

### Lecture worksheet III: coupled oscillators – finding normal mode solutions

**Key:** seek special same  $\omega$  “normal mode” solutions  $x_1 = C_1 e^{i\omega t}$   $x_2 = C_2 e^{i\omega t}$  and plug into

$$\begin{aligned} m_1 \ddot{x}_1 &= -k_1 x_1 + k_2 (x_2 - x_1) & \text{to get} & & m_1 (-\omega^2 C_1) &= -k_1 C_1 + k_2 (C_2 - C_1) \\ m_2 \ddot{x}_2 &= -k_3 x_2 - k_2 (x_2 - x_1) & & & m_2 (-\omega^2 C_2) &= -k_3 C_2 - k_2 (C_2 - C_1) \end{aligned}$$

IV a) Are these equivalent to  $\begin{pmatrix} k_1 + k_2 - m_1 \omega^2 & -k_2 \\ -k_2 & k_3 + k_2 - m_2 \omega^2 \end{pmatrix} \begin{pmatrix} C_1 \\ C_2 \end{pmatrix} = \begin{pmatrix} 0 \\ 0 \end{pmatrix}$ ? (yes or no)

b) What must be true about the matrix to have a chance of finding solutions with nonzero  $\frac{C_1}{C_2}$ ?

c) For the special case  $k_1 = k_3 = k$ ,  $m_1 = m_2 = m$   $\begin{pmatrix} k + k_2 - m\omega^2 & -k_2 \\ -k_2 & k + k_2 - m\omega^2 \end{pmatrix} \begin{pmatrix} C_1 \\ C_2 \end{pmatrix} = \begin{pmatrix} 0 \\ 0 \end{pmatrix}$

what specific values of  $\omega^2$  can give solutions with nonzero  $C_1$  and  $C_2$  here?

d) I'll call the higher value of  $\omega$  you just found  $\omega_H$ . Plug that larger  $\omega^2$  back in to find the allowed ratio  $\frac{C_2}{C_1}$  for that frequency. [Should you get same answer from both equations?]

e) If I pick a very specific set of initial conditions that makes the masses only oscillate at frequency  $\omega_H$  and in that case  $x_1(t) = A_H \cos(\omega_H t - \delta_H)$  for a specific  $A_H$  value and  $\delta_H$  value, what should  $x_2(t)$  be then in the real world

[Hint1: set  $C_1 = A_H e^{-i\delta_H}$ . If  $x_1 = C_1 e^{i\omega t}$   $x_2 = C_2 e^{i\omega t}$  are solutions, linear equations mean  $x_1 = \text{Re}[C_1 e^{i\omega t}]$   $x_2 = \text{Re}[C_2 e^{i\omega t}]$  are solutions too -- ones appropriate for REAL world.

Hint2: What is the required ratio  $\frac{C_2}{C_1}$ ? Feel free to make an easy choice for phase and then make the amplitude ratio right. Then take a real part to get a real solution for  $x_2$  for higher frequency.]

f) I'll call the lower value of  $\omega$  you found in (c)  $\omega_L$ . We plugged that lower  $\omega^2$  back in to find  $C_1 = C_2$  for THAT frequency.

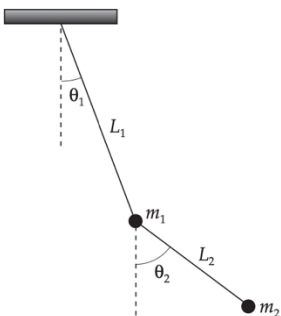
If I pick a specific set of initial conditions that makes the masses only oscillate at frequency  $\omega_L$  and in that case  $x_1(t) = A_L \cos(\omega_L t - \delta_L)$  for a specific  $A_L$  value and  $\delta_L$  value, what should  $x_2(t)$  be then in the real world?

h) To make a general solution, I try  $x_1(t) = A_L \cos(\omega_L t - \delta_L) + A_H \cos(\omega_H t - \delta_H)$ . What solution for  $x_2(t)$  has to go with it? [Plugging  $x_1(t) = A_L \cos(\omega_L t - \delta_L) + A_H \cos(\omega_H t - \delta_H)$  AND your answer for  $x_2(t)$  here into the equations at top has to work -- we found 2 pairs of solutions that work by themselves.]

i) Are there any OTHER frequencies that can have NONzero solutions for  $C_1, C_2$  above? Y or N

j) Will any trial solution where each mass moves at a single frequency DIFFERENT from the other's frequency be able to fit the equations at top of page 1 at **ALL** times? Y or N

k) For (h) answer, write down a set of possible  $\delta_L, \delta_H$  values that could fit initial conditions where the starting velocity of both masses is 0.



l) For SMALL angles this double pendulum follows  $\ddot{\theta}_1 + A\ddot{\theta}_2 + B\theta_1 = 0$  and (with  $A = \frac{m_2 L_2}{(m_1 + m_2) L_1}$   $B = \frac{g}{L_1}$   $D = \frac{L_2}{L_1}$ )  $\ddot{\theta}_1 + D\ddot{\theta}_2 + B\theta_2 = 0$

To find a normal mode where both angles oscillate at the same  $\omega$ , what matrix has to have a nonzero null space (and thus a special determinant value)?