Derivation of Numerical Slowing Down Equations for Homogenous Mixture of Fuel and Absorber

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We're interested in running deterministic, multigroup transport calculations on a reactor core, to determine it's reactivity as a function of state points - control rod positions, temperature, moderator flow rate, etc. These libraries can be generated by doing neutron slowing down calculations on a fuel assembly to determine neutron spectra across each coarse energy group. These relatively computationally cheap calculations can be repeated quickly for many state points, allowing reactor designers to generate a table of multigroup cross sections for each background cross sections $\sigma_b^i = \sum_{j\neq i} N_j \sigma_j$, where each background cross section corresponds to a reactor state point of interest. Using this library, whole core transport methods can be used that treat space and angle withna higher fidelity, by using a few-energy group calculation that still preserves energy features of interest, i.e. resonance self-shielding.

Let's take a look at a very simple example: a homogenized fuel assumby made up of just two nuclides, an absorber nuclide (e.g. U238) and a moderator nuclide (e.g. H1). This simple system allows us to use a simple numerical slowing down scheme to calculate the shape of the neutron flux in the slowing down range (below the fission neutron energy, above thermal energies) as a function of the number density ratio of the absorber to the moderator N_m/N_a; or equivalently, as a function of the background cross section $\sigma_b = \Sigma_p^m/N_a$, where Σ_p is a potential scattering cross section. This then allows us to generate a multigroup cross section library for a tabulated list of backround cross sections.

For simpicity's sake, we are going to treat our fuel assumbly as infinite - Eq. 1 is effectively is an integration of a full steady-state transport equation over space and solid angle. Further, we are confining ourselves to an energy region where fission can be ignored and scatterring energy distributions are uniform. It is possible to include an escape term in our slowing down equations to account for simple geomeetries - like an infinite fuel lattice, but we will keep things simple for now.

Our system is parameterized by:

- σ_s^a, σ_s^m, the absorber and moderator (respectively) microscopic scattering cross sections

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- σ_a^a, σ_a^m, the absorber and moderator (respectively) microscopic absorption cross sections,

- α_a, α_m, where $\alpha = \frac{(A-1)^2}{(A+1)^2}$, and A_a, A_m are the respective nuclear masses of the absorber and moderator, in units of neutron masses.

- N_m/N_a, the number density ratio if moderator to absorber

We have the total cross section for a species $\sigma_t = \sigma_a + \sigma_s$ Ignoring fission and other reactions is a good approximation in our energy range of interest.

We have the lethargy (u) dependent analytical nuetron slowing down euqation in a homogenous material, for multiple nuclides indexed by j:

$$\sum_j N_j \sigma_t^j(u)\phi(u) = \sum_j \int_{u-\epsilon_j}^u N_j \sigma_s^j(u')\phi(u') \frac{e^{-(u-u')}}{1-\alpha_j} du' \qquad (1)$$

Here $\epsilon_j = \log(1/\alpha_j)$ is the maximum lethargy gain from scattering on nuclide j. We omit any source distribution over our lethargy range. We will further exam our boundary conditions later.

Explicitly writing this out for our two nuclides, absorber a and moderator m, we have:

$$\left(N_a\sigma_t^a(u)+N_m\sigma_t^m(u)\right)\phi(u) = \int_{u-\epsilon_a}^u N_a\sigma_s^a(u')\phi(u') \frac{e^{-(u-u')}}{1-\alpha_a} du' + \int_{u-\epsilon_m}^u N_m\sigma_s^m(u')\phi(u') \frac{e^{-(u-u')}}{1-\alpha_m} du'$$
$$(2)$$

Dividing through by N_a, and denoting the ratio $N_m/N_a \equiv r$:

$$\left(r\sigma_t^a(u)+\sigma_t^m(u)\right)\phi(u) = \int_{u-\epsilon_a}^u \sigma_s^a(u')\phi(u') \frac{e^{-(u-u')}}{1-\alpha_a} du' + \int_{u-\epsilon_m}^u r\sigma_s^m(u')\phi(u') \frac{e^{-(u-u')}}{1-\alpha_m} du'.$$
$$(3)$$

This is the analytical equation we want to solve. To do so, let's define an equal lethargy grid, with $I-1$ groups, numbered from 1 through I. Each group has width $\Delta u << \epsilon_a < \epsilon_m$, and lethargy is tabulated at each of the I group edges, from $u_0 = 0; E_0 = E_{\max}$ to some $u_I = \log(E_{\max}/E_{\min})$. Lethargy group i starts at $u_{i-1} = (i-1)\Delta u$ and goes to $u_i = i\Delta u$. Let's also define a lethargy group averaging operator over lethargy group i: $\int_{(i-1)\Delta u}^{i\Delta u} (*) \, du$.

Applying this operator to 3:

$$\int_{(i-1)\Delta u}^{i\Delta u} \left(r\sigma_t^a(u) + \sigma_t^m(u)\right)\phi(u) \, du = \int_{(i-1)\Delta u}^{i\Delta u} \left(\int_{u-\epsilon_a}^u \sigma_s^a(u')\phi(u') \frac{e^{-(u-u')}}{1-\alpha_a} du'\right) du$$
$$(4)$$

$$+ \int_{(i-1)\Delta u}^{i\Delta u} \left(\int_{u-\epsilon_m}^u r\sigma_s^m(u')\phi(u') \frac{e^{-(u-u')}}{1-\alpha_m} du'.\right) du$$
$$(5)$$

Next, we make the approximation that our cross sections and fluxes are constant over the groups:

$$\phi(u) = \phi_i \qquad\qquad u_{i-1} < u < u_i \qquad\qquad (6)$$

$$\sigma_k^j(u) = \sigma_{ki}^j \qquad\qquad u_{i-1} < u < u_i, \;\; \forall \sigma_k \in \{\sigma_a, \sigma_s, \sigma_t\}. \qquad (7)$$

First, we write the collision term:

$$\int_{(i-1)\Delta u}^{i\Delta u} \left(r\sigma_t^a(u) + \sigma_t^m(u)\right)\phi(u)\,du = \left(r\sigma_{ti}^a + \sigma_{ti}^m\right)\phi_i\Delta u \qquad (8)$$

The in-scattering term into group i will be a sum of contributions from lesser lethargy groups l. Contributions from the absorber will run from same-group in-scattering to the group with lower lethargy bound a distance of the maximum lethargy gain away from the group in question: $l \in \{\, l\,|\,u_{l-1} \geq u_i - \epsilon_a\}$ Scattering off the moderator will have contributions from a larger set of groups, $l \in \{\, l\,|\,u_{l-1} \geq u_i - \epsilon_m\}$. Applying the lethargy group averaging operator to the RHS of Eq. 3:

$$\int_{u_{i-1}}^{u_i} \left(\int_{u-\epsilon_a}^{u} \sigma_s^a(u')\phi(u')\frac{e^{-(u-u')}}{1-\alpha_a}du'\right) du + \int_{(i-1)\Delta u}^{i\Delta u} \left(\int_{u-\epsilon_m}^{u} r\sigma_s^m(u')\phi(u')\frac{e^{-(u-u')}}{1-\alpha_m}du'.\right) du = \qquad (9)$$

$$\int_{u_{i-1}}^{u_i} \left(\int_{u-\epsilon_a}^{u_{i-1}} \sigma_s^a(u')\phi(u')\frac{e^{-(u-u')}}{1-\alpha_a}du' + \int_{u_{i-1}}^{u} \sigma_s^a(u')\phi(u')\frac{e^{-(u-u')}}{1-\alpha_a}du'\right) du + \qquad (10)$$

$$\int_{u_{i-1}}^{u_i} \left(\int_{u-\epsilon_m}^{u_{i-1}} r\sigma_s^m(u')\phi(u')\frac{e^{-(u-u')}}{1-\alpha_m}du' + \int_{u_{i-1}}^{u} r\sigma_s^m(u')\phi(u')\frac{e^{-(u-u')}}{1-\alpha_m}du'.\right) du \qquad (11)$$

The integral on the bounds $[u_{i-1}, u_i]$ represents scattering events that don't leave group i:

$$\int_{u_{i-1}}^{u_i} \left(\int_{u_{i-1}}^{u_i} r\sigma_s^m(u')\phi(u')\frac{e^{-(u-u')}}{1-\alpha_m}du'.\right) du \qquad (12)$$

$$= \frac{r\sigma_{si}^m\phi_i}{1-\alpha_m} \int_{u_{i-1}}^{u_i} due^{-u} \int_{u_{i-1}}^{u_i} du'e^{u'} \qquad (13)$$

$$= \frac{r\sigma_{si}^m\phi_i}{1-\alpha_m}\left(\Delta u - 1 + e^{-\Delta u_i}\right) \qquad (14)$$

The absorber same-group scattering term can be derived identically:

$$\frac{\sigma_{si}^m\phi_i}{1-\alpha_a}\left(\Delta u - 1 + e^{-\Delta u_i}\right) \qquad (15)$$

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We can rewrite the bounds of the other integral, $u \in [u_{i-1} - \epsilon_j, u_{i-1}]$, as $l \in [i - \frac{\epsilon_j}{\Delta u}, i - 1]$. We must be careful here, however, if our moderator is hydrogen, it is likely that $\frac{\epsilon_m}{\Delta u} > I$. In other words, as we aren't modeling thermal or fission spectrum energies, the lethargy range of our problem is likely less than the maximum lethargy gain from our presumably hydrogenous moderator. For the purposes of this analysis, we will assume such, and thus take the lower bound of this range to be group 1, as group 1 can scatter into every other group.

For the absorber however, most likely dozens of ϵ_a's will fit into the lethargy range of interest; each group will only have in-scatter from a limited subset of lower lethargy groups. The range $l \in [i - \frac{\epsilon_a}{\Delta u}, i - 1]$ will be useful, except for in the first $\frac{\epsilon_a}{\Delta u}$ groups. We can define

$$
n = \begin{cases} i - \frac{\epsilon_a}{\Delta u} & \frac{\epsilon_a}{\Delta u} + 1 < i \\ 0 & \frac{\epsilon_a}{\Delta u} + 1 > i \end{cases} \tag{16}
$$

and use the range $l \in [n, i - 1]$.

Using this new terminology, let's see how applying the group averaging operator to the in-scatter term looks for the moderator:

$$
\int_{u_{i-1}}^{u_i} \left(\int_{u-\epsilon_m}^{u_{i-1}} r\sigma_s^m(u')\phi(u')\frac{e^{-(u-u')}}{1-\alpha_m}du'. \right) du \tag{17}
$$

$$
= \int_{u_{i-1}}^{u_i} \left(\int_0^{u_{i-1}} r\sigma_s^m(u')\phi(u')\frac{e^{-(u-u')}}{1-\alpha_m}du'. \right) du \tag{18}
$$

$$
= \sum_{l=0}^{i-1} r\sigma_{sl}^m \phi_l \int_{u_{i-1}}^{u_i} du \int_{u_{l-1}}^{u_l} du' e^{-(u-u')} \tag{19}
$$

$$
= \sum_{l=0}^{i-1} r\sigma_{sl}^m \phi_l (e^{u_l} - e^{u_{l-1}})(e^{-u_{i-1}} - e^{-u_i}) \tag{20}
$$

We can derive the term for the absorber nearly identically:

$$
\int_{u_{i-1}}^{u_i} \left(\int_{u-\epsilon_a}^{u_{i-1}} \sigma_s^a(u')\phi(u')\frac{e^{-(u-u')}}{1-\alpha_a}du'. \right) du \tag{21}
$$

$$
= \sum_{l=n}^{i-1} \sigma_{sl}^a \phi_l (e^{u_l} - e^{u_{l-1}})(e^{-u_{i-1}} - e^{-u_i}) \tag{22}
$$

Putting our scattering term together:

$$\int_{u_{i-1}}^{u_i} \left(\int_{u-\epsilon_a}^{u_{i-1}} \sigma_s^a(u')\phi(u')\frac{e^{-(u-u')}}{1-\alpha_a}du' + \int_{u_{i-1}}^{u} \sigma_s^a(u')\phi(u')\frac{e^{-(u-u')}}{1-\alpha_a}du' \right) du$$
$$+\int_{u_{i-1}}^{u_i} \left(\int_{u-\epsilon_m}^{u_{i-1}} r\sigma_s^m(u')\phi(u')\frac{e^{-(u-u')}}{1-\alpha_m}du' + \int_{u_{i-1}}^{u} r\sigma_s^m(u')\phi(u')\frac{e^{-(u-u')}}{1-\alpha_m}du'. \right) du$$
$$=\frac{r\sigma_{si}^m\phi_i}{1-\alpha_m}\left(\Delta u - 1 + e^{-\Delta u_i}\right) + \frac{\sigma_{si}^m\phi_i}{1-\alpha_a}\left(\Delta u - 1 + e^{-\Delta u_i}\right)$$
$$+\sum_{l=n}^{i-1} \sigma_{sl}^a\phi_l(e^{u_l} - e^{u_{l-1}})(e^{-u_{i-1}} - e^{-u_i})$$
$$+\sum_{l=0}^{i-1} r\sigma_{sl}^m\phi_l(e^{u_l} - e^{u_{l-1}})(e^{-u_{i-1}} - e^{-u_i})$$

$$(23)$$

We have no source term, so at this point we are ready to write the neutron balance equation in a group:

$$\left(r\sigma_{ti}^a + \sigma_{ti}^m\right)\phi_i\Delta u$$
$$=\frac{r\sigma_{si}^m\phi_i}{1-\alpha_m}\left(\Delta u - 1 + e^{-\Delta u_i}\right) + \frac{\sigma_{si}^a\phi_i}{1-\alpha_a}\left(\Delta u - 1 + e^{-\Delta u_i}\right)$$
$$+\sum_{l=n}^{i-1} \sigma_{sl}^a\phi_l(e^{u_l} - e^{u_{l-1}})(e^{-u_{i-1}} - e^{-u_i})$$
$$+\sum_{l=0}^{i-1} r\sigma_{sl}^m\phi_l(e^{u_l} - e^{u_{l-1}})(e^{-u_{i-1}} - e^{-u_i}),$$

$$(24)$$

By grouping terms of the same group i, we can see this is a lower triangular system:

$$\left[\left(r\sigma_{ti}^a + \sigma_{ti}^m\right)\Delta u - \left(\frac{r\sigma_{si}^m}{1-\alpha_m} - \frac{\sigma_{si}^a}{1-\alpha_a}\right)\left(\Delta u - 1 + e^{-\Delta u_i}\right)\right]\phi_i$$
$$=\sum_{l=n}^{i-1} \sigma_{sl}^a\phi_l(e^{u_l} - e^{u_{l-1}})(e^{-u_{i-1}} - e^{-u_i})$$
$$+\sum_{l=0}^{i-1} r\sigma_{sl}^m\phi_l(e^{u_l} - e^{u_{l-1}})(e^{-u_{i-1}} - e^{-u_i}),$$

$$(25)$$

with solution:

$$\phi_i = \frac{\sum_{l=n}^{i-1} \sigma_{sl}^a \phi_l (e^{u_l} - e^{u_{l-1}})(e^{-u_{i-1}} - e^{-u_i}) + \sum_{l=0}^{i-1} r\sigma_{sl}^m \phi_l (e^{u_l} - e^{u_{l-1}})(e^{-u_{i-1}} - e^{-u_i})}{\left(r\sigma_{ti}^a + \sigma_{ti}^m\right)\Delta u - \left(\frac{r\sigma_{si}^m}{1-\alpha_m} - \frac{\sigma_{si}^a}{1-\alpha_a}\right)\left(\Delta u - 1 + e^{-\Delta u_i}\right)}$$
$$(26)$$

Now, we need a boundary condition. We will use an asymptotic scattering source , by assuming $\sigma_s(u) = \sigma_p$ and $\phi(u) = 1$ above E_0. We only care abou the flux shape, and we're well out of the fission range, so this is a good approximation.

This means, in lethargy group 1, we have the following neutron balance equation:

$$\left(r\sigma_{t1}^a + \sigma_{t1}^m\right)\phi_1 \Delta u = \left(\frac{r\sigma_p^m}{1-\alpha_m} - \frac{\sigma_p^a}{1-\alpha_a}\right)\left[-e^{-(u-u')}\right]_{u'=u_0}^{u'=-\infty}$$
$$+ \left(\frac{r\sigma_{s1}^m}{1-\alpha_m} - \frac{\sigma_{s1}^a}{1-\alpha_a}\right)\left(\Delta u - 1 + e^{-\Delta u_1}\right)$$
$$(27)$$

If we evaluate the integral at $u = u_0$, the lower lethargy boundary of group 1, we simply have:

$$\phi_1 = \frac{\frac{r\sigma_p^m}{1-\alpha_m} - \frac{\sigma_p^a}{1-\alpha_a}}{\left(r\sigma_{t1}^a + \sigma_{t1}^m\right)\Delta u - \left(\frac{r\sigma_{s1}^m}{1-\alpha_m} - \frac{\sigma_{s1}^a}{1-\alpha_a}\right)\left(\Delta u - 1 + e^{-\Delta u_1}\right)}$$
$$(28)$$

Now we have everything we need to solve this system numerically, given a set of cross sections.

Let's look at some parameters for an example problem. Using H-1 as our moderator, we have $A_m = 0.9992$, so we have a maximum fractional energy loss in our problem of $\alpha_m = 1.600e - 7$, and maximum lethargy gain of $\epsilon_m = 15.648$. For a typical absorber, U-238, we have $A_a = 11.293$, so we have a maximum fractional energy loss in our problem of $\alpha_a = 0.7011$, and maximum lethargy gain of $\epsilon_m = 0.355$, from scattering on the absorber.

The main constraint on the number of groups is that the lethargy groups must be small enough to have constant cross sections over the resolved resonance region of the absorber. Note that the above derivation does not require equal lethargy groups. Choosing 1 eV and 20 keV as our energy bounds, we have a maximum lethargy of 9.903. Choosing $\frac{\epsilon_a}{\Delta u} = 2e4$, meaning $2e4$ lethargy groups covering the range of the maximum lethargy gain from the absorber in the problem, gives us a group width of $\Delta u = 1.776e - 5$, which gives us 557,708 groups in the problem lethargy space, if we adjust to a minimum energy of 1.0004753 eV. I'm just going to use 1 eV, because percision is overrated.