

# Reduced terrestrial evaporation increases atmospheric water vapor by generating cloud feedbacks

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## Abstract

Reduced terrestrial evaporation directly warms the surface by reducing latent cooling, but also indirectly modifies surface climate by altering atmospheric processes. We use a global climate model to explore two end cases of terrestrial evaporation, comparing the climate of SwampLand, a world where land is always fully saturated with water, to that of DesertLand, where land is always completely lacking in soil moisture. When we suppress evaporation to create a desert-like planet, we find that temperatures increase and precipitation decreases in the global mean. We find an increase in atmospheric water vapor over both land and ocean in the DesertLand simulation. Suppressing evaporative cooling over the continents reduces continental cloud cover, allowing more energy input to the surface and increasing surface moist static energy over land. The residence time of atmospheric water vapor increases by about 50 percent. Atmospheric feedbacks such as changes in air temperatures and cloud cover contribute larger changes to the terrestrial surface energy budget than the direct effect of suppressed evaporation alone. Without the cloud feedback, the land surface still warms with suppressed land evaporation, but total atmospheric water vapor decreases, and the anomalous atmospheric circulations over the continents are much shallower than in simulations with cloud changes; that is, the cloud feedback changes the sign of the water vapor response. This highlights the importance of accounting for atmospheric feedbacks when exploring land surface change impacts on the climate system.

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## 29    1 Introduction

30    Changes in terrestrial evapotranspiration directly impact climate. Reducing evapo-  
 31    ration from the land surface has a direct warming effect by reducing energy loss from  
 32    the land surface (Shukla & Mintz, 1982; Fraedrich et al., 1999; Davin et al., 2010; Laguë  
 33    et al., 2019). However, in idealized continental configurations with large land masses, re-  
 34    ducing terrestrial evaporation can instead drive terrestrial cooling by reducing atmospheric  
 35    water vapor concentrations and the strength of the water vapor greenhouse effect (Laguë,  
 36    Pietschnig, et al., 2021). Changes in atmospheric temperatures and water vapor driven  
 37    by changes in terrestrial evaporation are not only important for terrestrial surface cli-  
 38    mate, but also for the global atmospheric energy budget, as changes in atmospheric tem-  
 39    peratures, moisture content, and cloud cover driven by terrestrial processes alter the global  
 40    radiative balance of the atmosphere (Swann et al., 2010, 2012; Boos & Korty, 2016; Laguë  
 41    & Swann, 2016; Laguë, Swann, & Boos, 2021).

42    Terrestrial processes are an integral part of the global water cycle. Water evapo-  
 43    rates from the oceans, is transported by the atmosphere, and falls as precipitation over  
 44    the land. Water on the land surface is evaporated or transpired to the atmosphere, stored  
 45    as ground water, or returned to the ocean as runoff. Terrestrial evapotranspiration is de-  
 46    termined by a combination of soil moisture, atmospheric demand for water, terrestrial  
 47    re-distribution of water and the evaporative properties of vegetation and soils (Monteith,  
 48    1965; Bonan, 2008; Eltahir & Bras, 1996), while large-scale climate features, topogra-  
 49    phy, and soil properties modulate soil water available for evapotranspiration (see Kot-  
 50    tek et al., 2006, and references therein).

51    Terrestrial evaporation can be limited both by the availability of water to evapo-  
 52    rate and by energy (Budyko, 1961; Vargas Zeppetello et al., 2019). Vegetation can di-  
 53    rectly modulate transpiration by opening and closing stomata (Jones, 1998; Sellers et  
 54    al., 1996; Pielke et al., 1998); transpiration can change with both water availability and  
 55    vegetation controls on leaf area (Bonan, 2008), stomatal properties (Ball et al., 1987; Med-  
 56    llyn et al., 2011), and root depth (Lai & Katul, 2000). Evapotranspiration (ET) has changed  
 57    over the historical period (Hobeichi et al., 2021), with some regions (e.g. tropical Africa)  
 58    exhibiting a negative trend in ET from 1980-2018, while other regions (e.g. Europe) show  
 59    a positive trend. Earth system models project more changes in the future (Swann et al.,  
 60    2016; Berg et al., 2016), though the response is complex and varies across models, with  
 61    stomatal closure in response to increased atmospheric CO<sub>2</sub> acting to reduce ET and in-  
 62    creased leaf area and atmospheric demand for water acting to increase ET. Changes in  
 63    evapotranspiration are an expected result of land use change (Wang et al., 2021), veg-  
 64    etation responses to increased atmospheric CO<sub>2</sub> (Field et al., 1995; Sellers et al., 1996;  
 65    Norby et al., 2010; Donohue et al., 2013; Lemordant et al., 2018), and climate change  
 66    (Swann et al., 2016; Berg et al., 2016; Collins et al., 2013).

67    How changes in terrestrial evaporation relate to the water cycle both regionally and  
 68    globally remains an area of active research (Swann et al., 2016; Berg et al., 2016; Koster  
 69    et al., 2006; Dirmeyer, 2011, 1994, 2006; Byrne & O’Gorman, 2015, 2016; Seneviratne  
 70    et al., 2010; Laguë, Pietschnig, et al., 2021). Total atmospheric water vapor is projected  
 71    to increase in simulations of global warming (Sherwood et al., 2010); while relative hu-  
 72    midity over the oceans is expected to remain roughly constant, relative humidity over  
 73    land is expected to decrease (O’Gorman & Muller, 2010a; Byrne & O’Gorman, 2016).  
 74    Independent of the radiative effects of CO<sub>2</sub>, plant responses to increased atmospheric  
 75    CO<sub>2</sub> are projected to reduce near-surface relative humidity on land (Swann et al., 2016).

76    In this study, we explore the effect of extreme end-cases of terrestrial evaporation  
 77    on global climate for the modern continental configuration: land that is always fully sat-  
 78    urated (all land looks like a swamp), versus land that is always fully desiccated (all land  
 79    looks like a desert). We find that fully suppressing terrestrial evaporation leads to in-  
 80    creased water vapor concentrations throughout the atmospheric column over most con-

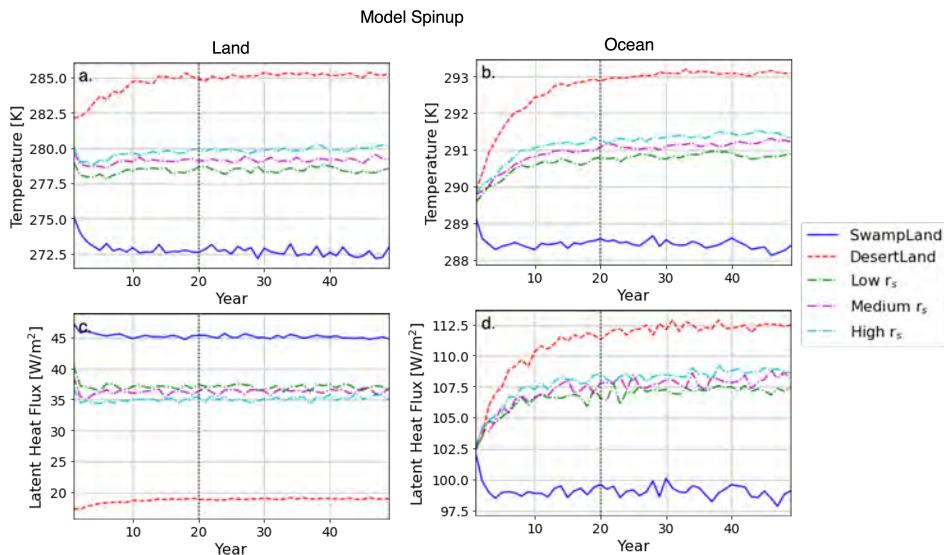
81 continental and ocean regions. While terrestrial relative humidity decreases with suppressed  
 82 evaporation, strong cloud feedbacks enhance the energy content and specific humidity  
 83 of air over land.

## 84 2 Methods

85 We conduct experiments using two climate models to study how changes in land  
 86 evaporation impact the atmosphere. We use a radiative kernel to decompose the impact  
 87 of temperature, moisture, cloud, and albedo changes on the atmospheric energy budget.

### 88 2.1 Models

89 We use a modified version of the Community Earth System Model (CESM) (Hurrell et al., 2013), comprised of the Community Atmosphere Model v. 5 (CAM5) coupled  
 90 to a slab ocean model (Neale et al., 2012), the CICE5 interactive sea ice model (Bailey  
 91 et al., 2018), and the Simple Land Interface Model (SLIM) (Laguë et al., 2019). Sim-  
 92 ulations are run at  $2.5^\circ$  resolution for 50 years, with the first 20 years discarded to al-  
 93 low for model spin-up. After 20 years, there is  $< 0.1$  K drift in global mean surface tem-  
 94 peratures (Fig. 1); the top of atmosphere energy imbalance in these near-equilibrium sim-  
 95 ulations is near-zero ( $\approx 0.3$  W/m $^2$ ).  
 96



**Figure 1.** Annual mean, spatially averaged surface temperature (top) and latent heat flux (bottom) over land (left) and ocean (right) regions for the five CESM simulations explored in this study. Simulations have reached equilibrium prior to year 20.

97 The slab ocean has prescribed, seasonally and spatially varying pre-industrial ocean  
 98 heat transport and heat capacity (slab or mixed-layer “depth”) which is identical between  
 99 simulations and from year-to-year within the same simulation. This allows sea surface  
 100 temperatures to evolve in response to forcings, but does not allow for changes in ocean  
 101 heat transport. The mixed-layer depth and heat transport values used in these sim-  
 102 ulations are calculated from the dynamic ocean component of the pre-industrial control sim-  
 103 ulation of the fully coupled CESM 1.2 model (as used in Garcia et al., 2016; Swann et  
 104 al., 2018; Laguë et al., 2019). We use slab ocean values from a pre-industrial vs. present-  
 105 day simulation because the pre-industrial climate is in equilibrium, while the ocean acts

as a net energy sink in the present day climate. The ocean is a large source of variability in the Earth system, and the use of a slab ocean model allows us to focus on the atmospheric response to land surface changes, as the internal variability of the slab ocean model is small, allowing us to use a single simulation for each experiment rather than an ensemble. To explicitly demonstrate this, we run the initial 10 years of the two simulations described below (“SwampLand” and “DesertLand”) with three ensemble members, each with initial land surface temperatures perturbed by  $1e-6$  K; the spread between ensemble members is very small compared to the difference between the two experiments (Fig. S1). The default ocean albedo parameterization is used, where ocean albedo varies with solar declination, with default values of 0.06 for direct and 0.07 for diffuse radiation.

SLIM is used to allow us to directly control the physical properties of the land surface in a way that is difficult with complex land surface models. Hydrology is represented using a bucket model, with the resistance to evaporation calculated as a function of how much water is in the bucket as well as an additional user-prescribed resistance. A simple snow model allows for snow-albedo feedbacks on the land surface. The idealized land surface model allows us to artificially control terrestrial water availability without altering other aspects of the land surface. In contrast, if we were to override soil moisture in a complex land surface model like CLM5 (Lawrence et al., 2019), this would have follow-on impacts on leaf area, plant hydraulics, and the carbon cycle, which would in turn have follow-on impacts on albedo and surface evaporative resistance. In order to isolate the effect of surface water availability on climate, we need to leverage an idealized land surface model like SLIM.

We also want to test the response of the climate system to changes in land evaporation without cloud responses, which are a large source of model uncertainty. We could force clouds in CESM to be transparent to radiation, but this would lead to an unreasonably dark top of atmosphere albedo (unless we *also* modified surface albedo) and a much hotter base-state climate than CESM produces with the normal cloud parameterization. Instead, we conduct two additional simulations using Isca (Vallis et al., 2018), an idealized global circulation model which has radiatively interactive water vapor and produces precipitation, but where clouds are “invisible” to radiation in the standard, tested configuration. Simulations use a T42 grid, a slab ocean with a 20m mixed layer depth and prescribed modern heat transport, a bucket model for land hydrology with a heat capacity equal to that of a 2m ocean mixed layer, no snow albedo feedbacks, the modern continental configuration, realistic topography, and the RRTM radiation scheme (Clough et al., 2005). The albedo of the surface is set to be much higher than realistic values (0.25 for ocean and 0.325 for land), to generate a reasonable climate and top of atmosphere albedo (Thomson & Vallis, 2019; Geen et al., 2018).

## 2.2 Simulations

We consider five CESM simulations in this study, primarily focusing on two end-members of wet/dry land. Simulations have globally uniform land surface properties: all non-glaciated points on the land surface have the same albedo, aerodynamic roughness, capacity to hold water, and evaporative resistance. Glaciated gridcells (Greenland and Antarctica) have surface properties associated with ice, and do not vary between simulations (see Lagu   et al., 2019, for details). The snow-free albedo of the land surface is set to 0.2 in visible wavelengths and 0.3 in the near-infrared, while the aerodynamic roughness is set to 0.1 m. While the base-state climate differs regionally in these simulations with globally uniform land surface properties compared to simulations with present-day land surface properties (see Lagu   et al., 2019), we choose to use uniform, idealized land surface properties in order to isolate the sensitivity of the atmosphere to only the change in water availability at the land surface, without introducing the added dimen-

157 sion of imposing the difference in water availability onto different land surface types (e.g.  
 158 land with different albedo).

159 In DesertLand, the first of our extreme simulations, the capacity of the land to hold  
 160 water is reduced to 20 mm everywhere (compared to a typical value of  $\approx$  200 mm), and  
 161 the resistance to evaporation is set to 100,000 s/m (compared to a typical value of  $\approx$  100  
 162 s/m); this effectively turns off evaporation from the land surface, regardless of precip-  
 163 itation or the atmospheric demand for water. DesertLand can physically be thought of  
 164 as land free of vegetation with extremely well draining soils, such that all precipitation  
 165 that falls on the land is quickly transferred into below ground aquifers or sub-surface runoff  
 166 and returned to the oceans. In SwampLand, the second extreme simulation, the land sur-  
 167 face is forced to be fully saturated with water at every time step. Land always has 200  
 168 mm of water available for evaporation, regardless of the precipitation or evaporation rates  
 169 at each point. SwampLand can be thought of as land with a high water table and un-  
 170 limited ground water supply. Physically, this is comparable to swampy regions on the  
 171 modern land surface, but in this idealized simulation, these swamps are imposed over  
 172 the entire non-glaciated land surface, regardless of elevation, slope, or distance from a  
 173 water body.

174 Three additional simulations with interactive surface hydrology are also briefly con-  
 175 sidered, differing in their prescribed evaporative resistance: 30 s/m (“low” for low re-  
 176 sistance), 100 s/m (“medium” for medium resistance), and 200 s/m (“high” for high re-  
 177 sistance). Each simulation has the capacity to hold 200 mm of water at each non-glaciated  
 178 land surface point, but the amount of water actually on the land surface is interactively  
 179 simulated by the model based on precipitation and evaporation at each location.

180 Two additional simulations, similar to DesertLand and SwampLand, are conducted  
 181 in the Isca model. In the DesertLand Isca simulation, the bucket capacity is set to 0.01  
 182 mm, while the bucket capacity in SwampLand is set to 150 mm and is “topped up” to  
 183 always have 150 mm of water available at all land points at each time step.

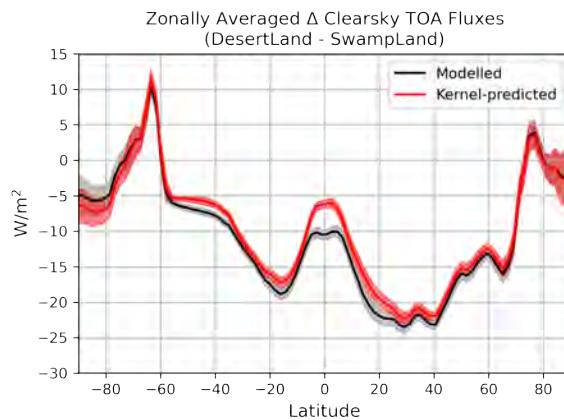
### 184 2.3 Radiative Kernel

185 Changing surface evaporation between simulations alters atmospheric and surface  
 186 temperatures, water vapor, cloud cover, and snow and ice extent. To isolate the individ-  
 187 ual contributions of each of these responses to the atmospheric energy budget, we use  
 188 a radiative kernel for the CAM5 atmospheric model (Pendergrass et al., 2018) and fol-  
 189 low the procedure introduced in Lagu  , Swann, & Boos (2021), which we summarize here.

190 The radiative kernel provides the change in top-of-atmosphere (TOA) net short-  
 191 wave and longwave radiation under both “full-sky” (including the effects of clouds) and  
 192 “clear-sky” (without the effects of clouds) conditions that result from independent changes  
 193 in the following quantities: surface albedo; surface temperature; air temperature at each  
 194 level of the atmosphere; and specific humidity given a unit change in air temperature at  
 195 constant relative humidity at each level of the atmosphere. The kernel also provides the  
 196 change in downwelling shortwave and longwave radiation at the surface associated with  
 197 each of these perturbations. Multiplying our simulated change in temperature, water va-  
 198 por, etc. by the kernel returns the effect of that change on TOA or surface radiative fluxes.

199 We mask out differences in the stratosphere between simulations (as in Pendergrass  
 200 et al., 2018; Shell et al., 2008), and multiply the water vapor kernel by the change in the  
 201 natural logarithm of the simulated change in water vapor. We apply the clear-sky lin-  
 202 earity test (Vial et al., 2013) to our most extreme simulations (SwampLand and Desert-  
 203 Land) and find generally excellent agreement between the change in TOA fluxes simu-  
 204 lated by the full model and those predicted by the radiative kernel (Fig. 2). The excep-  
 205 tion is in the deep tropics, where the kernel and model are qualitatively similar, but dis-  
 206 agree by a few W/m<sup>2</sup>; these changes are primarily driven by disagreements in the trop-

207      ical Atlantic and over tropical Africa (not shown). Note that all simulations here have  
 208      a net input of energy into the TOA in the low latitudes and a net removal of energy from  
 209      the TOA in the high latitudes under both clear-sky and full-sky conditions; Fig. 2 shows  
 210      the *difference* in the TOA energy balance between two simulations. In the high latitudes,  
 211      the positive values on this graph show that DesertLand is less negative than SwampLand,  
 212      i.e. DesertLand is losing less energy at this latitude than SwampLand is. The negative  
 213      values in the low latitudes show that DesertLand is less positive than SwampLand, i.e.  
 214      that SwampLand is absorbing more total net energy into the Earth System at low lat-  
 215      itudes compared to DesertLand. When we show results of calculations using the radia-  
 216      tive kernel, any non-linearities or residuals in the radiative kernel are necessarily included  
 217      in the cloud term; however, the excellent agreement in the clear-sky linearity test gives  
 218      us confidence that such residuals are small.



**Figure 2.** Zonal mean change (DesertLand - SwampLand) in TOA clear-sky radiation directly from the model (gray) and predicted by the clear-sky radiative kernel (red). Shading indicates  $\pm 1\sigma$  of inter-annual standard deviation.

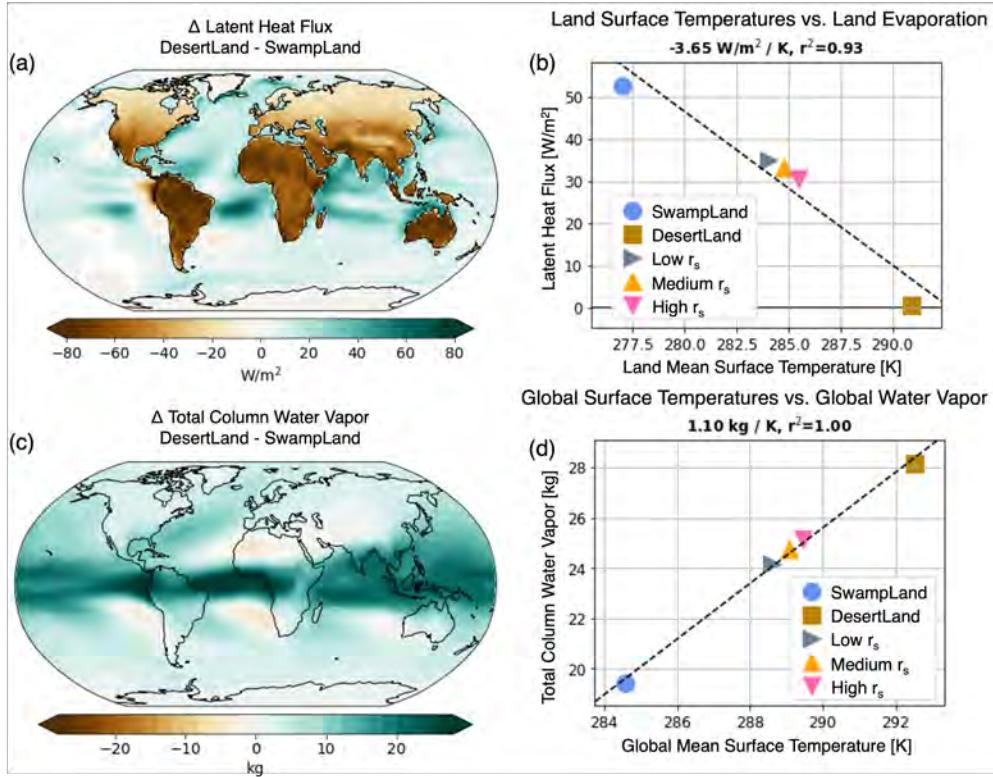
219      Analysis was conducted with the Python programming language, primarily with  
 220      the NumPy (Harris et al., 2020) and xarray (Hoyer & Hamman, 2017) packages, using  
 221      the JupyterHub (<https://jupyter.org/>) service on the Cheyenne computing system  
 222      (Computational and Information Systems Laboratory, 2019). When statistical signifi-  
 223      cance is shown on maps and vertical cross-sections, a value is deemed statistically sig-  
 224      nificant if the p-values calculated from a Student's t-test pass a false discovery rate of  
 225      0.15 (following Wilks, 2016). Uncertainty intervals indicate  $\pm 1\sigma$  of inter-annual stan-  
 226      dard deviation.

### 227      3 Results & Discussion

#### 228      3.1 Column water vapor increases with suppressed land evaporation

229      As expected, SwampLand—the simulation with perpetually saturated land—has  
 230      the largest terrestrial evaporation of the simulations considered, and the lowest average  
 231      land surface temperatures (Fig. 3). In contrast, DesertLand—the simulation with per-  
 232      petually suppressed land evaporation—has the lowest terrestrial evaporation (by design),  
 233      and the warmest surface temperatures both on land and globally (Fig. 3a,b). The three  
 234      simulations with intermediate values of terrestrial evaporative resistance lay between Swamp-  
 235      Land and DesertLand in terms of evaporation, surface temperatures, and total atmo-  
 236      spheric water vapor. There is a strong linear relationship across the simulations between

terrestrial evaporation and terrestrial surface temperature, and between global mean surface temperatures and total atmospheric water vapor (Fig. 3c,d).



**Figure 3.** Annual mean change in latent heat flux (a) and total column water vapor (c) for the DesertLand - SwampLand simulations. Scatter plots showing the relationship between annual mean (b) land surface temperature and terrestrial evaporation and (d) global mean surface temperature and global mean total column water vapor. In a/c, only changes that pass a statistical test are shown, where values are significant if the p-values calculated from a student's t-test pass a false discovery rate of 0.15.

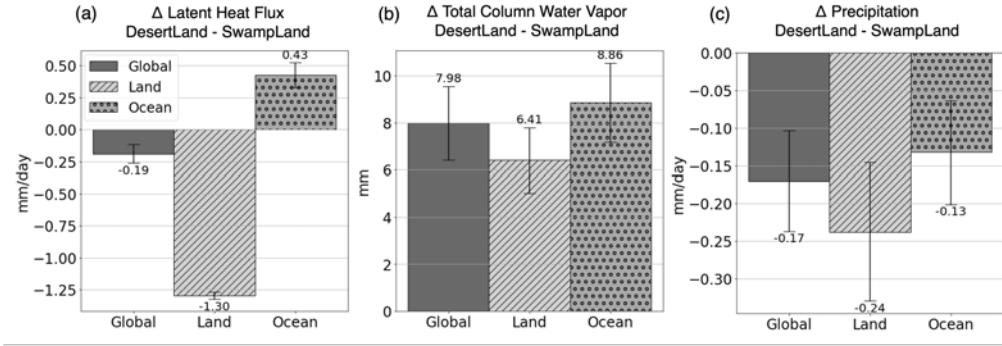
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DesertLand has the most atmospheric water vapor, despite having suppressed land evaporation (Figs. 3). The planet as a whole is not water limited in the modern continental configuration, so ocean evaporation increases in the DesertLand simulation (Fig. 4a). However, this only partially compensates for the reduction in land evaporation, so there is less surface evaporation in the global mean in DesertLand. Precipitation over both land and ocean is accordingly reduced in DesertLand, but the atmosphere has more total water vapor in DesertLand than SwampLand (Fig. 4c). This is true of both land and ocean regions in the lower and upper troposphere, except for a drying of some subtropical regions and in the lower troposphere of inland continental regions (Figs. 5, 6).

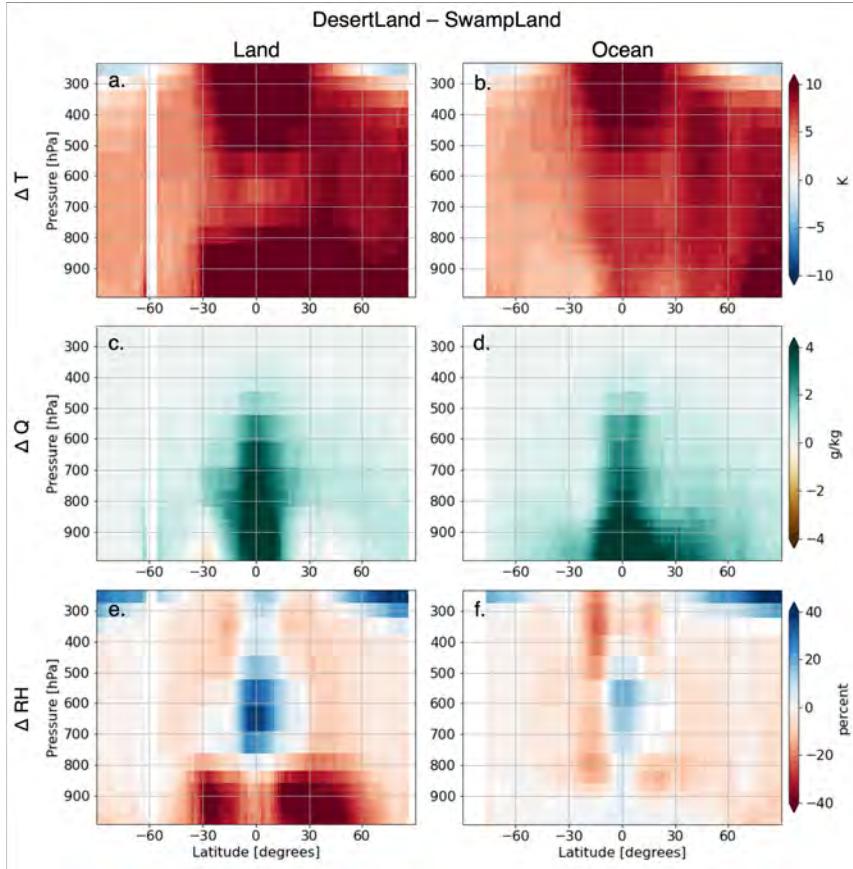
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250  
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The reduction in global mean precipitation and increase in global mean water vapor content together imply an increase in the residence time of atmospheric water vapor. Specifically, this residence time has been defined as the ratio of global mean precipitable water  $Q$  to global mean precipitation  $P$  (Trenberth, 1998),

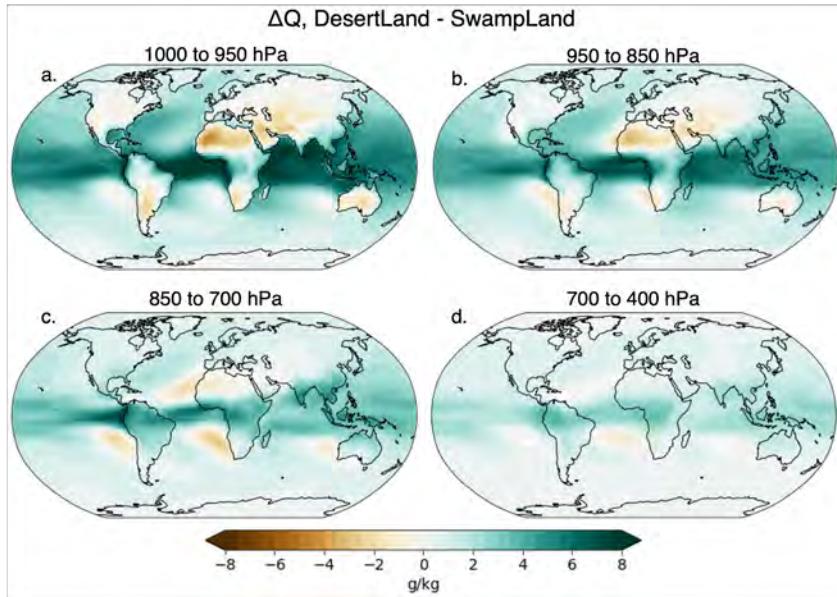
$$\tau \equiv \frac{Q}{P} \quad (1)$$



**Figure 4.** Annual mean change (DesertLand - SwampLand) in area-weighted mean (a) latent heat flux [mm/day], (b) column water vapor [mm], and (c) precipitation [mm/day], separated into the global (solid), land (hatched), and ocean (dotted) contributions. Error bars indicate  $\pm 1\sigma$  of inter-annual standard deviation. Note that as there is more ocean than land area, bars within each subplot do not sum directly.



**Figure 5.** Zonal mean vertical cross sections of the change in temperature (T [K], top), specific humidity (Q [g/kg], middle), and relative humidity (RH [%], bottom) over land regions (left) and ocean regions (right). Only changes that pass a statistical test are shown, where values are significant if the p-values calculated from a student's t-test pass a false discovery rate of 0.15.

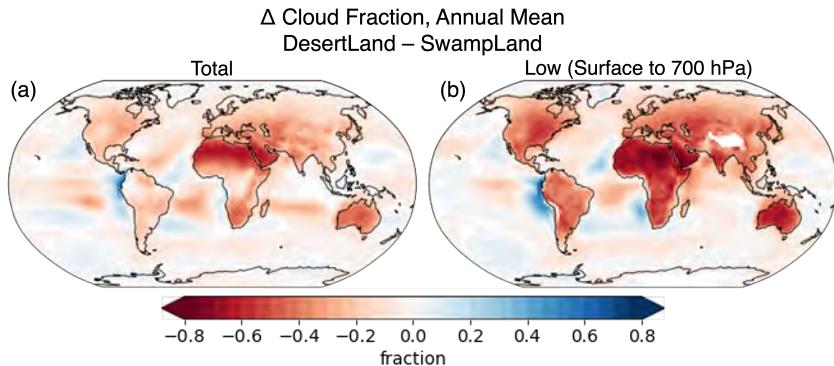


**Figure 6.** Change (DesertLand - SwampLand) in specific humidity [g/kg] from (a) 1000-950 hPa, (b) 950-850 hPa, (c) 850-700 hPa, and (d) 700-450 hPa. Only changes that pass a statistical test are shown, where values are significant if the p-values calculated from a student's t-test pass a false discovery rate of 0.15.

Here we find that  $\tau$  increases from 6.7 days in SwampLand to 10.2 days in DesertLand. One may alternatively interpret this change as a reduction in the convective mass flux that transports water vapor vertically until it condenses and precipitates (Held & Soden, 2006). Thus, while a reduction in land evaporation is expected to produce a transient reduction in local atmospheric water vapor, changes in the precipitating atmospheric circulation dominate to allow more water to accumulate in the atmosphere and then be maintained at that higher level. Locally, water vapor over land is maintained by a balance between local evaporation, local precipitation, and the convergence of water by large-scale winds; the latter two are typically large compared to the former, making possible indirect effects of a surface evaporative forcing. Furthermore, Sun & Wang (2022) found a reduction in precipitation intensity in hot weather in moisture limited regions (e.g. over land) as a result of an increased saturation deficit.

### 3.2 Cloud feedbacks enhance energy input over land

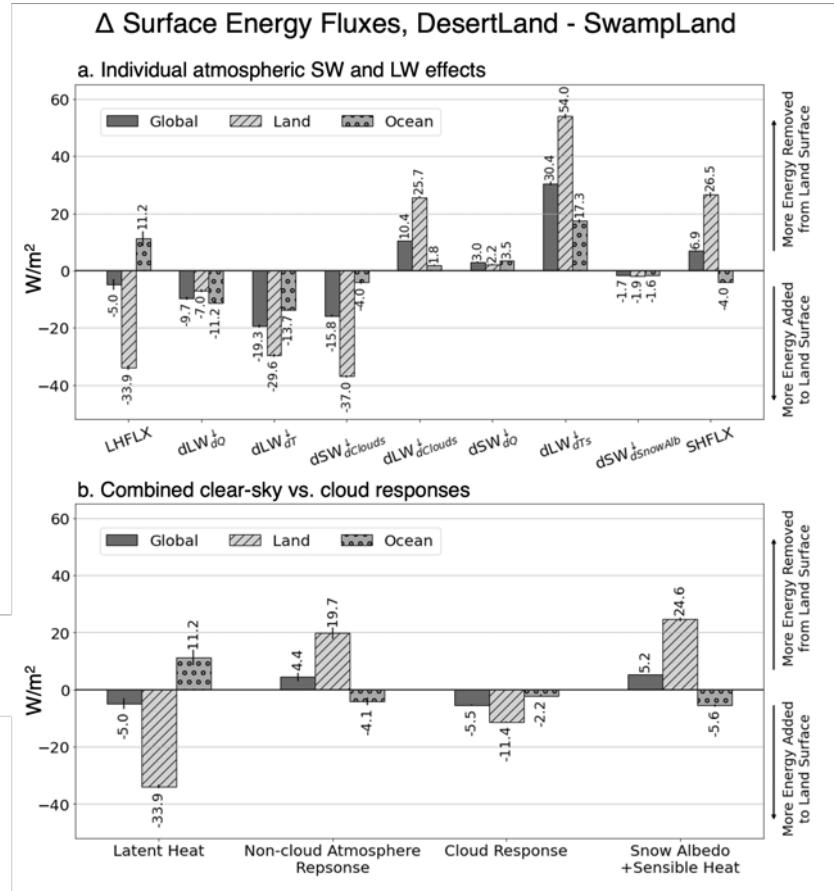
We now describe how suppression of surface evaporation produces a reduction in low cloud cover (Fig. 7), increasing the energy absorbed by land and the precipitating large-scale circulation driven by that energy source. Reductions in cloud cover driven by suppressed terrestrial evaporation lead to an additional  $\approx 37 \text{ W/m}^2$  of shortwave radiation absorbed by the surface (Fig. 8). The largest reductions occur in low clouds over land, and are thus consistent with a local response to the land evaporative forcing. The reduced cloud cover also makes it easier for longwave radiation emitted by land to exit the top of the atmosphere, but this is smaller in magnitude than the shortwave cloud effect (as expected given the low altitude of the cloud changes), yielding a net positive heating of land by cloud radiative effects of about  $11 \text{ W/m}^2$ . Other components of the surface energy budget are discussed below, when we use the model's radiative kernel to decompose the net change into different physical components.



**Figure 7.** Annual mean change (DesertLand - SwampLand) in total cloud fraction (left) and low cloud fraction (right) for the CESM simulations. Only changes that pass a statistical test are shown, where values are significant if the p-values calculated from a student's t-test pass a false discovery rate of 0.15.

277 Reductions in low cloud cover over land as a result of suppressing terrestrial evap-  
 278 oration also impact the TOA energy budget (Fig. 9). In particular, loss of low cloud cover  
 279 over land lowers the planetary albedo; this results in more energy absorbed at the TOA  
 280 and in turn should increase global mean temperatures. Increased temperatures should,  
 281 in turn, lead to an increase in water vapor following the Clausius-Clapeyron relationship.  
 282 That is, the reduction in low cloud cover alone should lead to increased atmospheric tem-  
 283 peratures and water vapor. Without the reduction in low cloud cover, suppressing ter-  
 284 restrial evaporation would not necessarily lead to any increased energy into the Earth  
 285 system at the top of the atmosphere, and would lead to weaker warming at the surface  
 286 as the only warming mechanism would be the reduction in latent cooling, with no ad-  
 287 dditional warming from increased solar radiation due to reduced cloud cover. Indeed, we  
 288 see this later when we conduct similar simulations in a model without cloud cover.

289 We roughly estimate the effect of the change in clouds on global mean tempera-  
 290 tures and water vapor as follows: in the global mean, the change in clouds leads to a 12.8  
 291 W/m<sup>2</sup> increase in energy into the Earth system at the TOA (14.8 W/m<sup>2</sup> from shortwave  
 292 radiation, -2 W/m<sup>2</sup> from longwave radiation; Fig. 9). Using a climate response param-  
 293 eter of 1.5 W/m<sup>2</sup>/K (from Gregory (2004)'s estimates for increased TOA insolation), this  
 294 would result in a roughly 8.5 K temperature increase, close to the actual global mean  
 295 temperature increase in our simulations of roughly 8 K (from a global mean of 283.5 to  
 296 291 K; Fig. 3d). Assuming that column water vapor scales with surface air temperature  
 297 at a rate of roughly 7%/K per the Clausius-Clapeyron relationship (O'Gorman & Muller,  
 298 2010b; Held & Soden, 2006), we would expect water vapor to increase roughly 70% (the  
 299 Clausius-Clapeyron equation is exponential, so the relative increase in saturation vapor  
 300 pressure for the relatively large warming of 8 K is substantially larger than the infinites-  
 301 imal rate of change of 7%/K); the actual increase in total atmospheric water vapor in  
 302 our simulations was roughly 40% (from a global mean of 19 to 28 kg; Fig. 3d), with re-  
 303 ductions in tropical lower-tropospheric relative humidity over land that peak near -40%  
 304 (Fig. 5e). Thus, this simple argument provides similar order of magnitude changes in tem-  
 305 peratures and water vapor as our full simulations, but with a nearly factor-of-two over-  
 306 estimate in the vapor increase because the radiative forcing (a reduction in low clouds)  
 307 is driven by lower-tropospheric drying. We note that there is a lot of uncertainty in the  
 308 value of the climate response parameter, both across models and across forcing mech-  
 309 anisms (e.g. CO<sub>2</sub>, insolation, etc.) within a single model (Gregory, 2004; Hansen et al.,



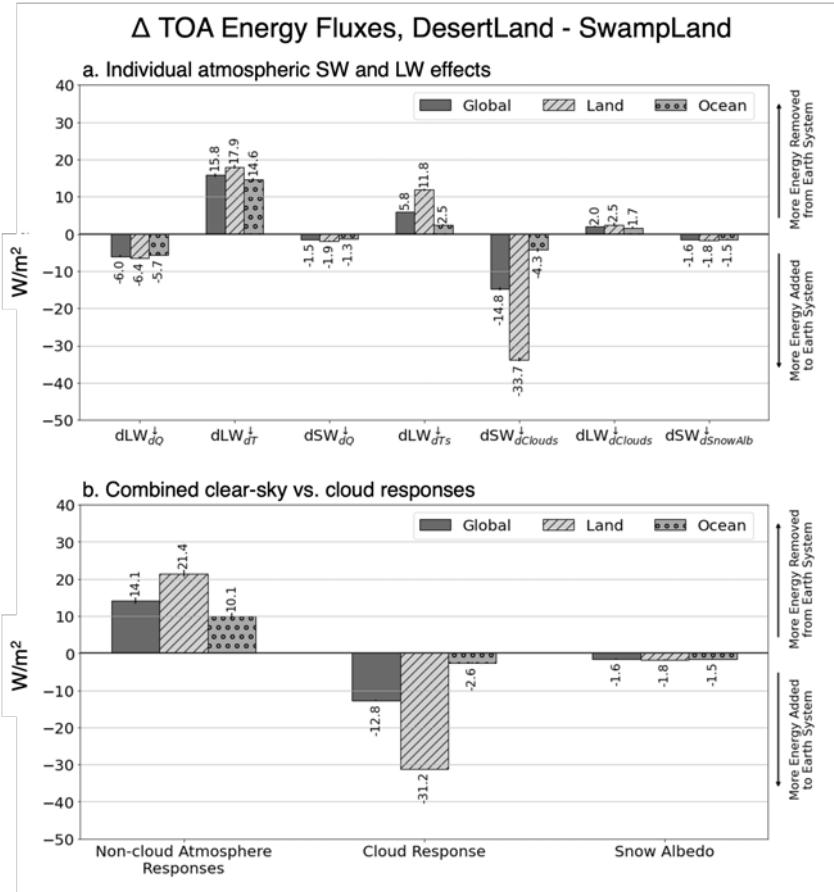
**Figure 8.** Change in global (solid), land (hatched), and ocean (dotted) annual mean surface fluxes for DesertLand - SwampLand. The breakdown of surface flux changes due to latent (LH-FLX) and sensible (SHFLX) heat, and shortwave (SW) and longwave (LW) radiation are shown. Panel (a) shows the changes in surface shortwave and longwave radiative fluxes decomposed using the radiative kernel into the contributions due to water vapor, atmospheric temperatures, surface temperatures, cloud cover, and surface albedo (i.e. snow changes), as well as the changes due to latent and sensible heat. Panel (b) combines the fluxes into those driven by latent heat flux, the cloud-free atmosphere and Plank response (air temperature, water vapor, and surface temperature), those driven by clouds, and other surface changes (sensible heat flux, albedo). The inter-annual standard deviation is marked by the vertical black lines capping each bar. Downwards (negative) values indicate that the change in the flux leads to more energy into the land surface (i.e. a warming effect), while upwards (positive) values indicate less energy into the land surface or more energy removed from the land surface.

1997), and that 7%/K is not an exact Clausius-Clapeyron scaling (O'Gorman & Muller, 2010b).

### 3.2.1 Decomposing the energy balance with a radiative kernel

The increase in near-surface MSE is driven by a combination of factors. In the absence of evaporative cooling (i.e. the DesertLand simulation), changes in the atmosphere which increase radiative fluxes into the land surface necessarily lead to warming. Over

land, suppressed terrestrial evaporation directly increases the energy that must be removed from the surface as sensible heat or longwave radiation (Fig. 8, spatial patterns shown in Fig. S2). Reducing evaporation results in excess energy available in the land surface ( $\approx 34 \text{ W/m}^2$  averaged over all land areas). The increase in energy into the land surface from increase downwelling longwave radiation as a result of suppressed land evaporation ( $29.6 \text{ W/m}^2$  from air temperatures and an additional  $7 \text{ W/m}^2$  from increased water vapor) is of a similar magnitude to the increase in energy into the land surface from suppressed latent heat flux ( $33.9 \text{ W/m}^2$ ).



**Figure 9.** Change in global (solid), land (hatched), and ocean (dotted) annual mean top of atmosphere (TOA) fluxes for DesertLand - SwampLand. The breakdown of toa flux changes due to shortwave (SW) and longwave (LW) radiation are shown. Panel (a) shows the changes in toa shortwave and longwave radiative fluxes decomposed using the radiative kernel into the contributions due to water vapor, atmospheric temperatures, surface temperatures, cloud cover, and surface albedo (i.e. snow changes). Panel (b) combines the fluxes into those driven by the cloud-free atmosphere and Plank response (air temperature, water vapor, and surface temperature), those driven by clouds, and surface albedo. The inter-annual standard deviation is marked by the vertical black lines capping each bar. Downwards (negative) values indicate changes that lead to more energy absorbed by the Earth system at the TOA while upwards (positive) values indicate energy removed from the Earth system at the TOA.

The longwave effects of increased water vapor act against the negative longwave cloud effect (Figs. 8, S2b,d). Increased water vapor itself drives a feedback on the sur-

face energy balance, with more water vapor leading to more longwave radiation into the surface ( $7.9 \text{ W/m}^2$ ); this term is of comparable magnitude to (though slightly smaller than) net cloud radiative effects ( $11.4 \text{ W/m}^2$ ) over land; Figs. 8, S2). Previous work has shown how changes in terrestrial evaporation modulate the water vapor greenhouse effect; specifically, Lagu  , Pietschnig, et al. (2021) show that while reducing land evaporation directly warms the surface, over very large idealized continents, reductions in land evaporation lead to reduced atmospheric water vapor and drive an overall cooling at the surface by reducing the water vapor greenhouse effect. In this study (with the modern Earth's continental configuration), our results show that land's control on water vapor is still an important contribution to the radiation budget, with the changes in surface and TOA fluxes driven by changes in water vapor of comparable magnitude to the combined shortwave and longwave effects of changes in cloud cover.

The combined effect of atmospheric responses to suppressed terrestrial evaporation is a slight increase in longwave energy into the surface, and a larger increase in shortwave radiation into the surface which, over land areas, is comparable in magnitude to the increase in surface energy coming from reduced latent cooling. Suppressing terrestrial evaporation leads to a net radiative flux of roughly  $6 \text{ W/m}^2$  out of the land surface (the sum of all the radiative fluxes in Fig. 8a), which is balanced by the combined changes in sensible and latent heat flux. While increased sensible heat flux (associated with higher surface temperatures in DesertLand vs. SwampLand) increases energy removed from the surface (by  $26.5 \text{ W/m}^2$  over land), suppressed latent heat flux over land means  $33.9 \text{ W/m}^2$  of energy is not removed from the land surface through evaporation. The response of the radiative terms of the surface energy budget over the oceans are of the same sign as the changes over land, though of different magnitude. In contrast, because ocean evaporation increases when terrestrial evaporation is suppressed, this removes energy from the surface, balanced by reductions in sensible heat flux. Over land areas, the radiative effects of cloud changes are larger than the combined radiative effects of water vapor and temperature changes (Fig. 8b).

The importance of the cloud feedback is reinforced by considering the change in top of atmosphere (TOA) energy fluxes. The shortwave effects of reductions in cloud cover in DesertLand-SwampLand are the single largest contributor to changes in the top of atmosphere energy balance over land (Fig. 9, spatial patterns in Fig. S3; the largest cloud reductions occur in low clouds over land areas (Fig. 7). The difference between the TOA and surface energy flux anomalies is the anomalous net energy input (NEI) to the atmosphere discussed below. Both the DesertLand and SwampLand simulations are in equilibrium, so the net TOA energy balance (the sum of the bars in Fig. 9) is near zero.

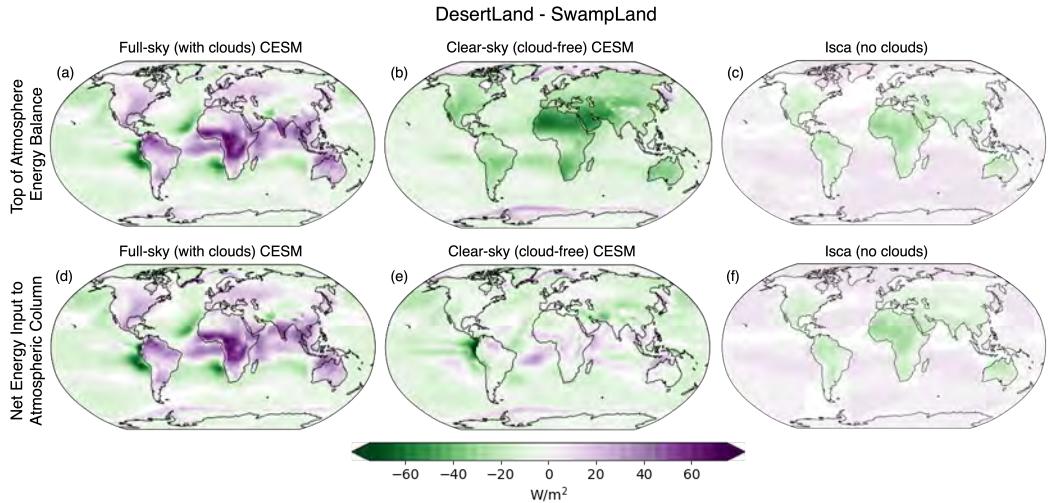
### 3.2.2 Enhanced energy sources drive tropical ascent

The reduction in low clouds in the Desertland vs. SwampLand simulation lead to an increase in the total amount of solar energy absorbed at the land surface—energy which is then emitted back to the atmosphere through the surface energy budget. Specifically, suppressed terrestrial evaporation leads to an increase over most land regions in the net energy input to the atmosphere (NEI), which is the sum of radiative and surface turbulent fluxes into the atmosphere through its top and bottom boundaries (Fig. 10d). It also leads to an increase in near-surface moist static energy (MSE) over most continental regions, especially in the tropics (Fig. 11). We calculate moist static energy as

$$MSE = c_p T + L_v Q + gZ, \quad (2)$$

the sum of the dry energy (the heat capacity  $c_p$  of dry air multiplied by the air temperature  $T$ ), potential energy (the gravitational constant of acceleration  $g$  multiplied by the geopotential height  $Z$ ), and the moist energy (the latent heat of vaporization  $L_v$  multiplied by water vapor  $Q$ ) of a static (non-dynamic) parcel of air. Note that while we show the NEI for clear-sky CESM, the latent and sensible heat fluxes which go into the NEI

376 calculation are from the full CESM model that includes cloud effects on the surface en-  
 377 ergy budget; as such, Fig. 10e should be treated with appropriate caution.



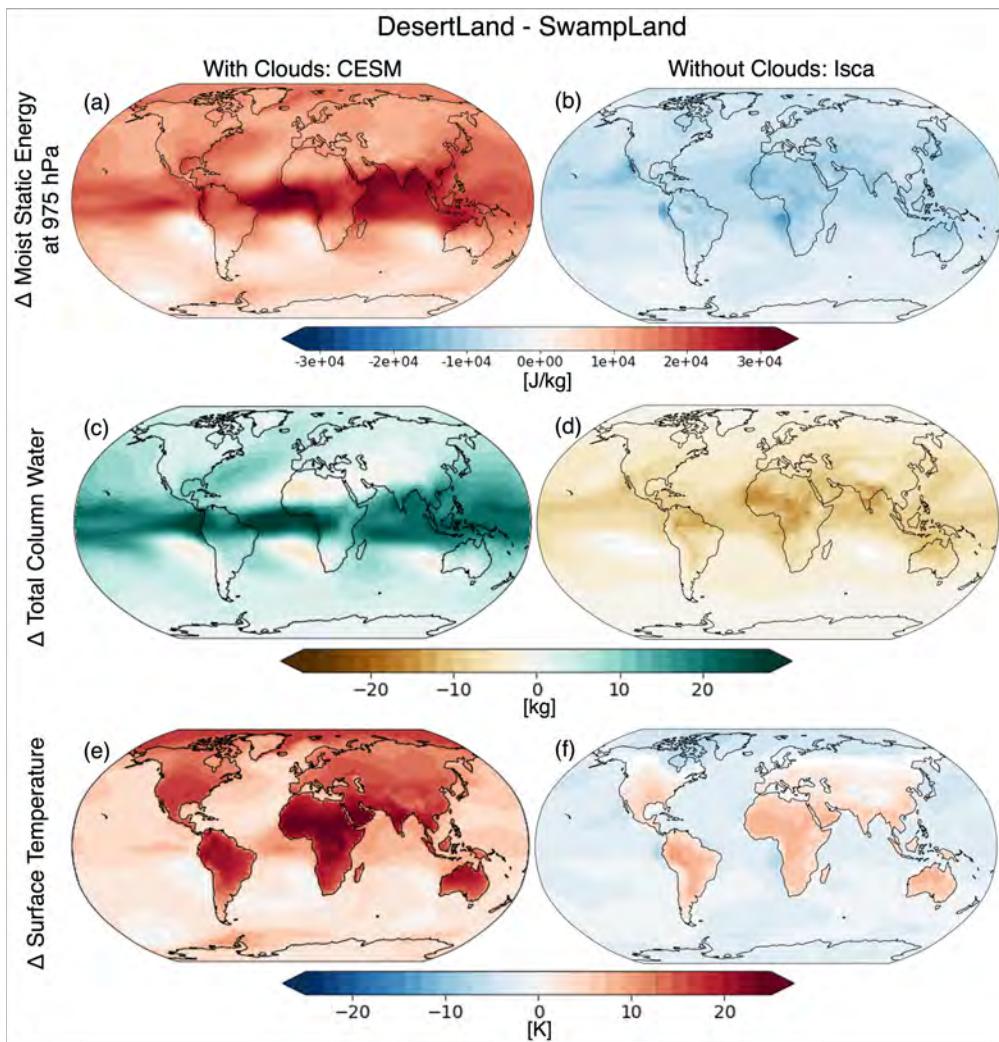
**Figure 10.** Annual mean change (DesertLand - SwampLand) in the top of atmosphere energy budget (net shortwave - outgoing longwave radiation; top row) and net energy input (NEI: net top of atmosphere - net surface fluxes; bottom row). The full CESM simulation, including cloud radiative effects, is shown in the left column; CESM where radiative fluxes ignore the influence of clouds (“clear-sky” conditions) are shown in the centre column, while Isca, which is always radiatively cloud-free, is shown in the right column. Only changes that pass a statistical test are shown, where values are significant if the p-values calculated from a student’s t-test pass a false discovery rate of 0.15.

378 Two complementary theoretical frameworks can then help in understanding the re-  
 379 sponse of the large-scale circulation to our surface evaporative forcing. First, by verti-  
 380 cally integrating the moist static energy budget, one can relate large-scale vertical winds  
 381 to the local input of energy through the top and bottom boundaries of the atmosphere,

$$-\omega_1 M = NEI \quad (3)$$

382 where  $\omega_1$  is the vertical motion at a characteristic level, and  $M$  is a coefficient known  
 383 as the gross moist stability (Neelin & Held, 1987; Sobel et al., 2007; Raymond et al., 2015).  
 384 Being a measure of the vertical energy stratification of the atmosphere,  $M$  is typically  
 385 positive in the time mean in deep-convection regions, expressing the fact that time-mean  
 386 ascent typically exports energy from the column in many tropical regions. Thus, the in-  
 387 crease in  $NEI$  due to reduced low cloud cover is accompanied by enhanced large-scale  
 388 ascent in many tropical land regions (Figs. 12a, S4). While this anomalous ascent is bottom-  
 389 heavy, it extends the full depth of the troposphere. There is an increase in precipitation  
 390 over many of the regions in which there is an increase in  $NEI$  (Fig. 12b). Outside of the  
 391 tropics, precipitation over land decreases (Fig. 12b). Though there is increased time-mean  
 392 upwards motion and increased  $NEI$  (Fig. S5), this does not lead to more precipitation,  
 393 consistent with the fact that precipitation there is generated primarily by the moisture  
 394 converged by transient motions, rather than time-mean flow.

395 In an alternate framework, one can consider the energy content of air (MSE) in-  
 396 stead of the source of atmospheric energy (the NEI); precipitating ascent in the low-latitude  
 397 atmosphere generally lies near the maximum in surface air MSE and increases in inten-  
 398 sity with horizontal gradients in that MSE. This is expected when surface air MSE is



**Figure 11.** Annual mean change (DesertLand - SwampLand) in near-surface (975 hPa) moist static energy (top), total column water vapor (middle), and surface temperature (bottom) for CESM (left) and Isca (right) simulations. Only changes that pass a statistical test are shown, where values are significant if the p-values calculated from a student's t-test pass a false discovery rate of 0.15.

convectively coupled to free-tropospheric temperatures (Emanuel et al., 1994; Neelin, 2007; Privé & Plumb, 2007; Shekhar & Boos, 2016), and has been shown to describe the observed tropical climatology and interannual variability (Nie et al., 2010; Hurley & Boos, 2013). Suppressed surface evaporation increases MSE (because of the atmospheric feedbacks discussed above), particularly over tropical regions which thus enhances the magnitude of the meridional gradient in MSE, and in turn the large-scale tropical overturning circulation (Fig. 11a, 12). Decreases in moisture in the lower atmosphere over most land areas (driven by reduced surface evaporation) would lead to a reduction in MSE, but this is more than compensated for by increases in temperatures, which lead to an overall increase in MSE at 975 hPa over land areas (Fig. S6); changes in geopotential heights are negligible.

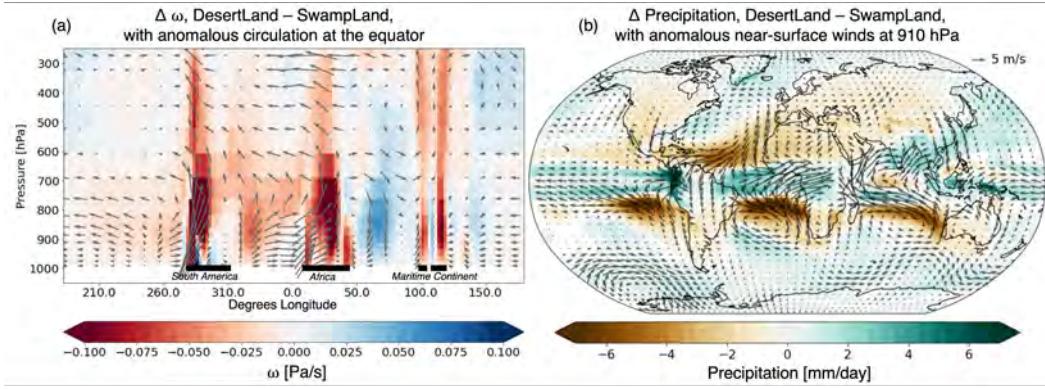
Neither of these frameworks provides a closure for the amount of water vapor in the atmosphere, but in Earth's low-latitude atmosphere, where most water vapor lies and where horizontal temperature gradients are weak, MSE generally scales like precipitable water (Charney, 1963; Sobel et al., 2001) given the general availability of the ocean water supply in the modern continental configuration. In other words, large-scale ascent will advect water vapor upward, humidifying the column over land regions despite suppressed terrestrial evaporation. The MSE increase aloft over land in DesertLand indeed consists of increases in both temperatures and specific humidities (specific humidity decreases near the surface, but increases aloft), though the latter increase slowly enough that RH over land decreases (Fig. 5e); the increases in relative humidity occurring near the poles at high altitudes occur in the stratosphere and are not considered in this study (indeed, changes in T and Q above the tropopause are masked out of all calculations involving the radiative kernel, as in Pendergrass et al. (2018); Shell et al. (2008); Laguë, Swann, & Boos (2021)). While the warming in our simulations is driven by changes in land evaporation, for CO<sub>2</sub>-driven warming, Byrne & O'Gorman (2016) also find near-surface continental relative humidity decreases. In simulations of future climate, we would expect both the radiative effects of CO<sub>2</sub> as well as changes in terrestrial evaporation to drive changes in atmospheric moisture. The only terrestrial regions where total column water vapor decreases is in the dry subtropics (e.g. the Sahara) where subsidence increases (Figs. S4, S5).

The discussion above differs greatly from arguments in which continental water vapor is treated as being set by transport from ocean regions. Anomalous near-surface winds do bring moist ocean air onto tropical land in DesertLand vs. SwampLand, e.g. in tropical South America, Africa, and Asia (Fig. 12b). However, such onshore winds need not produce ascent that spans most of the depth of the troposphere. Strong land-sea temperature gradients, such as those between a desert and ocean on Earth, often produce shallow, non-precipitating circulations; we will show below that, without cloud-radiative effects, suppressed land evaporation indeed leads to enhanced onshore flow that is shallow, non-precipitating, and that does not result in enhanced total column water vapor.

### 3.2.3 The response without cloud-radiative effects

When similar “desert” and “swamp” simulations are repeated with a cloud-free idealized global circulation model (Isca), surface temperatures still increase in response to suppressed terrestrial evaporation (Fig. 11f). However, both near-surface MSE (Fig. 11b) and atmospheric water vapor over continental regions (Fig. 11d) decrease. Unlike the CESM simulations, where the reduction in low cloud cover leads to an increase in energy into the Earth system at the top of the atmosphere, the Isca simulations do not have this additional energy source to the system as there are no clouds.

Suppressing terrestrial evaporation leads to a decrease in the atmospheric NEI and the net top of atmosphere energy budget over continents in the Isca simulations (Fig. 10). While the Isca simulations still show some anomalous upward motion over the continents in the lower atmosphere, and near-surface onshore flow from the oceans to the land in many regions (Figs. S7 & S8), these are shallow, non-precipitating anomalous circulations that contrast strongly with the deep, precipitating anomalous flow in CESM (c.f. S4 & S7). Similarly, the warming induced by suppressed terrestrial evaporation in Isca is restricted to the near-surface atmosphere over land regions, while CESM warms throughout the column (c.f. Fig. 5 & S9). While there are many differences between CESM and Isca, a key difference in the response to altered terrestrial evaporation is the response of cloud cover, which CESM includes and Isca does not. This highlights the importance of understanding cloud feedbacks—already a large source of climate uncertainty (Zelinka et al., 2017)—for determining how terrestrial evaporation changes alter the climate system.



**Figure 12.** (a) Vertical cross section (longitude vs. pressure) at the equator for the change in vertical motion  $\omega$  [Pa/s] for DesertLand – SwampLand is shown in colors, with red indicating negative  $\omega$ , i.e. positive vertical motion. Vectors show the direction of anomalous motion:  $u$  winds in [m/s] in the x-direction, and  $\omega \times -150$  in [Pa/s] in the y direction.  $\omega$  is multiplied by  $-1$  so arrows point in the direction of motion, and by  $150$  as an arbitrary scaling such that the vertical component is of comparable magnitude to the horizontal component, for ease of visualization. Horizontal black lines at the surface indicate land masses. (b) Change in precipitation (shading) for DesertLand – SwampLand, with vectors showing anomalous  $\langle u, v \rangle$  winds [m/s] at 910 hPa. Only changes that pass a statistical test are shown, where values are significant if the p-values calculated from a student's t-test pass a false discovery rate of 0.15.

While we do not have a CESM simulation without radiatively interactive clouds, we can crudely compare the radiative fluxes in the Isca simulations to the top of atmosphere radiative fluxes calculated using "clear-sky" conditions in CESM (Fig. 11b/e). The clear-sky fluxes are calculated at each time step ignoring the radiative effects of clouds; however the model is integrated forwards using the full sky (including the radiative effects of clouds) fluxes—that is, the temperature, moisture, dynamics, etc. of CESM are consistent with the full-sky radiative fluxes, not the clear-sky radiative fluxes. Both the clear-sky CESM simulations and the Isca simulations show a decrease in the TOA energy balance over land regions when terrestrial evaporation is suppressed, reflecting the fact that the increase in energy into the Earth system in the CESM simulations are due to changes in low cloud cover over land.

While both CESM and Isca allow us to test the response of the climate system to suppressed terrestrial evaporation, we note that there are limitations in comparing results between these two models, as they do not provide an apples-to-apples comparison. In future work, a perhaps more robust comparison would be to compare the full CESM-SLIM model to a version of CESM where the clouds have been modified to be transparent to radiation; however, this model would produce unrealistically high TOA albedos, thus an increase in the surface albedo (such as is done in Isca) would be necessary to replicate a similar climate to the modern Earth. Such a modified model configuration would need to be evaluated and benchmarked before using it for this kind of experiment. Similarly, a modified version of Isca that included radiatively interactive clouds could also be leveraged to understand interactions between terrestrial evaporation, cloud cover, and the global energy budget. In this study however, we present the results of both the CESM and Isca simulations as complimentary; both indicate a strong control of terrestrial evaporation on surface climate, and the model that allows for interactive cloud changes (CESM) suggests further investigation of the coupling of terrestrial evaporation and cloud cover is of merit. Additionally, we note that all the simulations shown in this study are highly

488 idealized; they are useful for improving our mechanistic understanding of interactions  
 489 between terrestrial processes and global climate, and can be used to inform more real-  
 490 istic studies of the effects of changes in terrestrial evapotranspiration driven by vegeta-  
 491 tion change, land use, agriculture, climate change, etc. on global climate.

## 492 4 Conclusions

493 In a global model with realistic continental geometry, reducing terrestrial evapo-  
 494 ration increases the total amount of atmospheric water vapor over most land and ocean  
 495 regions. The residence time of water vapor in the atmosphere increases by roughly 50%  
 496 from the simulation with fully saturated land to the simulation with desert land. Sup-  
 497 pressing land evaporation has a direct warming effect on the land surface by reducing  
 498 latent cooling of the surface, but also drives atmospheric feedbacks including reductions  
 499 in terrestrial cloud cover. The anomalous surface energy fluxes driven by atmospheric  
 500 cloud, water vapor, and temperature feedbacks are larger than the initial change in la-  
 501 tent heat flux driven directly by suppressed terrestrial evaporation. The cloud feedback  
 502 is critical for increasing near-surface moist static energy and generating anomalous at-  
 503 mospheric circulations throughout the depth of the troposphere. Simulations conducted  
 504 in a cloud-free model still show surface warming with suppressed terrestrial evaporation,  
 505 but also show a decrease, rather than an increase, in near surface MSE. Anomalous at-  
 506 mospheric circulations over the continents in cloud-free simulations are much shallower,  
 507 and the atmosphere shows reduced atmospheric water vapor with suppressed terrestrial  
 508 evaporation. This extreme experiment raises the question of how real-world changes to  
 509 the land surface (e.g. land use, agriculture) may be contributing to climate change by  
 510 altering atmospheric water vapor and cloud cover, and how terrestrial evaporation mod-  
 511 ulates climate on other planets or in past continental configurations of Earth's history.

## 512 Data Availability Statement

513 Source code for the models used are publicly available on github at [https://github](https://github.com/marysa/SimpleLand/releases/tag/slim0.1.006_release-cesm2.1.4)  
 514 .com/marysa/SimpleLand/releases/tag/slim0.1.006\_release-cesm2.1.4 for SLIM,  
 515 <https://github.com/ESCOMP/CESM/releases/tag/release-cesm2.1.1> for CESM, and  
 516 <https://github.com/ExeClim/Isca/releases/tag/v1.0> for Isca. The specific versions  
 517 used in this study are archived on zenodo in the following locations: SLIM, at <https://zenodo.org/badge/latestdoi/509227365> (DOI: 10.5281/zenodo.7352717); CESM, at  
 518 <https://doi.org/10.5281/zenodo.3895315> (DOI: 10.5281/zenodo.3895315); and Isca,  
 519 at <https://github.com/ExeClim/Isca/releases/tag/v1.0>, respectively. The mod-  
 520 ified version of Isca used in this study is available via zenodo at <https://doi.org/10.5281/zenodo.6800218>. Output from the simulations, analysis scripts, and model setup  
 521 files used in this study are available on Zenodo at <https://doi.org/10.5281/zenodo.7909213>.

## 525 Author Contributions Statement

526 MML developed the idea for the study, modified the models, performed the sim-  
 527 ulations, and conducted the analysis. All authors discussed the experimental design and  
 528 results. MML wrote the initial draft and all authors contributed to the manuscript.

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538 **References**

- Bailey, D., Hunke, E., DuVivier, A., Lipscomb, B., Bitz, C., Holland, M., ...  
 Schramm, J. (2018). *CESM CICE5 Users Guide* (Tech. Rep.).
- Ball, J. T., Woodrow, I. E., & Berry, J. A. (1987). A Model Predicting Stomatal Conductance and Its Contribution to the Control of Photosynthesis Under Different Environmental Conditions. *Progress in Photosynthesis Research* (January). doi: 10.1007/978-94-017-0519-6
- Berg, A., Findell, K., Lintner, B., Giannini, A., Seneviratne, S. I., Van Den Hurk, B., ... Milly, P. C. D. (2016). Land-atmosphere feedbacks amplify aridity increase over land under global warming. *Nature Climate Change*, 6(May), 869–874. doi: 10.1038/nclimate3029
- Bonan, G. B. (2008). *Ecological Climatology*. Cambridge Univ. Press, Cambridge, UK.
- Boos, W. R., & Korty, R. L. (2016). Regional energy budget control of the intertropical convergence zone and application to mid-Holocene rainfall. *Nature Geoscience*, 9(12), 892–897. doi: 10.1038/ngeo2833
- Budyko, M. I. (1961). The Heat Balance of the Earth's Surface. *Soviet Geography*, 2(4), 3–13. doi: 10.1080/00385417.1961.10770761
- Byrne, M. P., & O'Gorman, P. A. (2015, October). The Response of Precipitation Minus Evapotranspiration to Climate Warming: Why the "Wet-Get-Wetter, Dry-Get-Drier" Scaling Does Not Hold over Land\*. *Journal of Climate*, 28(20), 8078–8092. doi: 10.1175/JCLI-D-15-0369.1
- Byrne, M. P., & O'Gorman, P. A. (2016). Understanding decreases in land relative humidity with global warming: Conceptual model and GCM simulations. *Journal of Climate*, 29(24), 9045–9061. doi: 10.1175/JCLI-D-16-0351.1
- Charney, J. G. (1963, November). A Note on Large-Scale Motions in the Tropics. *Journal of the Atmospheric Sciences*, 20(6), 607–609. doi: 10.1175/1520-0469(1963)020<0607:ANOLSM>2.0.CO;2
- Clough, S. A., Shephard, M. W., Mlawer, E. J., Delamere, J. S., Iacono, M. J., Cady-Pereira, K., ... Brown, P. D. (2005). Atmospheric radiative transfer modeling: A summary of the AER codes. *Journal of Quantitative Spectroscopy and Radiative Transfer*, 91(2), 233–244. doi: 10.1016/j.jqsrt.2004.05.058
- Collins, M., Knutti, R., Arblaster, J., Dufresne, J. L., Fichefet, T., Friedlingstein, P., ... Wehner, M. (2013). Long-term climate change: Projections, commitments and irreversibility. *Climate Change 2013 the Physical Science Basis: Working Group I Contribution to the Fifth Assessment Report of the Intergovernmental Panel on Climate Change*, 9781107057, 1029–1136. doi: 10.1017/CBO9781107415324.024
- Computational and Information Systems Laboratory. (2019). *Cheyenne: HPE/SGI ICE XA System (University Community Computing)*. Boulder, CO: National Center for Atmospheric Research. doi: 10.5065/D6RX99HX
- Davin, E. L., de Noblet-Ducoudré, N., de Noblet-Ducoudre, N., & de Noblet-Ducoudré, N. (2010, January). Climatic Impact of Global-Scale Deforestation: Radiative versus Nonradiative Processes. *Journal of Climate*, 23(1), 97–112. doi: 10.1175/2009JCLI3102.1
- Dirmeyer, P. A. (1994). Vegetation Stress as a Feedback Mechanism in Midlatitude Drought. *Journal of Climate*, 7.
- Dirmeyer, P. A. (2006). The hydrologic feedback pathway for land-climate coupling.

- 585                   *Journal of Hydrometeorology*, 7(5), 857–867. doi: 10.1175/JHM526.1
- 586 Dirmeyer, P. A. (2011). The terrestrial segment of soil moisture-climate coupling.  
 587                   *Geophysical Research Letters*, 38(16), 1–5. doi: 10.1029/2011GL048268
- 588 Donohue, R. J., Roderick, M. L., McVicar, T. R., & Farquhar, G. D. (2013). Im-  
 589 pact of CO<sub>2</sub> fertilization on maximum foliage cover across the globe's warm,  
 590 arid environments.     *Geophysical Research Letters*, 40(12), 3031–3035.     doi:  
 591 10.1002/grl.50563
- 592 Eltahir, E. A., & Bras, R. L. (1996). Precipitation Recycling.   *Reviews of Geo-*  
 593 *physics*(96), 367–378.
- 594 Emanuel, K. A., David Neelin, J., & Bretherton, C. S. (1994, July). On large-scale  
 595 circulations in convecting atmospheres.   *Quarterly Journal of the Royal Meteorolo-*  
 596 *gical Society*, 120(519), 1111–1143. doi: 10.1002/qj.49712051902
- 597 Field, C. B., Jackson, R. B., & Mooney, H. A. (1995). Stomatal responses to in-  
 598 creased CO<sub>2</sub>: Implications from the plant to the global scale.   *Plant, Cell & Envi-*  
 599 *ronment*, 18(10), 1214–1225. doi: 10.1111/j.1365-3040.1995.tb00630.x
- 600 Fraedrich, K., Kleidon, A., & Lunkeit, F. (1999). A Green Planet versus a Desert  
 601 World: Estimating the Effect of Vegetation Extremes on the Atmosphere.   *Journal*  
 602 *of Climate*, 12(10), 3156–3163. doi: 10.1175/1520-0442(1999)012<3156:AGPVAD>2  
 603 .0.CO;2
- 604 Garcia, E. S., Swann, A. L. S., Villegas, J. C., Breshears, D. D., Law, D. J., Saleska,  
 605 S. R., & Stark, S. C. (2016). Synergistic ecoclimate teleconnections from forest  
 606 loss in different regions structure global ecological responses.   *PLoS ONE*, 11(11),  
 607 1–12. doi: 10.1371/journal.pone.0165042
- 608 Geen, R., Lambert, F. H., & Vallis, G. K. (2018). Regime change behavior during  
 609 Asian monsoon onset.   *Journal of Climate*, 31(8), 3327–3348. doi: 10.1175/JCLI-D  
 610 -17-0118.1
- 611 Gregory, J. M. (2004). A new method for diagnosing radiative forcing and cli-  
 612 mate sensitivity.   *Geophysical Research Letters*, 31(3), L03205.     doi: 10.1029/  
 613 2003GL018747
- 614 Hansen, J., Sato, M., & Ruedy, R. (1997, March). Radiative forcing and climate re-  
 615 sponse.   *Journal of Geophysical Research: Atmospheres*, 102(D6), 6831–6864. doi:  
 616 10.1029/96JD03436
- 617 Harris, C. R., Millman, K. J., van der Walt, S. J., Gommers, R., Virtanen, P., Cour-  
 618 napeau, D., ... Oliphant, T. E. (2020, September). Array programming with  
 619 NumPy.   *Nature*, 585(7825), 357–362. doi: 10.1038/s41586-020-2649-2
- 620 Held, I. M., & Soden, B. J. (2006, November). Robust Responses of the Hydrological  
 621 Cycle to Global Warming.   *Journal of Climate*, 19(21), 5686–5699. doi: 10.1175/  
 622 JCLI3990.1
- 623 Hobeichi, S., Abramowitz, G., & Evans, J. P. (2021). Robust historical evapotranspi-  
 624 ration trends across climate regimes.   *Hydrology and Earth System Sciences*, 25(7),  
 625 3855–3874. doi: 10.5194/hess-25-3855-2021
- 626 Hoyer, S., & Hamman, J. (2017). Xarray: N-D labeled arrays and datasets in  
 627 Python.   *Journal of Open Research Software*, 5(1). doi: 10.5334/jors.148
- 628 Hurley, J. V., & Boos, W. R. (2013, December). Interannual Variability of Mon-  
 629soon Precipitation and Local Subcloud Equivalent Potential Temperature.   *Journal*  
 630 *of Climate*, 26(23), 9507–9527. doi: 10.1175/JCLI-D-12-00229.1
- 631 Hurrell, J. W., Holland, M. M., Gent, P. R., Ghan, S., Kay, J. E., Kushner, P. J.,  
 632 ... Marshall, S. (2013). The Community Earth System Model: A Framework for  
 633 Collaborative Research.   *Bulletin of the American Meteorological Society*, 94(9),  
 634 1339–1360. doi: 10.1175/BAMS-D-12-00121.1
- 635 Jones, H. G. (1998). Stomatal control of photosynthesis and transpiration.   *Jour-*  
 636 *nal of Experimental Botany*, 39(SPECIAL ISSUE: STOMATAL BIOLOGY),  
 637 387–398.
- 638 Koster, R. D., Guo, Z., Dirmeyer, P. A., Bonan, G., Chan, E., Cox, P., ... Yamada,

- 639 T. (2006). GLACE: The Global Land-Atmosphere Coupling Experiment. Part I:  
 640 Overview. *Journal of Hydrometeorology*, 7(4), 590–610. doi: 10.1175/JHM510.1
- 641 Kottek, M., Grieser, J., Beck, C., Rudolf, B., & Rubel, F. (2006). World map of the  
 642 Köppen-Geiger climate classification updated. *Meteorologische Zeitschrift*, 15(3),  
 643 259–263. doi: 10.1127/0941-2948/2006/0130
- 644 Laguë, M. M., Bonan, G. B., & Swann, A. L. S. (2019). Separating the Impact of  
 645 Individual Land Surface Properties on the Terrestrial Surface Energy Budget in  
 646 both the Coupled and Uncoupled Land–Atmosphere System. *Journal of Climate*,  
 647 32(18), 5725–5744. doi: 10.1175/jcli-d-18-0812.1
- 648 Laguë, M. M., Pietschnig, M., Ragen, S., Smith, T. A., & Battisti, D. S. (2021).  
 649 Terrestrial evaporation and global climate: Lessons from Northland, a planet  
 650 with a hemispheric continent. *Journal of Climate*, 34(6), 2253–2276. doi:  
 651 10.1175/jcli-d-20-0452.1
- 652 Laguë, M. M., & Swann, A. L. (2016). Progressive midlatitude afforestation: Im-  
 653 pacts on clouds, global energy transport, and precipitation. *Journal of Climate*,  
 654 29(15), 5561–5573. doi: 10.1175/jcli-d-15-0748.1
- 655 Laguë, M. M., Swann, A. L. S., & Boos, W. R. (2021). Radiative feedbacks on  
 656 land surface change and associated tropical precipitation shifts. *Journal of Cli-  
 657 mate*(2012), 1–63. doi: 10.1175/jcli-d-20-0883.1
- 658 Lai, C.-T., & Katul, G. (2000, January). The dynamic role of root-water uptake  
 659 in coupling potential to actual transpiration. *Advances in Water Resources*, 23(4),  
 660 427–439. doi: 10.1016/S0309-1708(99)00023-8
- 661 Lawrence, D. M., Fisher, R. A., Koven, C. D., Oleson, K. W., Swenson, S. C., Bo-  
 662 nan, G., ... Zeng, X. (2019). The Community Land Model Version 5: De-  
 663 scription of New Features, Benchmarking, and Impact of Forcing Uncertainty.  
 664 *Journal of Advances in Modeling Earth Systems*, 11(12), 4245–4287. doi:  
 665 10.1029/2018MS001583
- 666 Lemordant, L., Gentine, P., Swann, A. S., Cook, B. I., & Scheff, J. (2018). Critical  
 667 impact of vegetation physiology on the continental hydrologic cycle in response to  
 668 increasing CO<sub>2</sub>. *Proceedings of the National Academy of Sciences of the United  
 669 States of America*, 115(16), 4093–4098. doi: 10.1073/pnas.1720712115
- 670 Medlyn, B. E., Duursma, R. A., Eamus, DEREK., Ellsworth, D. S., Prentice, I. C.,  
 671 Barton, C. V. M., ... Wingate, L. (2011). Reconciling the optimal and empirical  
 672 approaches to modelling stomatal conductance. *Global Change Biology*, 17(6),  
 673 2134–2144. doi: 10.1111/j.1365-2486.2010.02375.x
- 674 Monteith, J. (1965). Evaporation and environment. *Symposia of the Society for Ex-  
 675 perimental Biology*(19), 205–234.
- 676 Neale, R. B., Gettelman, A., Park, S., Chen, C.-C., Lauritzen, P. H., Williamson,  
 677 D. L., ... Others (2012). Description of the NCAR community atmosphere model  
 678 (CAM 5.0). *NCAR Tech. Note NCAR/TN-486+STR*.
- 679 Neelin, J. D. (2007). Moist dynamics of tropical convection zones in monsoons, tele-  
 680 connections, and global warming. In *The global circulation of the atmosphere* (pp.  
 681 267–301). doi: 10.2307/j.ctv1t1kg52
- 682 Neelin, J. D., & Held, I. M. (1987, January). Modeling tropical convergence based  
 683 on the moist static energy budget. *Monthly Weather Review*, 115(1), 3–12. doi: 10  
 684 .1175/1520-0493(1987)115<0003:MTCBOT>2.0.CO;2
- 685 Nie, J., Boos, W. R., & Kuang, Z. (2010). Observational evaluation of a convective  
 686 quasi-equilibrium view of monsoons. *Journal of Climate*, 23(16), 4416–4428. doi:  
 687 10.1175/2010JCLI3505.1
- 688 Norby, R. J., Warren, J. M., Iversen, C. M., Medlyn, B. E., & McMurtrie, R. E.  
 689 (2010). CO<sub>2</sub> enhancement of forest productivity constrained by limited nitrogen  
 690 availability. *Proceedings of the National Academy of Sciences of the United States  
 691 of America*, 107(45), 19368–19373. doi: 10.1073/pnas.1006463107
- 692 O’Gorman, P. A., & Muller, C. J. (2010a). How closely do changes in surface and

- 693 column water vapor follow Clausius-Clapeyron scaling in climate change simu-  
 694 lations? *Environmental Research Letters*, 5(2). doi: 10.1088/1748-9326/5/2/  
 695 025207
- 696 O'Gorman, P. A., & Muller, C. J. (2010b, April). How closely do changes in  
 697 surface and column water vapor follow Clausius-Clapeyron scaling in climate  
 698 change simulations? *Environmental Research Letters*, 5(2), 025207. doi:  
 699 10.1088/1748-9326/5/2/025207
- 700 Pendergrass, A. G., Conley, A., & Vitt, F. M. (2018). Surface and top-of-  
 701 Atmosphere radiative feedback kernels for cesm-cam5. *Earth System Science  
 702 Data*, 10(1), 317–324. doi: 10.5194/essd-10-317-2018
- 703 Pielke, R. A., Sr., Avissar, RonI., Raupach, M., Dolman, A. J., Zeng, X., & Den-  
 704 ning, A. S. (1998, June). Interactions between the atmosphere and terrestrial  
 705 ecosystems: Influence on weather and climate. *Global Change Biology*, 4(5),  
 706 461–475. doi: 10.1046/j.1365-2486.1998.t01-1-00176.x
- 707 Privé, N. C., & Plumb, R. A. (2007). Monsoon dynamics with interactive forcing.  
 708 Part II: Impact of eddies and asymmetric geometries. *Journal of the Atmospheric  
 709 Sciences*, 64(5), 1431–1442. doi: 10.1175/JAS3917.1
- 710 Raymond, D., Fuchs, Ž., Gjorgjevska, S., & Sessions, S. (2015, September). Bal-  
 711 anced dynamics and convection in the tropical troposphere. *Journal of Advances  
 712 in Modeling Earth Systems*, 7(3), 1093–1116. doi: 10.1002/2015MS000467
- 713 Sellers, P. J., Bounoua, L., Collatz, G. J., a. Randall, D., a. Dazlich, D., Los, S. O.,  
 714 ... Others (1996). Comparison of radiative and physiological effects of doubled  
 715 atmospheric CO<sub>2</sub> on climate. *SCIENCE-NEW YORK THEN WASHINGTON-*,  
 716 271(5254), 1402–1405. doi: 10.1126/science.271.5254.1402
- 717 Seneviratne, S. I., Corti, T., Davin, E. L., Hirschi, M., Jaeger, E. B., Lehner, I., ...  
 718 Teuling, A. J. (2010, May). Investigating soil moisture–climate interactions in  
 719 a changing climate: A review. *Earth-Science Reviews*, 99(3-4), 125–161. doi:  
 720 10.1016/j.earscirev.2010.02.004
- 721 Shekhar, R., & Boos, W. R. (2016, July). Improving Energy-Based Estimates  
 722 of Monsoon Location in the Presence of Proximal Deserts. *Journal of Climate*,  
 723 29(13), 4741–4761. doi: 10.1175/JCLI-D-15-0747.1
- 724 Shell, K. M., Kiehl, J. T., & Shields, C. A. (2008). Using the radiative kernel tech-  
 725 nique to calculate climate feedbacks in NCAR's Community Atmospheric Model.  
 726 *Journal of Climate*, 21(10), 2269–2282. doi: 10.1175/2007JCLI2044.1
- 727 Sherwood, S. C., Roca, R., Weckwerth, T. M., & Andronova, N. G. (2010). Tropo-  
 728 spheric water vapor, convection, and climate. *Reviews of Geophysics*, 48(2), 1–29.  
 729 doi: 10.1029/2009RG000301
- 730 Shukla, J., & Mintz, Y. (1982). Influence of Land-Surface Evapotranspiration on the  
 731 Earth's Climate. *Science*, 215(4539), 1498–1501. doi: 10.1126/science.215.4539  
 732 .1498
- 733 Sobel, A. H., Bellon, G., & Bacmeister, J. (2007, November). Multiple equilibria in  
 734 a single-column model of the tropical atmosphere. *Geophysical Research Letters*,  
 735 34(22), L22804. doi: 10.1029/2007GL031320
- 736 Sobel, A. H., Nilsson, J., & Polvani, L. M. (2001, December). The Weak Tem-  
 737 perature Gradient Approximation and Balanced Tropical Moisture Waves\*.  
 738 *Journal of the Atmospheric Sciences*, 58(23), 3650–3665. doi: 10.1175/  
 739 1520-0469(2001)058<3650:TWTGAA>2.0.CO;2
- 740 Sun, X., & Wang, G. (2022). Causes for the negative scaling of extreme precipitation  
 741 at high temperatures. *Journal of Climate*, 35(18), 6119–6134.
- 742 Swann, A. L. S., Fung, I. Y., & Chiang, J. C. H. (2012, January). Mid-latitude af-  
 743 forestation shifts general circulation and tropical precipitation. *Proceedings of the  
 744 National Academy of Sciences*, 109(3), 712–716. doi: 10.1073/pnas.1116706108
- 745 Swann, A. L. S., Fung, I. Y., Levis, S., Bonan, G. B., & Doney, S. C. (2010).  
 746 Changes in Arctic vegetation amplify high-latitude warming through the green-

- 747 house effect. *Proceedings of the National Academy of Sciences*, 107(4), 1295–1300.  
 748 doi: 10.1073/pnas.0913846107
- 749 Swann, A. L. S., Hoffman, F. M., Koven, C. D., & Randerson, J. T. (2016). Plant  
 750 responses to increasing CO<sub>2</sub> reduce estimates of climate impacts on drought severity.  
 751 *PNAS*, 113(36), 10019–10024. doi: 10.1073/pnas.1604581113
- 752 Swann, A. L. S., Laguë, M. M., Garcia, E. S., Breshears, D. D., Field, J., Moore, D.,  
 753 ... Minor, D. M. (2018). Local tree die-off impacts vegetation productivity across  
 754 the US regardless of location via continental-scale ecoclimate teleconnections.  
 755 *Environmental Research Letters*, Submitted.
- 756 Thomson, S. I., & Vallis, G. K. (2019). Hierarchical modeling of solar system planets  
 757 with Isca. *Atmosphere*, 10(12), 1–21. doi: 10.3390/ATMOS10120803
- 758 Trenberth, K. E. (1998). Atmospheric Moisture Residence Times and Cycling:  
 759 Implications for Rainfall Rates and Climate Change. , *Climatic Change*(39), 667–  
 760 694.
- 761 Vallis, G. K., Colyer, G., Geen, R., Gerber, E., Jucker, M., Maher, P., ... Thomson,  
 762 S. I. (2018, March). Isca, v1.0: A framework for the global modelling of the atmos-  
 763 pheres of Earth and other planets at varying levels of complexity. *Geoscientific  
 764 Model Development*, 11(3), 843–859. doi: 10.5194/gmd-11-843-2018
- 765 Vargas Zeppetello, L. R., Battisti, D. S., & Baker, M. B. (2019, October). The  
 766 origin of soil moisture evaporation “regimes”. *Journal of Climate*, 32(20), 6939–  
 767 6960. doi: 10.1175/JCLI-D-19-0209.1
- 768 Vial, J., Dufresne, J. L., & Bony, S. (2013). On the interpretation of inter-model  
 769 spread in CMIP5 climate sensitivity estimates. *Climate Dynamics*, 41(11-12),  
 770 3339–3362. doi: 10.1007/s00382-013-1725-9
- 771 Wang, Q., Cheng, L., Zhang, L., Liu, P., Qin, S., Liu, L., & Jing, Z. (2021). Quantify-  
 772 ing the impacts of land-cover changes on global evapotranspiration based on the  
 773 continuous remote sensing observations during 1982–2016. *Journal of Hydrology*,  
 774 598(February), 126231. doi: 10.1016/j.jhydrol.2021.126231
- 775 Wilks, D. S. (2016, March). “The stippling shows statistically significant grid-  
 776 points”: How Research Results are Routinely Overstated and Over-interpreted,  
 777 and What to Do About It. *Bulletin of the American Meteorological Society*,  
 778 160309141232001. doi: 10.1175/BAMS-D-15-00267.1
- 779 Zelinka, M. D., Randall, D. A., Webb, M. J., & Klein, S. A. (2017). Clear-  
 780 ing clouds of uncertainty. *Nature Climate Change*, 7(10), 674–678. doi:  
 781 10.1038/nclimate3402

1                   **Supporting Information for: Reduced terrestrial  
2 evaporation increases atmospheric water vapor by  
3 generating cloud feedbacks**

4                   **M. M. Lagu  <sup>1,2</sup>, G. R. Quetin<sup>3</sup>, W. R. Boos<sup>4,5</sup>**

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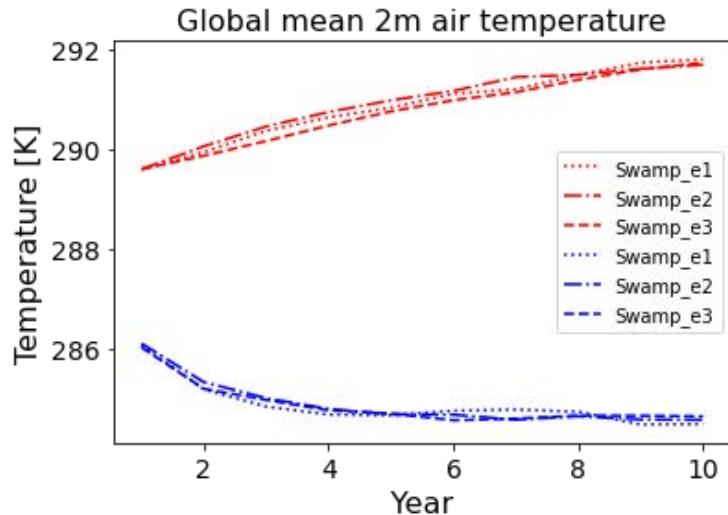
6                   <sup>2</sup>Department of Atmospheric Sciences, University of Utah, Salt Lake City, UT, USA

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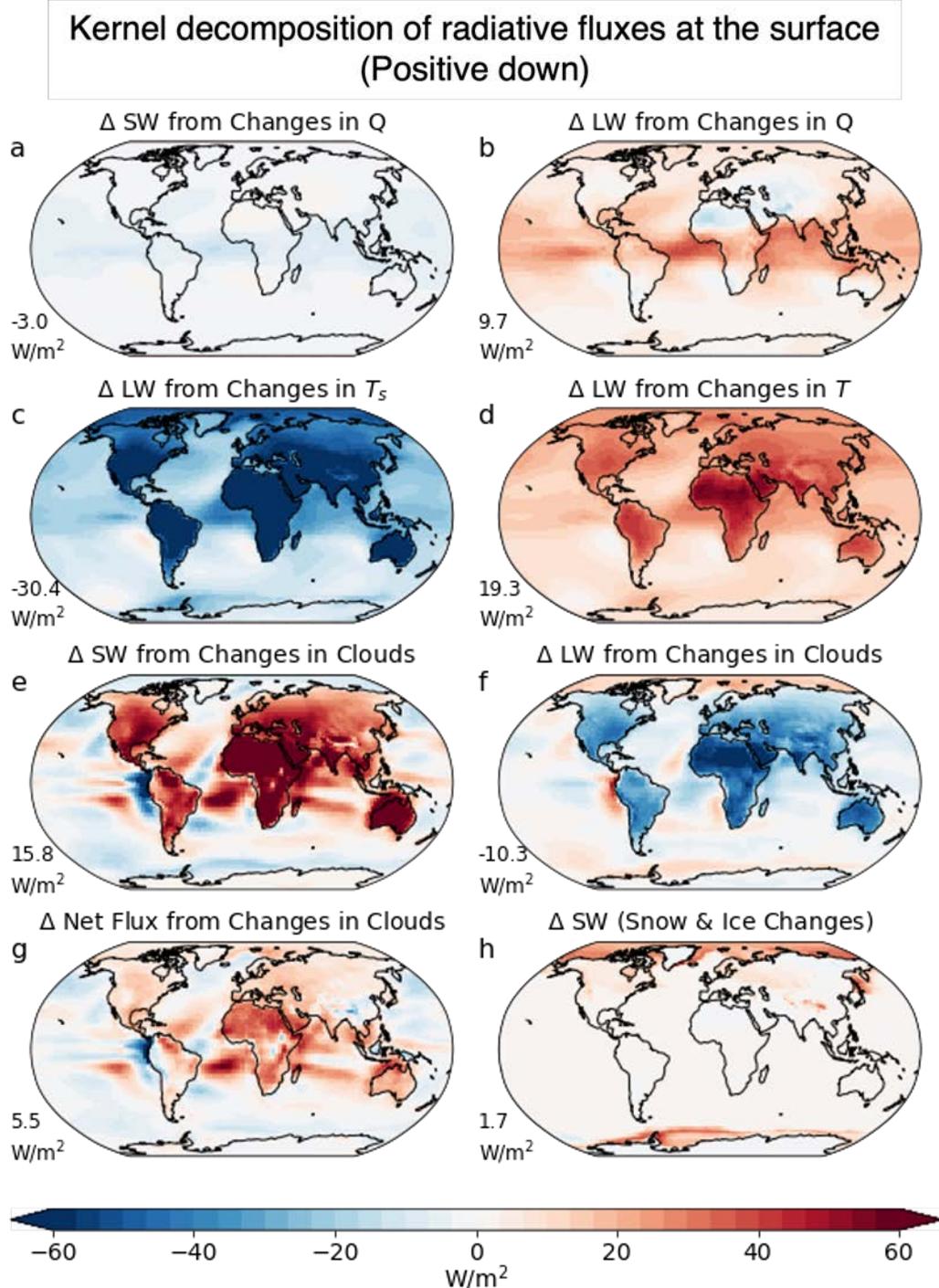
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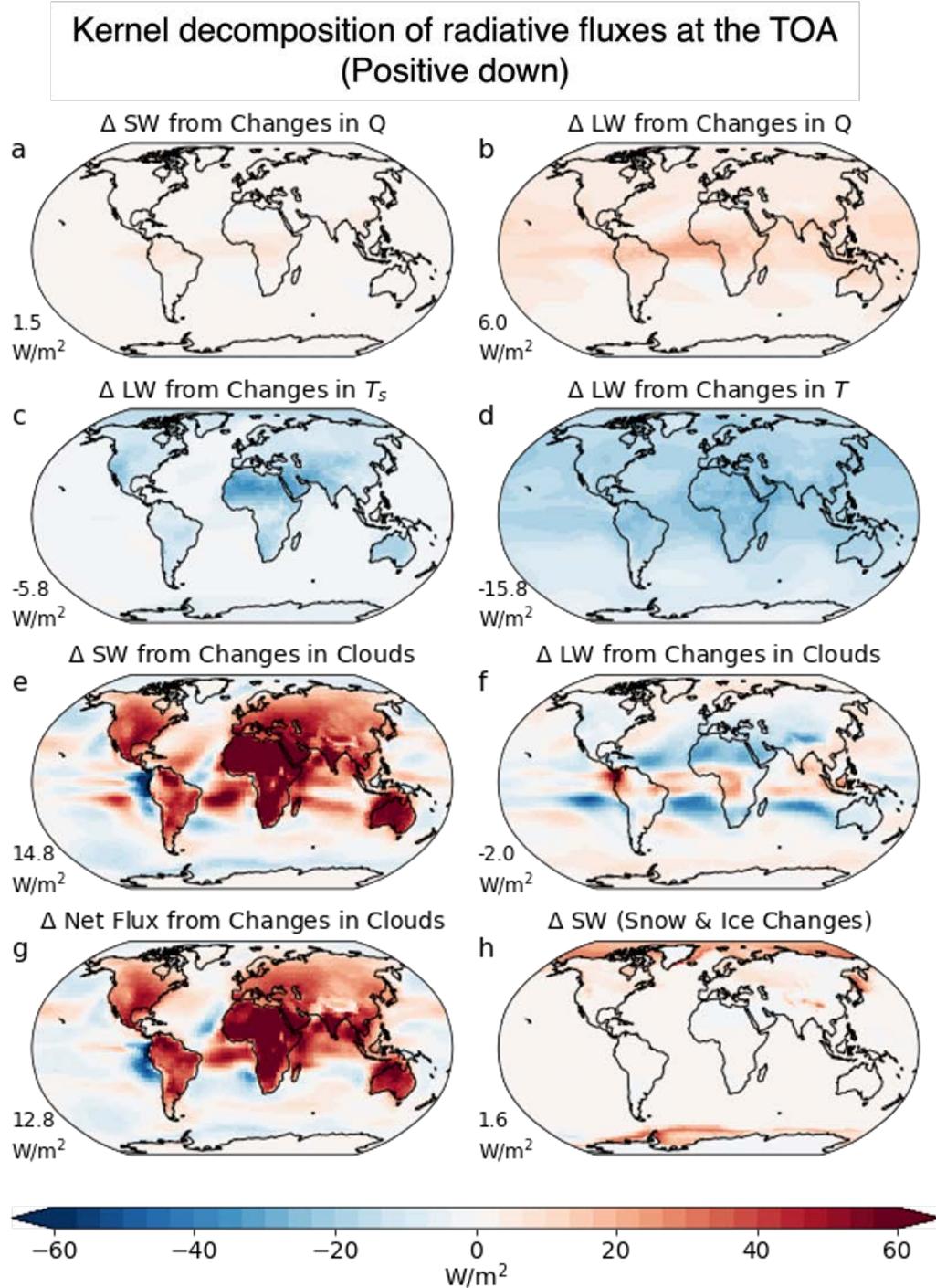
10                  **Supplemental Figures**



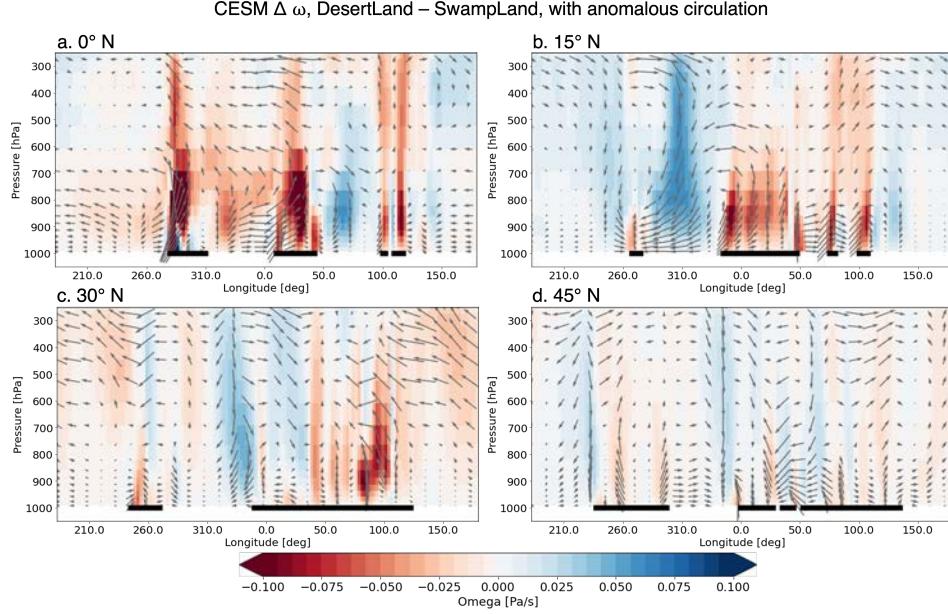
**Figure S1.** Spin-up of global mean 2m air temperatures across 3 ensemble members each of the Swamp and Desert simulations (swamp\_e1-swamp\_e3 and desert\_e1-desert\_e3, respectively), showing that the spread between ensemble members is small compared to inter-annual variability, and the difference between Swamp or Desert is much larger than the difference between ensemble members within either the Swamp or Desert simulation-type.



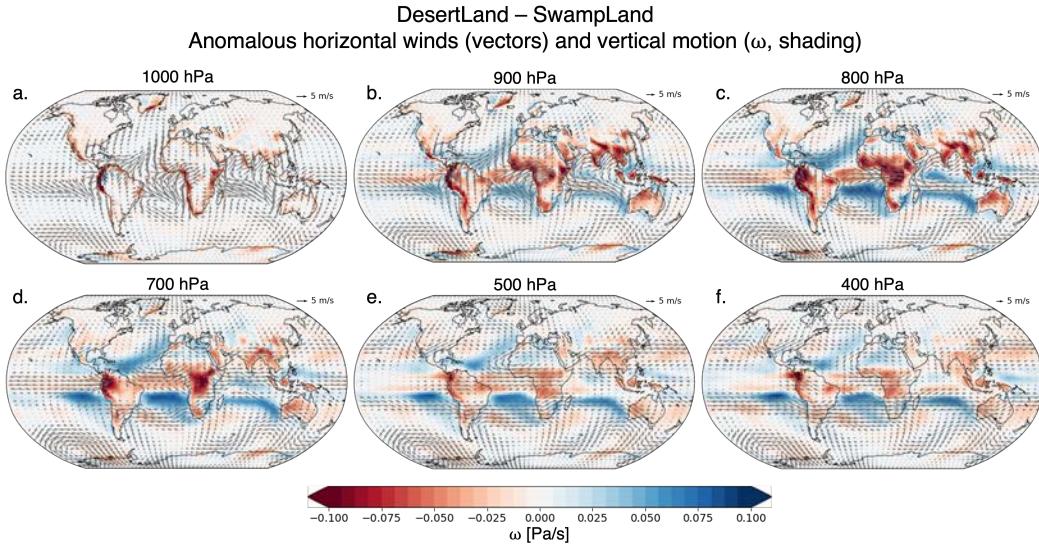
**Figure S2.** Decomposition of surface radiative fluxes using the radiative kernel for the DesertLand - SwampLand simulations. Fluxes are positive down into the surface, such that red regions indicate more energy into the surface attributable to that term in the DesertLand simulation than in the SwampLand simulation. The global mean value is noted to the lower left of each panel.



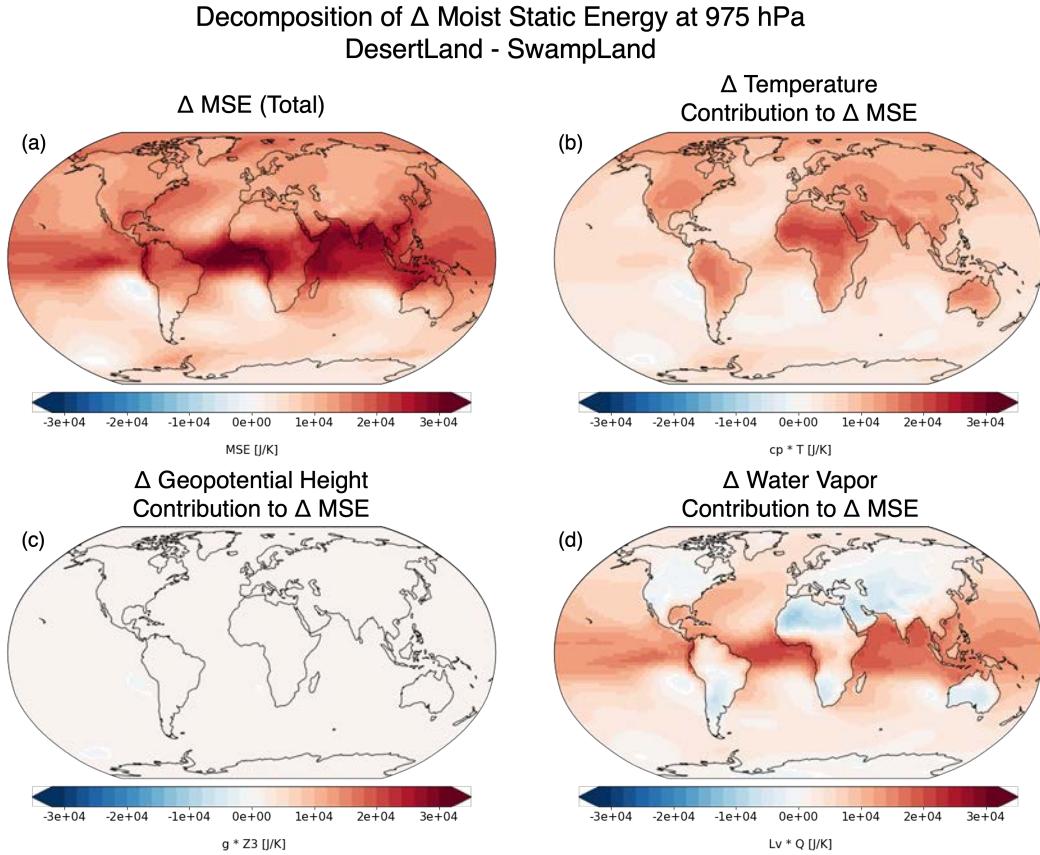
**Figure S3.** Decomposition of top of atmosphere radiative fluxes using the radiative kernel for the DesertLand - SwampLand simulations. Fluxes are positive down into the top of the atmosphere, such that red regions indicate more energy into the atmospheric column attributable to that term in the DesertLand simulation than in the SwampLand simulation. The global mean value is noted to the lower left of each panel.



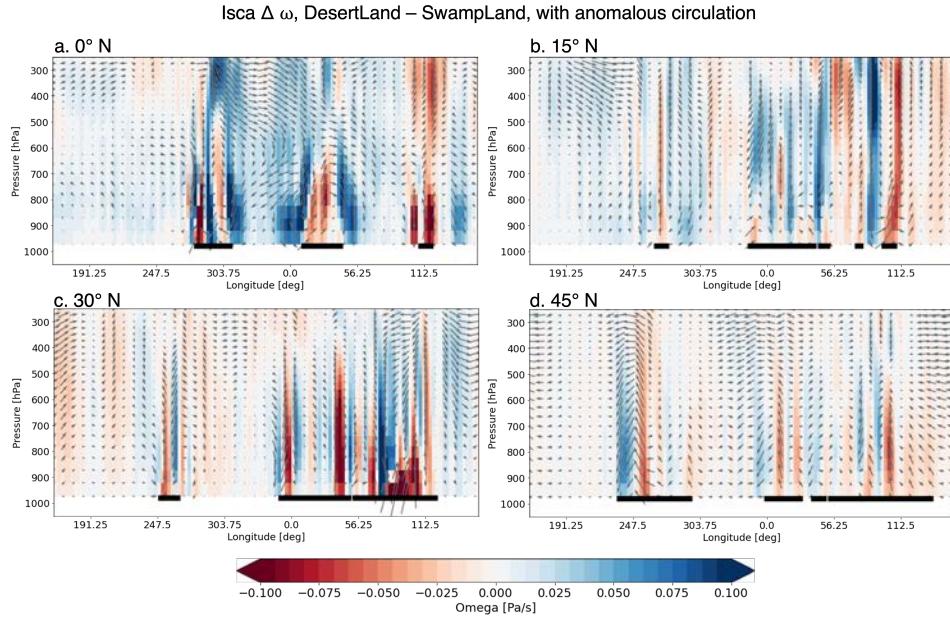
**Figure S4.** CESM vertical cross section (longitude vs. pressure) at various latitudes for the change in vertical motion  $\omega$  [Pa/s] for DesertLand – SwampLand is shown in colors, with red indicating negative  $\omega$ , i.e. positive vertical motion. Vectors show the direction of anomalous motion:  $u$  winds in [m/s] in the x-direction, and  $\omega \times -150$  in [Pa/s] in the y direction.  $\omega$  is multiplied by  $-1$  so arrows point in the direction of motion, and by  $50$  so the vertical component is of comparable magnitude to the horizontal component, for ease of visualization. Horizontal black lines at the surface indicate land masses.



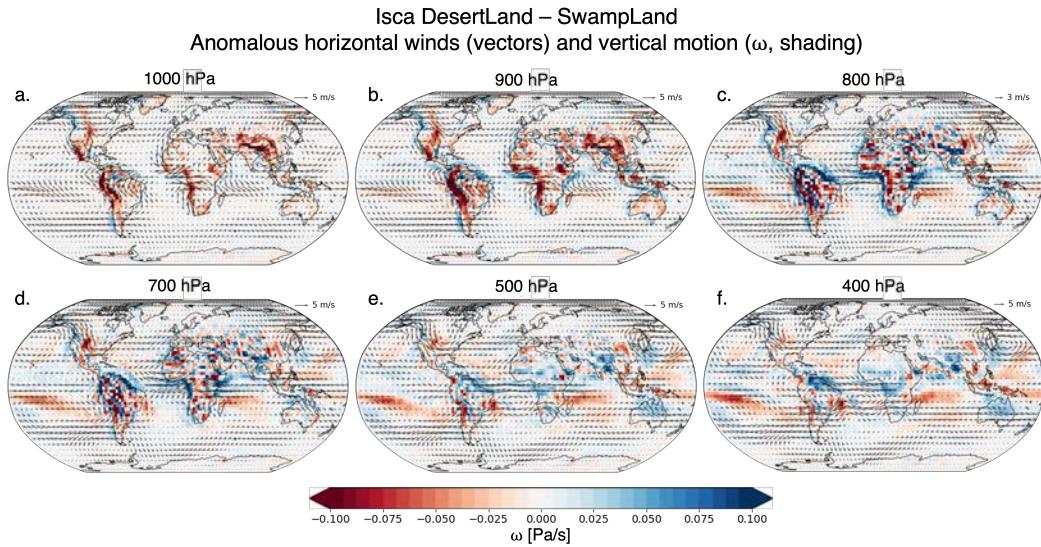
**Figure S5.** CESM change (DesertLand – SwampLand) in vertical motion  $\omega$  [Pa/s] (shading, with red indicating upwards motion) and horizontal winds  $\langle u, v \rangle$  [m/s] (vectors) at various levels of the atmosphere.



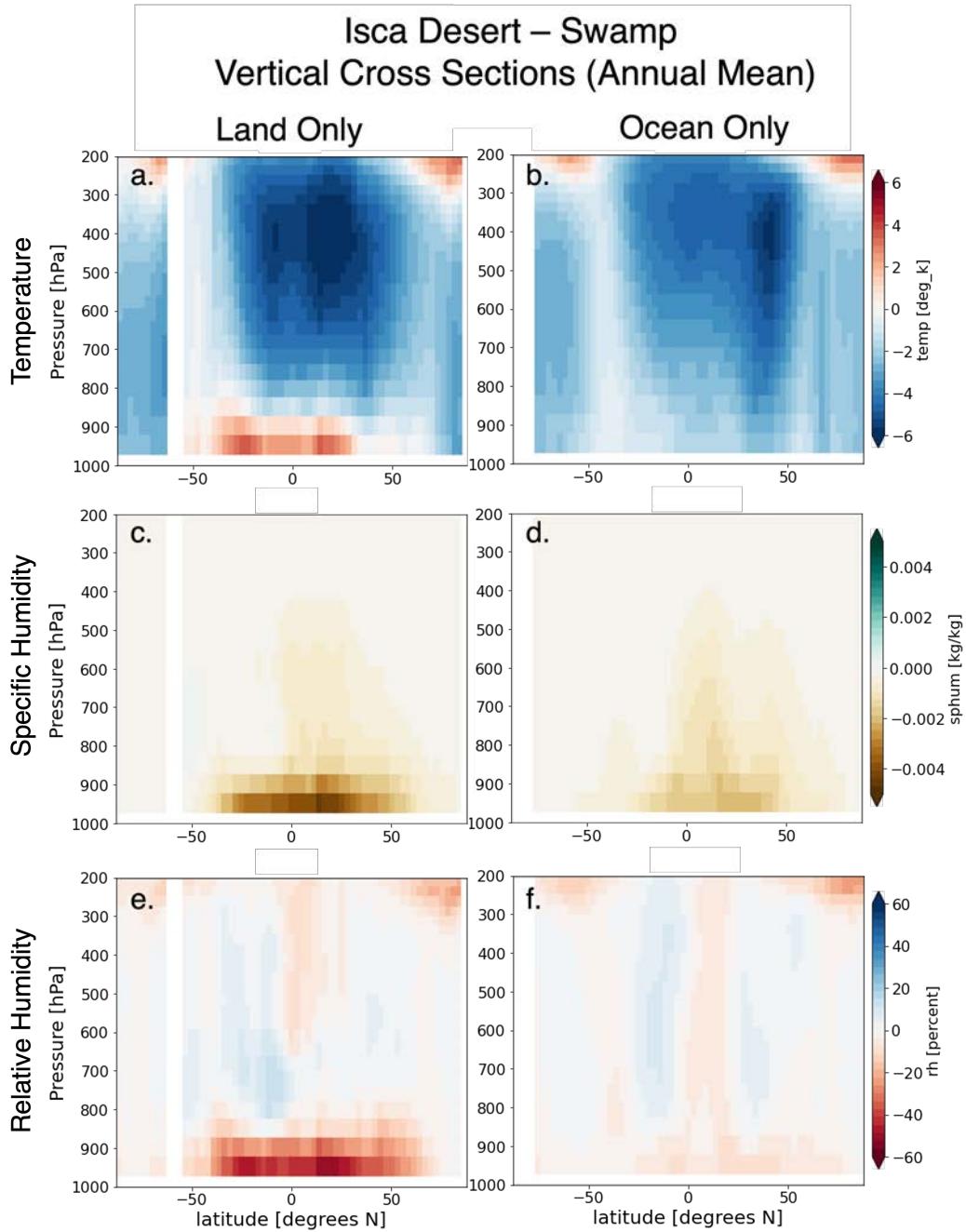
**Figure S6.** Decomposition of the change in moist static energy (MSE) into (a) the total and the components attributable to the change in (b) temperatures ( $c_p T$ ), (c) geopotential heights ( $gZ$ ), and (d) moisture ( $L_v Q$ ) at 975 hPa, for the difference between the Desert and Swamp CESM simulations. Only changes that pass a statistical test are shown, where values are significant if the p-values calculated from a student's t-test pass a false discovery rate of 0.15.



**Figure S7.** As in Fig. ??, but for Isca.



**Figure S8.** As in Fig. ??, but for Isca.



**Figure S9.** Change in the Isca simulations (DesertLand - SwampLand) in zonally averaged temperature (top), specific humidity (middle) and relative humidity (bottom) for land-only (left) and ocean-only (right).