

CEE 576: Nonlinear Finite Elements - Fall 2016

Homework 2

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Solⁿ 1:

The general nonlinear finite element formulation for 3-D Elasticity problem can be described as follows:

$$a(\mathbf{w}, \mathbf{u}; \mathbf{u}, \mathbf{x}) = (\mathbf{w}, \mathbf{f})_{\Omega} + (\mathbf{w}, \mathbf{h})_{\Gamma_h} \quad (1)$$

where

$$a(\mathbf{w}, \mathbf{u}; \mathbf{u}, \mathbf{x}) = \int_{\Omega} w_{i,j} \sigma_{ij}(\epsilon_{ij}, \mathbf{x}) d\Omega ; \quad (\mathbf{w}, \mathbf{f})_{\Omega} = \int_{\Omega} w_i f_i d\Omega ; \quad (\mathbf{w}, \mathbf{h})_{\Gamma_h} = \int_{\Gamma_h} w_i h_i d\Gamma$$

The corresponding 1-D version of Equation (1) be stated as follows:

$$a(w, u; u, x) = (w, f)_{\Omega} + (w, h)_{\Gamma}$$

The Galerkin form is similar to the above equation, except that the functions w^h , u^h and f^h are piece-wise continuous over each element. Moreover f^h can be thought of as the mapping of the forcing function of the finite element space of functions.

$$a(w^h, u^h; u^h, x) = \int_{\Omega} w_{,x}^h \sigma^h(u_{,x}^h) ; \quad (w^h, f^h)_{\Omega} = \int_{\Omega} w^h f^h d\Omega ; \quad (w^h, h) = \int_{\Gamma_h} w^h h d\Gamma \quad (2)$$

(a) Strong Form (S):

For the nonlinear 1D elasticity problem, given external forcing function $f(x)$ and material parameters $\{\Omega = (0, 1)\}$, find $u(x) : \Omega \rightarrow \mathbb{R}$ such that $u(x) \in \mathcal{U}$ and

$$\begin{aligned} \sigma_{,x}(u_{,x}; x) + f(x) &= 0 \\ u(1) &= g; \quad \sigma(0) = h \end{aligned} \quad (3)$$

where

$$\mathcal{U} = \{u(x) \mid u(x) \in C^2(0, 1), u(1) = g\}$$

The dependence on x (inhomogeneity) appears in the first term in the expression, which ultimately appears in the Consistent-tangent tensor. The forcing function $f(x)$ may (or may not) potentially depend on the spatial variable x .

(b) Weak Form (W):

We multiply Equation(3)₁ by a test function $w(x)$ and integrate by parts to get

$$\int_{\Omega} w_{,x}(x) \sigma(u_{,x}; x) d\Omega = w(x) \sigma(u_{,x}; x) |_{\Gamma_h} + \int_{\Omega} w(x) f(x) d\Omega$$

Thus the Weak form (W) can be stated as follows:

Given $f(x)$ and other material parameters $\{\Omega = (0, 1)\}$, find $u(x) : \Omega \rightarrow \mathbb{R}$ such that $u(x) \in \mathcal{U}_0(\Omega)$

$$\int_{\Omega} w_{,x}(x) \sigma(u_{,x}; x) d\Omega = w(x) \sigma(u_{,x}; x) |_{\Gamma_h} + \int_{\Omega} w(x) f(x) d\Omega \quad \forall w(x) \in \mathcal{V}(\Omega) \quad (4)$$

$$\begin{aligned} \mathcal{U}_0(\Omega) &= \left\{ u \mid u(x) \in \mathbf{H}^1(\Omega), u(x)|_{\Gamma_g} = g \right\} \\ \mathcal{V}(\Omega) &= \left\{ w \mid w(x) \in \mathbf{H}^1(\Omega), w(x)|_{\Gamma_g} = 0 \right\} \end{aligned}$$

(c) Galerkin-Form (G):

We now introduce the concept of discretization (mesh). The Galerkin-form (G) is, find $u^h(x) : \Omega_e \rightarrow \mathbb{R}$ such that $u^h(x) \in \mathcal{U}_0^h(\Omega_e)$

$$\sum_{e=1}^{n_{el}} \int_{\Omega_e} w^h_{,x}(x) \sigma(u^h_{,x}; x) d\Omega = w^h(x) \sigma(u^h_{,x}; x) |_{\Gamma_h} + \sum_{e=1}^{n_{el}} \int_{\Omega_e} w^h(x) f^h(x) d\Omega \quad \forall w^h(x) \in \mathcal{V}^h(\Omega_e) \quad (5)$$

$$\begin{aligned} \mathcal{U}_0^h(\Omega_e) &= \left\{ u^h(x) \mid u^h(x) \in \mathbf{H}^1(\Omega_e), u^h(x)|_{\Gamma_g} = g \right\} \\ \mathcal{V}^h(\Omega_e) &= \left\{ w^h(x) \mid w^h(x) \in \mathbf{H}^1(\Omega_e), w^h(x)|_{\Gamma_g} = 0 \right\} \end{aligned}$$

Choice of Shape Functions:

For the first part of the problem we consider linear shape functions to compute the external force vector, internal force vector, and consistent tangent i.e. for $-1 \leq \xi \leq 1$

$$N(\xi) = \frac{1}{2} \begin{Bmatrix} 1 - \xi \\ 1 + \xi \end{Bmatrix}$$

For the Part.B and Part.C of the problem, separate reference to quadratic shape functions is made. We use isoparametric elements throughout, and hence

$$x = \sum_{i=1}^{n_{en}} N_i^e(\xi) x_i$$

(d) Nonlinear Matrix Problem:

1. External Force Vector:

The external load vector, for an element is calculated as follows:

$$\begin{aligned}\mathbf{f}_a^e &= \int_{\Omega_e} N_a(\xi) \hat{f}(x) d\Omega \\ &= \int_{\Omega} N_a(\xi) N_b(\xi) f_b d\Omega\end{aligned}\tag{6}$$

where f_b represent the nodal values of $f^h(x)$. It can be interpreted as the projection of the actual forcing function on the finite element space of functions used to discretize the field. We use one point Gauss-Quadrature to evaluate the integral

$$\begin{aligned}\mathbf{f}_a^e &= \int_{\Omega_e} \mathbf{N}^T \mathbf{N} \mathbf{f}_{\text{nodal}} d\Omega \\ &= \int_{\Omega_e} \frac{1}{4} \begin{bmatrix} (1-\xi)^2 & (1-\xi^2) \\ (1-\xi^2) & (1+\xi)^2 \end{bmatrix} \begin{Bmatrix} f_1 \\ f_2 \end{Bmatrix} d\Omega \\ &= \int_{-1}^1 \frac{1}{4} \begin{bmatrix} (1-\xi)^2 & (1-\xi^2) \\ (1-\xi^2) & (1+\xi)^2 \end{bmatrix} \begin{Bmatrix} f_1 \\ f_2 \end{Bmatrix} \frac{1}{2} d\xi \\ &= \frac{1}{4} \begin{Bmatrix} f_1 + f_2 \\ f_1 + f_2 \end{Bmatrix} + \begin{Bmatrix} h \\ 0 \end{Bmatrix}\end{aligned}$$

2. Internal Force Vector:

The internal force vector for an element is calculated as follows:

$$\begin{aligned}\mathbf{n}_a^e &= \int_{\Omega_e} N_{a,x} \sigma(\epsilon^h; x) d\Omega \\ \text{where, } \epsilon^h &= u^h_{,x}(x) = \sum_{e=1}^{n_{en}} N_{e,x}(\xi) d_e \\ \mathbf{n}_a^e &= \int_{\xi=-1}^1 \frac{2}{h^e} \frac{1}{2} \begin{Bmatrix} -1 \\ 1 \end{Bmatrix} \sigma\left(\frac{1}{h^e} (d_2^e - d_1^e); x(\xi)\right) \frac{h^e}{2} d\xi\end{aligned}$$

Using one point Quadrature ($W = 2$ and $\xi = 0$):

$$\begin{aligned}\mathbf{n}_a^e &= 2 \cdot \frac{1}{2} \begin{Bmatrix} -1 \\ 1 \end{Bmatrix} \sigma\left(\frac{1}{h^e} (d_2^e - d_1^e); \frac{x_1^e + x_2^e}{2}\right) \\ &= \begin{Bmatrix} -1 \\ 1 \end{Bmatrix} \sigma\left(d_2^e - d_1^e; \frac{x_1^e + x_2^e}{2}\right)\end{aligned}$$

The Matrix Problem is just $\mathbb{A}_{e=1}^{n_{el}} \{\mathbf{n}_a^e\} = \mathbb{A}_{e=1}^{n_{el}} \mathbf{f}_a^e$, where the operator, $\mathbb{A}_{e=1}^{n_{el}}$, is the assembly operator over all elements in the mesh. Thus we get

$$\mathbf{f}_a^e = \int_{\Omega} \mathbf{N}_a \mathbf{N}_b f_b^{\text{nodal}} d\Omega + \mathbf{N}_{a_{x=0}} \cdot h$$

(e) Consistent Tangent Tensor:

The consistent tangent tensor can be obtained by taking the variational derivative (in this case the Taylor-Expansion) of the internal force vector:

$$\begin{aligned}\mathbb{D}\mathbf{n}_a^e &= \frac{\partial \mathbf{n}_a^e}{\partial d_b} \\ &= \frac{\partial}{\partial d_b} \int_{\Omega_e} N_{a,x} \sigma(\epsilon^h; x) d\Omega\end{aligned}$$

Applying the chain-rule of differentiation, with $(h^e = 1)$:

$$\begin{aligned}& \int_{\Omega_e} N_{a,x}(\xi) \frac{\partial}{\partial \epsilon^h} \sigma(\epsilon^h; x) \frac{\partial \epsilon^h}{\partial d_b} d\Omega \\ &= \int_{\Omega_e} N_{a,x}(\xi) \mathbb{C}(\epsilon^h; x) N_{e,x}(\xi) \delta_{eb} d\Omega \\ &= \int_{\Omega_e} N_{a,x}(\xi) \mathbb{C}(\epsilon^h; x) N_{b,x}(\xi) d\Omega \\ &= \int_{-1}^1 \frac{2}{h^e} N_{a,\xi}(\xi) \mathbb{C}(\epsilon^h(\xi); x(\xi)) N_{b,\xi}(\xi) d\xi \\ &= 2 \cdot 2 \frac{1}{4} \begin{Bmatrix} -1 \\ 1 \end{Bmatrix} \mathbb{C}\left(\frac{d_2^e - d_1^e}{h^e}; \frac{x_2^e + x_1^e}{2}\right) \begin{Bmatrix} -1 & 1 \end{Bmatrix} \\ &= \begin{bmatrix} 1 & -1 \\ -1 & 1 \end{bmatrix} \mathbb{C}\left(d_2^e - d_1^e; \frac{x_2^e + x_1^e}{2}\right)\end{aligned} \tag{7}$$

(f) Internal Force:

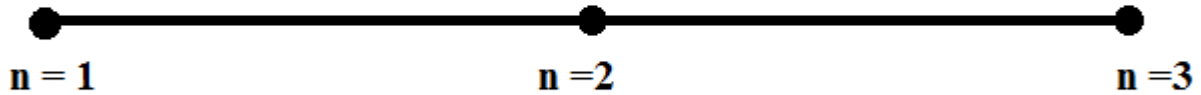
For this part, and the next, calculations are presented with respect to quadratic shape functions and 2-point numerical quadrature.

Choice of Shape Functions:

1-D Quadratic Shape Functions, listed for the given element as shown in the figure:

$$N(\xi) = \frac{1}{2} \begin{Bmatrix} \xi(\xi - 1) \\ 2(1 + \xi)(1 - \xi) \\ \xi(\xi + 1) \end{Bmatrix}$$

Assumption: The internal node is equidistant from the other two nodes.



Calculations:

$$\mathbf{n}_a^e = \int_{\Omega_e} N_{a,x} \sigma(\epsilon^h; x) d\Omega \quad (8)$$

Again using isoparametric elements:

$$x = \sum_{i=1}^{n_{en}} N_i^e(\xi) x_i$$

Thus

$$J = \frac{dx}{d\xi} = \sum_{i=1}^{n_{en}} \frac{dN_i^e}{d\xi}(\xi) x_i^e = \frac{(2\xi - 1)}{2} x_1^e - 2\xi x_2^e + \frac{(2\xi + 1)}{2} x_3^e = \frac{1}{2}$$

$$\mathbf{n}_a^e = \int_{-1}^1 \frac{1}{2} \begin{bmatrix} 2\xi - 1 \\ -4\xi \\ 2\xi + 1 \end{bmatrix} \sigma(\epsilon^h(\xi); x(\xi)) d\xi$$

Now using 2-point Gauss-Quadrature to integrate the above polynomial, we get:

$$\begin{aligned} & \frac{1}{2} \begin{bmatrix} 2\xi - 1 \\ -4\xi \\ 2\xi + 1 \end{bmatrix} \sigma(\epsilon^h(\xi); x(\xi)) \Big|_{\xi=-\frac{1}{\sqrt{3}}} + \frac{1}{2} \begin{bmatrix} 2\xi - 1 \\ -4\xi \\ 2\xi + 1 \end{bmatrix} \sigma(\epsilon^h(\xi); x(\xi)) \Big|_{\xi=\frac{1}{\sqrt{3}}} \\ &= \begin{bmatrix} -1.07735 \\ 1.155 \\ -0.07735 \end{bmatrix} \sigma\left(\epsilon^h\left(\frac{-1}{\sqrt{3}}\right); x\left(\frac{-1}{\sqrt{3}}\right)\right) + \begin{bmatrix} 0.07735 \\ -1.155 \\ 1.07735 \end{bmatrix} \sigma\left(\epsilon^h\left(\frac{1}{\sqrt{3}}\right); x\left(\frac{1}{\sqrt{3}}\right)\right) \end{aligned} \quad (9)$$

where,

$$\begin{aligned} \epsilon^h\left(\frac{-1}{\sqrt{3}}\right) &= -2.1547d_1^e + 2.311d_2^e - 0.1547d_3^e; & x\left(\frac{-1}{\sqrt{3}}\right) &= 0.455x_1^e + 0.667x_2^e - 0.122x_3^e = 0.2115 \\ \epsilon^h\left(\frac{1}{\sqrt{3}}\right) &= 0.1547d_1^e - 2.311d_2^e + 2.1547d_3^e; & x\left(\frac{1}{\sqrt{3}}\right) &= -0.122x_1^e + 0.667x_2^e + 0.455x_3^e = 0.7885 \end{aligned}$$

(g) Consistent Tangent Tensor:

Using the calculation carried out in (7) we have

$$\begin{aligned} \mathbb{D}\mathbf{n}_a^e &= \int_{-1}^1 \frac{2}{h^e} N_{a,\xi}(\xi) \mathbb{C}(\epsilon^h(\xi); x(\xi)) N_{b,\xi}(\xi) d\xi \\ &= \int_{-1}^1 \frac{2}{h^e} \frac{1}{4} \begin{bmatrix} 2\xi - 1 \\ -4\xi \\ 2\xi + 1 \end{bmatrix} \mathbb{C}(\epsilon^h(\xi); x(\xi)) \begin{bmatrix} 2\xi - 1 & -4\xi & 2\xi + 1 \end{bmatrix} d\xi \\ &= \int_{-1}^1 \frac{1}{2} \begin{bmatrix} (2\xi - 1)^2 & -4\xi(2\xi - 1) & 4\xi^2 - 1 \\ -4\xi(2\xi - 1) & 16\xi^2 & -4\xi(2\xi + 1) \\ 4\xi^2 - 1 & -4\xi(2\xi + 1) & (2\xi + 1)^2 \end{bmatrix} \mathbb{C}(\epsilon^h(\xi); x(\xi)) d\xi \end{aligned} \quad (10)$$

Using 2-point Gauss-Quadrature to calculate the above integral we get,

$$\frac{1}{2} \left[\begin{bmatrix} 4.642 & -4.976 & 0.333 \\ -4.976 & 5.33 & -0.357 \\ 0.333 & -0.357 & 0.0239 \end{bmatrix} \mathbb{C} \left(\epsilon^h \left(\frac{-1}{\sqrt{3}} \right); x \left(\frac{-1}{\sqrt{3}} \right) \right) + \begin{bmatrix} 0.0239 & -0.357 & 0.333 \\ -0.357 & 5.33 & -4.976 \\ 0.333 & -4.976 & 4.642 \end{bmatrix} \mathbb{C} \left(\epsilon^h \left(\frac{1}{\sqrt{3}} \right); x \left(\frac{1}{\sqrt{3}} \right) \right) \right] \quad (11)$$

Thus the Consistent Tangent Tensor, evaluated using the 2-point Gauss-Quadrature would be as follows:

$$\begin{bmatrix} 2.321 & -2.488 & 0.167 \\ -2.488 & 2.667 & -0.179 \\ 0.167 & -0.179 & 0.012 \end{bmatrix} \mathbb{C} \left(\epsilon^h \left(\frac{-1}{\sqrt{3}} \right); x \left(\frac{-1}{\sqrt{3}} \right) \right) + \begin{bmatrix} 0.012 & -0.179 & 0.167 \\ -0.179 & 2.667 & -2.488 \\ 0.167 & -2.488 & 2.321 \end{bmatrix} \mathbb{C} \left(\epsilon^h \left(\frac{1}{\sqrt{3}} \right); x \left(\frac{1}{\sqrt{3}} \right) \right) \quad (12)$$

where,

$$\begin{aligned} \epsilon^h \left(\frac{-1}{\sqrt{3}} \right) &= -2.1547d_1^e + 2.311d_2^e - 0.1547d_3^e ; & x \left(\frac{-1}{\sqrt{3}} \right) &= 0.455x_1^e + 0.667x_2^e - 0.122x_3^e = 0.2115 \\ \epsilon^h \left(\frac{1}{\sqrt{3}} \right) &= 0.1547d_1^e - 2.311d_2^e + 2.1547d_3^e ; & x \left(\frac{1}{\sqrt{3}} \right) &= -0.122x_1^e + 0.667x_2^e + 0.455x_3^e = 0.7885 \end{aligned}$$

Part C.:

No, 1-point Gauss-Quadrature cannot integrate this problem exactly.

- A general m -point Quadrature can exactly integrate a polynomial of order $2m - 1$, thus 1-point Quadrature can only integrate linear polynomials.
- For the consistent tangent tensor there are quadratic polynomials to be integrated (at-least considering \mathbb{C} remains constant) or higher.
- Thus, we would need higher order quadrature, potentially 2 (in case of constant \mathbb{C}) or even more depending on the dependance of \mathbb{C} on x , to exactly integrate this problem.

Some remarks about internal force vector:

- For the internal force vector, we have linear polynomials to be integrated at least, (if σ is constant), but in case of non-homogeneous materials $\sigma(\epsilon; x)$ and hence one point quadrature would not be able to exactly integrate the problem.

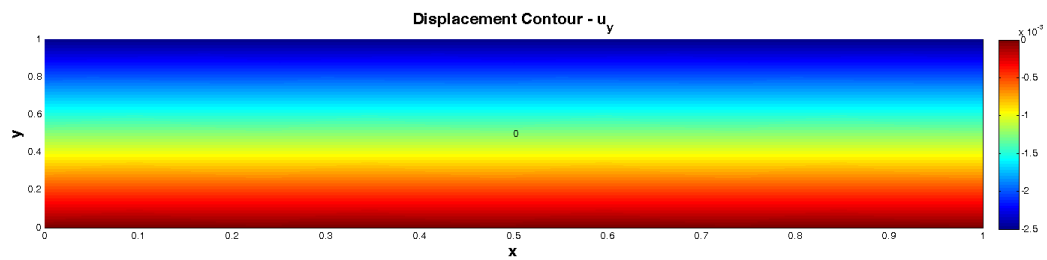
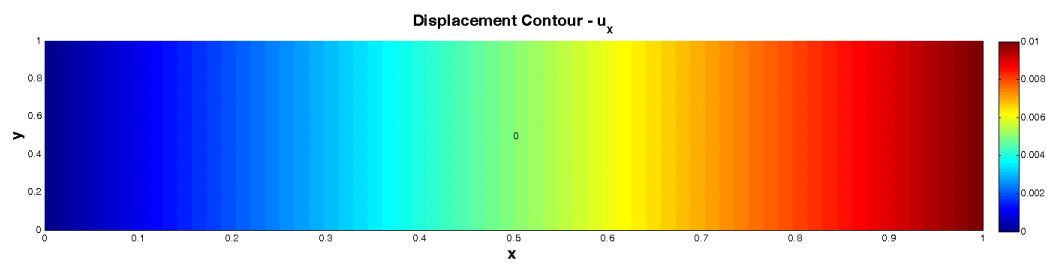
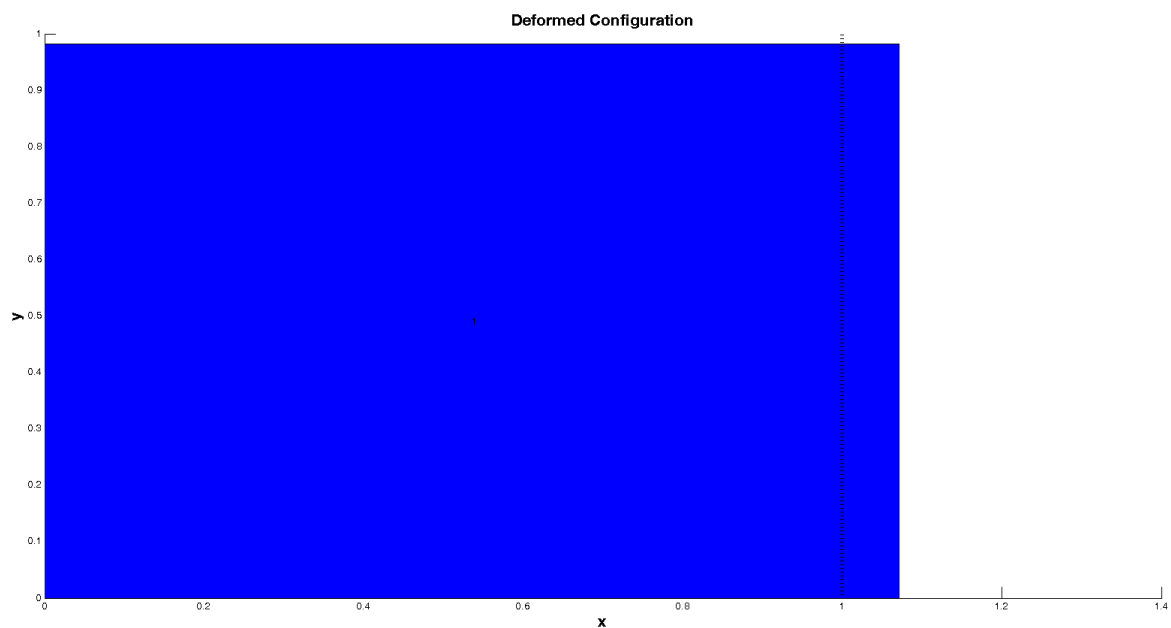
Some remarks about external force vector:

- In case of the external force vector we will have the product of shape functions $N_i N_j$ and the nodal values of the forcing function.
- This would lead to a biquadratic polynomial and hence, even 2-point Gauss-Quadrature would not be sufficient to exactly integrate the external load vector.

Solⁿ 2:

The given problem is solved using the provided Matlab-Code:

(a) Results: Plots



(b) Stress-Strain values

The values for stresses and strains obtained at the integration points are described below:

Table 1: Stress at Integration points

Integration Point \ Stress	σ_{xx}	σ_{yy}	σ_{xy}
1	1	4.547E-13	0
2	1	4.547E-13	0
3	1	4.547E-13	0
4	1	4.547E-13	0

Table 2: Strain at Integration Points

Integration Point \ Strain	ϵ_{xx}	ϵ_{yy}	ϵ_{xy}
1	0.01	-0.0025	0
2	0.01	-0.0025	0
3	0.01	-0.0025	0
4	0.01	-0.0025	0

Table 3: Nodal Displacements

Node	Ux	Uy
1	0	0
2	0.01	0
3	0.01	-0.0025
4	0	-0.0025

Matlab Code:

```
% Input File: One Quadrilateral Element Under Axial Load
%
% Copyright (C) Arif Masud and Tim Truster
%
% This input file should be run prior to executing the FEA_Program routine.
%
% Format of required input:
%
%   numnp:           = number of nodes in the mesh (length(NodeTable))
%
%   numel:           = number of elements in the mesh
%
%   nen:             = maximum number of nodes per element (4)
%
%   PSPS:            = flag for plane stress ('s') or plane strain ('n')
%
%   NodeTable:       = table of mesh nodal coordinates defining the
```



```

% geometry of the mesh; format of the table is as
% follows:
%
%      Nodes |      x-coord  y-coord
%      n1    |      NodeTable = [x1      y1
%      n2    |                  x2      y2
%      ...   |                  ..      ..
%      nnumnp |                  xnumnp ynumnp];
%
% ix:      = table of mesh connectivity information, specifying
%           how nodes are attached to elements and how materials
%           are assigned to elements; entries in the first nen
%           columns correspond to the rows of NodeTable
%           representing the nodes attached to element e;
%           entries in the last nen+1 column are rows from MateT
%           signifying the material properties assigned to
%           element e; format of the table is as follows:
%           Elements |      n1      n2      n3      n4      mat
%           e1       |      ix = [e1n1  e1n2  e1n3  e1n4  e1mat
%           e2       |                  e2n1  e2n2  e2n3  e2n4  e2mat
%           ...      |                  ..      ..      ..      ..      ..
%           enumel   |                  values for element numel    ];
%
% MateT:    = table of mesh material properties for each distinct
%           set of material properties; these sets are
%           referenced by element e by setting the value of
%           ix(e,nen+1) to the row number of the desired
%           material set; format of the table is as follows:
%           Materials |      E      v      t
%           mat1      |      MateT = [E1   v1   t1
%           mat2      |                  E2   v2   t2
%           ...       |                  ..   ..   ..];
%
% BCLIndex: = list of the number of boundary conditions and loads
%           applied to the mesh; first entry is the number of
%           prescribed displacements at nodes; second entry is
%           the number of nodal forces
%
% NodeBC:   = table of prescribed nodal displacement boundary
%           conditions; it contains lists of nodes, the
%           direction of the displacement prescribed (x=1, y=2),
%           and the value of the displacement (set 0 for fixed
%           boundary); the length of the table must match the
%           entry in BCLIndex(1), otherwise an error will result
%           if too few conditions are given or extra BCs will be
%           ignored in the model input module; format of the
%           table is as follows:
%           BCs      |      nodes direction value
%           bc1      |      NodeBC = [bc1n   bc1dir  bc1u
%           bc2      |                  bc2n   bc2dir  bc2u
%           ...      |                  ..      ..      .. ];

```

```
%
% NodeLoad:          = table of prescribed nodal forces; it contains lists
%                    of nodes, the direction of the force prescribed
%                    (x=1, y=2), and the value of the force; the length
%                    of the table must match the entry in BCLIndex(2),
%                    otherwise an error will result if too few conditions
%                    are given or extra loads will be ignored in the
%                    model input module; format of the table is as
%                    follows:
%
%                    Loads |      nodes   direction   value
%                    P1   |      P1n     P1dir      P1P
%                    P2   |      P2n     P2dir      P2P
%                    ...  |      ..      ..         .. ];
```

The diagram illustrates two basic finite element shapes. On the left is a square element defined by four nodes labeled 1, 2, 3, and 4. Node 1 is at the bottom-left corner, node 2 is at the bottom-right, node 3 is at the top-right, and node 4 is at the top-left. Solid lines connect the nodes to form the square boundary. On the right is a triangular element defined by three nodes labeled 1, 2, and 3. Node 1 is at the bottom-right corner, node 2 is at the top vertex, and node 3 is at the bottom-left corner. The triangle's edges are shown with solid lines.

```
% Arbitrary data for assistance in defining the mesh
```

```
% Mesh Nodal Coordinates
```

```
NodeTable = [0 0
              1 0
              1 1
              0 1];
numnp = length(NodeTable);
```

```
% Mesh Element Connectivities
```

$$\mathbf{i}_X = [1 \ 2 \ 3 \ 4 \ 1];$$

```
nen = 4;
numel = 1;
```

```
% Mesh Boundary Conditions and Loads
```

```

BCIndex = [6 0]';
NodeBC = [1 1 0
           1 2 0
           2 1 0.01
           2 2 0
           3 1 0.01
           4 1 0
           ];
NodeLoad = 0;

```

```

% Mesh Material Properties
young = 100;
pois = .25;
thick = 1;
PSPS = 's';
MateT = [young pois thick];
FEA_Program

function [strain, stress] = CompStrainStress_Elem_Cee570(xl, ul, mateprop, nel,
    ndf, PSPS)
%
% Subroutine to compute strain and stress for linear
% 2-dimensional elasticity element. Element currently supports bilinear
% quadrilateral elements with the following node and shape function
% labelling scheme:
%
%      (-1, 1)  4  -----  3  ( 1, 1)
%               |           |
%               |    s     |
%               |    ^     |
%               |    |     |
%               |    .-> r  |
%               |           |
%      (-1,-1)  1  -----  2  ( 1,-1)
%
% Element local coordinates (r,s) are defined by a coordinate axis with the
% origin at the center of the element; the corners of the element have
% local coordinate values as shown in the figure.
%
% Definitions for input:
%
%   xl:                = local array containing (x,y) coordinates of nodes
%                       forming the element; format is as follows:
%
%                       Nodes   |       n1  n2  n3  n4
%                       x-coord |   xl = [x1  x2  x3  x4
%                       y-coord |       y1  y2  y3  y4];
%
%   mateprop:          = vector of material properties:
%                       mateprop = [E v t];
%                               = [(Young's Modulus) (Poisson's Ratio)
%                               (thickness)];
%
%   nel:               = number of nodes on current element (4)
%
%   ndf:               = max number of DOF per node (2)
%
%   ndm:               = space dimension of mesh (2)
%
%   PSPS:              = flag for plane stress ('s') or plane strain ('n')

```

```

%
% Definitions for output:
%
%   strain:           = strain array containing strain components
%                       at integration points:
%                       xx    yy    xy
%   int1   strain[ .    .    .
%   int2           .    .    .
%   int3           .    .    .
%   int4           .    .    .   ];
%
%   stress:           = stress array containing stress components
%                       at integration points:
%                       xx    yy    xy
%   int1   stress[ .    .    .
%   int2           .    .    .
%   int3           .    .    .
%   int4           .    .    .   ];
%
% Definitions of local constants:
%
%   nst:               = size of element arrays (ndf*nel)
%
%

```

```

% Set Material Properties

```

```

ElemE = mateprop(1);
Elemv = mateprop(2);
thick = mateprop(3);

```

```

if PSPS == 's' %Plane Stress

```

```

    Dmat = ElemE/(1-Elemv^2)*[1      Elemv    0
                             Elemv    1      0
                             0        0      (1-Elemv)/2];

```

```

else %Plane Strain

```

```

    %   Dmat =

```

```

end

```

```

% Load Guass Integration Points

```

```

if nel == 3
    lint = 4;
else
    lint = 4;
end

```

```

% Initialize Matrix and Vector

nst = nel*ndf;
ul_elem = reshape(ul,ndf*nel,1);
strain = zeros(lint,3);
stress = zeros(lint,3);
strain_temp = zeros(3,1);
stress_temp = zeros(3,1);

% Loop over integration points
for l = 1:lint

    if nel == 3
        [Wgt,r,s] = intpntt(l,lint,0);
    else
        [Wgt,r,s] = intpntq(l,lint,0);
    end

    % Evaluate local basis functions at integration point
    shp = shpl_2d(r,s,nel);

    % Evaluate first derivatives of basis functions at int. point
    [Qxy, Jdet] = shpg_2d(shp,xl,nel);

    % Form B matrix
    if nel == 3
        Bmat = [Qxy(1,1) 0 Qxy(1,2) 0 Qxy(1,3) 0
                0 Qxy(2,1) 0 Qxy(2,2) 0 Qxy(2,3)
                Qxy(2,1) Qxy(1,1) Qxy(2,2) Qxy(1,2) Qxy(2,3) Qxy(1,3)];
    else
        Bmat = [Qxy(1,1) 0 Qxy(1,2) 0 Qxy(1,3) 0 Qxy
                  (1,4) 0
                  0 Qxy(2,1) 0 Qxy(2,2) 0 Qxy(2,3) 0
                  Qxy(2,4)
                  Qxy(2,1) Qxy(1,1) Qxy(2,2) Qxy(1,2) Qxy(2,3) Qxy(1,3) Qxy
                  (2,4) Qxy(1,4)];
    end

    % Compute strain
    strain_temp = Bmat*ul_elem;

    strain(1,1) = strain_temp(1,1);
    strain(1,2) = strain_temp(2,1);
    strain(1,3) = strain_temp(3,1);

    % Compute stress
    stress_temp = Dmat*strain_temp;

```

```

stress(1,1) = stress_temp(1,1);
stress(1,2) = stress_temp(2,1);
stress(1,3) = stress_temp(3,1);

```

```
end
```

```
function [ElemK, ElemF] = Elast2d_Elem(xl, mateprop, nel, ndf, ndm, PSPS)
```

```
%
```

```
% Copyright (C) Arif Masud and Tim Truster
```

```
%
```

```
% Subroutine to compute stiffness matrix and force vector for linear
% 2-dimensional elasticity element. Element currently supports bilinear
% quadrilateral elements with the following node and shape function
% labelling scheme:
```

```
%
```

```
% (-1, 1) 4 ----- 3 ( 1, 1)
```

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%
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```
% Element local coordinates (r,s) are defined by a coordinate axis with the
% origin at the center of the element; the corners of the element have
% local coordinate values as shown in the figure.
```

```
%
```

```
% Definitions for input:
```

```
%
```

```
% xl: = local array containing (x,y) coordinates of nodes
% forming the element; format is as follows:
```

```
%
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```

% Definitions for output:
%
%   ElemK:           = element stiffness matrix containing stiffness
%                     entries in the following arrangement, where
%                     wij corresponds to weighting function (i),
%
%                     coordinate
%
%                     direction (j), and ukl corresponds to displacement
%                     function (k), coordinate direction (l):
%
%                     ulx  uly  u2x  u2y  u3x  u3y  u4x  u4y
%   w1x  ElemK[  .    .    .    .    .    .    .    .
%   w1y      .    .    .    .    .    .    .    .
%   w2x      .    .    .    .    .    .    .    .
%   w2y      .    .    .    .    .    .    .    .
%   w3x      .    .    .    .    .    .    .    .
%   w3y      .    .    .    .    .    .    .    .
%   w4x      .    .    .    .    .    .    .    .
%   w4y      .    .    .    .    .    .    .    . ];
%
%   ElemF:           = element force vector containing force entries in the
%                     following arrangement:
%   w1x  ElemF[  .
%   w1y      .
%   w2x      .
%   w2y      .
%   w3x      .
%   w3y      .
%   w4x      .
%   w4y      . ];
%
% Definitions of local constants:
%
%   nst:              = size of element arrays (ndf*nel)
%
%

```

```

% Set Material Properties

```

```

ElemE = mateprop(1);
Elemv = mateprop(2);
thick = mateprop(3);

```

```

if PSPS == 's' %Plane Stress

```

```

    Dmat = ElemE/(1-Elemv^2)*[1      Elemv  0
                             Elemv  1      0
                             0      0      (1-Elemv)/2];

```

```

else %Plane Strain

```

```

end

% Initialize Matrix and Vector

nst = nel*ndf;
ElemK = zeros(nst);
ElemF = zeros(nst,1);

% Load Guass Integration Points

if nel == 3
    lint = 4;
else
    lint = 4;
end

% Loop over integration points
for l = 1:lint

    if nel == 3
        [Wgt,r,s] = intpntt(l,lint,0);
    else
        [Wgt,r,s] = intpntq(l,lint,0);
    end

    % Evaluate local basis functions at integration point
    shp = shpl_2d(r,s,nel);

    % Evaluate first derivatives of basis functions at int. point
    [Qxy, Jdet] = shpg_2d(shp,xl,nel);

    % Form B matrix
    if nel == 3
        Bmat = [Qxy(1,1) 0          Qxy(1,2) 0          Qxy(1,3) 0
                0          Qxy(2,1) 0          Qxy(2,2) 0          Qxy(2,3)
                Qxy(2,1) Qxy(1,1) Qxy(2,2) Qxy(1,2) Qxy(2,3) Qxy(1,3)];
    else
        Bmat = [Qxy(1,1) 0          Qxy(1,2) 0          Qxy(1,3) 0          Qxy
                  (1,4) 0
                  0          Qxy(2,1) 0          Qxy(2,2) 0          Qxy(2,3) 0
                  Qxy(2,4)
                  Qxy(2,1) Qxy(1,1) Qxy(2,2) Qxy(1,2) Qxy(2,3) Qxy(1,3) Qxy
                  (2,4) Qxy(1,4)];
    end

    % Update integration weighting factor
    W = Wgt*Jdet*thick;

```



```
ElemK = ElemK + W*Bmat'*Dmat*Bmat;
```

```
end %je
```