

# Reduced basis model order reduction of dissipatively perturbed constrained Hamiltonian systems

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## 1 Introduction

In this paper, we consider ~~the~~ model order reduction (MOR) of the differential algebraic equations of the form

$$\left. \begin{aligned} \dot{\mathbf{q}} &= \nabla_{\mathbf{p}} H(\mathbf{q}, \mathbf{p}), \\ \dot{\mathbf{p}} &= -\nabla_{\mathbf{q}} H(\mathbf{q}, \mathbf{p}) - \gamma \mathbf{p} - \mathbf{G}(\mathbf{q})^\top \boldsymbol{\lambda}, \\ \mathbf{0} &= \mathbf{g}(\mathbf{q}). \end{aligned} \right\} \quad (1)$$

Here  $\mathbf{q}(t), \mathbf{p}(t) : \mathbb{R} \rightarrow \mathbb{R}^d$  are referred to as the position and conjugate momenta, respectively, and constitute the solution of the system. The smooth function  $H : \mathbb{R}^{2d} \rightarrow \mathbb{R}$  is the Hamiltonian,  $\gamma > 0$  is the dissipation parameter,  $\mathbf{g} \in \mathbb{R}^m$  are the holonomic constraints (depend only on  $\mathbf{q}$ ),  $\mathbf{G}$  is the Jacobian matrix of  $\mathbf{g}$ , i.e.,  $\mathbf{G}(\mathbf{q}) = \mathbf{g}_{\mathbf{q}}(\mathbf{q})$ , and  $\boldsymbol{\lambda} \in \mathbb{R}^m$  are the Lagrange multipliers. By differentiating the constraint  $\mathbf{g}(\mathbf{q}) = \mathbf{0}$ , we see that the solution of eq. (1) evolves on the manifold defined by the phase-space

$$\mathfrak{P} = \{(\mathbf{q}, \mathbf{p}) \in \mathbb{R}^{2d} : \mathbf{g}(\mathbf{q}) = \mathbf{0}, \mathbf{G}(\mathbf{q})\nabla_{\mathbf{p}} H(\mathbf{q}, \mathbf{p}) = \mathbf{0}\}.$$

Here and henceforth we have dropped the explicit dependence of variables on  $t$  to keep the notation concise. The constraints defining  $\mathfrak{P}$  are also known as the *weak invariants* of the system. Equation (1) can be thought of as a dissipative perturbation of a constrained Hamiltonian system. Constrained Hamiltonian systems find application in many areas including mechanics, chemical processes, and molecular dynamics.

*Constrained Hamiltonian system* is an important special case of eq. (1) when  $\gamma = 0$ . In this special case, the Hamiltonian  $H$  and symplecticness are constants of motion of the constrained Hamiltonian system. The *Hamiltonian system* is obtained from eq. (1) when  $\mathbf{g} = \mathbf{0}$ ,  $\gamma = 0$ . Structure-preserving techniques, preserving the symplecticness through MOR, have been developed in [1, 3, 6].

A generalization of the Hamiltonian system is *port-Hamiltonian system*, which is the aggregate system model obtained when the core dynamics of the subsystem components are described by variational principles. MOR techniques preserving the port-Hamiltonian structure of the port-Hamiltonian systems are developed in [4].

Another relevant DAE that finds a lot of mention in the literature is the *descriptor system*

$$\left. \begin{aligned} \mathbf{E}\dot{\mathbf{x}}(t) &= \mathbf{A}\mathbf{x}(t) + \mathbf{B}\mathbf{u}(t), \\ \mathbf{y}(t) &= \mathbf{C}\mathbf{x}(t) + \mathbf{D}\mathbf{u}(t), \end{aligned} \right\}$$

where  $\mathbf{x} \in \mathbb{R}^n$ ,  $\mathbf{u} \in \mathbb{R}^m$ ,  $\mathbf{y} \in \mathbb{R}^p$  are states, inputs, outputs respectively. Matrix  $\mathbf{E} \in \mathbb{R}^{n \times n}$  is a singular matrix, and  $\mathbf{A} \in \mathbb{R}^{n \times n}$ ,  $\mathbf{B} \in \mathbb{R}^{n \times m}$ ,  $\mathbf{C} \in \mathbb{R}^{p \times n}$ ,  $\mathbf{D} \in \mathbb{R}^{p \times m}$ . The descriptor systems have been reduced using interpolatory projection methods in [5] and using balanced truncation in [2, 7].

The research on MOR of DAEs so far has mainly focused on ~~the~~ reduced basis MOR of ~~the~~ port-Hamiltonian systems and on interpolatory and balanced truncation MOR of ~~the~~ descriptor systems. In this paper, we take a reduced-basis approach to MOR of the nonlinear DAE (1). Our objective is twofold:

- To reduce eq. (1) using the reduced basis method such that the conformal symplecticness of the equation is preserved through MOR
- To accelerate the simulation of the reduced order model (ROM) using a nonlinearity reduction technique, such as DEIM, by structure-preserving hyper-reduction of the conservative and constraining forces in eq. (1)

In the end, we conduct numerical experiments on the ROM using a conformal symplectic numerical method to solve it.

! Paper with long progress: Comp. of reduced techniques w/ (pH system) (time integration)

## 2 MOR of Eq. (1)

We begin by addressing our first objective – structure-preserving reduction (SPR) of Eq. (1) using a reduced basis method. To this end, let us consider the following generalization of Eq. (1).

$$\begin{cases} \dot{z} = J \nabla_z \tilde{H}(z) - \gamma z, \\ 0 = g(q). \end{cases} \quad (2)$$

Here  $z(t) : \mathbb{R} \rightarrow \mathbb{R}^{2d}$  is the solution of the system, and the structure matrix  $J$  is skew-symmetric and non-singular. The smooth function  $\tilde{H} = H + g^\top \lambda$  is the augmented Hamiltonian,  $\gamma > 0$  is the dissipation parameter, and  $g \in \mathbb{R}^m$  are the constraints. The separable system (1) can be deduced from eq. (2) replacing  $\gamma$  by  $\gamma/2$ ,  $H(z)$  by  $H(q, p) + \frac{\gamma}{2} z$ , and the structure matrix  $J$  by

$$\begin{bmatrix} 0 & I \\ -I & 0 \end{bmatrix}.$$

The system (2) is *conformal symplectic* since

$$dz \wedge J^{-1} dz = -2\gamma dz \wedge J^{-1} dz,$$

where  $dz$  is the solution of the variational equation associated with (2). We refer to Eq. (2) as *constrained perturbed Hamiltonian system* (CPHS).

Let us assume that we have obtained the projection bases  $V_r, W_r$  such that

$$z(t) \approx V_r z_r(t), \quad \nabla_z \tilde{H}(z(t)) \approx W_r f_r(t),$$

respectively, for some  $f_r \in \mathbb{R}^r$ . Assuming  $V_r^\top W_r = I$ , it follows that

$$\begin{aligned} \nabla_z \tilde{H}(V_r z_r) &\approx \nabla_z \tilde{H}(z) \\ &\approx W_r f_r(t) \\ &= W_r V_r^\top W_r f_r(t) \\ &\approx W_r V_r^\top \nabla_z \tilde{H}(V_r z_r) \\ &= W_r \nabla_{z_r} \tilde{H}_r(z_r), \end{aligned}$$

where  $\tilde{H}_r(z_r) = \tilde{H}(V_r z_r)$ .

Substituting  $z(t) \approx V_r z_r(t)$ ,  $\nabla_z \tilde{H}(V_r z_r) \approx W_r \nabla_{z_r} \tilde{H}_r(z_r)$  in Eq. (2), one gets the reduced system

$$\begin{cases} \dot{z}_r = J_r \nabla_{z_r} \tilde{H}_r(z_r) - \gamma z_r, \\ 0 = g(q_r), \end{cases} \quad (3)$$

where  $J_r = W_r^\top J W_r$  is skew-symmetric, and  $q_r$  denotes the first half of the vector  $V_r z_r$ . The reduced system (3) is again *conformal symplectic* because the *Jacobian*  $\partial_{z_r} \tilde{H}_r$  is symmetric.

We further reduce the nonlinearities in Eq. (3) with the discrete empirical interpolation method (DEIM). To this end, let us assume that the Hamiltonian  $H$  can be expressed as

$$H = \frac{1}{2} z^\top Q z + h(z)$$

with  $Q$  being a positive-definite matrix and  $h$  being the nonlinear part of  $H$ . Let us define  $\tilde{h} = h + g^\top \lambda$ , and assume that  $\nabla_z \tilde{h}(z) \approx U_l k_l(t)$  and that  $V_r^\top Q V_r = I$ ,  $U_l^\top Q U_l = I$ . Let us further assume the DEIM projection  $\mathbb{P} = U_l (E^\top U_l)^{-1} E^\top$  for the hyper-reduction of  $\tilde{h}$ . Then the hyper-reduced system reads

$$\begin{cases} \dot{\hat{z}}_r = J_r \nabla_{\hat{z}_r} \hat{H}_r(\hat{z}_r) - \gamma \hat{z}_r, \\ 0 = g(\hat{q}_r). \end{cases} \quad (4)$$

Here  $\hat{H} = \frac{1}{2} \hat{z}^\top Q \hat{z} + \mathbb{P} \tilde{h}(\mathbb{P}^\top \hat{z})$ , and  $\nabla_{\hat{z}_r} \hat{H}_r(\hat{z}_r) = z_r + V_r^\top \mathbb{P} \nabla_z \tilde{h}(\mathbb{P}^\top V_r \hat{z}_r)$ . The hyper-reduced system (4) is also conformal symplectic because the Jacobian of  $\hat{H}_r$  is symmetric. Indeed

$$\partial_{\hat{z}_r} \hat{H}_r(\hat{z}_r) = I + V_r^\top \mathbb{P} \partial_{z^2} \tilde{h}(\mathbb{P}^\top V_r \hat{z}_r) \mathbb{P}^\top V_r.$$

Q: status Babak preprint?

! Basis generator not yet mentioned! !

This implies that the Jacobian  $\partial_{\hat{\mathbf{z}}_r} \hat{H}_r(\hat{\mathbf{z}}_r)$  is symmetric and, therefore,

what eqn.?

$$\begin{aligned} \frac{1}{2} d\hat{\mathbf{z}}_r \wedge \mathbf{J}_r^{-1} d\hat{\mathbf{z}}_r &= d\hat{\mathbf{z}}_r \wedge \partial_{\hat{\mathbf{z}}_r} \hat{H}_r(\hat{\mathbf{z}}_r) - \gamma d\hat{\mathbf{z}}_r \wedge \mathbf{J}_r^{-1} d\hat{\mathbf{z}}_r \\ &= -\gamma d\hat{\mathbf{z}}_r \wedge \mathbf{J}_r^{-1} d\hat{\mathbf{z}}_r \end{aligned}$$

(?)

where  $d\hat{\mathbf{z}}_r$  is the solution of the variational equation associated with eq. (4). Equation (4) can now be solved with a conformal symplectic integrator to obtain the hyper-reduced solution  $\hat{\mathbf{z}}_r$ .

We summarize the process of obtaining Eq. (4) in Table 1 (c.f. Algorithm 5 of [4]).

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Table 1: POD-DEIM structure-preserving reduction of CPHS (2).

1. Select a positive-definite  $\mathbf{Q}$  such that

$$H = \frac{1}{2} \mathbf{z}^\top \mathbf{Q} \mathbf{z} + h(\mathbf{z}).$$

2. Generate the solutions  $\mathbf{z}(t)$  and construct the snapshot matrices

$$\begin{aligned} \mathbb{Z} &= [\mathbf{z}(t_0), \mathbf{z}(t_1), \dots, \mathbf{z}(t_N)] \\ \mathbb{F} &= [\nabla_{\mathbf{z}} \tilde{H}(\mathbf{z}(t_0)), \nabla_{\mathbf{z}} \tilde{H}(\mathbf{z}(t_1)), \dots, \nabla_{\mathbf{z}} \tilde{H}(\mathbf{z}(t_N))] \\ \mathbb{K} &= [\nabla_{\mathbf{z}} \tilde{h}(\mathbf{z}(t_0)), \nabla_{\mathbf{z}} \tilde{h}(\mathbf{z}(t_1)), \dots, \nabla_{\mathbf{z}} \tilde{h}(\mathbf{z}(t_N))] \end{aligned}$$

(Q: no symplecticity of  $V_r, W_r$  required?)

3. Truncate the SVD of  $\mathbb{Z}, \mathbb{F}, \mathbb{K}$  to obtain the POD bases  $\mathbf{V}_r, \mathbf{W}_r$ , and  $\mathbf{U}_l$ , respectively, such that

$$\mathbf{z}(t) \approx \mathbf{V}_r \mathbf{z}_r(t), \nabla_{\mathbf{z}} \tilde{H}(\mathbf{z}) \approx \mathbf{W}_r \mathbf{f}_r(t), \nabla_{\mathbf{z}} \tilde{h}(\mathbf{z}) \approx \mathbf{U}_l \mathbf{k}_l(t).$$

Ensure that the bases satisfy  $\mathbf{W}_r^\top \mathbf{V}_r = \mathbf{I}$ ,  $\mathbf{V}_r^\top \mathbf{Q} \mathbf{V}_r = \mathbf{I}$ , and  $\mathbf{U}_l^\top \mathbf{Q} \mathbf{U}_l = \mathbf{I}$ .

4. Construct the DEIM projection  $\mathbb{P} = \mathbf{U}_l (\mathbf{E}^\top \mathbf{U}_l)^{-1} \mathbf{E}^\top$ .

5. The POD-DEIM reduced system reads as eq. (4).

discons with previous

### 3 Numerical results

In this section, we consider eq. (1) with

$$H(t) = \frac{1}{2} \|\mathbf{p}(t)\|^2 - \sum_{i=1}^d \cos(\mathbf{q}_i(t)), \quad d=3, \gamma=0.1, \quad \mathbf{g}(\mathbf{q}) = \sum_{i=1}^d (\alpha_i q_i^2) - 1, \quad \|\boldsymbol{\alpha}\| = 1. \quad (5)$$

We will reduce and simulate eq. (5) in structure-preserving manner.

Work in progress ..

### References

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