

Part 3: The Soft/Collinear Limit of QCD

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Last time: Jet Algorithms, workhorse of hadronic final states

Jet as proxy for parton \rightarrow Fundamentally ambiguous
but still essential

Anti- k_t is current standard \rightarrow Circular jets, area πR^2
IRC safe

Today: a jet is more than \vec{p}_{jet} !

Jets have constituents and you can measure its substructure

Example: Using soft/collinear limit of QCD to
understand quark/gluon discrimination

Pedagogical choice: I'm going to tell you a key fact,
not derive it, and go from there.

Glosses over many cool aspects of
QCD, but limited time, and I
think this is best for lectures on jets.

Exercise: Show that in soft & collinear limit

$$\sum_{\text{polarizations}} \int dT_{\text{soft}} \left| \text{Diagram} \right|^2 \quad (\text{with } q \text{ and } g \text{ in the same jet})$$

$$\approx \sum_{\text{polarizations}} \int dT_n \left| \text{Diagram} \right|^2$$

$$\times \int_0^1 dz \int_0^R d\theta \frac{2\alpha_s}{\pi} C_F \frac{1}{z} \frac{1}{\theta}$$

color factor,
Swap to C_A for gluons
soft/collinear singularities

Energy fraction : $z = \frac{E_{\text{gluon}}}{E_{\text{jet}}}$ $z \rightarrow 0$ soft limit

$$\int_0^z \theta \rightarrow 1-z$$

Spl. tting angle : $\Theta = \Theta_{qg}$ $\Theta \rightarrow 0$ collinear limit

Color factors:

$$SU(N_c) \rightarrow SU(3)$$

Quarks: $\sum_{a,j} t_{ij}^a t_{jk}^a = C_F \delta_{ik}$, $C_F = \frac{N_c^2 - 1}{2N_c} \rightarrow \frac{4}{3}$

Gluon $\sum_{a,b} f^{abc} f^{abd} = C_A \delta^{cd}$, $C_A = N_c \rightarrow 3$

Note: Above only true for soft & collinear limit

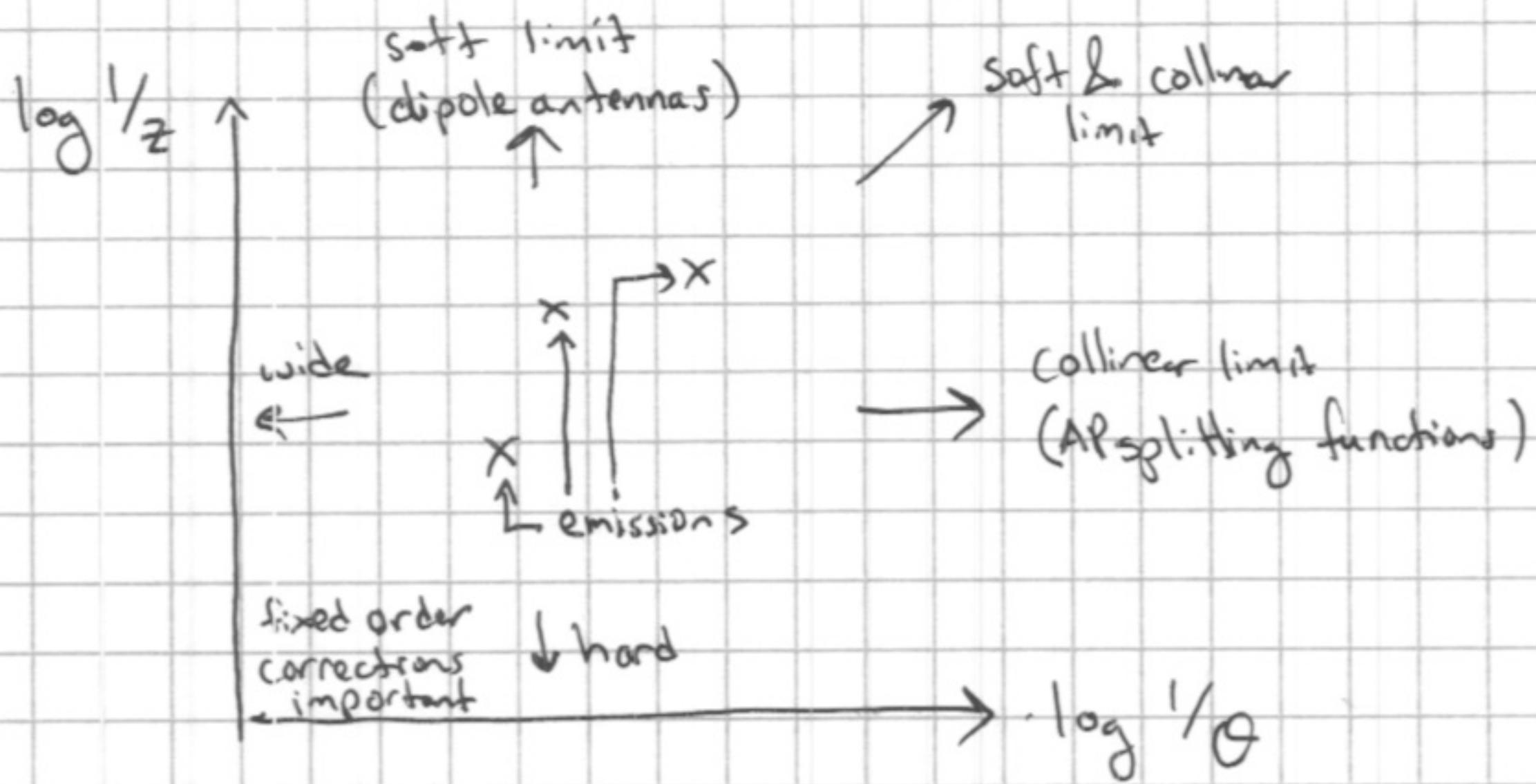
In this extreme limit, splitting probability is:

$$dP_{i \rightarrow i+g} = \frac{2 ds}{\pi} \frac{dz}{z} \frac{d\theta}{\Theta}$$

↑
soft gluon

(30)

Uniform emissions in the $(\log \frac{1}{\Theta}, \log \frac{1}{z})$ plane

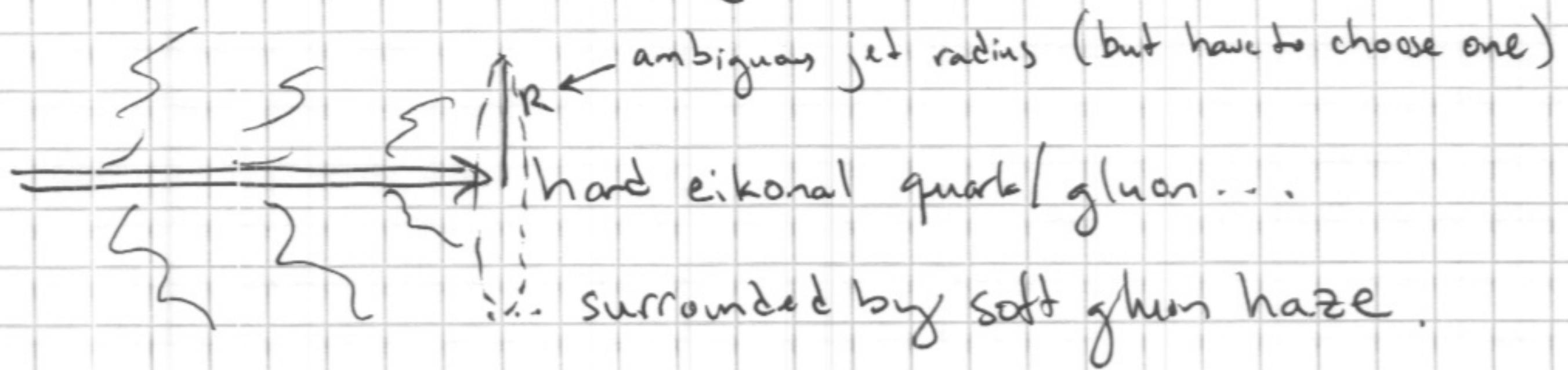


In "strongly-ordered" (or "leading log" or "double log") limit,
we can do simple but insightful calculations.

Very low accuracy, but gives a flavor for
questions worth asking.

Many approaches to systematically improving this picture,
e.g. SCET, direct resummation, ...

The (simplest) picture of a jet



In fact, basis for parton shower, which recursively applies $P_i \rightarrow jk$ (for all parton flavors)

Benefit: Some information at all orders in α_s .

Often more realistic qualitatively than fixed-order calculations.

Caution: Have to think carefully about quantitative accuracy.

Hard work to convince yourself that it is systematically improvable.

Insight: Jet is clearly not a single parton

(Wilson-line-wrapped eikonal parton)

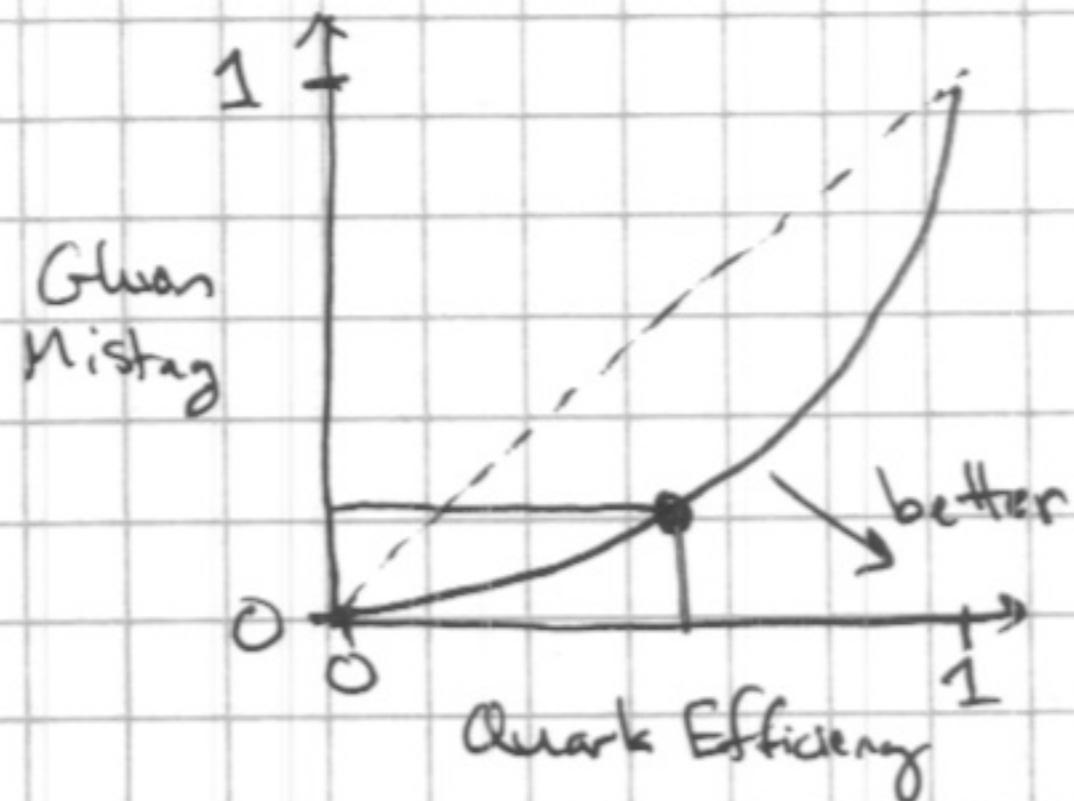
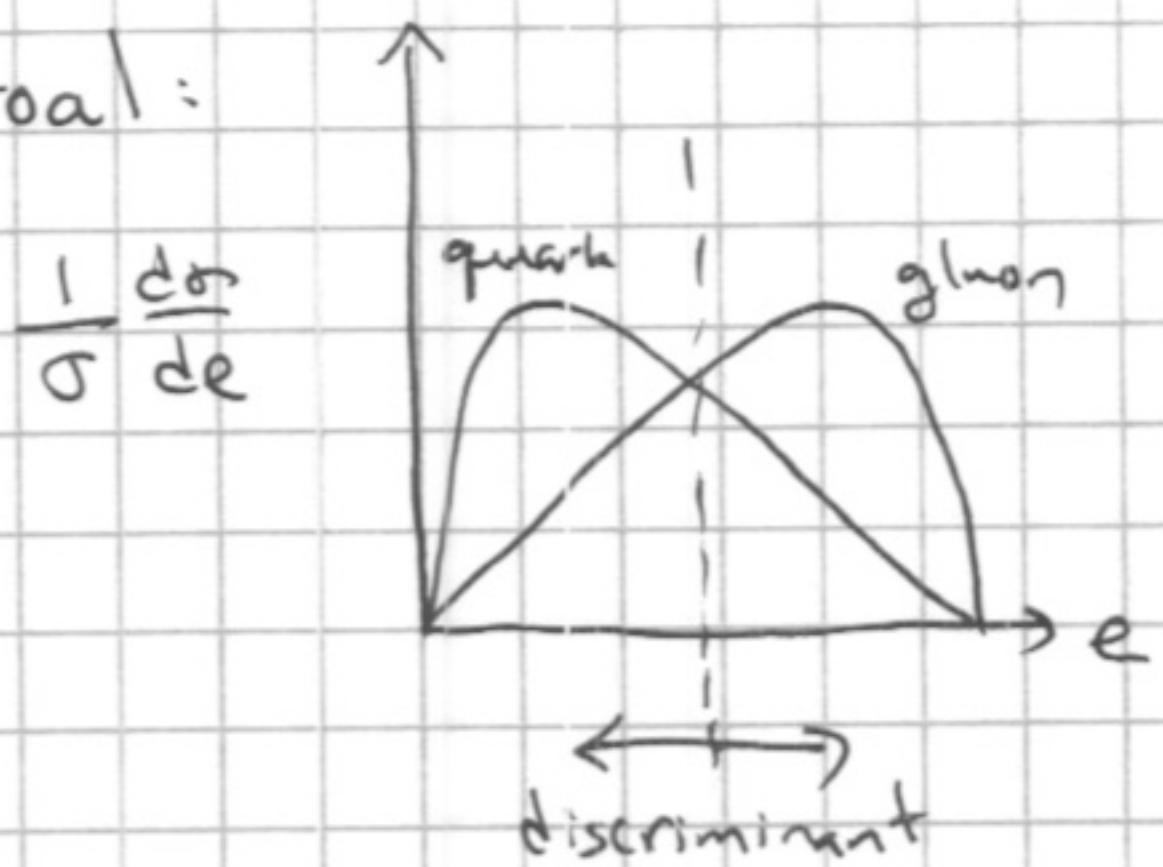
Very instructive to always visualize the haze.

Concrete question: Quark vs. Gluon Jets?

$C_F < C_A$, so gluon jets should be "fatter"

Need to define observables sensitive to this difference

Goal:



Ideally: Predict from first principles QCD

Validate in Monte Carlo parton showers

Test in LHC data

Why do we care?

Experimentally: Different calibrations for different jet types

QCD Theory: Observables sensitive to color structure of QCD

BSM Theory: Often, BSM signals are quark-rich while backgrounds are gluon-rich.

Choice of discriminant observable?

→ IRC Safe (in order to perform perturbative calculations)

→ Not overly sensitive to hadronization or other non-perturbative effects.

Want $\delta e \sim \left(\frac{\Delta QCD}{E_{jet}} \right)^{\beta}$

Simplest example I know: Energy-Energy Correlation Functions

$R_j \rightarrow E$ $R \rightarrow \theta$ for simplicity

$$e_2^{(\beta)} = \frac{\sum_{i \neq j} E_i E_j \theta_{ij}^{\beta}}{\left(\sum_i E_i \right)^2}$$

\uparrow

2-point
correlator for
1-prong testing

$\beta > 0$ for IRC safety
Also known as $C_1^{(\beta)}$.
 \uparrow
1-prong testing
with 2-point
correlator

Soft safe? Yes, e_2 doesn't change if $E_i = 0$

Collinear safe? Yes, e_2 is additive, so $E \rightarrow E_1 + E_2$
has no effect (for $\beta > 0$)

What is quark/gluon discrimination power for e_2 ?

Go to extreme soft & collinear & strongly-ordered limit of QCD (borderline physical)

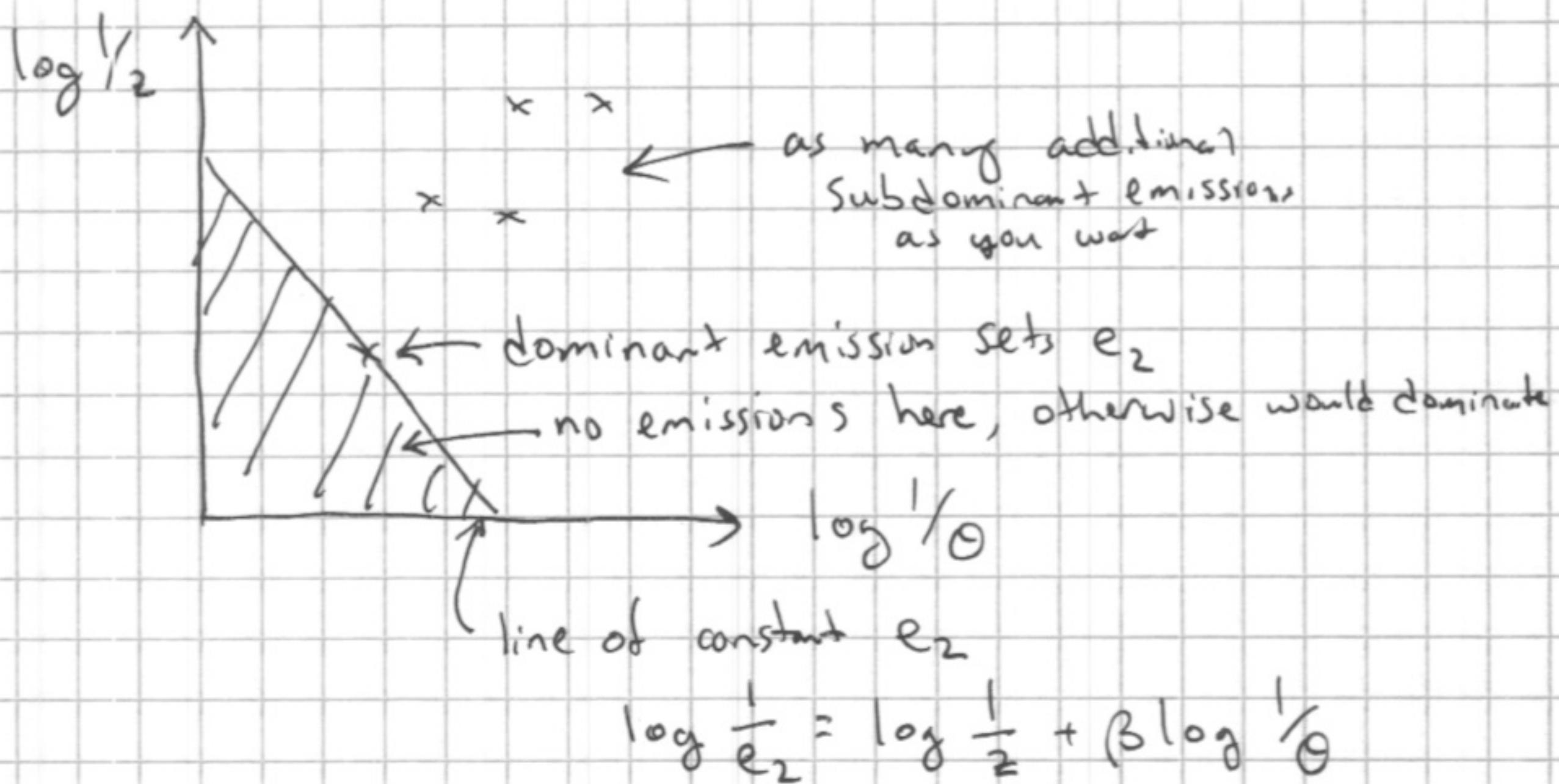
$$\frac{g}{q} \frac{E_2}{\theta_{12}} \frac{E_1}{q}$$

$$e_2 = \frac{E_1 E_2 \theta_{12}^\beta}{(E_1 + E_2)^2}$$

$$\approx z \theta^\beta \quad \text{for } E_2 \ll E_1$$

$$\text{where } z = \frac{E_2}{E_1 + E_2}$$

$$\text{Recall : } d\sigma_{i \rightarrow ig} = \frac{2 ds}{\pi} C_i \frac{dz}{z} \frac{d\theta}{\theta}$$



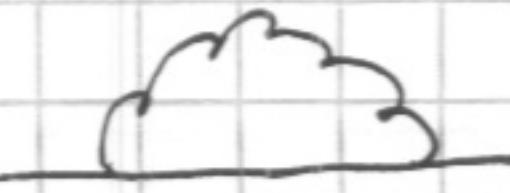
How do you calculate cross section for e_2 ?

In particular, $d\sigma_{i \rightarrow ig}$ is for real emissions.

Where are all the virtual diagrams?

Real : 

at finite z, θ

Virtual: 

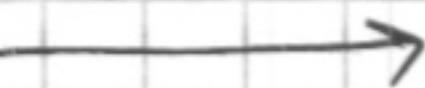
effectively at $(\log \frac{1}{\theta}, \log \frac{1}{z}) \rightarrow (\infty, \infty)$

Easiest to think of in terms of probabilities.

$$P_{\text{emit}} + P_{\text{no-emit}} = 1$$

$\Theta(\alpha_s^0)$

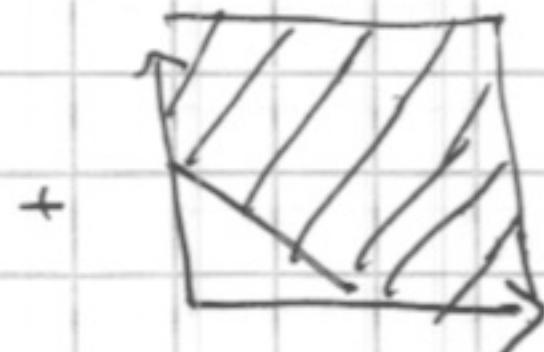
$$P_{\text{no-emit}}^{(0)} = 1$$

just 

$\Theta(\alpha_s')$

$$P_{\text{emit}}^{(1)} + P_{\text{no-emit}}^{(1)} = 0$$

↑ weird!



has to be zero.

Fun exercise in probability: chance to get value of e_2 less than e_2^{\max} is

$$\sum_q (e_2^{\max}) = \int_0^{e_2^{\max}} p(e_2) de_2$$

$$= \exp \left[-\frac{2\alpha_s}{\pi} C_F \left(\text{area of } \begin{array}{c} \text{cone} \\ \diagdown \end{array} \right) \right]$$

$$\left(= 1 - \begin{array}{c} \text{no emission} \\ @ \alpha_s^0 \end{array} + \frac{1}{2} \begin{array}{c} \text{veto emission} \\ @ \alpha_s^1 \end{array} - \dots \right)$$

subtleties
@ α_s^2
from $\Delta = -$ 

I assume you've never done a calculation like that before? Very standard in parton shower picture. Basis for MC generators.

Final result:

$$\Sigma_q(e_2^{\text{max}}) = \exp \left[\frac{-\alpha_s}{\pi} \frac{C_F}{\beta} \log^2 \frac{R^\beta}{e_2^{\text{max}}} \right]$$

↑
For gluons swap
 $C_F \rightarrow C_A$.

Called Sudakov double-logarithmic form factor.

Has some information (double-log enhanced) at all orders in α_s .

Systematically improvable in approaches to resummation like SCET

For simplicity: set $R=1$ (or equivalently, rescale e_2)

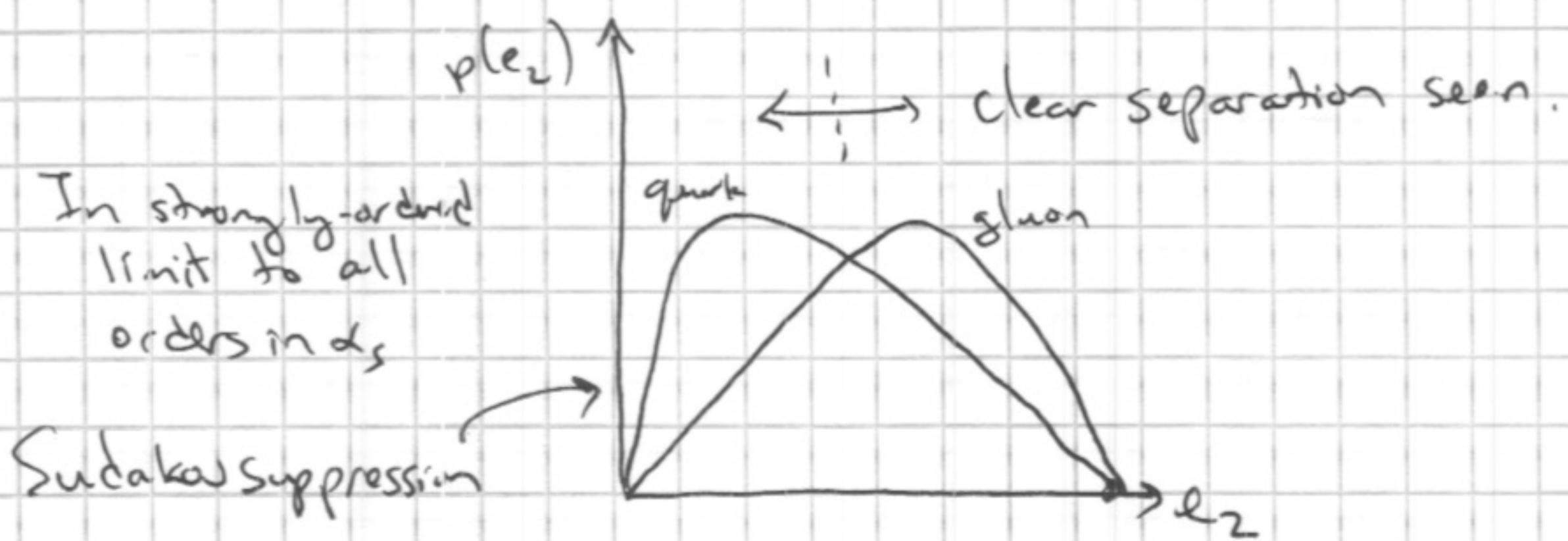
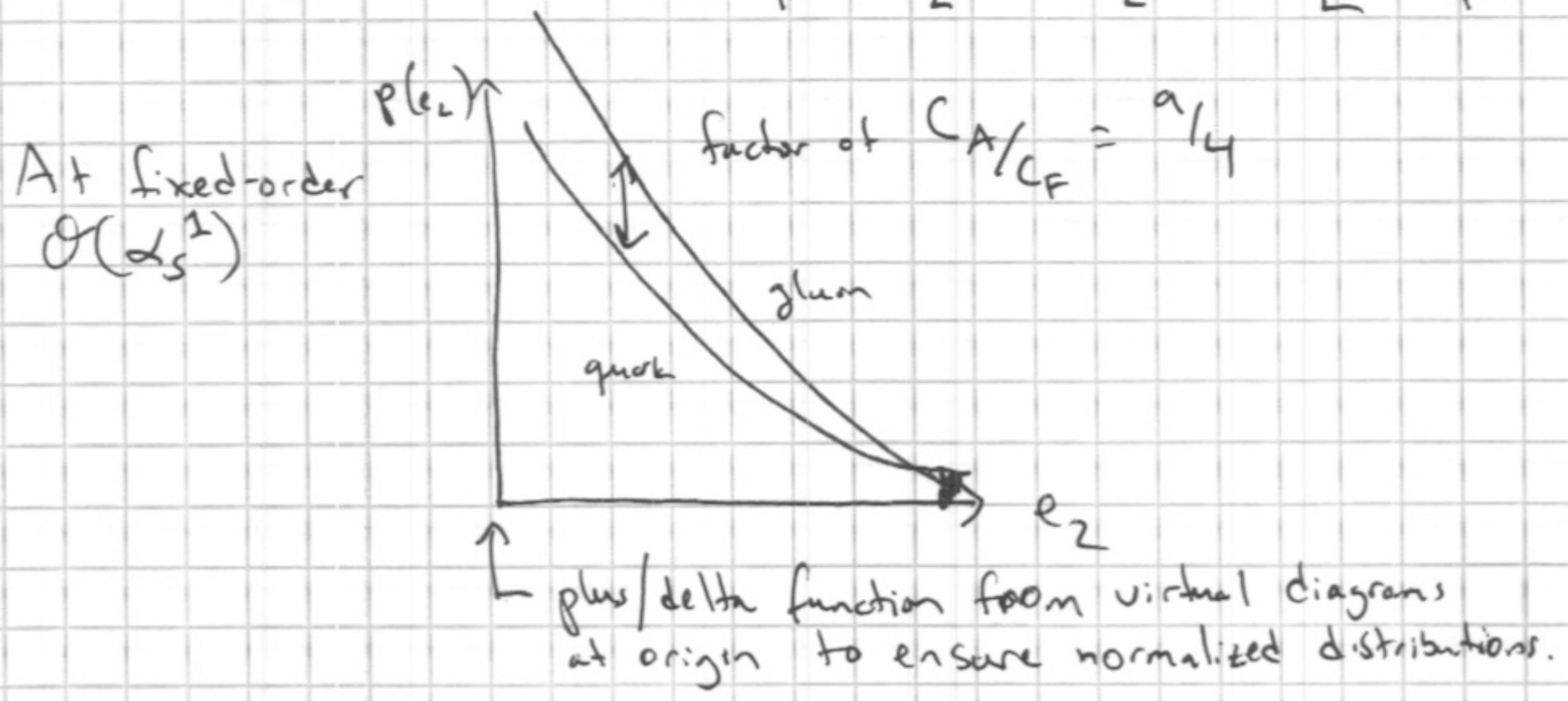
To go from cumulative distribution Σ to probability distribution, just take a derivative!

$$P(e_2) = \frac{1}{\sigma} \frac{d\sigma}{de_2} = \frac{\partial}{\partial e_2} \Sigma(e_2)$$

Time to interpret our result!

1) What is cross section?

$$\rho_{q\bar{q}}(e_2) = \frac{2\alpha_s}{\pi} \frac{C_F}{\beta} \frac{1}{e_2} \log \frac{1}{e_2} \exp \left[-\frac{\alpha_s}{\pi} \frac{C_F}{\beta} \log \frac{1}{e_2} \right]$$



This resummed calculation is much closer to what is seen experimentally.

2) Can we quantify discrimination power?

Place a cut of e_2^{cut}

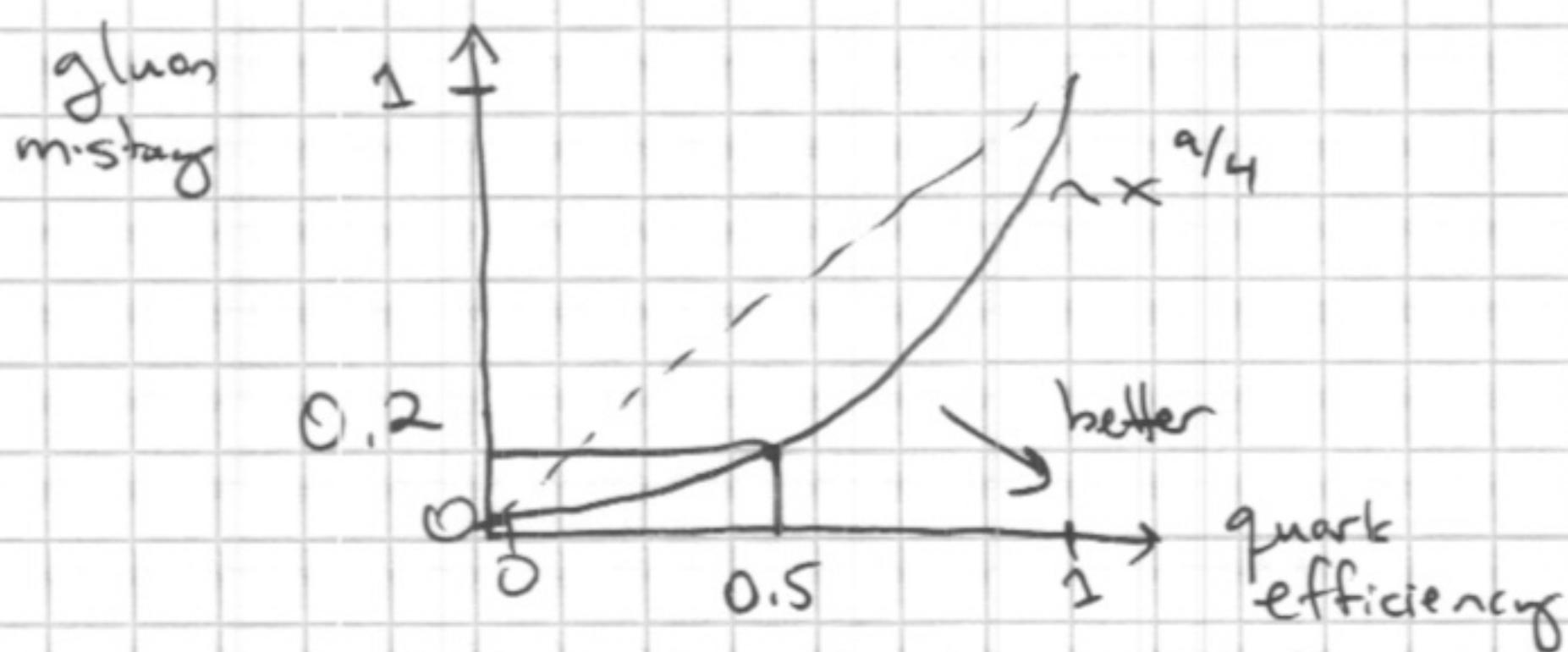
$e_2 < e_2^{\text{cut}} \Rightarrow$ "quark"

$e_2 > e_2^{\text{cut}} \Rightarrow$ "gluon"

Want to know mistag rate ($\sum_g (e_2^{\text{cut}})$) as a function of efficiency ($\sum_q (e_2^{\text{cut}})$).

Key: $\Sigma_g = (\Sigma_q)^{C_A/C_F \sim 9/4}$

Called "Casimir Scaling" (violated at higher orders)



In strongly-ordered limit, independent of β !

All* quark/gluon discriminants alike, because only difference at this order is C_F vs C_A .

Open question: How best to separate quarks from gluons using higher-order information?

Many important higher-order effects

- Multiple emissions
- Color coherence
- Subleading terms in splitting function
- Fixed-order corrections
- Running α_s
- Non-global logarithms
- Hadronization Effects
- ...

Hope you've gained some appreciation for
value of working to all orders in α_s .

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Next time: The Substructure (R) evolution.

Multi-prong Jets & Jet Cleaning

(more complicated
observables for
 $w/z/H/\text{top}$)

(dealing with
challenges of
jet contamination.)

General Strategy for Set Substructure

→ Figure out your goal

"Discriminate Quarks vs. Gluons"

→ Identify Underlying Physics

" C_F vs. C_A "

→ Construct "clean" observable to probe that physics

$$\sim e_2^{(\beta)} = \sum_{i \neq j} z_i z_j \Theta_{ij}^{\beta}$$

→ Use analytics / Monte Carlo generators to estimate distributions and discrimination power.

→ Convince experimentalists to apply your method to data (or if desperate, do yourself)

Ask me about
CMS Open Data