

Far-from-equilibrium dynamics of molecules in ^4He nanodroplets: a quasiparticle perspective

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Institute of Science and Technology Austria

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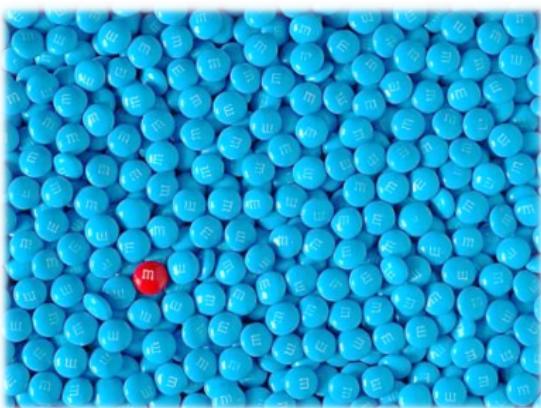
Quantum impurities

One particle (or a few particles)
interacting with a many-body
environment.

- Condensed matter
- Chemistry
- Ultracold atoms

How are the properties of the
particle modified by the interaction?

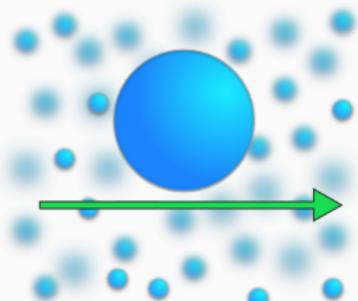
$\mathcal{O}(10^{23})$ degrees of freedom.



Quantum impurities

Structureless impurity: translational degrees of freedom/linear momentum exchange with the bath.

Most common cases: electron in a solid, atomic impurities in a BEC.



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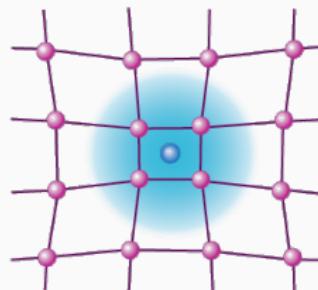


Image from: F. Chevy, Physics 9, 86.

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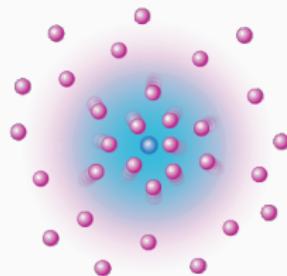


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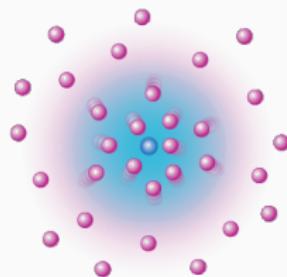
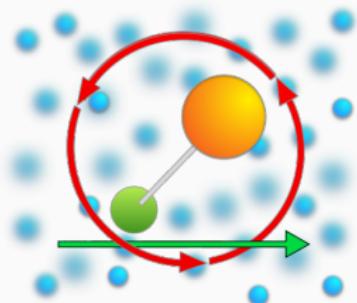


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Composite impurity (e.g. a molecule): translational *and* rotational degrees of freedom/linear and angular momentum exchange.

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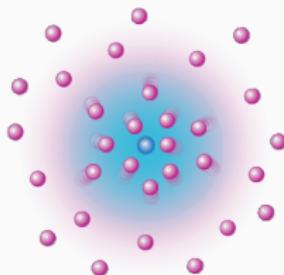
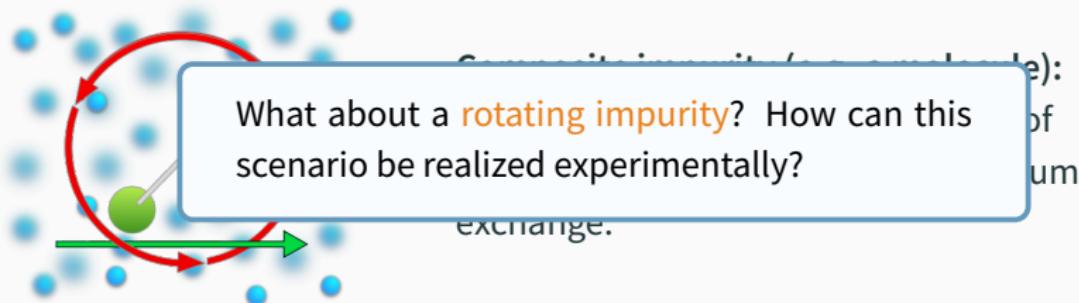


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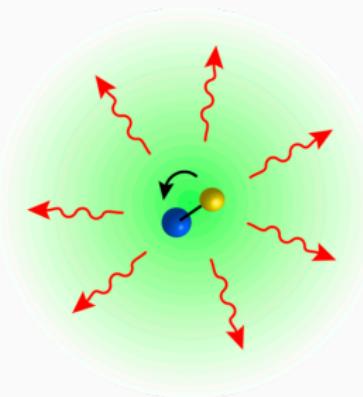


exchange.

Composite impurities: where to find them

Strong motivation for the study of composite impurities comes from many different fields. Composite impurities can be realized as:

- Ultracold molecules and ions.

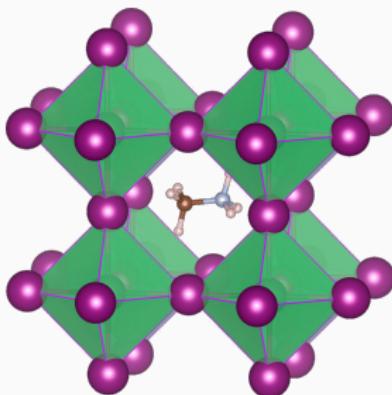


B. Midya, M. Tomza, R. Schmidt, and M. Lemeshko, Phys. Rev. A 94, 041601(R) (2016).

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- Rotating molecules inside a 'cage' in **perovskites**.



T. Chen et al., PNAS **114**, 7519 (2017).

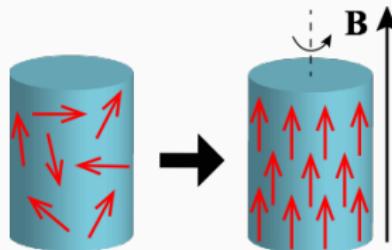
J. Lahnsteiner et al., Phys. Rev. B **94**, 214114 (2016).

Image from: C. Eames et al, Nat. Comm. **6**, 7497 (2015).

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- Rotating molecules inside a ‘cage’ in **perovskites**.
- Angular momentum transfer from the **electrons** to a **crystal lattice**.



J.H. Mentink, M.I. Katsnelson, M. Lemeshko, “*Quantum many-body dynamics of the Einstein-de Haas effect*”, Phys. Rev. B **99**, 064428 (2019).

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- Ultracold molecules and ions.
- Rotating molecules inside a 'cage' in **perovskites**.
- Angular momentum transfer from the **electrons** to a **crystal lattice**.
- **Molecules** embedded into **helium nanodroplets**.

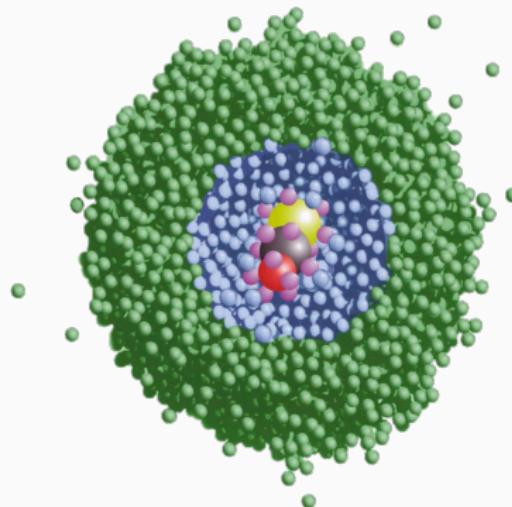
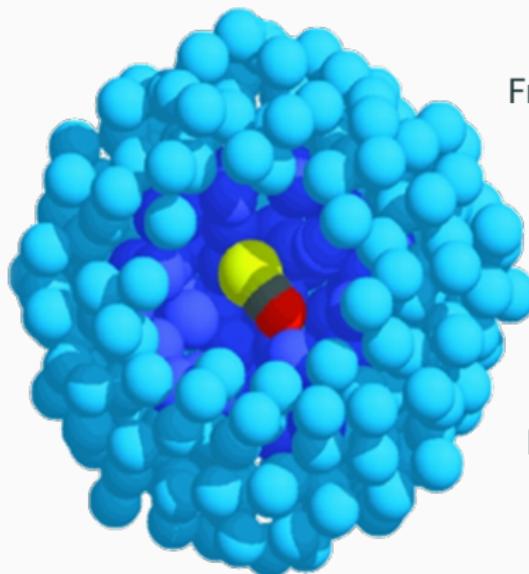


Image from: J. P. Toennies and A. F. Vilesov, Angew. Chem. Int. Ed. **43**, 2622 (2004).

Molecules in helium nanodroplets

A molecular impurity embedded into a helium nanodroplet: a controllable system to explore angular momentum redistribution in a many-body environment.



Temperature $\sim 0.4\text{K}$

Droplets are superfluid

Easy to produce

Free of perturbations

Only rotational degrees of freedom

Easy to manipulate by a laser

Image from: S. Grebenev *et al.*,
Science **279**, 2083 (1998).

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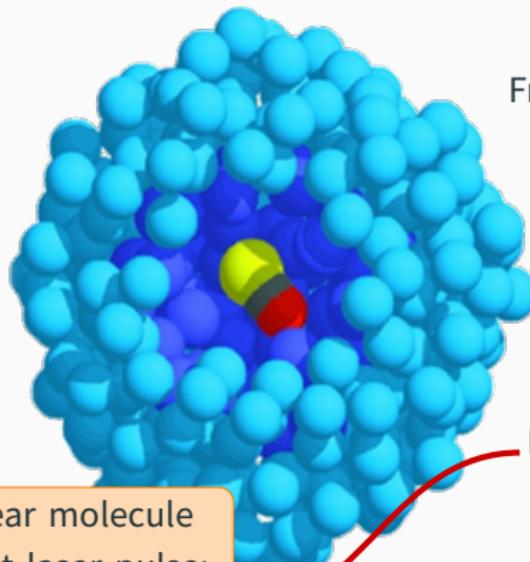
A molecular impurity embedded into a helium nanodroplet: a controllable system to explore angular momentum redistribution in a many-body environment.

Temperature $\sim 0.4\text{K}$

Droplets are superfluid

Easy to produce
Interaction of a linear molecule with an off-resonant laser pulse:

$$\hat{H}_{\text{laser}} = -\frac{1}{4}\Delta\alpha E^2(t) \cos^2 \hat{\theta}$$



Free of perturbations

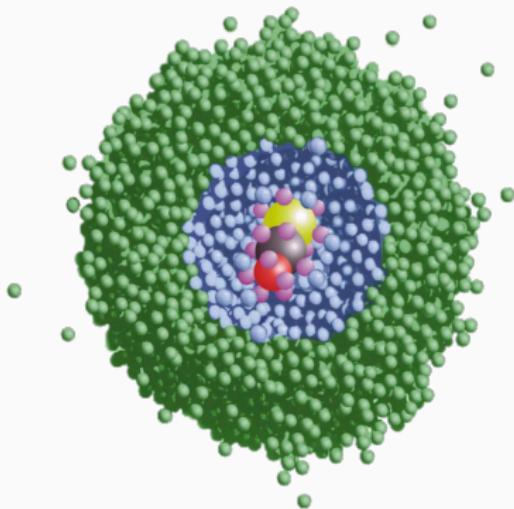
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Rotational spectrum of molecules in He nanodroplets

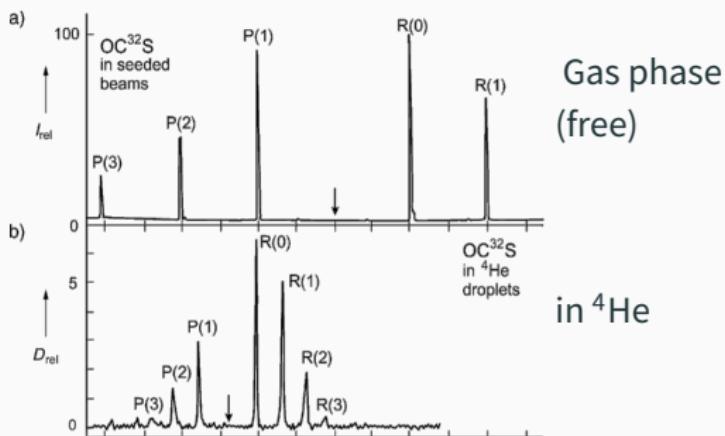
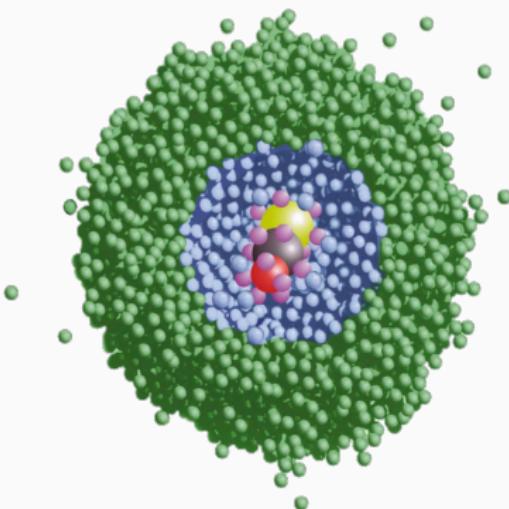
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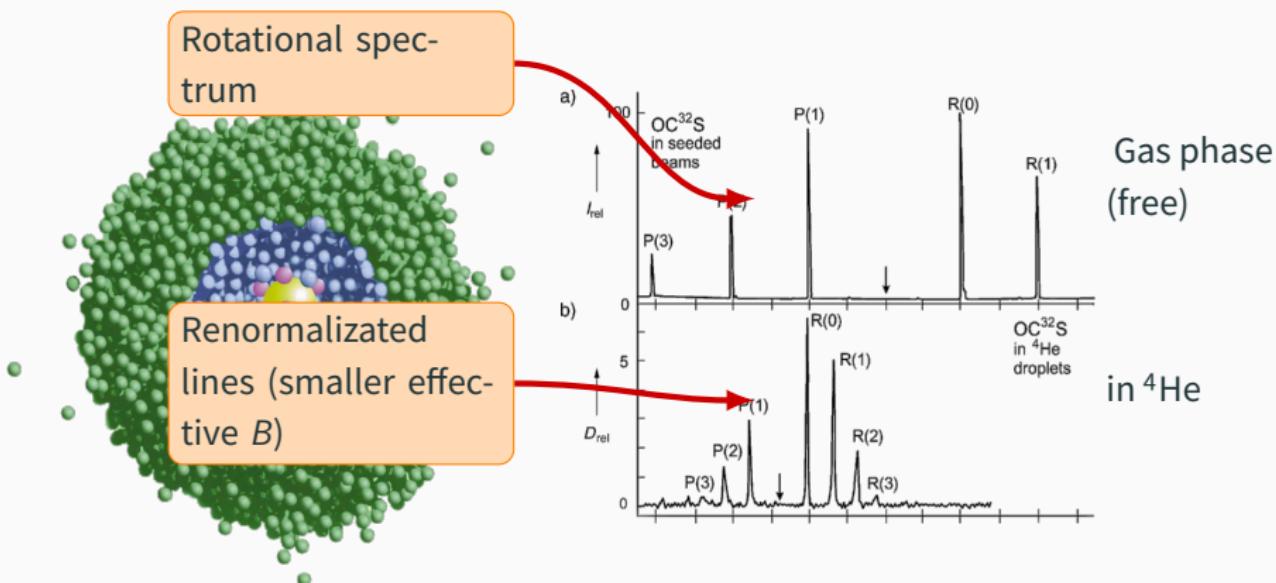
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Dynamical alignment of molecules in He nanodroplets

Dynamical alignment experiments

(Stapelfeldt group, Aarhus University):

- **Kick** pulse, aligning the molecule.
- **Probe** pulse, destroying the molecule.
- Fragments are imaged, reconstructing alignment as a function of time.
- Averaging over multiple realizations, and varying the time between the two pulses, one gets

$$\langle \cos^2 \hat{\theta}_{2D} \rangle(t)$$

with:

$$\cos^2 \hat{\theta}_{2D} \equiv \frac{\cos^2 \hat{\theta}}{\cos^2 \hat{\theta} + \sin^2 \hat{\theta} \sin^2 \hat{\phi}}$$

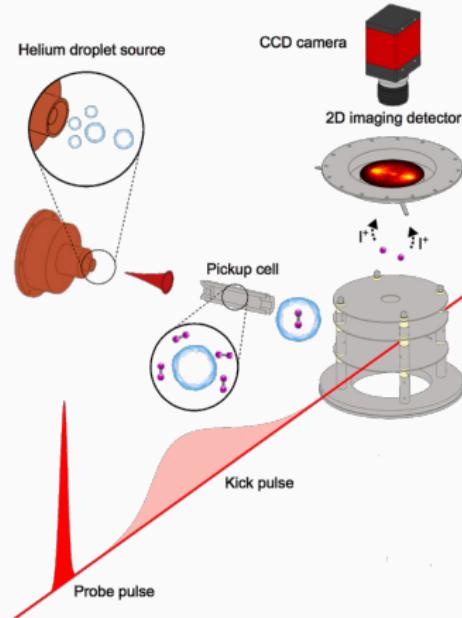
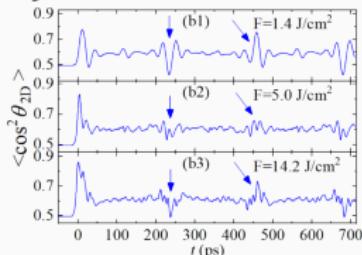


Image from: B. Shepperson *et al.*, Phys. Rev. Lett. 118, 203203 (2017).

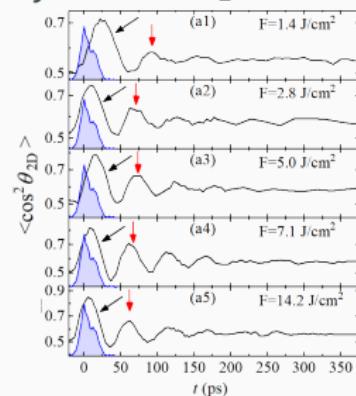
Dynamical alignment of molecules in He nanodroplets

Dynamics of isolated I₂ molecules



Experiment: Henrik Stapelfeldt, Lars Christiansen,
Anders Vestergaard Jørgensen (Aarhus University)

Dynamics of I₂ molecules in helium



Effect of the environment is substantial:

- The peak of **prompt alignment** doesn't change its shape as the fluence $F = \int dt I(t)$ is changed.
- The revival structure differs from the gas-phase: revivals with a 50ps period of **unknown origin**.
- The oscillations appear weaker at **higher fluences**.
- An intriguing **puzzle**: not even a qualitative understanding. Monte Carlo?
He-DFT?

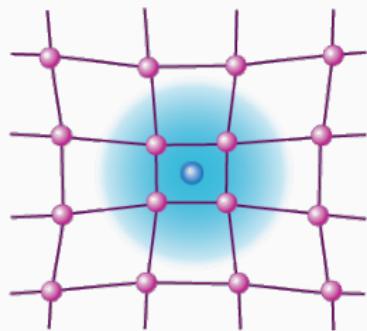
Quasiparticle approach

The quantum mechanical treatment of many-body systems is always challenging. How can one simplify the quantum impurity problem?

Quasiparticle approach

The quantum mechanical treatment of many-body systems is always **challenging**. How can one simplify the **quantum impurity** problem?

Polaron: an electron dressed by a field of many-body excitations.



Angulon: a quantum rotor dressed by a field of many-body excitations.

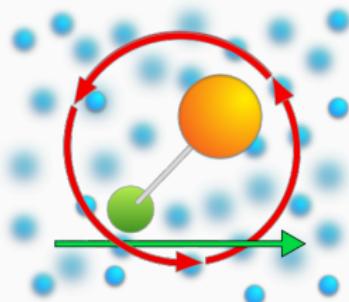


Image from: F. Chevy, Physics 9, 86.

The Hamiltonian

A **rotating linear molecule** interacting with a bosonic bath can be described in the frame co-rotating with the molecule by the following Hamiltonian:

$$\hat{\mathcal{H}} = B(\hat{\mathbf{L}} - \hat{\mathbf{\Lambda}})^2 + \sum_{k\lambda\mu} \omega_k \hat{b}_{k\lambda\mu}^\dagger \hat{b}_{k\lambda\mu} + \sum_{k\lambda} V_{k\lambda} (\hat{b}_{k\lambda 0}^\dagger + \hat{b}_{k\lambda 0}),$$

Notation:

- $\hat{\mathbf{L}}$ the total angular-momentum operator of the combined system, consisting of a molecule and helium excitations.
- $\hat{\mathbf{\Lambda}}$ is the angular-momentum operator for the bosonic helium bath, whose excitations are described by $\hat{b}_{k\lambda\mu}/\hat{b}_{k\lambda\mu}^\dagger$ operators.
- $k\lambda\mu$: angular momentum basis. k the magnitude of linear momentum of the boson, λ its angular momentum, and μ the z-axis angular momentum projection.
- ω_k gives the dispersion relation of superfluid helium.
- $V_{k\lambda}$ encodes the details of the molecule-helium interactions.

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Compare with the Lee-Low-Pines Hamiltonian

$$\hat{H}_{LLP} = \frac{(\mathbf{P} - \sum_{\mathbf{k}} \mathbf{k} \hat{b}_{\mathbf{k}}^\dagger \hat{b}_{\mathbf{k}})^2}{2m_I} + \sum_{\mathbf{k}} \omega_{\mathbf{k}} \hat{b}_{\mathbf{k}}^\dagger \hat{b}_{\mathbf{k}} + \frac{g}{V} \sum_{\mathbf{k}, \mathbf{k}'} \hat{b}_{\mathbf{k}'}^\dagger \hat{b}_{\mathbf{k}'}$$

Dynamics: time-dependent variational Ansatz

We describe dynamics using a **time-dependent variational** Ansatz, including excitations up to one phonon:

$$|\psi_{LM}(t)\rangle = \hat{U}(\mathbf{g}_{LM}(t) |0\rangle_{\text{bos}} |LM0\rangle + \sum_{k\lambda n} \alpha_{k\lambda n}^{LM}(t) b_{k\lambda n}^\dagger |0\rangle_{\text{bos}} |LMn\rangle)$$

Lagrangian on the variational manifold defined by $|\psi_{LM}\rangle$:

$$\mathcal{L} = \langle \psi_{LM} | i\partial_t - \hat{\mathcal{H}} | \psi_{LM} \rangle$$

Euler-Lagrange **equations of motion**:

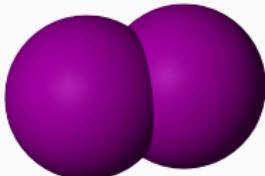
$$\frac{d}{dt} \frac{\partial \mathcal{L}}{\partial \dot{x}_i} - \frac{\partial \mathcal{L}}{\partial x_i} = 0$$

where $x_i = \{g_{LM}, \alpha_{k\lambda n}^{LM}\}$. We obtain a **differential system**

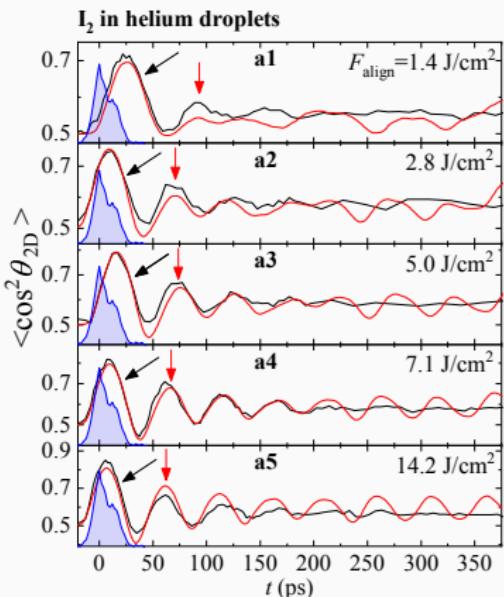
$$\begin{cases} \dot{g}_{LM}(t) = \dots \\ \dot{\alpha}_{k\lambda n}^{LM}(t) = \dots \end{cases}$$

to be solved numerically; in $\alpha_{k\lambda n}$ the momentum k needs to be discretized.

Theory vs. experiments: I₂



Comparison with experimental data from Stapelfeldt group, Aarhus University, for different molecules: I₂.

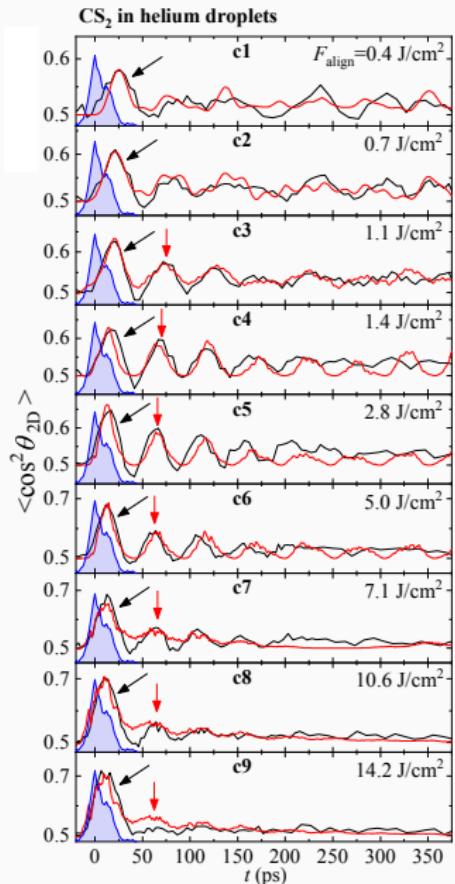


Generally good agreement for the main features in experimental data:

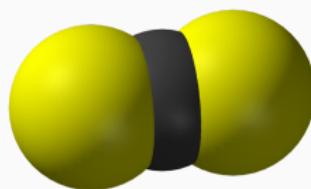
- Oscillations with a period of 50ps, growing in amplitude as the laser fluence is increased.
- Oscillations decay: at most 4 periods are visible.
- The width of the first peak does not change much with fluence.

— Experiment ■ Laser pulse
— Angulon theory

Theory vs. experiments: CS_2



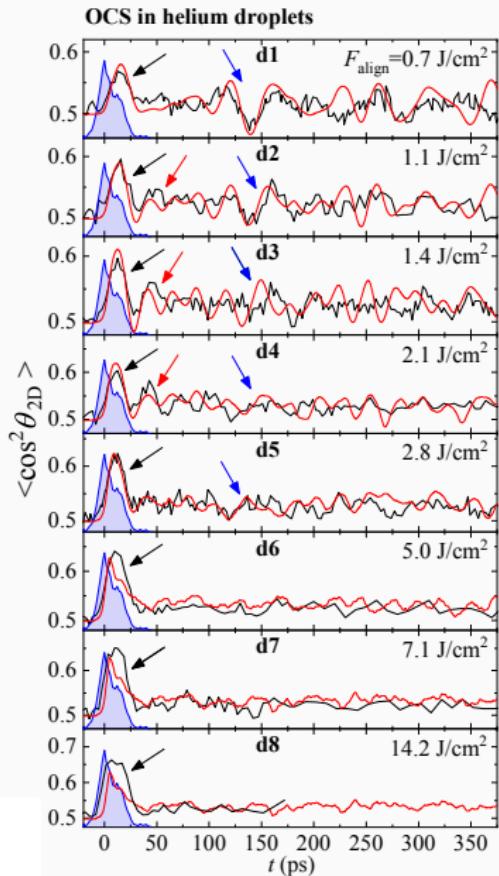
Comparison with experimental data from Stapelfeldt group, Aarhus University, for different molecules: CS_2 .



- Again, a persistent oscillatory pattern.
- For higher values of the fluence the oscillatory pattern disappears.

— Experiment — Laser pulse
— Angulon theory

Theory vs. experiments: OCS



Comparison with experimental data from Stapelfeldt group, Aarhus University, for different molecules: OCS.



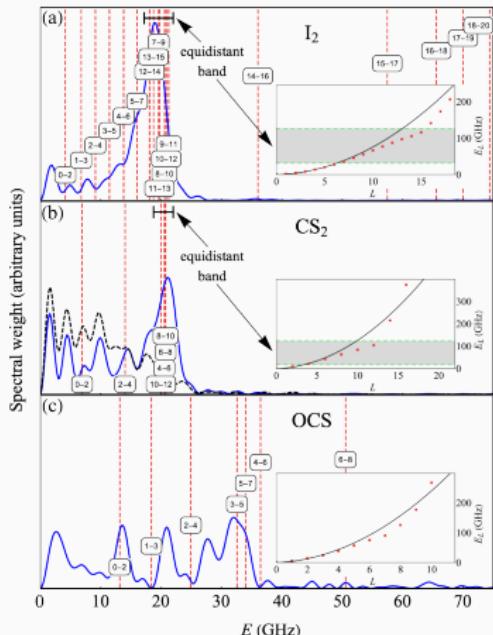
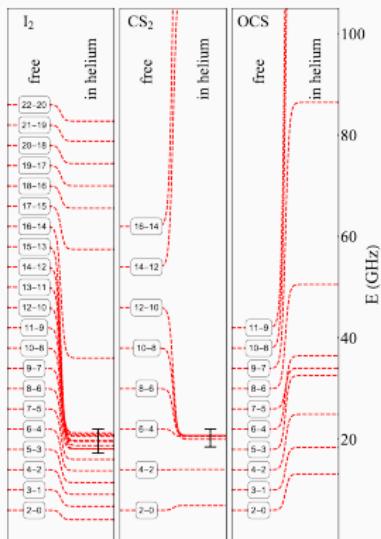
- Unfortunately the data is noisier.
- Oscillatory pattern not present, except in a couple of cases where one weak oscillation might be identified.

— Experiment ■ Laser pulse
— Angulon theory

Experiments vs. theory: spectrum

The Fourier transform of the measured alignment cosine $\langle \cos^2 \hat{\theta}_{2D} \rangle(t)$ is dominated by $(L) \leftrightarrow (L + 2)$ interferences. How is it affected when the level structure changes?

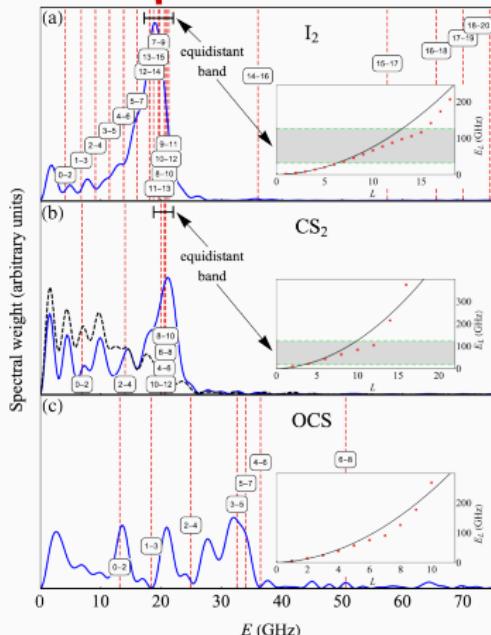
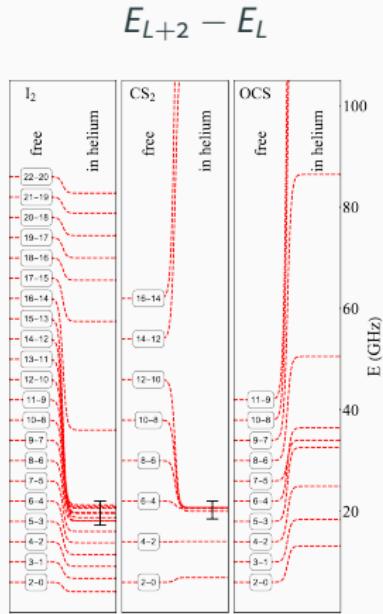
$$E_{L+2} - E_L$$



Experiments vs. theory: spectrum

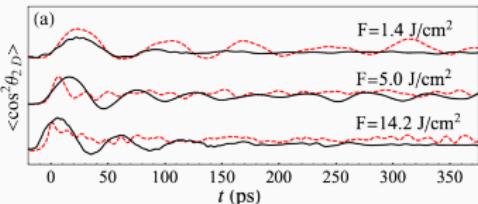
The Fourier transform of the measured alignment cosine $\langle \cos^2 \hat{\theta}_{2D} \rangle(t)$ is dominated by $(L) \leftrightarrow (L + 2)$ interferences. How is it affected when the level structure changes?

20Ghz corresponds to 50ps



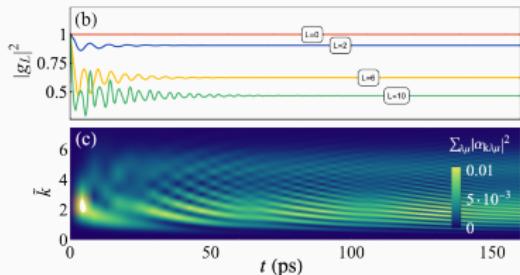
Many-body dynamics of angular momentum

i) Is this the full story? Can the observed dynamics be explained **only by means of renormalised rotational levels?**



Red dashed lines (only renormalised levels) vs. solid black line (full many-body treatment).

ii) How long does it take for a molecule to **equilibrate** with the helium environment and form an angulon quasiparticle? This requires tens of ps; which is also the **timescale of the laser!**

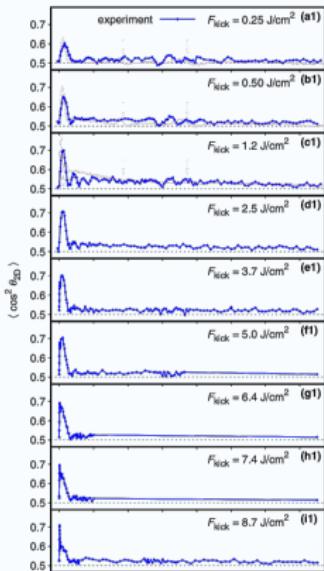


Approach to equilibrium of the quasiparticle weight $|g_{LM}|^2$ and of the phonon populations $\sum_k |\alpha_{k\lambda\mu}|^2$.

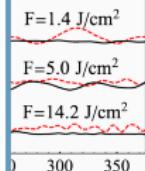
Many-body dynamics of angular momentum

i) Is this the fundamental dynamics being renormalised

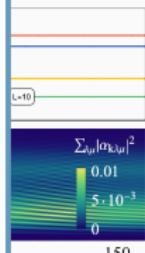
With a shorter 450 fs pulse, same molecule (I_2), the strong oscillatory pattern is absent:



ii) How long does it take for the system to equilibrate with the environment and form an angular momentum state? This requires tens of picoseconds, which is a timescale of the order of the cavity lifetime.



Population levels (red/blue) vs. time (blue) after treatment).



Biparticle populations

Image from: B. Shepperson *et al.*, Phys. Rev. Lett. **118**, 203203 (2017).

Conclusions

- A novel kind of pump-probe spectroscopy, based on **impulsive molecular alignment** in the laboratory frame, providing access to the structure of highly excited rotational states.
- Superfluid bath leads to formation of **robust long-wavelength oscillations** in the molecular alignment; an explanation requires a **many-body theory** of angular momentum redistribution.
- Our theoretical model allows us to interpret this behavior in terms of the dynamics of angulon quasiparticles, shedding light onto many-particle **dynamics of angular momentum at femtosecond timescales**.
- Future perspectives:
 - All molecular geometries (spherical tops, asymmetric tops).
 - Optical centrifuges and superrotors.
 - Can a rotating molecule create a vortex?
- For more details: arXiv:1906.12238



Institute of Science and Technology



Lemeshko group @ IST Austria:



Misha
Lemeshko

Dynamics in He



Enderalp
Yakaboylu



Xiang Li



Igor
Cherepanov



Wojciech
Rządkowski



Dynamical alignment
experiments

Collaborators:



Henrik
Stapelfeldt
(Aarhus)



Richard
Schmidt
(MPI Garching)

Thank you for your attention.



Institute of Science and Technology



Der Wissenschaftsfonds.

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These slides at <http://bigh.in>