Groove size	75
	9.11	Rayleig	gh-Huygens model	75
	9.12	Compa	risons of simulationas and analitic solution	76
	9.13	Simula	tion of Spectral defects and aberrations using Rayleigh-Huygens model	76
	9.14	implem	nentation	76
		9.14.1	Design	77
		9.14.2	Operating wavelength	77
		9.14.3	Dispersion	77
		9.14.4	Grating size	77
		9.14.5	Diffraction order and input, output angles	78
		9.14.6	Number of grooves	78
		9.14.7	Linear Dipersion	78
		9.14.8	input and output angle	79
		9.14.9	Output waveguide positioning	80
		9.14.10	Fabrication	80
		9.14.11	Testing	81
10	Ring	r enhai	nced spectrometer	83
10	`		Theory	
	10.1		Crosstalk	
	10.2			
			ation	
	10.0	Tabile	AUIOII	00

CONTENTS	5
CONTENTS	6

10.4	Testing	85
10.5	Increasing Channel Density	86
10.6	Device Comparison	88
10.7	Effect of transmission spectrum lineshape	88
10.8	Numerical calculation	91
10.9	Approximations	91

# Part I Optical Coherence Tomography

## Chapter 1

## Time Domain OCT

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#### 1.1 Theory

In OCT, the kind of interferometry used is the low-coherence type. In this section, we are going to describe how low-coherence interferometry is used to measure the distance of an object with more the one reflecting surface.

#### 1.1.1 Michelson Interferometer

Consider the interferometer shown in figure 1.1.1. To simplify the analisys, ignore the light polarization and dispersion effects. A light source emits a radiation whose electric field shine upon a beamspliter e described by  $E_0(t)$ . This light is devided and a fraction D of its power is directed

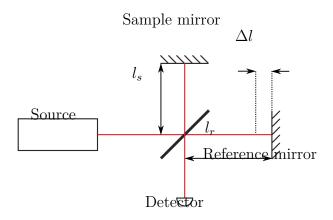


Figure 1.1.1: Michelson interferometer

to mirror S with reflectivity  $R_s$ . The light is reflected and returns to the beam spliter, where a fraction (1-D) of its power goes to the detector. Similarly, in the other interferometer arm, (1-D) of the light emited by the light source is redirected to the mirror  $R_r$ , which is reflected and returns to the beamspliter, where a fraction D of this light is redirected back to the detector. The electric field at the detector can be described by

$$E_{s} = \sqrt{K_{0}}E_{0}(t - \tau_{a}) + \sqrt{K_{r}}E_{0}(t - \tau_{r}), \qquad (1.1.1)$$

where  $K_i = R_i D (1 - D)$  is the light power fraction that is shined to the detector and  $\tau_i$  is the time the light takes to go from the beamspliter to the mirror i and back.

The relationship between the output light intensity  $\langle I \rangle$  and the mean electric field is:

$$\langle I \rangle = \left\langle \frac{|E(t)|^2}{2\eta_0} \right\rangle,$$
 (1.1.2)

where  $\eta_0 = \sqrt{\frac{\mu_0}{\varepsilon_0}}$  is the free space impedance. Using 1.1.1 and 1.1.2, we get the intensity at the detector as been

$$\langle I \rangle = \langle I_{\rm DC} \rangle + \sqrt{K_s K_r} \Re \left( \Gamma \left( \tau_{s,a} \right) \right)$$
 (1.1.3)

where  $\tau_{r,s} = \tau_s - \tau_r$ . We define the autocorrelation function for beam electric field as

$$\Gamma\left(\tau\right) = \frac{E\left(t - \tau\right)E^*\left(t\right)}{\eta_0}$$

and

$$\langle I_{\rm DC} \rangle = \sum_{j} K_{j} \langle I_{0} \rangle ,$$

where  $\langle I_0 \rangle = \left\langle \frac{|E_0(t)|^2}{2\eta_0} \right\rangle$  is the mean light intensity that leaves the source, which in this case is constant. More specifically, it is of interest that the light source should be stationary in the wide sense, which means

- 1.  $\langle E(t) \rangle$  is independent of t.
- 2.  $\langle E(t_1) E^*(t_2) \rangle$  depends only on  $\tau = t_2 t_1$ .

#### 1.1.2 Optical interferometry with coherent light

If the light is perfectly coherent (i.e. monochromatic), then its electric field can be described as  $E(t) = E_0 e^{-i2\pi\nu t}$  where  $\nu$  is the optical frequency and  $E_0$  is the field amplitude. Replacing  $E(t) = E_0 e^{-i2\pi\nu t}$  in equation and using, we get