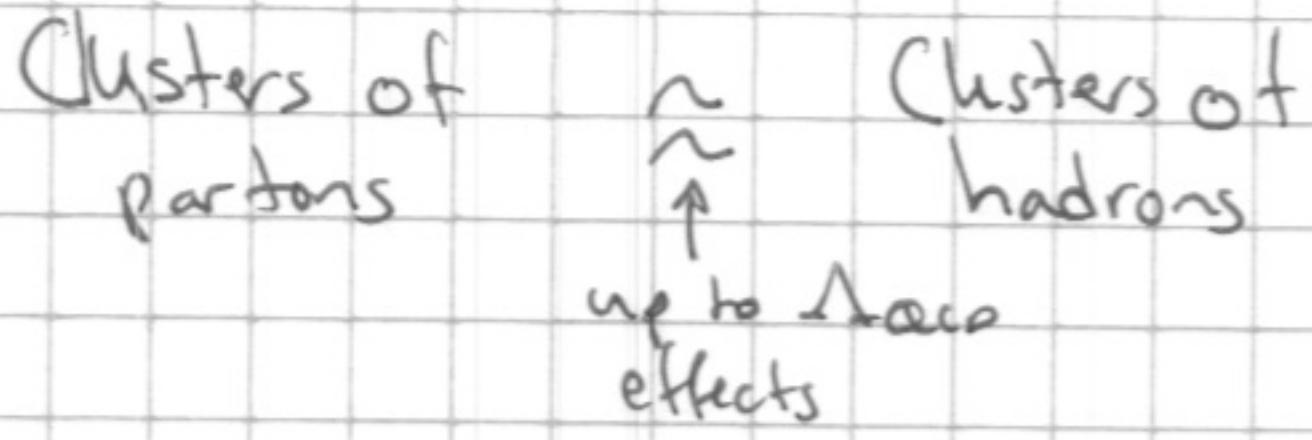


Part 2: Jet Algorithms

Last time:



- Jets are proxies for underlying quarks/gluons
 Essential for new physics searches
 Fundamentally ambiguous, but opportunity to optimize.

Today: Jet Algorithms as one way to study QCD processes

(Not only way! E.g. event shapes, probabilistic methods)

Offers unique event interpretation.

Generic Jet Algorithm

$$\{p_1, p_2, p_3 \dots p_n\}_{\text{hadrons}} \Rightarrow \{p_1, p_2 \dots p_N\}_{\text{jets}}$$

Usually:

$$p_{\text{jet}} = \sum_{i \in \text{jet}} p_i \quad (\text{E-Scheme recombination})$$

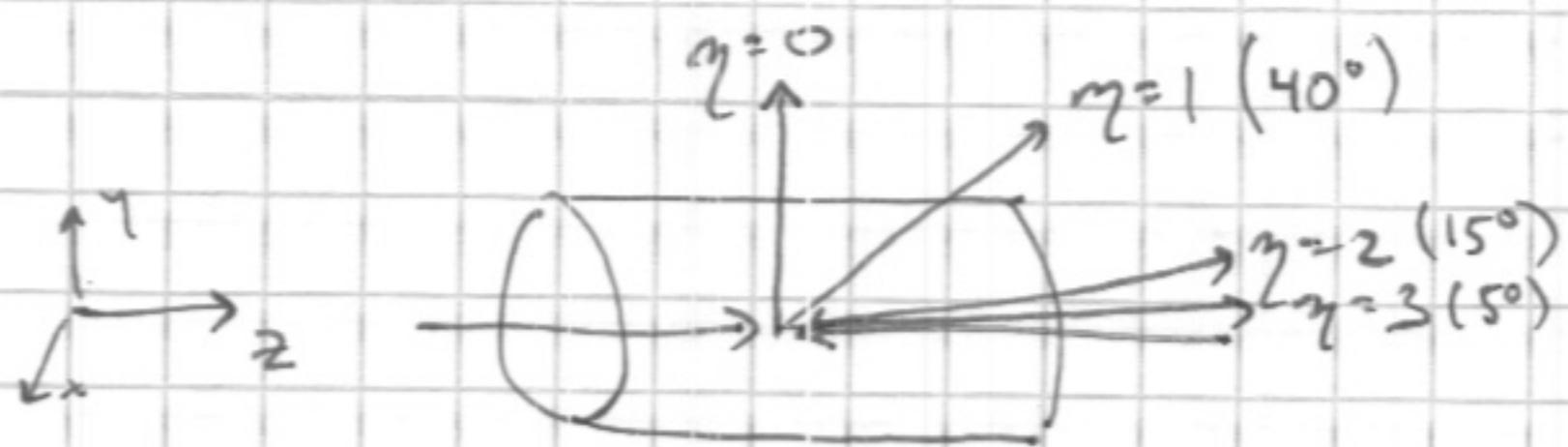
Note:

$$p_{\text{jet}}^2 = \left(\sum_i p_i \right)^2 \geq \sum_i m_i^2$$

Jets are full fourvectors with non-zero mass.

Mass has to be included for correct resonance reconstruction. (e.g. $Z \rightarrow \text{dijets}$)

Coordinate system for hadron colliders.



Typically:
longitudinally
boost invariance

$$p_T = \sqrt{p_x^2 + p_y^2}$$

$$\phi = \arctan \frac{p_y}{p_x}$$

$$\eta = \text{arctanh} \frac{p_z}{|\vec{p}|}$$

$$y = \frac{1}{2} \ln \frac{E + p_z}{E - p_z}$$

true rapidity

(longitudinal)

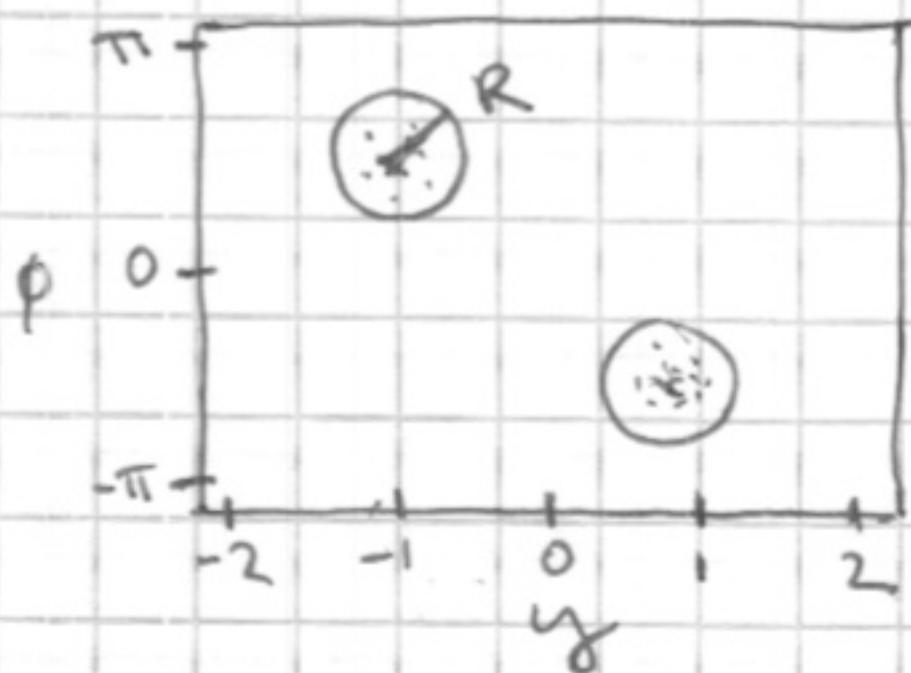
Boost-invariant distances

$$\Delta R = \sqrt{(\Delta\phi)^2 + (\Delta y)^2}$$

pseudorapidity
(only boost invariant
difference for massless particles)

$$= -\ln \tan \frac{\theta}{2}$$

Sets from cone algorithms?



Punch holes of radius R.

Immediate challenges;
which holes?
Split/Merge ambiguity

Cone algorithms are intuitive: flow of energy in a given direction within a given angle

Subtleties encountered in Tevatron era
(e.g. TPC unsafe seeds)

Should probably be revisited (more later)

(16)

What is ΔR ? (Beyond being just boost invariant.)

Standard opening angle: $d\Omega^2 = d\theta^2 + \sin^2\theta d\phi^2$

Exercise: Using $\sin\theta = \frac{1}{\cosh\eta}$ for massless particles

$$\text{Show: } d\Omega^2 = \frac{1}{\cosh^2\eta} \left(d\eta^2 + d\phi^2 \right) \approx dR^2$$

So ΔR is just like opening angle, rescaled by factors of $\cosh\eta$. ($\cosh(\eta=0)=1$)

For massless particles: $E = p_T \cosh\eta$

$$\Rightarrow E d\Omega \approx p_T dR \quad (\text{nice thing to remember})$$

\Rightarrow Jets are smaller in forward ($|\eta|$) direction.

Q : Why $\underline{p_T}, \Delta R$ instead of $\underline{E}, \Delta \eta$?

longitudinally
boost invariant

used extensively
in e^+e^- colliders

Ok answer: PDFs mean we don't know longitudinal momentum fraction of colliding quark/gluons.

Better answer: Rapidity Plateau: in "minimum bias" collisions, number of particles roughly constant as a function of η .

Q: Why fixed R ?

Actually, there is interest in variable R ,
dynamical R , multiple R ...

Right now, $R=0.4$ anti- k_T jets very standard

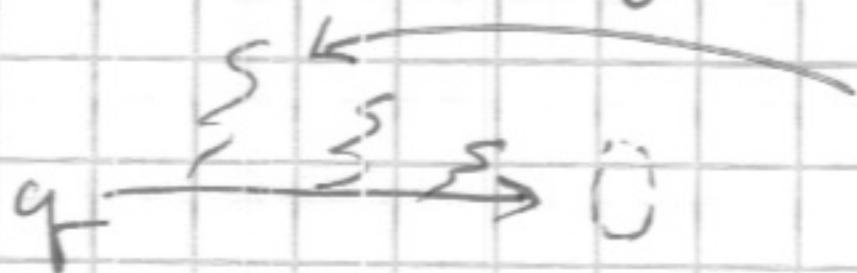
\hookrightarrow ($\sim 22^\circ$ half angle, $\sim 45^\circ$ full angle)

vs. 0.1×0.1 HCAL granularity

Fixed $R \Rightarrow$ roughly uniform jet contamination
as a function of η .

Q: Why not very small R ?

Perturbative $\log R$ resummation needed
Sensitivity to hadronization, though

 "out of cone" corrections.

The way I think of jet algorithms: Optimization problem

What do you want to optimize? (many choices...)

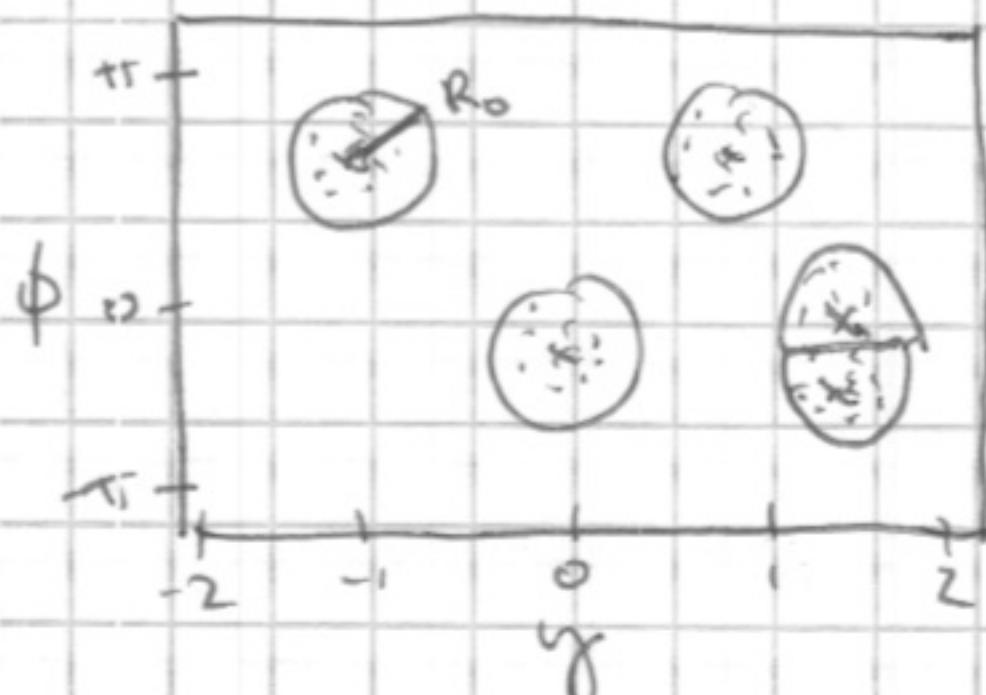
→ Hard partons should be clustered into jets

→ Jet axis \approx jet momentum direction

→ Roughly equal area (for contamination)

All of these can / should be revisited.

E.g. I want exactly N jets with axes: $\hat{n}_1, \hat{n}_2, \dots, \hat{n}_N$



Minimize τ_N

\Rightarrow axes that align
with hardest (by p_T)
clusters

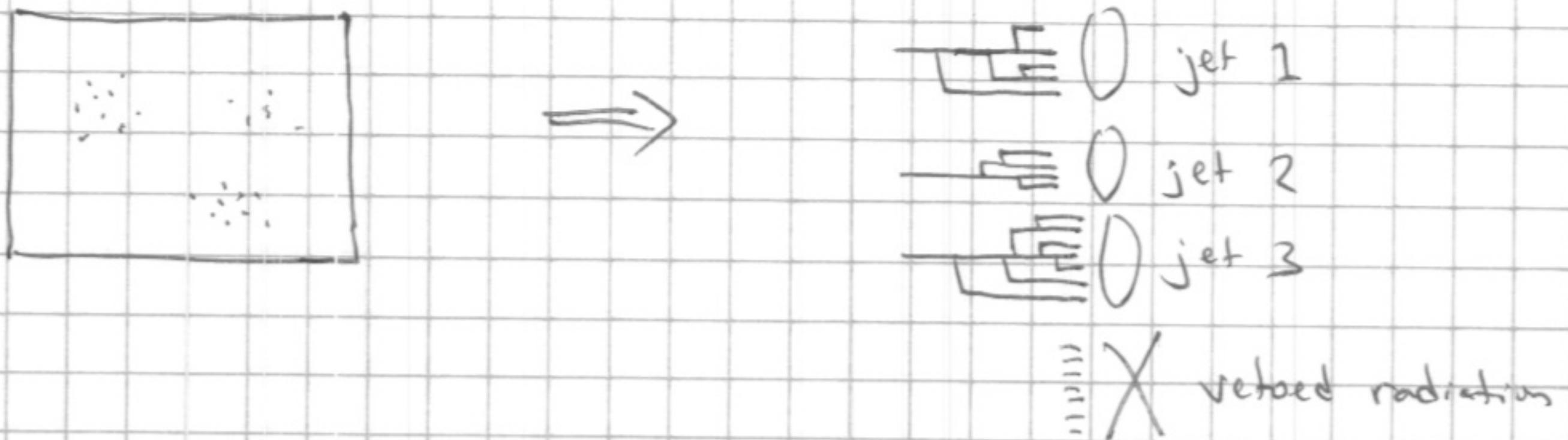
Basis for XCone jet algorithm.

Challenge: Computational expensive to fully optimize
Have to choose N in advance.

(Standard problems for cluster optimization
in Computer Science.)

Workhorse at the LHC: Recursive Jet Algorithm

(a.k.a hierarchical clustering trees)



Instead of global optimization, "local" optimization
↳ in tree-space.

Should I combine two protojets together or
should I declare a protojet finished?

Recursively
calculate
and take smallest

d_{ij}
Merge

d_i
Stop

(Long history,
going back to
JADE in 1980s)

Nice properties: Works the same on partons / hadrons / calo cells

Computationally very efficient ($n \log n$)

Defined perturbatively to all α_S orders
(Mostly) pathology-free.

Note: $P_{ij} = P_i + P_j$ in E-scheme

Let's make anti- k_t ! (Almost every LHC jet is anti- k_t)

Goal 1: Prioritize hardest jets

$$d_i = \frac{1}{P_{T,i}^2}$$

First remove jets of highest P_T

Goal 2: Jets should soak up all radiation within radius R_0 .

Want $d_{ij} < d_i$ if $R_{ij} < R_0$

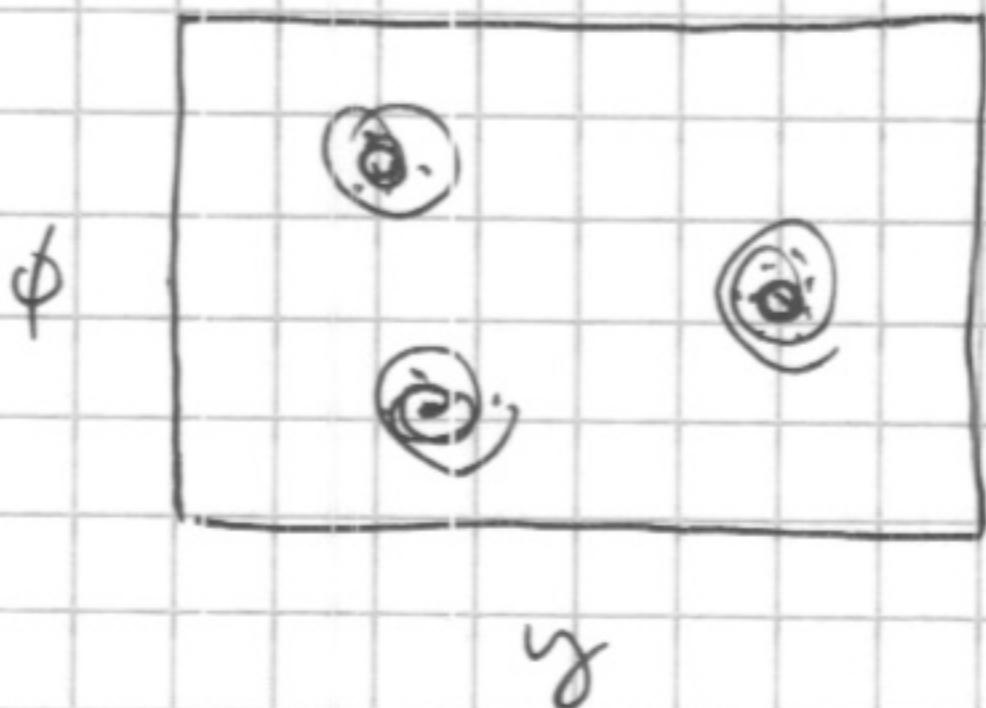
$$d_{ij} = \min \{d_i, d_j\} \left(\frac{R_{ij}}{R_0} \right)^2$$

$$d_{ij} = \frac{1}{\max \{P_{T,i}^2, P_{T,j}^2\}} \left(\frac{R_{ij}}{R_0} \right)^2$$

Done!

That's anti- k_t

Anti- k_t : Builds jets from hard to soft



Roughly: take hardest particle and merge with nearest neighbor until R_0

For me: A revelation when I read this paper in 2008
 (Why did it take so long?)

Different goals?

k_T

Cluster according
to singularity
structure of
QCD

$$d_{ij} = \min \left\{ p_{Ti}^2, p_{Tj}^2 \right\} \left(\frac{R_{ij}}{R_0} \right)^2$$

$$d_i = p_{Ti}^2$$

C/A

Cluster by
angle alone

$$d_{ij} = \frac{R_{ij}^2}{R_0^2}$$

$$d_i = 1$$

Both of these yield "amoeba jets"
with irregular shapes

Still, very important for
jet substructure.

Choice of jet algorithm depends on physics of interest!

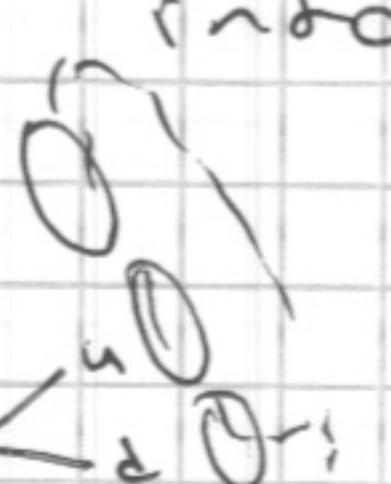
E.g.

Anti- k_T

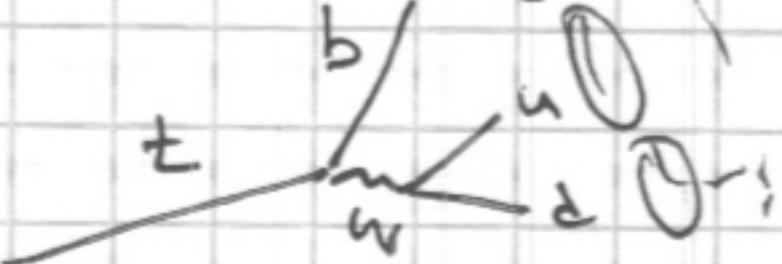
$R=1.0 \Rightarrow$ reclustered \rightarrow subjets

jets

3 k_T



for



boosted top
identification.

E.g. Anti- k_T
 $R=0.4$

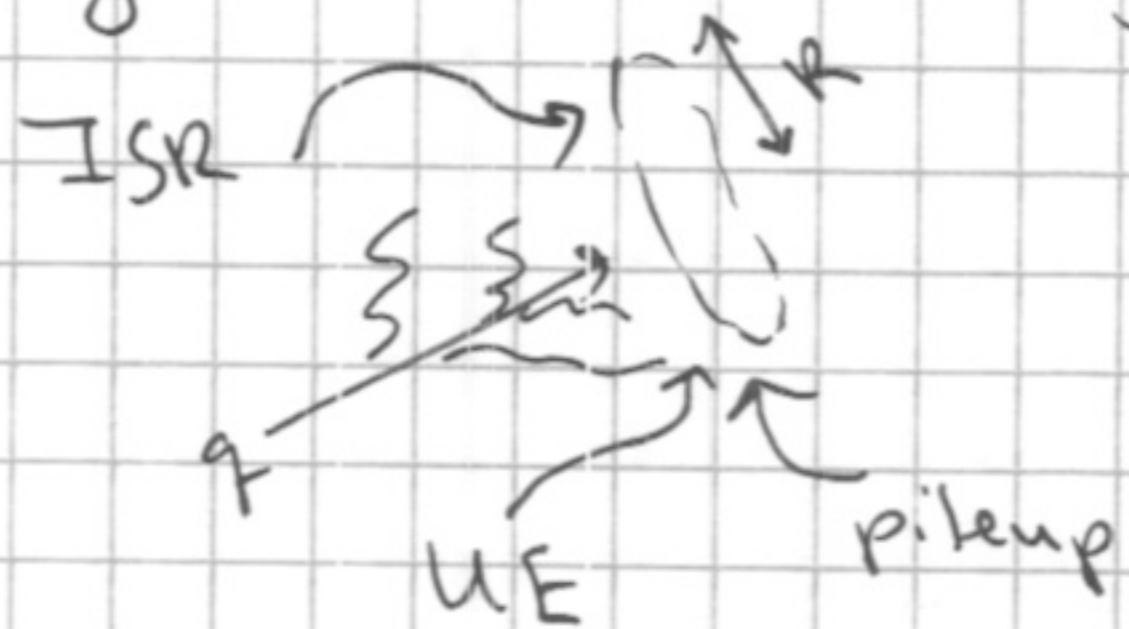
\Rightarrow Good balance between in-and-out-
of-cone corrections.
Well defined jet area (usually)

Periodically useful to revisit standard jet lore

- Longitudinal boost invariance?
- Manifest energy / momentum conservation? ($P_{12} = p_1 + p_2$)
- Fixed jet radius?
- Seedless jet algorithms?
- Unique assignment of hadrons to clusters?
- Unique jet reconstruction per event?

Your job is to question your elders!

E.g. Well-defined jet area?



Want to subtract contamination from pileup

$$P_{\text{jet}}^{\text{corr}} = P_{\text{jet}} - \rho A_{\text{jet}}$$

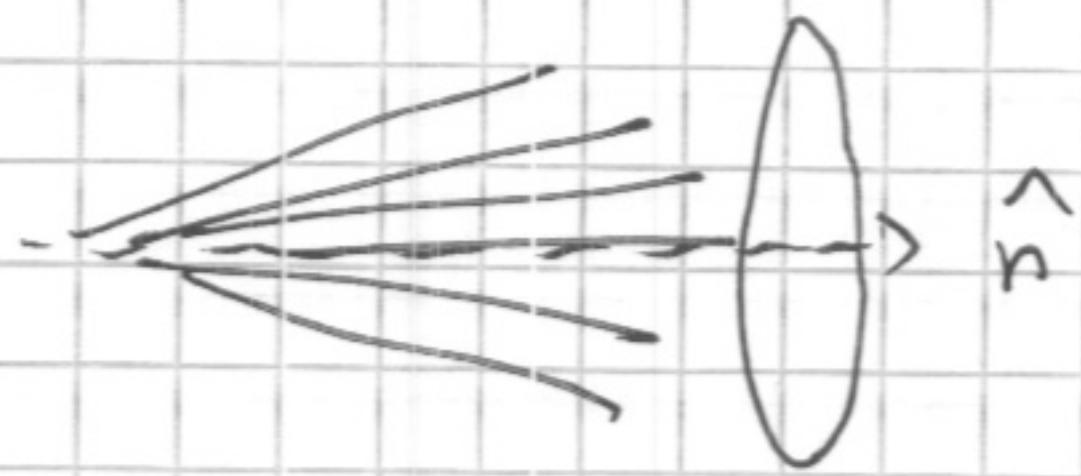
↑
jet area
contamination
per unit θ, ϕ
(rough 1 GeV per
unit per PV)

But why not build an algorithm that refuses to cluster soft particles to begin with?

"Zero Area Jet Algorithms?"

↳ under active investigation

E.g. Jet axis parallel to jet momentum?



$$\hat{n}_{\text{jet}} = \frac{\vec{p}_{\text{jet}}}{|\vec{p}_{\text{jet}}|}$$

In terms of optimization, very natural:

$$\chi_1 = \sum_i p_{T,i} \left((\phi_i - \phi_A)^2 + (y_i - y_A)^2 \right)$$

\sum_i \sum_i

$\sum_i p_{T,i}$

$$\frac{\partial \chi_1}{\partial \phi_A} = 0 = \sum_i 2 p_{T,i} (\phi_i - \phi_A)$$

\sum_i

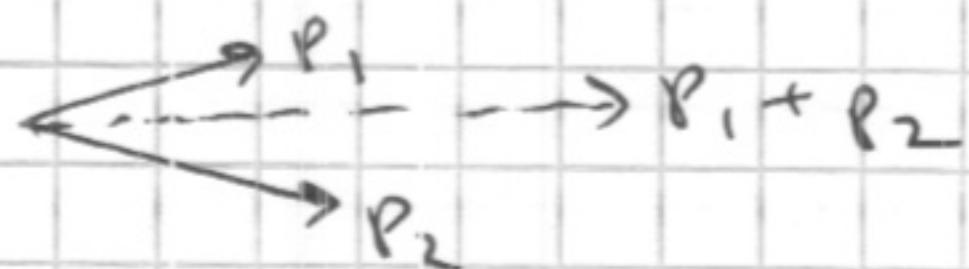
\hookrightarrow solved for $\phi_A = \frac{\sum_i p_{T,i} \phi_i}{\sum_i p_{T,i}}$

mostly the same.

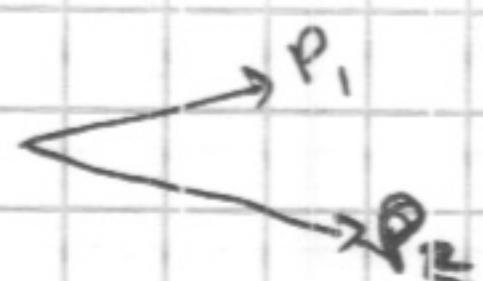
Same as finding "mean"

Why not "median"? $\chi_i = \sum_i p_{T,i} R_{i,A}^{-1}$

In language of clustering:



E-scheme axis. ("mean")



Winner-take-all axis ("median")

$$\frac{\vec{p}_2}{|\vec{p}_2|} (|\vec{p}_1| + |\vec{p}_2|)$$

Like median being robust to outliers,
WTA axis robust to contamination.

One lore that you challenge at your peril!

Infrared / Collinear Safety

(more next lecture)

(By the way, I've broken IRC safety with some success. Still, beware!)

Why IRC Safety?

Key success of QCD:

Fixed-order (and resummed) calculations
for jet cross sections.

Multi-loop, multi-leg, multi-log. (A whole separate
lecture series.)

Fair warning: Loop diagrams have UV divergences
(fix with renormalization)

and IR divergences

(cancel against real
emission diagrams)

$$\frac{d\sigma}{de} = \sum_N \int d\mathbf{I}_N \frac{d\sigma}{d\mathbf{I}_N} \delta(e - \hat{e}(\mathbf{I}_N))$$

↑
N-body
phase space

↑
measurement
function.

↑
IR divergences
at every α_s
order

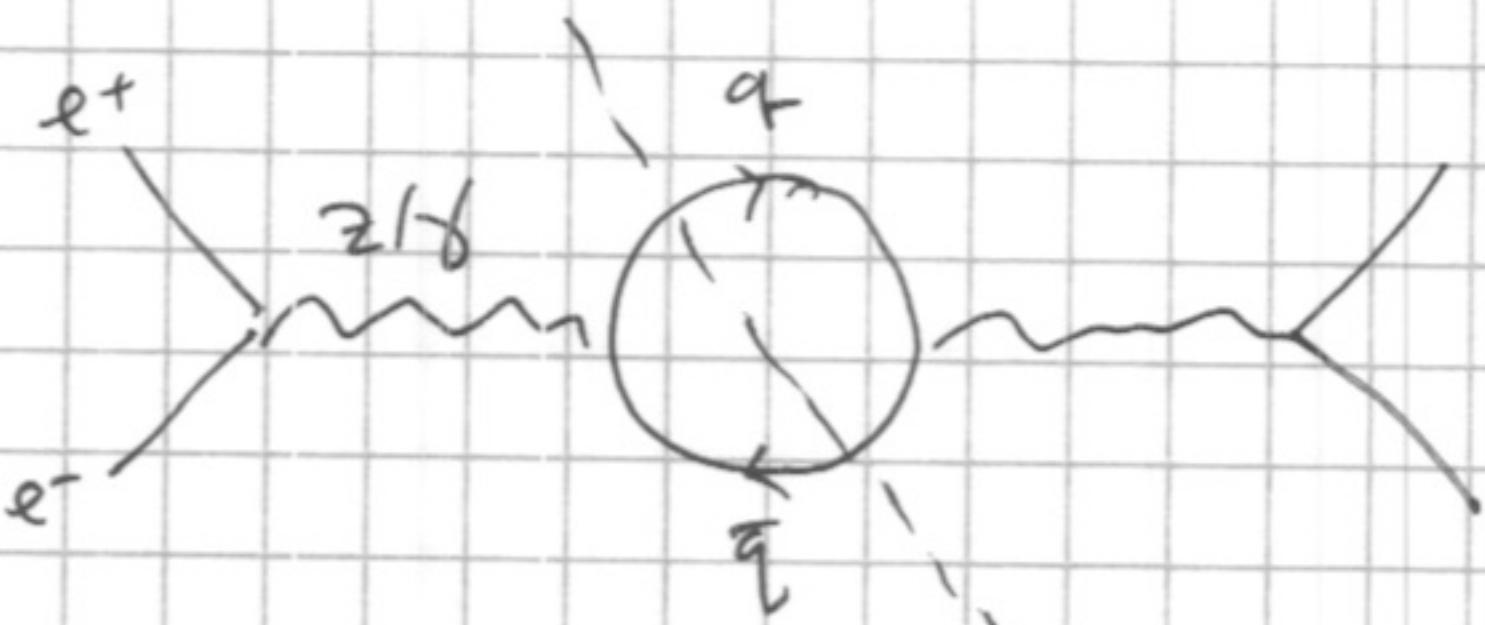
Finite if e
is IRC safe
observable

IRC Safety required for sensible perturbative calculations.
→ well, not quite (e.g. Sudakov safety)

For total cross sections, it's easy to understand by Optical theorem.

$$\sigma_{2 \rightarrow \text{anything}} \iff \text{Im } M_{\text{forward}}$$

e.g. $e^+e^- \rightarrow \text{hadrons} \iff e^+e^- \rightarrow e^+e^-$



$$\mathcal{O}(\alpha_s)$$

Imaginary part
from cutting rules.



$$\mathcal{O}(\alpha_s^2)$$

$\dots e^+e^- \rightarrow q\bar{q}\gamma$ (tree)

$\dots e^+e^- \rightarrow q\bar{q}$ (one-loop)

singularities
have to cancel!

Somewhat counterintuitively!

$2 \rightarrow 2$ tree-level cross section is good approximation to total cross section
(not exclusive 2 jet xsec)

$2 \rightarrow 2$
loop

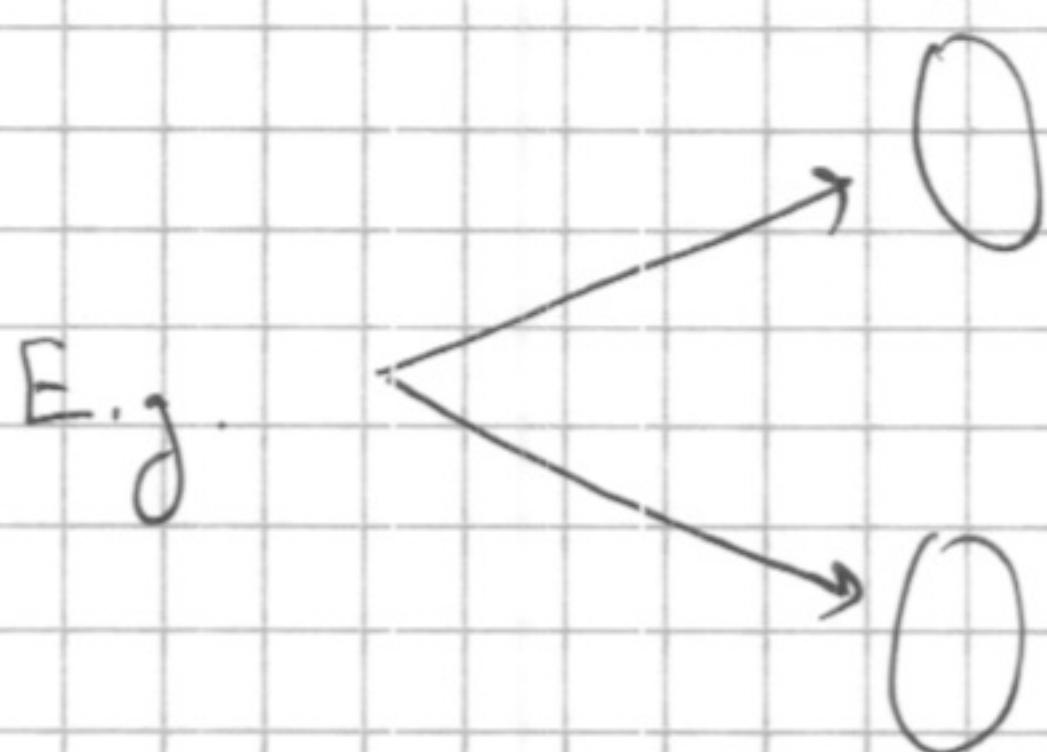
$2 \rightarrow 3$
tree

largely cancel.

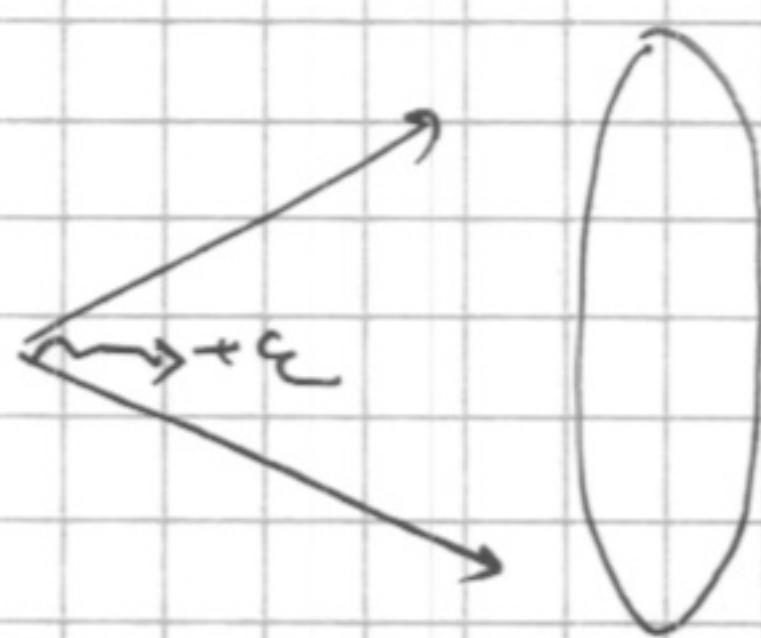
For jet cross sections, JRC safety is a bit non-trivial, and people have gotten in trouble over the years.

Heuristic for JRC Safety (almost always works)

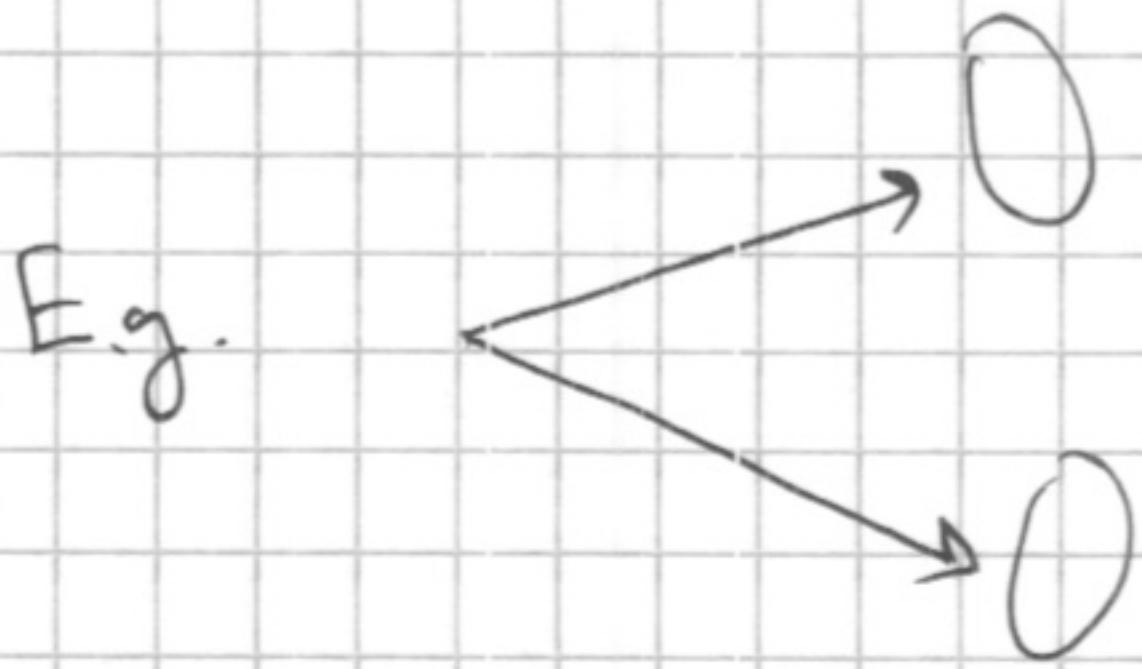
Soft safety: Answer should not change if you add infinitesimally soft particle



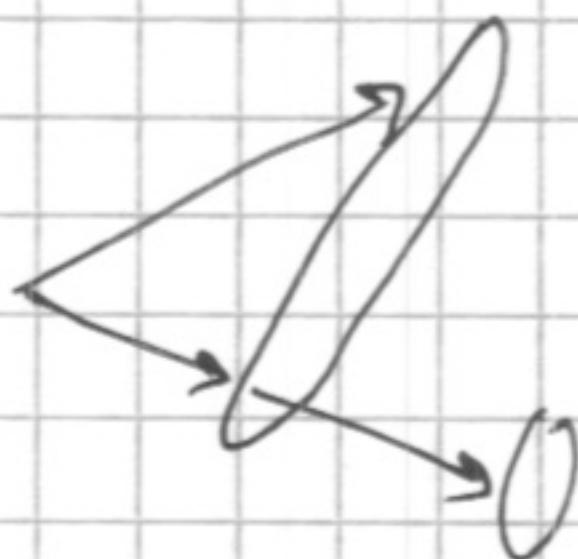
should not
= = =



Collinear safety: Answer should not change if you split a particle exactly in two



should not
= = =



Among jet aficionados, JRC Safety is the first thing we check about any new idea.

Not Safe? Not calculable in fixed-order perturbative QCD!

(Perhaps other strategies available, though)

Is anti- k_T safe?

$$d_{ij} = \frac{1}{\max\{P_{T,i}^2, P_{T,j}^2\}} \frac{R_i j^2}{R_0^2}$$

$$d_i = \frac{1}{P_{T,i}^2}$$

Soft emission? $d_i^{\text{soft}} \rightarrow \infty$ (no effect)

d_{ij}^{soft} could be important!

but infinitesimally soft particle
does not change jet P_T

Depends on question!

Anti- k_T jet P_T ? Soft safe

Anti- k_T jet multiplicity of constituents? Soft unsafe.

Collinear split?

$d_i^{\text{collinear}}$ changes!

But $d_{ij}^{\text{collinear}} = 0$, so first thing you do is put split back together.

So yes, anti- k_T is IRC safe.

NNLO jet cross sections being actively developed.

Next time: What exactly are soft / collinear singularities, and can we use them to understand quark/gluon tagging.