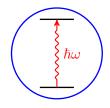
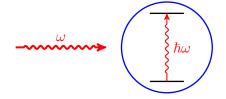
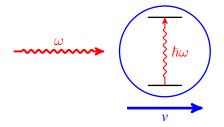
Lamb-Dicke regime/approximation

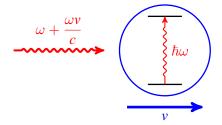
Yichao Yu

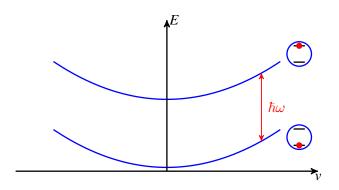
Journal Club

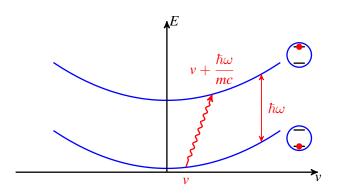


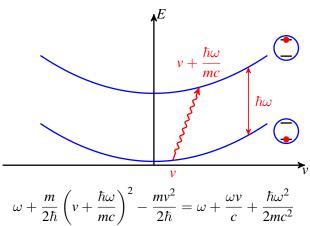






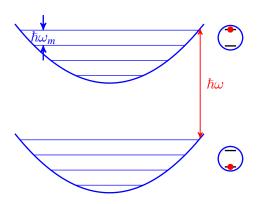




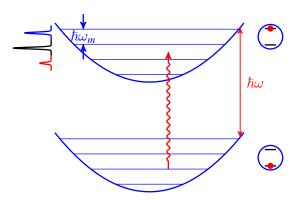


$$\omega + \frac{m}{2\hbar} \left(v + \frac{\hbar \omega}{mc} \right)^2 - \frac{mv^2}{2\hbar} = \omega + \frac{\omega v}{c} + \frac{\hbar \omega^2}{2mc^2}$$

Sideband

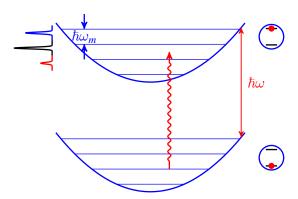


Sideband



Frequency: $\omega + n\omega_m$

Sideband



Frequency: $\omega + n\omega_m$

Strength: $\langle n|e^{ik\hat{x}}|n+\Delta n\rangle$

$$\langle n|\mathrm{e}^{\mathrm{i}k\hat{x}}|n+\Delta n\rangle$$

$$\langle n|\mathrm{e}^{\mathrm{i}k\hat{x}}|n+\Delta n\rangle$$

$$\hat{x} = \sqrt{\frac{\hbar}{2m\omega}} \left(a + a^{\dagger} \right) = z_0 \left(a + a^{\dagger} \right)$$

$$\langle n|e^{ik\hat{x}}|n+\Delta n\rangle$$
 $\hat{x}=\sqrt{\frac{\hbar}{2m\omega}}\left(a+a^{\dagger}\right)=z_{0}\left(a+a^{\dagger}\right)$ $k\hat{x}=\eta\left(a+a^{\dagger}\right)$ $\eta\equiv kz_{0}=k\sqrt{\frac{\hbar}{2m\omega}}$

$$\langle n|e^{ik\hat{x}}|n+\Delta n\rangle$$
 $\hat{x}=\sqrt{\frac{\hbar}{2m\omega}}\Big(a+a^{\dagger}\Big)=z_0\Big(a+a^{\dagger}\Big)$
 $\eta=\frac{2\pi z_0}{\lambda}$
 $k\hat{x}=\eta\Big(a+a^{\dagger}\Big)$
 $\eta=\sqrt{\frac{\omega_R}{\omega_m}}$
 $\eta\equiv kz_0=k\sqrt{\frac{\hbar}{2m\omega}}$

Sideband strength

$$\langle n|e^{ik\hat{x}}|n+\Delta n\rangle$$

$$=e^{-\eta^2/2}\eta^{\Delta n}\sqrt{\frac{n_-!}{n_+!}}L_{n_-}^{\Delta n}(\eta^2)$$

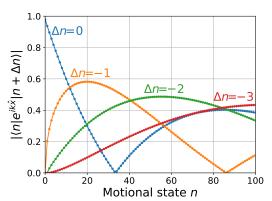
$$n_- \equiv \min(n, n+\Delta n), \quad n_+ \equiv \max(n, n+\Delta n)$$

Sideband strength

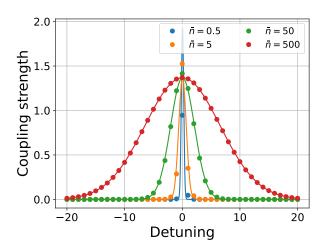
$$\langle n|e^{ik\hat{x}}|n+\Delta n\rangle$$

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Sideband spectrum



$$e^{i\eta(a+a^{\dagger})} = 1 + i\eta(a+a^{\dagger}) - \frac{\eta^2}{2}(a+a^{\dagger})^2 + \mathcal{O}(\eta^3)$$

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When $\eta \ll 1$: $e^{i\eta(a+a^{\dagger})} \approx 1 + i\eta(a+a^{\dagger})$

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When $\eta \ll 1$: $e^{i\eta(a+a^{\dagger})} \approx 1 + i\eta(a+a^{\dagger})$

To the first order,

$$\langle n|e^{i\eta(a+a^{\dagger})}|n\rangle \approx 1$$
$$\langle n+1|e^{i\eta(a+a^{\dagger})}|n\rangle \approx i\eta\sqrt{n}$$
$$\langle n-1|e^{i\eta(a+a^{\dagger})}|n\rangle \approx i\eta\sqrt{n+1}$$

$$e^{i\eta(a+a^{\dagger})} = 1 + i\eta(a+a^{\dagger}) - \frac{\eta^2}{2}(a+a^{\dagger})^2 + \mathcal{O}(\eta^3)$$
When $\eta \ll 1$: $e^{i\eta(a+a^{\dagger})} \approx 1 + i\eta(a+a^{\dagger})$

To the second order,

$$\langle n|e^{i\eta(a+a^{\dagger})}|n\rangle \approx 1 - \frac{\eta^2(2n+1)}{2}$$

$$\langle n|e^{\mathrm{i}\eta\left(a+a^{\dagger}\right)}|n\rangle\approx 1-\frac{\eta^{2}(2n+1)}{2}$$

$$\langle n|e^{\mathrm{i}\eta\left(a+a^{\dagger}\right)}|n\rangle\approx 1-\frac{\eta^{2}(2n+1)}{2}$$

Wavefunction spread,

$$\langle n|x^2|n\rangle = z_0^2(2n+1)$$

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Wavefunction spread,

$$\langle n|x^2|n\rangle = z_0^2(2n+1)$$

$$\eta_{\it eff} \equiv \eta \sqrt{2n+1}$$