

Reduced basis model order reduction of dissipatively perturbed constrained Hamiltonian systems

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1 Introduction

In this paper, we consider the model order reduction (MOR) of the differential algebraic equations of the form

$$\left. \begin{aligned} \dot{\mathbf{q}} &= \nabla_{\mathbf{p}} H(\mathbf{q}, \mathbf{p}), \\ \dot{\mathbf{p}} &= -\nabla_{\mathbf{q}} H(\mathbf{q}, \mathbf{p}) - \gamma \mathbf{p} - \mathbf{G}(\mathbf{q})^\top \boldsymbol{\lambda}, \\ \mathbf{0} &= \mathbf{g}(\mathbf{q}). \end{aligned} \right\} \quad (1)$$

Here $\mathbf{q}(t), \mathbf{p}(t) : \mathbb{R} \rightarrow \mathbb{R}^d$ are referred to as the position and conjugate momenta, respectively, and constitute the solution of the system. The smooth function $H : \mathbb{R}^{2d} \rightarrow \mathbb{R}$ is the Hamiltonian, $\gamma > 0$ is the dissipation parameter, $\mathbf{g} \in \mathbb{R}^m$ are the holonomic constraints (depend only on \mathbf{q}), \mathbf{G} is the Jacobian matrix of \mathbf{g} , i.e., $\mathbf{G}(\mathbf{q}) = \mathbf{g}_{\mathbf{q}}(\mathbf{q})$, and $\boldsymbol{\lambda} \in \mathbb{R}^m$ are the Lagrange multipliers. By differentiating the constraint $\mathbf{g}(\mathbf{q}) = \mathbf{0}$, we see that the solution of eq. (1) evolves on the manifold defined by the phase-space

$$\mathfrak{P} = \{(\mathbf{q}, \mathbf{p}) \in \mathbb{R}^{2d} : \mathbf{g}(\mathbf{q}) = \mathbf{0}, \mathbf{G}(\mathbf{q})\nabla_{\mathbf{p}} H(\mathbf{q}, \mathbf{p}) = \mathbf{0}\}.$$

Here and henceforth we have dropped the explicit dependence of variables on t to keep the notation concise. The constraints defining \mathfrak{P} are also known as the *weak invariants* of the system. Equation (1) can be thought of as a dissipative perturbation of a constrained Hamiltonian system. Constrained Hamiltonian systems find application in many areas including mechanics, chemical processes, and molecular dynamics. Constrained Hamiltonian system is an important special case of eq. (1) when $\gamma = 0$. In this special case, the Hamiltonian H and symplecticness are constants of motion of the constrained Hamiltonian system. The Hamiltonian system is obtained from eq. (1) when $\mathbf{g} \equiv \mathbf{0}$, $\gamma = 0$. Structure-preserving techniques, preserving the symplecticness through MOR, have been developed in [1, 3, 6].

A generalization of the Hamiltonian system is port-Hamiltonian system, which is the aggregate system model obtained when the core dynamics of the subsystem components are described by variational principles. MOR techniques preserving the port-Hamiltonian structure of the port-Hamiltonian systems are developed in [4].

Another relevant DAE that finds a lot of mention in the literature is the *descriptor system*

$$\left. \begin{aligned} \mathbf{E}\dot{\mathbf{x}}(t) &= \mathbf{A}\mathbf{x}(t) + \mathbf{B}\mathbf{u}(t), \\ \mathbf{y}(t) &= \mathbf{C}\mathbf{x}(t) + \mathbf{D}\mathbf{u}(t), \end{aligned} \right\}$$

where $\mathbf{x} \in \mathbb{R}^n$, $\mathbf{u} \in \mathbb{R}^m$, $\mathbf{y} \in \mathbb{R}^p$ are states, inputs, outputs respectively. Matrix $\mathbf{E} \in \mathbb{R}^{n \times n}$ is a singular matrix, and $\mathbf{A} \in \mathbb{R}^{n \times n}$, $\mathbf{B} \in \mathbb{R}^{n \times m}$, $\mathbf{C} \in \mathbb{R}^{p \times n}$, $\mathbf{D} \in \mathbb{R}^{p \times m}$. The descriptor systems have been reduced using interpolatory projection methods in [5] and using balanced truncation in [2, 7].

The research on MOR of DAEs so far has mainly focused on the reduced basis MOR of the port-Hamiltonian systems and on interpolatory and balanced truncation MOR of the descriptor systems. In this paper, we take a reduced-basis approach to MOR of the nonlinear DAE (1). Our objective is twofold:

- To reduce eq. (1) using the reduced basis method such that the conformal symplecticness of the equation is preserved through MOR
- To accelerate the simulation of the reduced order model (ROM) using a nonlinearity reduction technique, such as DEIM, by structure-preserving hyper-reduction of the conservative and constraining forces in eq. (1)

In the end, we conduct numerical experiments on the ROM using a conformal symplectic numerical method to solve it.

2 MOR of eq. (1)

We begin by addressing our first objective – structure-preserving reduction (SPR) of eq. (1) using a reduced basis method. To this end, let us consider the following generalization of eq. (1).

$$\begin{cases} \dot{\mathbf{z}} = \mathbf{J} \nabla_{\mathbf{z}} \tilde{H}(\mathbf{z}) - \gamma \mathbf{z}, \\ \mathbf{0} = \mathbf{g}(\mathbf{q}). \end{cases} \quad (2)$$

Here $\mathbf{z}(t) : \mathbb{R} \rightarrow \mathbb{R}^{2d}$ is the solution of the system, and the structure matrix \mathbf{J} is skew-symmetric and non-singular. The smooth function $\tilde{H} = H + \mathbf{g}^\top \boldsymbol{\lambda}$ is the augmented Hamiltonian, $\gamma > 0$ is the dissipation parameter, and $\mathbf{g} \in \mathbb{R}^m$ are the constraints. The separable system (1) can be deduced from eq. (2) replacing γ by $\gamma/2$, $H(\mathbf{z})$ by $H(\mathbf{q}, \mathbf{p}) + \frac{\gamma}{2} \mathbf{z}$, and the structure matrix \mathbf{J} by

$$\mathbf{J} = \begin{bmatrix} \mathbf{0} & \mathbf{I} \\ -\mathbf{I} & \mathbf{0} \end{bmatrix}.$$

The system (2) is conformal symplectic since

$$d\mathbf{z} \wedge \mathbf{J}^{-1} d\mathbf{z} = -2\gamma d\mathbf{z} \wedge \mathbf{J}^{-1} d\mathbf{z},$$

where $d\mathbf{z}$ is the solution of the variational equation associated with (2). We refer to eq. (2) as *constrained perturbed Hamiltonian system* (CPHS).

Let us assume that we have obtained the projection bases \mathbf{V}_r , \mathbf{W}_r such that

$$\mathbf{z}(t) \approx \mathbf{V}_r \mathbf{z}_r(t), \quad \nabla_{\mathbf{z}} \tilde{H}(\mathbf{z}(t)) \approx \mathbf{W}_r \mathbf{f}_r(t),$$

respectively, for some $\mathbf{f}_r \in \mathbb{R}^r$. Assuming $\mathbf{V}_r^\top \mathbf{W}_r = \mathbf{I}$, it follows that

$$\begin{aligned} \nabla_{\mathbf{z}} \tilde{H}(\mathbf{V}_r \mathbf{z}_r) &\approx \nabla_{\mathbf{z}} \tilde{H}(\mathbf{z}) \\ &\approx \mathbf{W}_r \mathbf{f}_r(t) \\ &= \mathbf{W}_r \mathbf{V}_r^\top \mathbf{W}_r \mathbf{f}_r(t) \\ &\approx \mathbf{W}_r \mathbf{V}_r^\top \nabla_{\mathbf{z}} \tilde{H}(\mathbf{V}_r \mathbf{z}_r) \\ &= \mathbf{W}_r \nabla_{\mathbf{z}_r} \tilde{H}_r(\mathbf{z}_r), \end{aligned}$$

where $\tilde{H}_r(\mathbf{z}_r) = \tilde{H}(\mathbf{V}_r \mathbf{z}_r)$.

Substituting $\mathbf{z}(t) \approx \mathbf{V}_r \mathbf{z}_r(t)$, $\nabla_{\mathbf{z}} \tilde{H}(\mathbf{V}_r \mathbf{z}_r) \approx \mathbf{W}_r \nabla_{\mathbf{z}_r} \tilde{H}_r(\mathbf{z}_r)$ in eq. (2), one gets the reduced system

$$\begin{cases} \dot{\mathbf{z}}_r = \mathbf{J}_r \nabla_{\mathbf{z}_r} \tilde{H}_r(\mathbf{z}_r) - \gamma \mathbf{z}_r, \\ \mathbf{0} = \mathbf{g}(\mathbf{q}_r), \end{cases} \quad (3)$$

where $\mathbf{J}_r = \mathbf{W}_r^\top \mathbf{J} \mathbf{W}_r$ is skew-symmetric, and \mathbf{q}_r denotes the first half of the vector $\mathbf{V}_r \mathbf{z}_r$. The reduced system (3) is again conformal symplectic because the Jacobian $\partial_{\mathbf{z}_r} \tilde{H}_r$ is symmetric.

We further reduce the nonlinearities in eq. (3) with the discrete empirical interpolation method (DEIM). To this end, let us assume that the Hamiltonian H can be expressed as

$$H = \frac{1}{2} \mathbf{z}^\top \mathbf{Q} \mathbf{z} + h(\mathbf{z})$$

with \mathbf{Q} being a positive-definite matrix and h being the nonlinear part of H . Let us define $\tilde{h} = h + \mathbf{g}^\top \boldsymbol{\lambda}$, and assume that $\nabla_{\mathbf{z}} \tilde{h}(\mathbf{z}) \approx \mathbf{U}_l \mathbf{k}_l(t)$ and that $\mathbf{V}_r^\top \mathbf{Q} \mathbf{V}_r = \mathbf{I}$, $\mathbf{U}_l^\top \mathbf{Q} \mathbf{U}_l = \mathbf{I}$. Let us further assume the DEIM projection basis $\mathbb{P} = \mathbf{U}_l (\mathbf{E}^\top \mathbf{U}_l)^{-1} \mathbf{E}^\top$ for the hyper-reduction of \tilde{h} . Then the hyper-reduced system reads

$$\begin{cases} \dot{\hat{\mathbf{z}}}_r = \mathbf{J}_r \nabla_{\hat{\mathbf{z}}_r} \hat{H}_r(\hat{\mathbf{z}}_r) - \gamma \hat{\mathbf{z}}_r \\ \mathbf{0} = \mathbf{g}(\hat{\mathbf{q}}_r). \end{cases} \quad (4)$$

Here $\hat{H}_r = \frac{1}{2} \hat{\mathbf{z}}_r^\top \mathbf{Q} \hat{\mathbf{z}}_r + \mathbb{P} \tilde{h}(\mathbb{P}^\top \hat{\mathbf{z}}_r)$, and $\nabla_{\hat{\mathbf{z}}_r} \hat{H}_r(\hat{\mathbf{z}}_r) = \mathbf{z}_r + \mathbf{V}_r^\top \mathbb{P} \nabla_{\mathbf{z}} \tilde{h}(\mathbb{P}^\top \mathbf{V}_r \hat{\mathbf{z}}_r)$. The hyper-reduced system (4) is also conformal symplectic because the Jacobian of \hat{H}_r is symmetric. Indeed

$$\partial_{\hat{\mathbf{z}}_r} \hat{H}_r(\hat{\mathbf{z}}_r) = \mathbf{I} + \mathbf{V}_r^\top \mathbb{P} \partial_{\mathbf{z}} \tilde{h}(\mathbb{P}^\top \mathbf{V}_r \hat{\mathbf{z}}_r) \mathbb{P}^\top \mathbf{V}_r.$$

$$\hat{H}_r(\mathbf{z}_r) = \hat{H}(\mathbf{V}_r \mathbf{z}_r)$$

How is (2) a generalization of (1)?

$$\left. \begin{aligned} \dot{z} &= J \nabla_z \tilde{H}(z) - \tilde{\gamma} z, \\ 0 &= g(q). \end{aligned} \right\} \quad (2)$$

$$\left. \begin{aligned} \dot{q} &= \nabla_p H(q, p), \\ \dot{p} &= -\nabla_q H(q, p) - \gamma p - G(q)^\top \lambda, \\ 0 &= g(q). \end{aligned} \right\} \quad (1)$$

$$\tilde{H}(z) = H(z) + g(q)^\top \lambda + \gamma/4 \ z^\top \begin{bmatrix} I & I \end{bmatrix} z$$

$$\tilde{\gamma} = \gamma/2$$

$$\nabla_z \tilde{H}(z) = \nabla_z H(z) + \begin{bmatrix} g(q)^\top \lambda \\ 0 \end{bmatrix} - \gamma/2 \underbrace{\begin{bmatrix} I & I \end{bmatrix}}_{= \begin{bmatrix} p \\ q \end{bmatrix}} z$$

$$\begin{bmatrix} -I & I \end{bmatrix} \nabla_z \tilde{H}(z) + \tilde{\gamma} z = \begin{bmatrix} \nabla_p H(z) - \gamma/2 \ q + \gamma/2 \ q \\ -(\nabla_q H(z) + g(q)^\top \lambda - \gamma/2 \ p) + \gamma/2 \ p \end{bmatrix} \quad \checkmark$$

→ maybe that is wrong in the numerics yet?

not sure how much novelty is in this method compared to Cecilia DEIM, DEIM in [4]

This implies that the Jacobian $\partial_{\hat{\mathbf{z}}_r^2} \hat{H}_r(\hat{\mathbf{z}}_r)$ is symmetric and, therefore,

$$\begin{aligned} \frac{1}{2} d\hat{\mathbf{z}}_r \wedge \mathbf{J}_r^{-1} d\hat{\mathbf{z}}_r &= d\hat{\mathbf{z}}_r \wedge \partial_{\hat{\mathbf{z}}_r^2} \hat{H}_r(\hat{\mathbf{z}}_r) - \gamma d\hat{\mathbf{z}}_r \wedge \mathbf{J}_r^{-1} d\hat{\mathbf{z}}_r \\ &= -\gamma d\hat{\mathbf{z}}_r \wedge \mathbf{J}_r^{-1} d\hat{\mathbf{z}}_r \end{aligned}$$

where $d\hat{\mathbf{z}}_r$ is the solution of the variational equation associated with eq. (4). Equation (4) can now be solved with a conormal symplectic integrator to obtain the hyper-reduced solution $\hat{\mathbf{z}}_r$.

We summarize the process of obtaining eq. (4) in Table 1 (c.f. Algorithm 5 of [4]).

rather use algorithm environment

Table 1: POD-DEIM structure-preserving reduction of CPHS (2).

1. Select a positive-definite \mathbf{Q} such that

$$H = \frac{1}{2} \mathbf{z}^\top \mathbf{Q} \mathbf{z} + h(\mathbf{z}).$$

2. Generate the solutions $\mathbf{z}(t)$ and construct the snapshot matrices

$$\begin{aligned} \mathbb{Z} &= [\mathbf{z}(t_0), \mathbf{z}(t_1), \dots, \mathbf{z}(t_N)] \\ \mathbb{F} &= [\nabla_z \tilde{H}(\mathbf{z}(t_0)), \nabla_z \tilde{H}(\mathbf{z}(t_1)), \dots, \nabla_z \tilde{H}(\mathbf{z}(t_N))] \\ \mathbb{K} &= [\nabla_z \tilde{h}(\mathbf{z}(t_0)), \nabla_z \tilde{h}(\mathbf{z}(t_1)), \dots, \nabla_z \tilde{h}(\mathbf{z}(t_N))] \end{aligned}$$

3. Truncate the SVD of \mathbb{Z} , \mathbb{F} , \mathbb{K} to obtain the POD bases \mathbf{V}_r , \mathbf{W}_r , and \mathbf{U}_l , respectively, such that

$$\mathbf{z}(t) \approx \mathbf{V}_r \mathbf{z}_r(t), \quad \nabla_z \tilde{H}(\mathbf{z}) \approx \mathbf{W}_r \mathbf{f}_r(t), \quad \nabla_z \tilde{h}(\mathbf{z}) \approx \mathbf{U}_l \mathbf{k}_l(t).$$

Ensure that the bases satisfy $\mathbf{W}_r^\top \mathbf{V}_r = \mathbf{I}$, $\mathbf{V}_r^\top \mathbf{Q} \mathbf{V}_r = \mathbf{I}$, and $\mathbf{U}_l^\top \mathbf{Q} \mathbf{U}_l = \mathbf{I}$.

4. Construct the DEIM projection $\mathbb{P} = \mathbf{U}_l (\mathbf{E}^\top \mathbf{U}_l)^{-1} \mathbf{E}^\top$.
5. The POD-DEIM reduced system reads as eq. (4).

3 Numerical results

In this section, we consider eq. (1) with

$$H(t) = \frac{1}{2} \|\mathbf{p}(t)\|^2 - \sum_{i=1}^d \cos(\mathbf{q}_i(t)), \quad d=3, \gamma=0.1, \quad \mathbf{g}(\mathbf{q}) = \sum_{i=1}^d (\alpha_i q_i^2) - 1, \quad \|\boldsymbol{\alpha}\| = 1. \quad (5)$$

We will reduce and simulate eq. (5) in structure-preserving manner.

Work in progress ..

what is the motivation behind this equation?

non-fat q?

is α given or a variable?

isn't that very low-dimensional?

maybe good as dummy example?

References

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