

## 1 Reduced models of short term synaptic plasticity

Let  $c$ ,  $p$ , and  $x$  represent the intracellular concentration of  $\text{Ca}^{2+}$  at the presynaptic terminal, the proportion of released vesicles per unit time, and the normalized readily releasable vesicles. Assume that the dynamics are given by

$$\partial_t c = \frac{c_b - c}{\tau_b} - k_r \alpha_r (1 - r) c + f(t) \quad (1)$$

$$\partial_t r = \alpha_r c (1 - r) - \beta_r r \quad (2)$$

$$\partial_t q = \frac{q_\infty - q}{\tau_q} - \delta q r \quad (3)$$

where  $c_b$  and  $\tau_b$  represent the steady state and time constant for presynaptic calcium due to buffering in the absence of any other perturbation. The term  $f(c)$  represents the flux of  $\text{Ca}^{2+}$  into the terminal via voltage-dependent channels. The term  $k_c \alpha_r c (1 - r)$  represents the decrease in the calcium concentration due to binding to the release machinery.

**Table 1:** Parameters and values

$c_b$	$\mu\text{M}$	steady state constant for presynaptic $\text{Ca}^{2+}$ due to buffering
$\alpha_r$	$\text{ms}^{-1} \mu\text{M}^{-1}$	Activation rate of the release machinery in the presence of $\text{Ca}^{2+}$
$\beta_r$	$\text{ms}^{-1}$	Deactivation rate of the release machinery
$k_r$	$\mu\text{M}$	Impact of the $\text{Ca}^{2+}$ activating the release machinery on the intracellular $\text{Ca}^{2+}$ concentration in the terminal
$q_\infty$		Steady state for que normalized quanta in the readily releasable pool
$\tau_b$	$\text{ms}$	time constant for presynaptic $\text{Ca}^{2+}$ due to buffering
$\tau_q$	$\text{ms}$	Time constant for the recovery of the readily releasable pool
$\delta$	$\text{ms}^{-1}$	Conversion rate

### 1.1 $\text{Ca}^{2+}$ dynamics

It is worth noticing that the dynamics for  $c$  are linear when  $f(t) = 0$ . To see this, equation (1) can be transformed to

$$\partial_t c = f(t) + \frac{c_\infty - c}{\tau_c}$$

where

$$\tau_c = \frac{\tau_b}{1 + k_r \alpha_r (1 - r) \tau_b}, \quad (4)$$

$$c_\infty = \frac{c_b}{1 + k_r \alpha_r (1 - r) \tau_b}. \quad (5)$$

That is, if  $f(t) = 0$ , then the dynamics for  $c$  from an initial condition  $c_o = c(t_o)$  behave like

$$c(t) = c_\infty - (c_\infty - c_o) \exp\left(\frac{t_o - t}{\tau_c}\right).$$

which means that in the absence of any  $\text{Ca}^{2+}$  influx, the concentration of  $\text{Ca}^{2+}$  in the presynaptic terminal goes to a steady state  $c_\infty$  that becomes the steady state for buffering only for fast buffering (small  $\tau_b$ ), or if the rate of binding to the activation machinery is slow enough (small  $\alpha_r$ ).

**Calcium influx.** To better understand the  $\text{Ca}^{2+}$  dynamics, this time in the presence of calcium influx, assume  $f(t)$  is a Dirac comb given by

$$f(t) = h \sum_{k=1}^n \delta(t - t_k) \quad (6)$$

where  $t_0, \dots, t_n$  represent stimulus times (e.g. presynaptic action potential-driven fluxes). Also, suppose that there has been a long enough interval of time so that before the first pulse, presynaptic terminal  $\text{Ca}^{2+}$  concentration is at steady state; that is, that  $c(t) = c_\infty$  for  $t \leq t_0$ . Also, suppose that, at the  $k$ th pulse time, the value of  $c(t)$  changes to  $c_k = c(t_k) + h$ , for  $k \in \{0, \dots, n\}$ . Then, at  $t = t_0$ , let  $c_0 = c_\infty + h$  be the new value of  $c$  after the jump. The dynamics for  $c$  before the next pulse are given by,

$$c(t) = c_\infty + h \exp\left(\frac{t_0 - t}{\tau_c}\right)$$

for  $t \in (t_0, t_1)$ . At  $t = t_1$  the calcium concentration changes again, to  $c_1 = c(t_1) + h$ , which becomes a new initial condition. Then, for  $t \in (t_1, t_2)$ ,

$$\begin{aligned} c(t) &= c_\infty - (c_\infty - c_1) \exp\left(\frac{t_1 - t}{\tau_c}\right), \\ &= c_\infty + h \left( \exp\left(\frac{t_1 - t}{\tau_c}\right) + \exp\left(\frac{t_0 - t}{\tau_c}\right) \right). \end{aligned}$$

In general, for  $t \in (t_{n-1}, t_n)$ ,

$$c(t) = c_\infty + h \sum_{k=0}^{n-1} \exp\left(\frac{t_k - t}{\tau_c}\right)$$

At  $t_n$  the value of  $c$  jumps to

$$c(t_n) = c_\infty + h \sum_{k=0}^n \exp\left(\frac{t_k - t_n}{\tau_c}\right)$$

If pulses are periodic, with  $d = t_{k+1} - t_k$ , for all  $k \in \{1, \dots, n\}$  then  $t_n - t_k = (n - k)d$ . In this case, the calcium concentration becomes

$$\begin{aligned} c(t_n) &= c_\infty + h \sum_{k=0}^n u^{(n-k)} = c_\infty + h \sum_{l=0}^n u^l, \\ &= c_\infty + h \left( \frac{1 - u^{n+1}}{1 - u} \right), \end{aligned}$$

where

$$u = \exp\left(-\frac{d}{\tau_c}\right). \quad (7)$$

Explicitly,

$$c(t_n) = c_\infty + h \left[ \frac{1 - \exp\left(-(n+1)\frac{d}{\tau_c}\right)}{1 - \exp\left(-\frac{d}{\tau_c}\right)} \right]. \quad (8)$$

The asymptotic behavior as  $n \rightarrow \infty$  is then

$$c_* = c_\infty + h \left( \frac{1}{1 - \exp\left(-\frac{d}{\tau_c}\right)} \right). \quad (9)$$

which written explicitly in terms of the parameters for  $\text{Ca}^{2+}$  buffering and activation of the release machinery, becomes

$$c_* = c_\infty + h \left( \frac{1}{1 - \exp\left(-d \frac{1+k_r\alpha_r(1-r)\tau_b}{\tau_b}\right)} \right). \quad (10)$$

## 1.2 $p$ dynamics and the subsystem $c - p$

The dynamics of the subsystem  $c, p$  are given by

$$\partial_t c = \frac{c_\infty - c}{\tau_c} + f(t) \quad (11)$$

$$\partial_t p = \alpha c (1 - p) - \beta p \quad (12)$$

The time constant and steady state for the  $p$  are

$$\tau_p(c) = \frac{1}{\alpha c + \beta} \quad (13)$$

$$p_\infty(c) = \frac{\alpha c}{\alpha c + \beta} = \frac{c}{c + \frac{\beta}{\alpha}} \quad (14)$$

$c$  increases as action potentials arrive, then  $\tau_p(c)$  decreases and  $p_\infty(c)$  increases. So  $p$  increases with  $c$ , but the dynamics for  $p$  become faster as that happens.

if  $\tau_c \gg \tau_p$ , then the dynamics for  $p$  are fast enough to substitute  $p$  with  $p_\infty(c)$ . Explicitly, it would be required that

$$\tau_c > \frac{1}{\alpha c + \beta}.$$

## 2 Binary classification trains of pulses.

$a_0, \dots, a_n$  pulse amplitudes. The change in amplitude during the  $k$ th interval is  $s_k = \frac{a_k - a_{k-1}}{\delta_k}$ , for  $k \in \{1, \dots, n\}$ . Similarly, the change in amplitude relative to the first pulse in a train is given by  $r_k = \frac{a_k - a_0}{\sum_{i=1}^k \delta_i}$ , for  $k \in \{1, \dots, n\}$

The heavy side function defined by  $H(x) = 1$  if  $x > 0$ ,  $0$  if  $x \leq 0$  can be used to determine whether there was an increase in the amplitudes between any two pulses. In particular,

One way to describe the history of changes during a train of postsynaptic responses is to record whether there was an increase or decrease in the amplitudes for each pair of responses using a binary code, and add all the results. To do so, let

$$b_{pp}(s_1, \dots, s_n) = \sum_{k=1}^n \frac{H(s_k)}{2^k} \quad (15)$$

for the slopes of pulses taken by pairs. Similarly, for the slopes relative to the first pulse, let

$$b_{fp}(r_1, \dots, r_n) = \sum_{k=1}^n \frac{H(r_k)}{2^k} \quad (16)$$

### 3 Short term synaptic plasticity from a whole cell approach

The model can be thought of in terms of 3 main components. First, a system capable of producing spiking dynamics, which can be 2- or 3-dimensional in case bursting or spike frequency adaptation are required. Second, a 2- or 3-dimensional system that describes the dynamics of presynaptic release, and third, at least one more equation that describes the dynamics of postsynaptic activation upon release.

$$\partial_t c = \frac{c_\infty - c}{\tau_c} + f(t) \quad (17)$$

$$\partial_t p = \alpha c (1 - p) - \beta p \quad (18)$$

$$\partial_t x = x \left( \frac{x_\infty - x}{\tau_x} \right) - x p \quad (19)$$

where  $f(c)$  represents the flux of  $\text{Ca}^{2+}$  into the terminal.