

Jacobian-Free Newton- Krylov (JFNK) Methods for Nonlinear Neutronics/Thermal-Hydraulic Equations

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Outline

1 Introduction

1 Governing Equations

1 Solvers

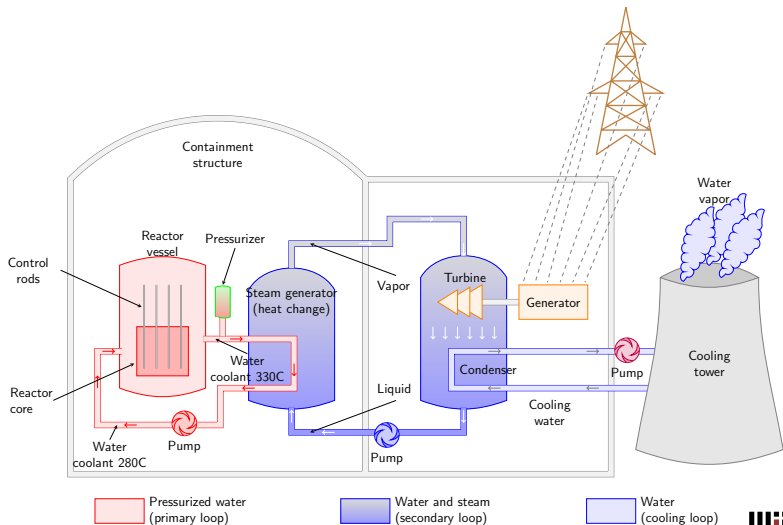
1 Results

1 Conclusions

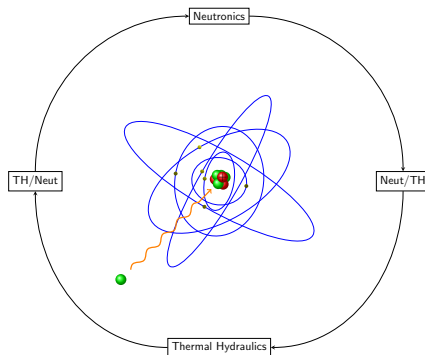
Motivation

- Eventually will be part of thesis work
- JFNK method not currently used in nuclear reactor analysis
- Incorporates a lot of ideas from 2.29 class
- Coupled physics is fun!

Nuclear Reactor Plant

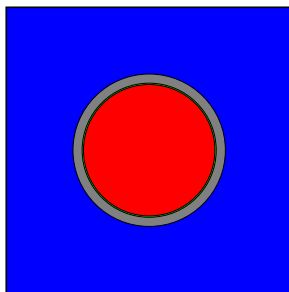


Nuclear Feedback



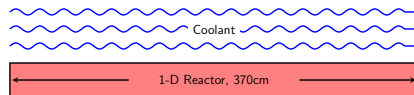
- Fuel Temperature Feedback - $T_f \uparrow$, U-238 Capture \uparrow , Fission Rate \downarrow , Power \downarrow
- Coolant Density Feedback - $\rho \downarrow$, $E_n \uparrow$, Fission Rate \downarrow , Power \downarrow

1-D Slab Reactor Geometry



■ Coolant ■ Clad ■ Gas Gap ■ Fuel

Top-view
Fuel Rod Unit-Cell



Side-view (vertical)
1-D Model of Reactor

Neutronics

■ Basic Neutron Conservation:

$$\textit{Change} + \textit{Leakage} + \textit{Interactions} = \textit{Scattering} + \textit{Fission}$$

■ Neutron Transport Equation:

$$\underbrace{\frac{1}{v} \frac{\partial \varphi}{\partial t}}_{\text{time-dependent}} + \underbrace{\boldsymbol{\Omega} \cdot \nabla \varphi(\mathbf{r}, E, \boldsymbol{\Omega}, t)}_{\text{neutron leakage}} + \underbrace{\Sigma_t(\mathbf{r}, E, t) \varphi(\mathbf{r}, E, \boldsymbol{\Omega}, t)}_{\text{interaction of neutrons with medium}} = \underbrace{Q(\mathbf{r}, E, \boldsymbol{\Omega}, t)}_{\text{neutron source}}$$

■ Neutron Diffusion Equation (1-D Energy Integrated)

$$\frac{1}{v} \frac{\partial \phi}{\partial t} - D(x, t) \frac{\partial^2 \phi}{\partial^2 x} + \Sigma_a(x, t) \phi(x, t) = \frac{1 - \beta}{k_{eff}} \nu \Sigma_f(x, t) \phi(x, t) + \lambda_d c(x, t)$$

$$\frac{\partial c}{\partial t} = \frac{\beta}{k_{eff}} \nu \Sigma_f(x, t) \phi(x, t) - \lambda_d c(x, t)$$

Discretization of Neutronics Equations

Assumptions:

- 1 One-dimensional finite volume spatial discretization
- 2 Central difference scheme for diffusion term
- 3 No-incoming current of neutrons at boundaries
- 4 Implicit Euler time discretization

- Discretized neutronics equation for interior cell

$$\frac{1}{v} \frac{d\bar{\phi}_i}{dt} + -\frac{2}{\Delta x^2} \frac{D_i D_{i-1}}{D_i + D_{i-1}} \bar{\phi}_{i-1} + \left(\frac{2}{\Delta x^2} \frac{D_{i+1} D_i}{D_{i+1} + D_i} + \frac{2}{\Delta x^2} \frac{D_i D_{i-1}}{D_i + D_{i-1}} + \Sigma_{a,i} \right) \bar{\phi}_i - \frac{2}{\Delta x^2} \frac{D_{i+1} D_i}{D_{i+1} + D_i} \bar{\phi}_{i+1} = \frac{1-\beta}{k_{eff}} \nu \Sigma_{f,i} \bar{\phi}_i + \lambda_d \bar{c}_i$$

- Matrix-form of neutronics equations

$$\bar{\Phi}^{n+1} - \bar{\Phi}^n + v\Delta t (\mathbb{M}\bar{\Phi}^{n+1} - (1-\beta)\lambda\mathbb{F}\bar{\Phi}^{n+1} - \lambda_d\bar{\mathbf{c}}^{n+1}) = 0$$

- Matrix-form of precursors

$$\bar{\mathbf{c}}^{n+1} - \bar{\mathbf{c}}^n + \Delta t (\lambda_d \bar{\mathbf{c}}^{n+1} - \beta \lambda \mathbb{F} \bar{\Phi}^{n+1}) = 0$$

Thermal Hydraulics

- Energy Equation - single phase fluid and inviscid fluid

$$\frac{\partial(\rho h)}{\partial t} + \nabla \cdot (\rho h \mathbf{u}) = -\nabla \cdot \mathbf{q}'' + q'''$$

- Assuming fissions are a volumetric heat source in 1-D

$$\rho A \frac{\partial h}{\partial t} + \dot{m} \frac{\partial h}{\partial x} = q'$$

- For an incompressible fluid, $dh = c_p dT$

$$\rho A c_p \frac{\partial T}{\partial t} + \dot{m} c_p \frac{\partial T}{\partial x} = q'$$

Discretization of Energy Equation

- 1 One-dimensional finite volume spatial discretization
- 2 Upwind difference scheme for diffusion term
- 3 Specify inlet conditions and mass flow rate
- 4 Implicit Euler time discretization

■ spatial discretization

$$\frac{\rho A \Delta x}{w} \frac{d\bar{T}_i}{dt} + \bar{T}_i - \bar{T}_{i-1} = \frac{1}{2w c_p} (Q_{i-1} + Q_i)$$

■ Matrix-form with time discretization

$$\bar{\mathbf{T}}^{n+1} - \bar{\mathbf{T}}^n + \frac{w \Delta t}{\mathcal{P}^{n+1} A \Delta x} (\mathbf{S} \bar{\mathbf{T}}^{n+1} - \mathbf{R} \mathbf{Q}^{n+1}) = 0$$

Note: \mathcal{P} is a vector of cell-averaged coolant densities

Physics Coupling

Neutronics - Thermal Hydraulics

- Neutrons cause fissions
- Large portion of fission energy deposited in coolant
- This is represented by

$$\mathbf{Q} = \tilde{c} \mathbb{E} \bar{\Phi} \Delta x$$

where $\mathbb{E} = \text{diag} \{ \kappa \Sigma_f \}$ characterizes energy per fission and \tilde{c} is flux-power normalization constant

Thermal-Hydraulics to Neutronics

- Diffusion theory parameters depend on coolant density
- This dependence is determined with a transport theory code
- D , Σ_a , $\nu \Sigma_f$, $\kappa \Sigma_f$ are all affected by coolant density variations
- Data is fitted with a linear regression of the form:

$$\Sigma = \Sigma^{ref} + \frac{\partial \Sigma}{\partial \rho} (\rho - \rho^{ref})$$

The Steady-State Eigenvalue Problem

- The steady state equations must be solved first
- Reducing the neutronics equation to steady-state:

$$\mathbb{M}\bar{\Phi} = \lambda\mathbb{F}\bar{\Phi}, \quad \lambda = \frac{1}{k_{eff}}$$

- Eigenvalue, λ , and eigenvector, Φ , must be determined
- Flux-power normalization constant determined from reactor power

$$Q_R = \tilde{c} \int_0^L dx \kappa \Sigma_f(x) \phi(x) = \tilde{c} \sum_i \kappa \Sigma_{f,i} \bar{\phi}_i \Delta x = \tilde{c} \kappa \Sigma_f^T \bar{\Phi} \Delta x$$

- λ and \tilde{c} are input as constants for time-dependent calculations

Newton's Method

Krylov Subspace Methods

Generalized Minimal RESidual Method

Inexact Newton's Method

Jacobian-Free Approximation

Steady State

Conclusions