Lecture Note #7

MTH653: Advanced Numerical Analysis

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Preconditioners for discontinuous Galerkin discretizations

In the previous lecture, we analyzed the stability and accuracy of DG methods for the Poisson problem

$$-\Delta u = f$$
.

We now turn our attention to the development of preconditioners for the resulting linear system

$$Ax = b$$
.

Since the interior penalty bilinear form is symmetric and coercive, the resulting matrix A is symmetric and positive-definite. Recall that the interior penalty bilinear form is defined by

$$a(u,v) := \int_{\Omega} \nabla_h u \cdot \nabla_h v \, dx - \int_{\Gamma} \llbracket u \rrbracket \cdot \{ \nabla_h v \} \, ds - \int_{\Gamma} \{ \nabla_h u \} \cdot \llbracket v \rrbracket \, ds + \int_{\Gamma} \alpha \llbracket u \rrbracket \cdot \llbracket v \rrbracket \, ds.$$

Since we require $\alpha \gtrsim h^{-1}$, we write

$$\alpha := \eta h^{-1} \tag{1}$$

for some constant η bounded below.

We continue to use the same definition of the DG norm,

$$|||u|||^2 := ||\nabla_h u||^2_{L^2(\Omega)} + h^{-1}||[u]||^2_{L^2(\Gamma)}.$$

We recall the continuity and coercivity properties.

$$a(u, v) \lesssim ||u|| ||v||,$$

 $a(u, u) \gtrsim ||u||^2.$

First, we wish to establish some condition number estimates. A useful tool will be the following Poincaré-type inequality.

Lemma 1. Let $u \in H^1(\mathcal{T})$. Then,

$$||u||_{L^{2}(\Omega)}^{2} \lesssim ||\nabla_{h}u||_{L^{2}(\Omega)}^{2} + h^{-1}||[u]||_{L^{2}(\Gamma)}^{2} = ||u||^{2}.$$
(2)

Proof. Let ψ solve the auxiliary problem

$$-\Delta \psi = u.$$

Then, $\|\psi\|_{H^2(\Omega)} \lesssim \|u\|_{L^2(\Omega)}$. Furthermore,

$$\begin{split} \|u\|_{L^{2}(\Omega)}^{2} &= (u, -\Delta\psi) \\ &= (\nabla_{h}u, \nabla\psi) - \int_{\Gamma} \nabla\psi \cdot [\![u]\!] \, ds \\ &\leq \left(\|\nabla_{h}u\|_{L^{2}(\Omega)}^{2} + h^{-1} \|[\![u]\!]\|_{L^{2}(\Gamma)}^{2} \right)^{1/2} \left(\|\nabla\psi\|_{L^{2}(\Omega)}^{2} + h \|\partial\psi/\partial n\|_{L^{2}(\Gamma)}^{2} \right)^{1/2} \end{split}$$

A scaling argument shows that

$$\|\partial \psi/\partial n\|_{L^2(\Gamma)}^2 \lesssim h^{-1} \|\psi\|_{H^2(\Omega)}^2,$$

from which we conclude that

$$\|\nabla \psi\|_{L^{2}(\Omega)}^{2} + h\|\partial \psi/\partial n\|_{L^{2}(\Gamma)}^{2} \lesssim \|\psi\|_{H^{2}(\Omega)}^{2} \lesssim \|u\|_{L^{2}(\Omega)}^{2}.$$

Dividing through by $||u||_{L^2(\Omega)}$, we obtain the desired result.

We will also make use the following **inverse inequality** and **trace inequality**, which both follow from a scaling argument, and equivalence of finite-dimensional norms. For any $u_h \in V_h$, we have that

$$||u_h||_{H^1(\Omega)} \lesssim h^{-1} ||u_h||_{L^2(\Omega)},$$

 $||u_h||_{L^2(\partial \kappa)} \lesssim h^{-1/2} ||u_h||_{L^2(\kappa)}.$

These estimates can be combined to give the following result, which can be used to estimate the eigenvalues of the matrix A.

Proposition 1. For any $u_h \in V_h$, it holds that

$$||u_h||_{L^2(\Omega)}^2 \lesssim a(u_h, u_h) \lesssim \eta h^{-2} ||u_h||_{L^2(\Omega)}^2.$$

Proof. The lower bound follows directly from (1) and coercivity of $a(\cdot \cdot)$ in the $\| \cdot \|$ norm. The upper bound can be seen by noting

$$a(u_h, u_h) \lesssim \eta |||u_h|||^2 \qquad \text{(continuity)}$$

$$= \eta \left(||\nabla_h u||^2_{L^2(\Omega)} + h^{-1} ||[u]||^2_{L^2(\Gamma)} \right) \qquad \text{(definition of } ||| \cdot |||)$$

$$\lesssim \eta \left(h^{-2} ||u||^2_{L^2(\Omega)} + h^{-1} ||[u]||^2_{L^2(\Gamma)} \right) \qquad \text{(inverse inequality)}.$$

On a given edge e adjacent to elements κ^{\pm} with traces u_h^{\pm} ,

$$||[u_h]||_{L^2(e)}^2 \le 2\left(||u_h^-||_{L^2(e)}^2 + ||u_h^+||_{L^2(e)}^2\right)$$
 (Young's inequality)
$$\lesssim h^{-1}||u_h||_{L^2(\kappa^-)}^2 + h^{-1}||u_h||_{L^2(\kappa^+)}^2$$
 (trace inequality)

and therefore

$$\|[u]\|_{L^2(\Gamma)}^2 \lesssim h^{-1} \|u_h\|_{L^2(\Omega)}^2$$
.

From this, the upper bound follows.