Response of single bit rotation to amplitude noise

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1 Goal

Try to derive the pulse sequence to achieve robustness against (DC) amplitude noise in a brute force/generic way.

2 Pulse sequence with three rotations

1. $\sigma_1 \equiv \sigma_x \cos\theta_1 + \sigma_y \sin\theta_1$ by angle ψ_1

2. $\sigma_2 \equiv \sigma_x$ by angle ψ_2

3. $\sigma_3 \equiv \sigma_x \cos\theta_3 + \sigma_y \sin\theta_3$ by angle ψ_3

Full rotation.

$$U = \exp\left(i\psi_1\sigma_1/2\right)\exp\left(i\psi_2\sigma_2/2\right)\exp\left(i\psi_3\sigma_3/2\right)$$

When ψ_1, ψ_2, ψ_3 are changed by a small fraction 2ε to $\psi_1(1+2\varepsilon)$, $\psi_2(1+2\varepsilon)$, $\psi_3(1+2\varepsilon)$ respectively.

The derivative of U w.r.t. ε is.

$$\begin{aligned}
\frac{\partial U}{\partial\varepsilon} =& i\psi_1 \exp\left(i\psi_1\sigma_1/2\right)\sigma_1 \exp\left(i\psi_2\sigma_2/2\right)\exp\left(i\psi_3\sigma_3/2\right)+ \\
& i\psi_2 \exp\left(i\psi_1\sigma_1/2\right)\sigma_2 \exp\left(i\psi_2\sigma_2/2\right)\exp\left(i\psi_3\sigma_3/2\right)+ \\
& i\psi_3 \exp\left(i\psi_1\sigma_1/2\right)\exp\left(i\psi_2\sigma_2/2\right)\sigma_3 \exp\left(i\psi_3\sigma_3/2\right)
\end{aligned}$$

Simplifying it by changing coordinates

$$\begin{aligned}
& \exp\left(-i\psi_1\sigma_1/2\right)\frac{\partial U}{\partial\varepsilon}\exp\left(-i\psi_2\sigma_2/2\right)\exp\left(-i\psi_3\sigma_3/2\right) \\
=& i\psi_1\sigma_1 + i\psi_2\sigma_2 + i\psi_3 \exp\left(i\psi_2\sigma_2/2\right)\sigma_3 \exp\left(-i\psi_2\sigma_1/2\right) \\
=& i\psi_1\sigma_1 + i\psi_2\sigma_2 + i\psi_3(\sigma_3 \cos\psi_2 - \sigma_z \sin\theta_3 \sin\psi_2 + \sigma_2 \cos\theta_3(1-\cos\psi_2))
\end{aligned}$$

For this to be 0, the coefficient for σ_z has to be 0. This means $\psi_3 = 0$ (two rotation only), $\sin\theta_3 = 0$ (third rotation on the same axis as the second one, so effectivel no third rotation), or $\sin\psi_2 = 0$. For a true three rotation sequence, the only option is $\sin\psi_2 = 0$, which leave us with two possibilities $\cos\psi_2 = \pm 1$.

For $\cos\psi_2 = 1$,

$$\begin{aligned}
& \exp\left(-i\psi_1\sigma_1/2\right)\frac{\partial U}{\partial\varepsilon}\exp\left(-i\psi_2\sigma_2/2\right)\exp\left(-i\psi_3\sigma_3/2\right) \\
=& i\psi_1\sigma_1 + i\psi_2\sigma_2 + i\psi_3\sigma_3
\end{aligned}$$

For $\cos\psi_2 = -1$,

$$\exp\left(-i\psi_1\sigma_1/2\right)\frac{\partial U}{\partial\varepsilon}\exp\left(-i\psi_2\sigma_2/2\right)\exp\left(-i\psi_3\sigma_3/2\right)$$
$$=i\psi_1\sigma_1 + i\psi_2\sigma_2 + i\psi_3(-\sigma_3 + 2\sigma_2\cos\theta_3)$$
$$=i\psi_1\sigma_1 + i\psi_2\sigma_2 + i\psi_3(\sigma_x\cos\theta_3 - \sigma_y\sin\theta_3)$$

For both of these cases, the robust condition is for the sum of the three Pauli vectors to be 0 (in the first case these are the Pauli vector that corresponds to the rotation and in the second case the last one is replaced with its reflection about the x axis). The robust condition for the first one is also the same as the cross talk calculation condition which allows the pulse sequence to cancel cross talk and overrotation error to the first order at the same time.

As for the actual rotation, for the first case, the second pulse is a full rotation so the equivalent final operation is simply the combination of the first and the third one. Unless one of them is a π or 2π rotation, the combined result will have a component that is a z rotation. This operation cannot be removed by changing the coordinate system (i.e. by making the second rotation along another axis rather than the x axis). This could be cancelled out for single-qubit gate by adjusting the phase but may not be as easy for two-qubit gates.

For the second case,

$$U = \exp\left(i\psi_1\sigma_1/2\right)\sigma_x\exp\left(i\psi_3\sigma_3/2\right)$$
$$=\left(\cos\frac{\psi_1}{2} + i\sigma_x\cos\theta_1\sin\frac{\psi_1}{2} + i\sigma_y\sin\theta_1\sin\frac{\psi_1}{2}\right)\sigma_x\left(\cos\frac{\psi_3}{2} + i\sigma_x\cos\theta_3\sin\frac{\psi_3}{2} + i\sigma_y\sin\theta_3\sin\frac{\psi_3}{2}\right)$$
$$=\sigma_x\cos\frac{\psi_1}{2}\cos\frac{\psi_3}{2} + i\cos\frac{\psi_1}{2}\cos\theta_3\sin\frac{\psi_3}{2} - \sigma_z\cos\frac{\psi_1}{2}\sin\theta_3\sin\frac{\psi_3}{2}$$
$$+ i\cos\theta_1\sin\frac{\psi_1}{2}\cos\frac{\psi_3}{2} - \sigma_x\cos\theta_1\cos\theta_3\sin\frac{\psi_1}{2}\sin\frac{\psi_3}{2} - \sigma_y\cos\theta_1\sin\theta_3\sin\frac{\psi_1}{2}\sin\frac{\psi_3}{2}$$
$$+ \sigma_z\sin\theta_1\sin\frac{\psi_1}{2}\cos\frac{\psi_3}{2} - \sigma_y\sin\theta_1\cos\theta_3\sin\frac{\psi_1}{2}\sin\frac{\psi_3}{2} + \sigma_x\sin\theta_1\sin\theta_3\sin\frac{\psi_1}{2}\sin\frac{\psi_3}{2}$$
$$=i\cos\theta_3\cos\frac{\psi_1}{2}\sin\frac{\psi_3}{2} + i\cos\theta_1\sin\frac{\psi_1}{2}\cos\frac{\psi_3}{2}$$
$$+ \sigma_x\cos\frac{\psi_1}{2}\cos\frac{\psi_3}{2} - \sigma_x\cos\theta_1\cos\theta_3\sin\frac{\psi_1}{2}\sin\frac{\psi_3}{2} + \sigma_x\sin\theta_1\sin\theta_3\sin\frac{\psi_1}{2}\sin\frac{\psi_3}{2}$$
$$- \sigma_y\cos\theta_1\sin\theta_3\sin\frac{\psi_1}{2}\sin\frac{\psi_3}{2} - \sigma_y\sin\theta_1\cos\theta_3\sin\frac{\psi_1}{2}\sin\frac{\psi_3}{2}$$
$$- \sigma_z\sin\theta_3\cos\frac{\psi_1}{2}\sin\frac{\psi_3}{2} + \sigma_z\sin\theta_1\sin\frac{\psi_1}{2}\cos\frac{\psi_3}{2}$$

(Seems possible to achieve pure $x-y$ plane rotation though not sure if there's a way to express the condition in a concise way.)

2.1   Maximizing the final rotation angle

Given the same total pulse lens (i.e. total rotation angle for all the pulses), we'd like to maximize the total effective rotation angle to make the pulse sequence the most efficient.

Fig. 1 shows the result of numerically optimizing the pulse sequnece by computing the maximum achievable rotation angle for a fixed total rotation angle (time). Two cases that corresponds to the minimum length second pulse are shown i.e. $\psi_2 = \pi$ and $\psi_2 = 2\pi$. The dashed line shows the comparison to the corresponding SK1(-like) pulse sequence. Since the total pulse sequence length
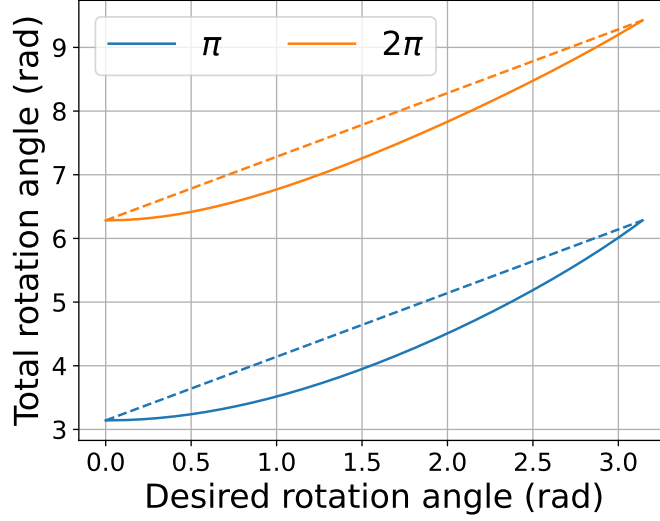
Figure 1: Minimum angle required to implement a total rotation angle. Rotation around z is ignored in the final rotation angle and only rotation around an axis in the xy plane is shown. The exact axis is also omitted from the plot since it can be changed to any desired direction by changing the axis of all the rotations at the same time. The legend shows the length of the middle pulse which is chosen to be π for the $\cos\psi_2 = -1$ case and 2π for the $\cos\psi_2 = 1$. The dashed line shows the corresponding total pulse length for the SK1(-like) pulse sequence where the first (or the last) pulse has the desired rotation angle.

must be at least twice the length of the second pulse, the improvement on top of the SK1 sequence is not very significant and is at most about 10%.

A Decompose an arbitrary single qubit rotation into an xy rotation and a z rotation

Arbitrary single qubit rotation,

$$U = a_0 + x_0\sigma_x + y_0\sigma_y + z_0\sigma_z$$

xy rotation followed by z rotation with global phase

$$
\begin{aligned}
U &= U_{xy}U_z \mathrm{e}^{\mathrm{i}\phi}\\
&= (a_1 + x_1\sigma_x + y_1\sigma_y)(\cos\theta + \mathrm{i}\sin\theta\sigma_z)\mathrm{e}^{\mathrm{i}\phi}\\
&= (a_1\cos\theta + \mathrm{i}a_1\sin\theta\sigma_z + \cos\theta x_1\sigma_x + \mathrm{i}\sin\theta x_1\sigma_x\sigma_z + \cos\theta y_1\sigma_y + \mathrm{i}\sin\theta y_1\sigma_y\sigma_z)\mathrm{e}^{\mathrm{i}\phi}\\
&= (a_1\cos\theta + (\cos\theta x_1 - \sin\theta y_1)\sigma_x + (\cos\theta y_1 + \sin\theta x_1)\sigma_y + \mathrm{i}a_1\sin\theta\sigma_z)\mathrm{e}^{\mathrm{i}\phi}
\end{aligned}
$$

where a_1, x_1, y_1 are all real numbers.

$$
\begin{aligned}
a_0\mathrm{e}^{-\mathrm{i}\phi} &= a_1\cos\theta\\
x_0\mathrm{e}^{-\mathrm{i}\phi} &= x_1\cos\theta - y_1\sin\theta\\
y_0\mathrm{e}^{-\mathrm{i}\phi} &= y_1\cos\theta + x_1\sin\theta\\
z_0\mathrm{e}^{-\mathrm{i}\phi} &= \mathrm{i}a_1\sin\theta
\end{aligned}
$$

so we have

$$
\begin{aligned}
a_1 &= \sqrt{|a_0|^2 + |z_0|^2}\\
\cos\theta &= \frac{|a_0|}{\sqrt{|a_0|^2 + |z_0|^2}}\\
\sin\theta &= \frac{|z_0|}{\sqrt{|a_0|^2 + |z_0|^2}}\\
x_1 &= \frac{a_0^* x_0 - \mathrm{i}z_0 y_0^*}{\sqrt{|a_0|^2 + |z_0|^2}}\\
y_1 &= \frac{a_0^* y_0 + \mathrm{i}z_0 x_0^*}{\sqrt{|a_0|^2 + |z_0|^2}}
\end{aligned}
$$