

1. Warm-up Quiz
2. HW4 due HW5 out \nearrow δ -function potentials
3. Reading Griffiths Sec. 2.6, 2.7 \rightarrow Finite well
4. Today:
 - momentum space, x, p
 - Maybe start Finite well

Last time: Free particle ($V=0$)

Plane waves

$\psi(x) = e^{ikx}$ \rightarrow eigenstate but not a real state

$$\psi(x) = \frac{1}{\sqrt{2\pi}} \int_{-\infty}^{\infty} \phi(k) e^{ikx} dk$$

Integral over plane waves of different k

$$\phi(k) = \frac{1}{\sqrt{2\pi}} \int_{-\infty}^{\infty} \psi(x) e^{-ikx} dx$$

Integral over plane waves w/ different positions.

What does well-behaved mean

$$\int_{-\infty}^{\infty} |\psi(x)|^2 dx = \int_{-\infty}^{\infty} |\phi(k)|^2 dk = 1$$

proof:

$$|\psi(x)|^2 = \left| \frac{1}{\sqrt{2\pi}} \int_{-\infty}^{\infty} \phi(k) e^{ikx} dk \right|^2$$

$$= \frac{1}{2\pi} \int_{-\infty}^{\infty} \phi^*(k') e^{-ik'x} dk' \int_{-\infty}^{\infty} \phi(k) e^{ikx} dk$$

$$= \frac{1}{2\pi} \iint \phi^*(k') \phi(k) e^{i(k-k')x} dk dk'$$

Integrate both sides over real space

$$\int_{-\infty}^{\infty} |\psi(x)|^2 dx = \frac{1}{2\pi} \iint \phi^*(k') \phi(k) \underbrace{\int_{-\infty}^{\infty} e^{i(k-k')x} dx}_{2\pi \delta(k-k')} dk dk'$$

See k&W Ch. 5

use

$$\int F(k) \delta(k-k') dk' = F(k)$$

$$= \int |\phi(k)|^2 dk \quad QED$$

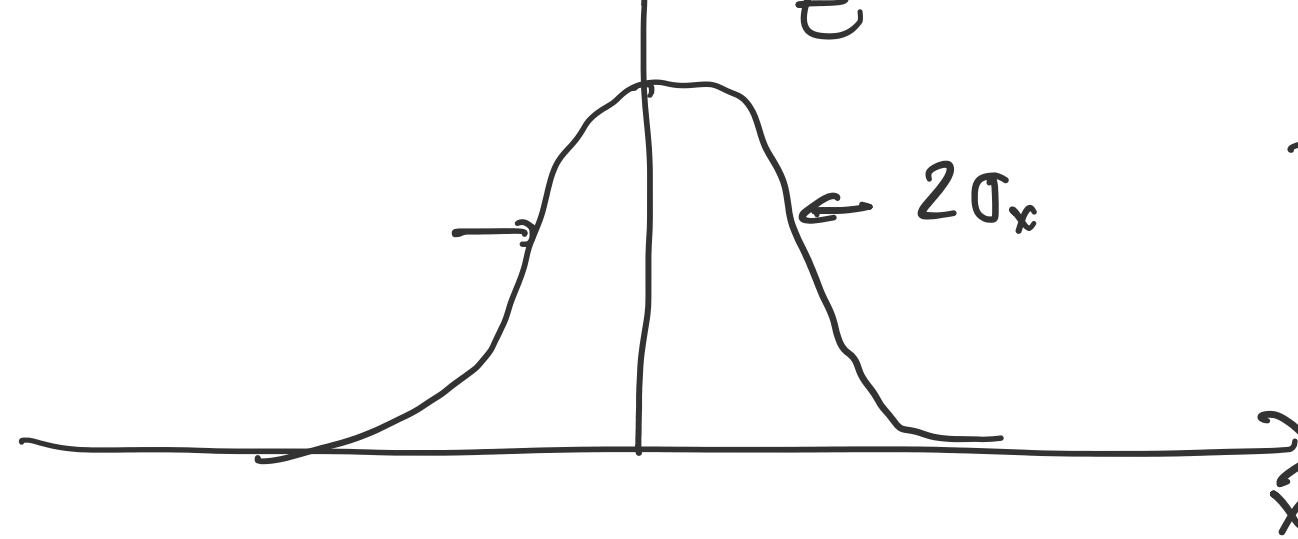
Aside: δ -function

$$\delta(x) = \begin{cases} 0 & x \neq 0 \\ \infty & x = 0 \end{cases}$$

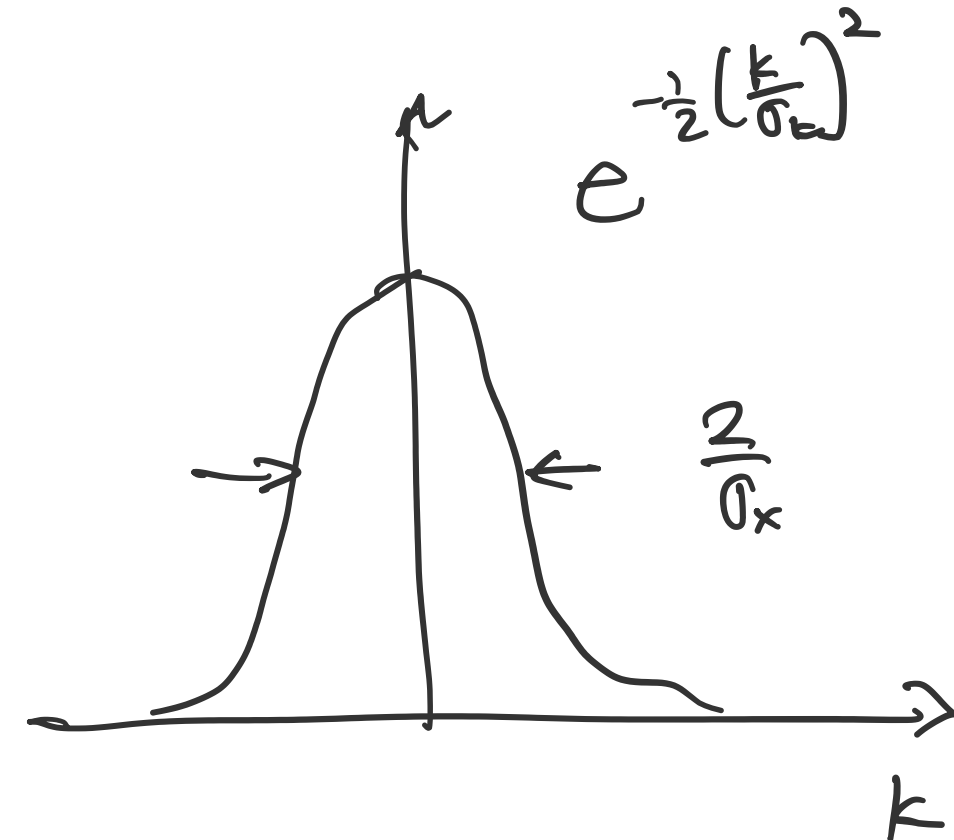
$$\int \delta(x) dx = 1$$

Example

Gaussian
 $e^{-\frac{1}{2}(\frac{x}{\sigma_x})^2}$



FT \rightarrow



- Note that $\sigma_x \sigma_k = 1$
- in general $\sigma_x \sigma_k \geq 1$ for well-behaved functions
- If a function is well-localized in one space it must be delocalized in the reciprocal space.

Time evolution of a wave packet $\psi(x, t=0)$
Follows Ex. 2.6 in Gr.

① Find eigenstates \rightarrow Free particle \rightarrow plane waves

② Fourier Transform into k -space $\phi(k) = \frac{1}{\sqrt{2\pi}} \int_{-\infty}^{\infty} \psi(x, 0) e^{-ikx} dx$

③ Transform back w/ Time dependence added

$$\psi(x, t) = \frac{1}{\sqrt{2\pi}} \int_{-\infty}^{\infty} \phi(k) e^{i(kx - \omega t)} dk$$

were $\omega = \frac{\hbar k^2}{2m}$

Operators in k -space

$$\hat{x} \text{ in } x\text{-space} \rightarrow x$$

$$\hat{p} \text{ in } x\text{-space} \rightarrow -i\hbar \frac{\partial}{\partial x}$$

$$\hat{p} \psi(x) = p \cdot \psi(x)$$

For plane waves

$$\hat{p} \psi(x) = -i\hbar \frac{\partial}{\partial x} e^{ikx} = \hbar k e^{ikx} = \underbrace{\hbar k}_p e^{ikx}$$

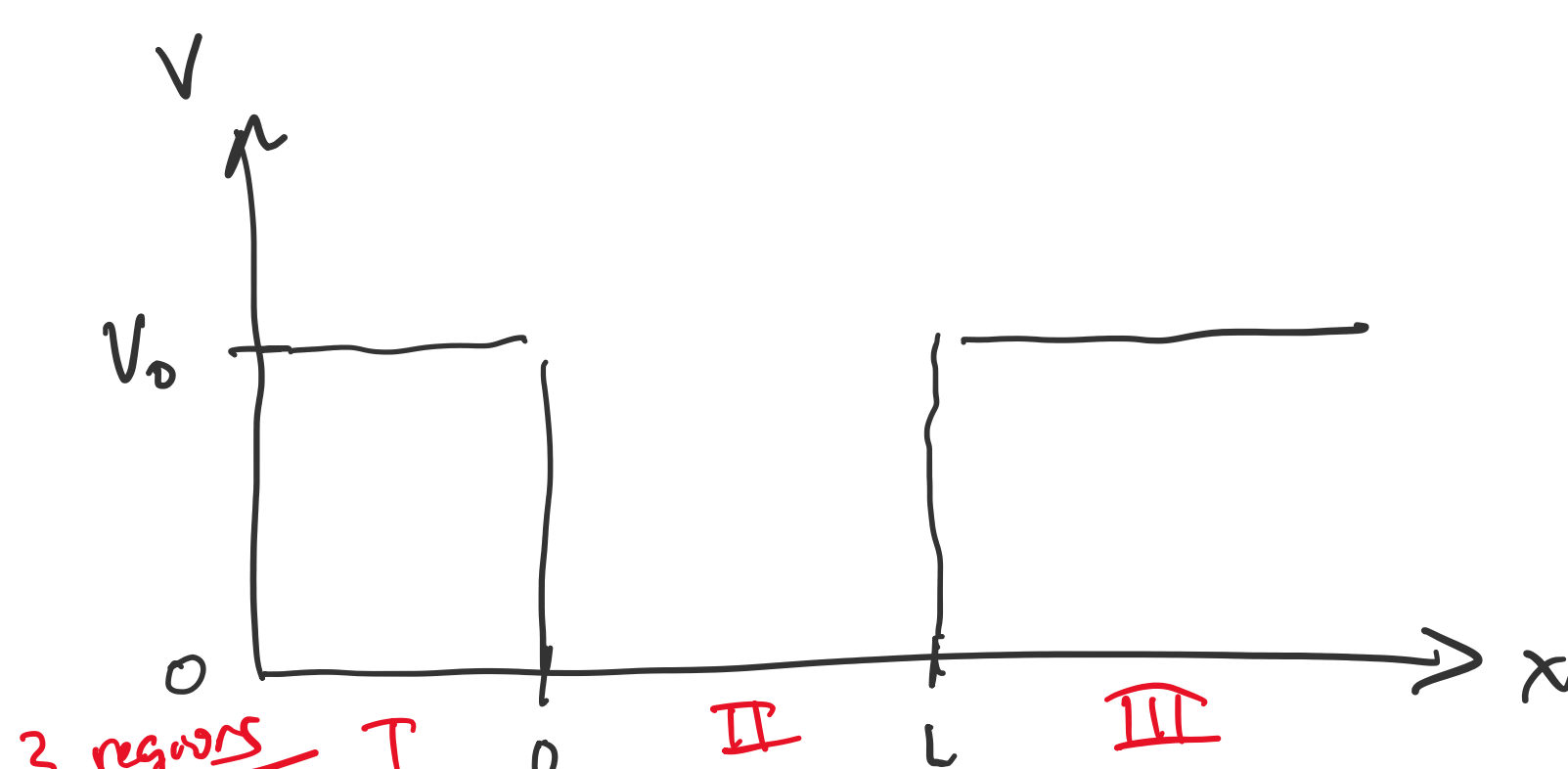
Now in k -space

$$\boxed{\hat{p} \rightarrow \hbar k}$$

$\hat{x} \rightarrow$ looking for $\hat{x} \phi(k) = x_0 \phi(k)$

so $\frac{\partial}{\partial k} \phi = -i x_0 e^{ikx_0} \Rightarrow \boxed{\hat{x} \rightarrow i \frac{\partial}{\partial k}}$

Next up: Finite Quantum Well in 1-D



$$V(x) = \begin{cases} V_0 & x < 0, x > L \\ 0 & 0 < x < L \end{cases}$$

- Expect ψ_n that look similar to the well
- ψ_n may extend out of the well.

Solve $-\frac{\hbar^2}{2m} \frac{d^2 \psi}{dx^2} + V(x) \psi = E \psi$

Start by looking for bound states ($E < V_0, E > 0$)
Later look @ "scattering states" ($E > V_0$)

Region II $V=0$ $E > 0$ $A' e^{ikx} + B' e^{-ikx}$

$$\frac{d^2 \psi}{dx^2} = \underbrace{-\frac{2mE}{\hbar^2}}_{-k^2} \psi \Rightarrow A \sin(kx) + B \cos(kx) = \psi$$

$$k = \frac{\sqrt{2mE}}{\hbar}$$

Region I, III ($x < 0, x > L$) $E > 0, E < V_0$

$$\frac{d^2 \psi}{dx^2} = \underbrace{+\frac{2m}{\hbar^2}}_{-k^2} (V_0 - E) \psi$$