

Chapter 2 Rigid Body Motion

Lecture Notes for A Geometrical Introduction to Robotics and Manipulation

Richard Murray and Zexiang Li and Shankar S. Sastry
CRC Press

Zexiang Li¹ and Yuanqing Wu¹

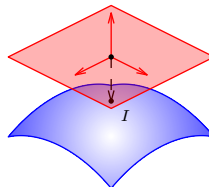
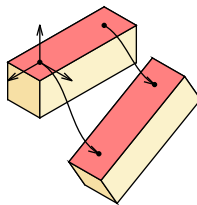
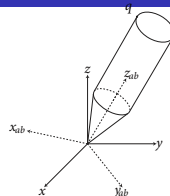
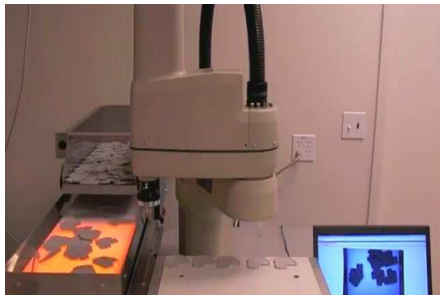
¹ECE, Hong Kong University of Science & Technology

August 30, 2012

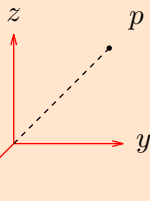
Table of Contents

Chapter 2 Rigid Body Motion

- 1 Rigid Body Transformations
- 2 Rotational motion in \mathbb{R}^3
- 3 Rigid Motion in \mathbb{R}^3
- 4 Velocity of a Rigid Body
- 5 Wrenches and Reciprocal Screws



Notations



$$p = \begin{bmatrix} p_x \\ p_y \\ p_z \end{bmatrix} \text{ or } p = \begin{bmatrix} p_1 \\ p_2 \\ p_3 \end{bmatrix}$$

For $p \in \mathbb{R}^n$, $n = 2, 3$ (2 for planar, 3 for spatial)

$$\text{Point: } p = \begin{bmatrix} p_1 \\ p_2 \\ \vdots \\ p_n \end{bmatrix}, \|p\| = \sqrt{p_1^2 + \cdots + p_n^2}$$

$$\text{Vector: } v = p - q = \begin{bmatrix} p_1 - q_1 \\ p_2 - q_2 \\ \vdots \\ p_n - q_n \end{bmatrix} = \begin{bmatrix} v_1 \\ v_2 \\ \vdots \\ v_n \end{bmatrix}, \|v\| = \sqrt{v_1^2 + \cdots + v_n^2}$$

$$\text{Matrix: } A \in \mathbb{R}^{n \times m}, A = \begin{bmatrix} a_{11} & a_{12} & \cdots & a_{1m} \\ \vdots & & & \vdots \\ a_{n1} & a_{n2} & \cdots & a_{nm} \end{bmatrix}$$

Description of point-mass motion

$$p(0) = \begin{bmatrix} x(0) \\ y(0) \\ z(0) \end{bmatrix} : \text{initial position}$$

$$p(t) = \begin{bmatrix} x(t) \\ y(t) \\ z(t) \end{bmatrix}, t \in (-\varepsilon, \varepsilon)$$

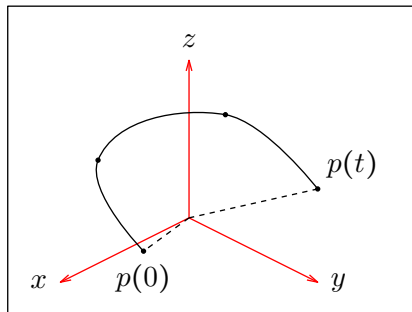


Figure 2.1

Definition: Trajectory

A **trajectory** is a curve $p : (-\varepsilon, \varepsilon) \mapsto \mathbb{R}^3, p(t) = \begin{bmatrix} x(t) \\ y(t) \\ z(t) \end{bmatrix}$

Rigid Body Motion

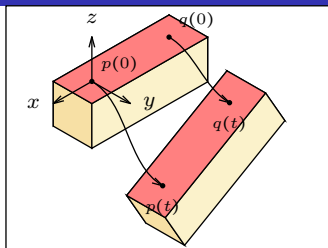


Figure 2.2

$$\|p(t) - q(t)\| = \|p(0) - q(0)\| = \text{constant}$$

Definition: Rigid body transformation

$$g : \mathbb{R}^3 \mapsto \mathbb{R}^3$$

s.t.

- ① Length preserving: $\|g(p) - g(q)\| = \|p - q\|$
- ② Orientation preserving: $g_*(v \times \omega) = g_*(v) \times g_*(\omega)$

† End of Section †

Rotational Motion in \mathbb{R}^3

- 1 Choose a reference frame A (spatial frame)
- 2 Attach a frame B to the body (body frame)

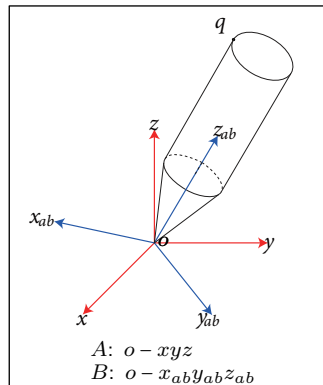


Figure 2.3

$x_{ab} \in \mathbb{R}^3$: coordinates of x_b in frame A
 $R_{ab} = [x_{ab} \ y_{ab} \ z_{ab}] \in \mathbb{R}^{3 \times 3}$: Rotation (or orientation) matrix of B w.r.t. A

Property of a Rotation Matrix

Let $R = [r_1 \ r_2 \ r_3]$ be a rotation matrix

$$\Rightarrow r_i^T \cdot r_j = \begin{cases} 0 & i \neq j \\ 1 & i = j \end{cases}$$

or

$$R^T \cdot R = \begin{bmatrix} r_1^T \\ r_2^T \\ r_3^T \end{bmatrix} [r_1 \ r_2 \ r_3] = I$$

or $R \cdot R^T = I$

We have:

$$\det(R^T R) = \det R^T \cdot \det R = (\det R)^2 = 1, \det R = \pm 1$$

As $\det R = r_1^T (r_2 \times r_3) = 1 \Rightarrow \det R = 1$

Definition:

$$SO(3) = \{R \in \mathbb{R}^{3 \times 3} \mid R^T R = I, \det R = 1\}$$

and

$$SO(n) = \{R \in \mathbb{R}^{n \times n} \mid R^T R = I, \det R = 1\}$$

◇ Review: Group

(G, \cdot) is a group if:

- ❶ $g_1, g_2 \in G \Rightarrow g_1 \cdot g_2 \in G$
- ❷ $\exists! e \in G, \text{ s.t. } g \cdot e = e \cdot g = g, \forall g \in G$
- ❸ $\forall g \in G, \exists! g^{-1} \in G, \text{ s.t. } g \cdot g^{-1} = g^{-1} \cdot g = e$
- ❹ $g_1 \cdot (g_2 \cdot g_3) = (g_1 \cdot g_2) \cdot g_3$

Examples of group

- ❶ $(\mathbb{R}^3, +)$
- ❷ $(\{0, 1\}, + \text{ mod } 2)$
- ❸ (\mathbb{R}, \times) Not a group (Why?)
- ❹ $(\mathbb{R}_* : \mathbb{R} - \{0\}, \times)$
- ❺ $S^1 \triangleq \{z \in \mathbb{C} \mid |z| = 1\}$

Property 1: $SO(3)$ is a group under matrix multiplication.

Proof :

- ❶ If $R_1, R_2 \in SO(3)$, then $R_1 \cdot R_2 \in SO(3)$, because
 - $(R_1 R_2)^T (R_1 R_2) = R_2^T (R_1^T R_1) R_2 = R_2^T R_2 = I$
 - $\det(R_1 \cdot R_2) = \det(R_1) \cdot \det(R_2) = 1$
- ❷ $e = I_{3 \times 3}$
- ❸ $R^T \cdot R = I \Rightarrow R^{-1} = R^T$



Configuration and rigid transformation

- $R_{ab} = [x_{ab} \ y_{ab} \ z_{ab}] \in SO(3)$
Configuration Space

- Let $q_b = \begin{bmatrix} x_b \\ y_b \\ z_b \end{bmatrix} \in \mathbb{R}^3$: coordinates of q in B .

$$q_a = x_{ab} \cdot x_b + y_{ab} \cdot y_b + z_{ab} \cdot z_b$$

$$= [x_{ab} \ y_{ab} \ z_{ab}] \begin{bmatrix} x_b \\ y_b \\ z_b \end{bmatrix} = R_{ab} \cdot q_b$$

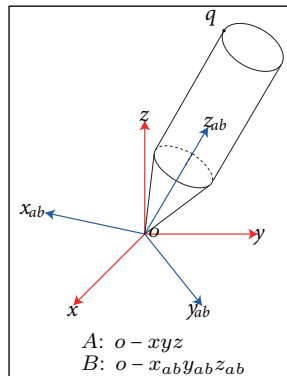


Figure 2.3

- A configuration $R_{ab} \in SO(3)$ is also a transformation:

$$R_{ab} : \mathbb{R}^3 \rightarrow \mathbb{R}^3, R_{ab}(q_b) = R_{ab} \cdot q_b = q_a$$

A config. \Leftrightarrow A transformation in $SO(3)$

Property 2: R_{ab} preserves distance between points and orientation.

$$\textcircled{1} \quad \|R_{ab} \cdot (p_b - q_b)\| = \|p_b - q_b\|$$

$$\textcircled{2} \quad R(v \times \omega) = (Rv) \times R\omega$$

Proof :

For $a \in \mathbb{R}^3$, let $\hat{a} = \begin{bmatrix} 0 & -a_3 & a_2 \\ a_3 & 0 & -a_1 \\ -a_2 & a_1 & 0 \end{bmatrix}$

Note that $\hat{a} \cdot b = a \times b$

$$\begin{aligned} \textcircled{1} \text{ follows from } \|R_{ab}(p_b - q_b)\|^2 &= (R_{ab}(p_b - q_b))^T R_{ab}(p_b - q_b) \\ &= (p_b - q_b)^T R_{ab}^T R_{ab}(p_b - q_b) \\ &= \|p_b - p_q\|^2 \end{aligned}$$

$$\textcircled{2} \text{ follows from } R\hat{v}R^T = (Rv)^\wedge \text{ (prove it yourself)}$$



Parametrization of $SO(3)$ (the exponential coordinate)

◇ **Review:** $S^1 = \{z \in \mathbb{C} \mid |z| = 1\}$

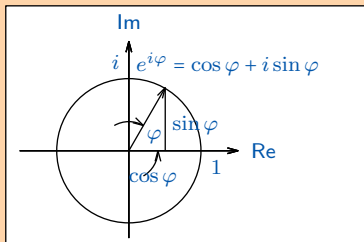


Figure 2.4

Euler's Formula

“One of the most remarkable, almost astounding, formulas in all of mathematics.”

R. Feynman

◇ **Review:**

$$\begin{cases} \dot{x}(t) = ax(t) \\ x(0) = x_0 \end{cases} \Rightarrow x(t) = e^{at}x_0$$

(Continues next slide)

$$R \in SO(3), R = \begin{bmatrix} r_{11} & r_{12} & r_{13} \\ r_{21} & r_{22} & r_{23} \\ r_{31} & r_{32} & r_{33} \end{bmatrix}$$

$$r_i \cdot r_j = \begin{cases} 0, & i \neq j \\ 1, & i = j \end{cases} \leftarrow 6 \text{ constraints}$$

$\Rightarrow 3$ independent parameters!

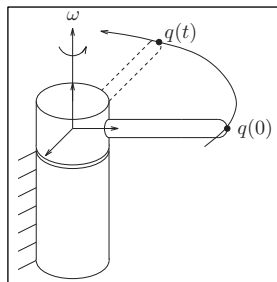


Figure 2.5

$$\begin{cases} \dot{q}(t) = \omega \times q(t) = \hat{\omega} q(t) \\ q(0): \text{Initial coordinates} \end{cases}$$

$$\Rightarrow q(t) = e^{\hat{\omega}t} q_0 \text{ where } e^{\hat{\omega}t} = I + \hat{\omega}t + \frac{(\hat{\omega}t)^2}{2!} + \frac{(\hat{\omega}t)^3}{3!} + \dots$$

By the definition of rigid transformation, $R(\omega, \theta) = e^{\hat{\omega}\theta}$. Let $so(3) = \{\hat{\omega} | \omega \in \mathbb{R}^3\}$ or $so(n) = \{S \in \mathbb{R}^{n \times n} | S^T = -S\}$ where $\wedge : \mathbb{R}^3 \mapsto so(3) : \omega \mapsto \hat{\omega}$, we have:

Property 3: $\exp : so(3) \mapsto SO(3), \hat{\omega}\theta \mapsto e^{\hat{\omega}\theta}$

Rodrigues formula

Rodrigues' formula ($\|\omega\| = 1$):

$$e^{\hat{\omega}\theta} = I + \hat{\omega} \sin \theta + \hat{\omega}^2 (1 - \cos \theta)$$

Proof :

Let $a \in \mathbb{R}^3$, write

$$a = \omega\theta, \omega = \frac{a}{\|a\|} \text{ (or } \|\omega\| = 1), \text{ and } \theta = \|a\|$$

$$e^{\hat{\omega}\theta} = I + \hat{\omega}\theta + \frac{(\hat{\omega}\theta)^2}{2!} + \frac{(\hat{\omega}\theta)^3}{3!} + \dots$$

As

$$\hat{a}^2 = aa^T - \|a\|^2 I, \hat{a}^3 = -\|a\|^2 \hat{a}$$

we have:

$$\begin{aligned} e^{\hat{\omega}\theta} &= I + \left(\theta - \frac{\theta^3}{3!} + \frac{\theta^3}{5!} - \dots\right)\hat{\omega} + \left(\frac{\theta^2}{2!} - \frac{\theta^4}{4!} + \dots\right)\hat{\omega}^2 \\ &= I + \hat{\omega} \sin \theta + \hat{\omega}^2 (1 - \cos \theta) \end{aligned}$$



Rodrigues formula

Rodrigues' formula for $\|\omega\| \neq 1$:

$$e^{\hat{\omega}\theta} = I + \frac{\hat{\omega}}{\|\omega\|} \sin \|\omega\|\theta + \frac{\hat{\omega}^2}{\|\omega\|^2} (1 - \cos \|\omega\|\theta)$$

Proof for Property 3:

Let $R \triangleq e^{\hat{\omega}\theta}$, then:

$$\begin{aligned} (e^{\hat{\omega}\theta})^{-1} &= e^{-\hat{\omega}\theta} = e^{\hat{\omega}^T\theta} = (e^{\hat{\omega}\theta})^T \\ \Rightarrow R^{-1} &= R^T \Rightarrow R^T R = I \Rightarrow \det R = \pm 1 \end{aligned}$$

From $\det \exp(0) = 1$, and the continuity of \det function w.r.t. θ , we have $\det e^{\hat{\omega}\theta} = 1, \forall \theta \in \mathbb{R}$ □

Property 4: The exponential map is onto.

Proof :

Given $R \in SO(3)$, to show $\exists \omega \in \mathbb{R}^3, \|\omega\| = 1$ and θ s.t. $R = e^{\hat{\omega}\theta}$

Let

$$R = \begin{bmatrix} r_{11} & r_{12} & r_{13} \\ r_{21} & r_{22} & r_{23} \\ r_{31} & r_{32} & r_{33} \end{bmatrix}$$

and

$$v_\theta = 1 - \cos \theta, c_\theta = \cos \theta, s_\theta = \sin \theta$$

By Rodrigues' formula

$$e^{\hat{\omega}\theta} = \begin{bmatrix} \omega_1^2 v_\theta + c_\theta & \omega_1 \omega_2 v_\theta - \omega_3 s_\theta & \omega_1 \omega_3 v_\theta + \omega_2 s_\theta \\ \omega_1 \omega_2 v_\theta + \omega_3 s_\theta & \omega_2^2 v_\theta + c_\theta & \omega_2 \omega_3 v_\theta - \omega_1 s_\theta \\ \omega_1 \omega_3 v_\theta - \omega_2 s_\theta & \omega_2 \omega_3 v_\theta + \omega_1 s_\theta & \omega_3^2 v_\theta + c_\theta \end{bmatrix}$$

(continues next slide)

Taking the trace of both sides,

$$\text{tr}(R) = r_{11} + r_{22} + r_{33} = 1 + 2 \cos \theta = \sum_{i=1}^3 \lambda_i$$

where λ_i is the eigenvalue of $R, i = 1, 2, 3$

Case 1: $\text{tr}(R) = 3$ or $R = I, \theta = 0 \Rightarrow \omega\theta = 0$

Case 2: $-1 < \text{tr}(R) < 3,$

$$\theta = \arccos \frac{\text{tr}(R) - 1}{2} \Rightarrow \omega = \frac{1}{2s_\theta} \begin{bmatrix} r_{32} - r_{23} \\ r_{13} - r_{31} \\ r_{21} - r_{12} \end{bmatrix}$$

Case 3: $\text{tr}(R) = -1 \Rightarrow \cos \theta = -1 \Rightarrow \theta = \pm\pi$

(continues next slide)

Following are 3 possibilities:

$$R = \begin{bmatrix} 1 & 0 & 0 \\ 0 & -1 & 0 \\ 0 & 0 & -1 \end{bmatrix} \Rightarrow \omega = \begin{bmatrix} 1 \\ 0 \\ 0 \end{bmatrix},$$

$$R = \begin{bmatrix} -1 & 0 & 0 \\ 0 & 1 & 0 \\ 0 & 0 & -1 \end{bmatrix} \Rightarrow \omega = \begin{bmatrix} 0 \\ 1 \\ 0 \end{bmatrix},$$

$$R = \begin{bmatrix} -1 & 0 & 0 \\ 0 & -1 & 0 \\ 0 & 0 & 1 \end{bmatrix} \Rightarrow \omega = \begin{bmatrix} 0 \\ 0 \\ 1 \end{bmatrix}$$

Note that if $\omega\theta$ is a solution, then $\omega(\theta \pm n\pi), n = 0, \pm 1, \pm 2, \dots$ is also a solution. □

Definition: Exponential coordinate

$\omega\theta \in \mathbb{R}^3$, with $e^{\hat{\omega}\theta} = R$ is called the exponential coordinates of R

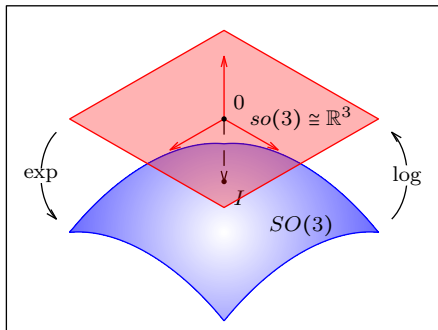


Figure 2.6

Property 5: \exp is 1-1 when restricted to an open ball in \mathbb{R}^3 of radius π .

Euler's rotation theorem

Theorem 1 (Euler):

Any orientation is equivalent to a rotation about a fixed axis $\omega \in \mathbb{R}^3$ through an angle $\theta \in [-\pi, \pi]$.

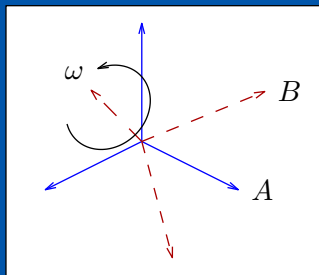


Figure 2.7



1707–1783

$SO(3)$ can be visualized as a solid ball of radius π .

Other Parametrizations of $SO(3)$

- XYZ fixed angles (or Roll-Pitch-Yaw angle)

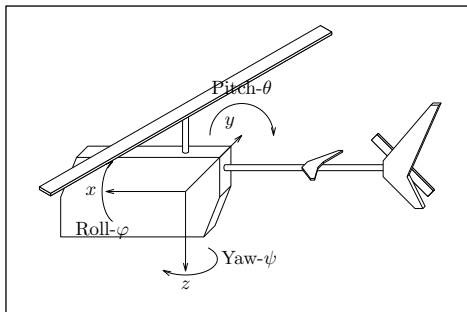


Figure 2.8

(continues next slide)

Other Parametrizations of $SO(3)$

- XYZ fixed angles (or Roll-Pitch-Yaw angle) Continued

$$R_x(\varphi) := e^{\hat{x}\varphi} = \begin{bmatrix} 1 & 0 & 0 \\ 0 & \cos \varphi & -\sin \varphi \\ 0 & \sin \varphi & \cos \varphi \end{bmatrix}$$

$$R_y(\theta) := e^{\hat{y}\theta} = \begin{bmatrix} \cos \theta & 0 & \sin \theta \\ 0 & 1 & 0 \\ -\sin \theta & 0 & \cos \theta \end{bmatrix}$$

$$R_z(\psi) := e^{\hat{z}\psi} = \begin{bmatrix} \cos \psi & -\sin \psi & 0 \\ \sin \psi & \cos \psi & 0 \\ 0 & 0 & 1 \end{bmatrix}$$

$$R_{ab} = R_x(\varphi)R_y(\theta)R_z(\psi)$$

$$= \begin{bmatrix} c_\theta c_\psi & -c_\theta s_\psi & s_\theta \\ s_\varphi s_\theta c_\psi + c_\varphi s_\psi & -s_\varphi s_\theta s_\psi + c_\varphi c_\psi & -s_\varphi c_\theta \\ -c_\varphi s_\theta c_\psi + s_\varphi s_\psi & c_\varphi s_\theta s_\psi + s_\varphi c_\psi & c_\varphi c_\theta \end{bmatrix}$$

Other Parametrizations of $SO(3)$

• ZYX Euler angle

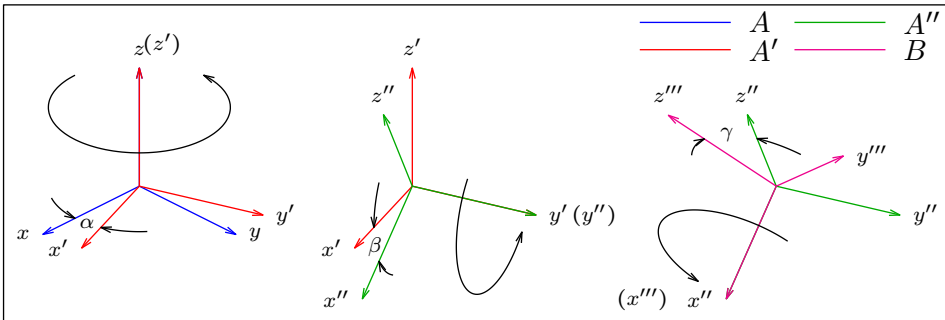


Figure 2.9

$$R_{aa'} = R_z(\alpha)$$

$$R_{a'a''} = R_y(\beta)$$

$$R_{a''b} = R_x(\gamma)$$

$$R_{ab} = R_z(\alpha)R_y(\beta)R_x(\gamma)$$

(continues next slide)

Other Parametrizations of $SO(3)$

• ZYX Euler angle (continued)

$$R_{ab}(\alpha, \beta, \gamma) = \begin{bmatrix} c_\alpha c_\beta & -s_\alpha c_\gamma + c_\alpha s_\beta s_\gamma & s_\alpha c_\gamma + c_\alpha s_\beta c_\gamma \\ s_\alpha c_\beta & c_\alpha c_\gamma + s_\alpha s_\beta s_\gamma & -c_\alpha c_\gamma + s_\alpha s_\beta c_\gamma \\ -s_\beta & c_\beta s_\gamma & c_\beta c_\gamma \end{bmatrix}$$

Note: When $\beta = \frac{\pi}{2}$, $\cos \beta = 0$, $\alpha + \gamma = \text{const} \Rightarrow \text{singularity!}$

$$\beta = \text{atan2}(-r_{31}, \sqrt{r_{32}^2 + r_{33}^2})$$

$$\alpha = \text{atan2}(r_{21}/c_\beta, r_{11}/c_\beta)$$

$$\gamma = \text{atan2}(r_{32}/c_\beta, r_{33}/c_\beta)$$

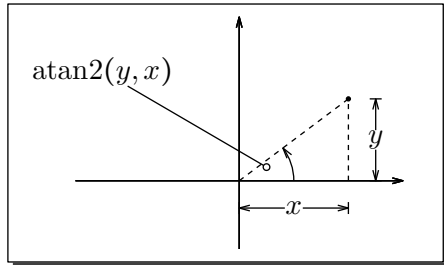


Figure 2.10

Other Parametrizations of $SO(3)$

§ Quaternions:

$$Q = q_0 + q_1 i + q_2 j + q_3 k$$

$$\text{where } i^2 = j^2 = k^2 = -1, i \cdot j = k, j \cdot k = i, k \cdot i = j$$

Property 1: Define $Q^* = (q_0, q)^* = (q_0, -q)$, $q_0 \in \mathbb{R}, q \in \mathbb{R}^3$

$$\|Q\|^2 = QQ^* = q_0^2 + q_1^2 + q_2^2 + q_3^2$$

Property 2: $Q = (q_0, q), P = (p_0, p)$

$$QP = (q_0 p_0 - q \cdot p, q_0 p + p_0 q + q \times p)$$

Property 3: (a) The set of unit quaternions forms a group

(b) If $R = e^{\hat{\omega}\theta}$, then $Q = \left(\cos \frac{\theta}{2}, \omega \sin \frac{\theta}{2}\right)$

(c) Q acts on $x \in \mathbb{R}^3$ by QXQ^* , where $X = (0, x)$

Other Parametrizations of $SO(3)$

□ Unit Quaternions:

Given $Q = (q_0, q)$, $q_0 \in \mathbb{R}$, $q \in \mathbb{R}^3$, the vector part of QXQ^* is given by $R(Q)x$, recall that

$$q_0 = \cos \frac{\theta}{2}, q = \omega \sin \frac{\theta}{2}$$

and the Rodrigues' formula:

$$e^{\hat{\omega}\theta} = I + \hat{\omega} \sin \theta + \hat{\omega}^2 (1 - \cos \theta)$$

then

$$\begin{aligned} R(Q) &= I + 2q_0\hat{q} + 2\hat{q}^2 \\ &= \begin{bmatrix} 1 - 2(q_2^2 + q_3^2) & -2q_0q_3 + 2q_1q_2 & 2q_0q_2 + 2q_1q_3 \\ 2q_0q_3 + 2q_1q_2 & 1 - 2(q_1^2 + q_3^2) & -2q_0q_1 + 2q_2q_3 \\ -2q_0q_2 + 2q_1q_3 & 2q_0q_1 + 2q_2q_3 & 1 - 2(q_1^2 + q_2^2) \end{bmatrix} \end{aligned}$$

where $\|Q\| \triangleq q_0^2 + q_1^2 + q_2^2 + q_3^2 = 1$

(continues next slide)

Other Parametrizations of $SO(3)$

□ Quaternions (continued):

Conversion from Roll-Pitch-Yaw angle to unit quaternions:

$$Q = \left(\cos \frac{\varphi}{2}, x \sin \frac{\varphi}{2}\right) \left(\cos \frac{\theta}{2}, y \sin \frac{\theta}{2}\right) \left(\cos \frac{\psi}{2}, z \sin \frac{\psi}{2}\right) \Rightarrow$$

$$q_0 = \cos \frac{\varphi}{2} \cos \frac{\theta}{2} \cos \frac{\psi}{2} - \sin \frac{\varphi}{2} \sin \frac{\theta}{2} \sin \frac{\psi}{2}$$

$$q = \begin{bmatrix} \cos \frac{\varphi}{2} \sin \frac{\theta}{2} \sin \frac{\psi}{2} + \sin \frac{\varphi}{2} \cos \frac{\theta}{2} \cos \frac{\psi}{2} \\ \cos \frac{\varphi}{2} \sin \frac{\theta}{2} \cos \frac{\psi}{2} - \sin \frac{\varphi}{2} \cos \frac{\theta}{2} \sin \frac{\psi}{2} \\ \cos \frac{\varphi}{2} \cos \frac{\theta}{2} \sin \frac{\psi}{2} + \sin \frac{\varphi}{2} \sin \frac{\theta}{2} \cos \frac{\psi}{2} \end{bmatrix}$$

Conversion from unit quaternions to roll-pitch-yaw angles (?)

† End of Section †

Rigid motion in \mathbb{R}^3

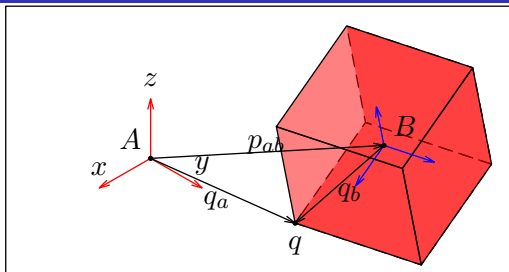


Figure 2.11

$p_{ab} \in \mathbb{R}^3$: Coordinates of the origin of B
 $R_{ab} \in SO(3)$: Orientation of B relative to A
 $SE(3) : \left\{ \begin{bmatrix} R & p \\ 0 & 1 \end{bmatrix} \middle| p \in \mathbb{R}^3, R \in SO(3) \right\}$: Orientation of B relative to A
 Or...as a transformation:

$$g_{ab} = (p_{ab}, R_{ab}) : \mathbb{R}^3 \mapsto \mathbb{R}^3$$

$$q_b \mapsto q_a = p_{ab} + R_{ab} \cdot q_b$$

Homogeneous Representation

Points:

$$q = \begin{bmatrix} q_1 \\ q_2 \\ q_3 \end{bmatrix} \in \mathbb{R}^3$$



$$\bar{q} = \begin{bmatrix} q_1 \\ q_2 \\ q_3 \\ 1 \end{bmatrix} \in \mathbb{R}^4$$

Vectors:

$$v = p - q = \begin{bmatrix} p_1 - q_1 \\ p_2 - q_2 \\ p_3 - q_3 \end{bmatrix} = \begin{bmatrix} v_1 \\ v_2 \\ v_3 \end{bmatrix}$$



$$\bar{v} = \bar{p} - \bar{q} = \begin{bmatrix} p_1 \\ p_2 \\ p_3 \\ 1 \end{bmatrix} - \begin{bmatrix} q_1 \\ q_2 \\ q_3 \\ 1 \end{bmatrix} = \begin{bmatrix} v_1 \\ v_2 \\ v_3 \\ 0 \end{bmatrix}$$

- ① Point-Point = Vector
- ② Vector+Point = Point
- ③ Vector+Vector = Vector
- ④ Point+Point: Meaningless

(continues next slide)

Homogeneous Representation

$$q_a = p_{ab} + R_{ab} \cdot q_b$$

$$\begin{bmatrix} q_a \\ 1 \end{bmatrix} = \underbrace{\begin{bmatrix} R_{ab} & p_{ab} \\ 0 & 1 \end{bmatrix}}_{\bar{g}_{ab}} \begin{bmatrix} q_b \\ 1 \end{bmatrix}$$

$$\bar{q}_a = \bar{g}_{ab} \cdot \bar{q}_b$$

□ **Composition Rule:**

$$\bar{q}_b = \bar{g}_{bc} \cdot \bar{q}_c$$

$$\bar{q}_a = \bar{g}_{ab} \cdot \bar{q}_b = \underbrace{\bar{g}_{ab} \cdot \bar{g}_{bc}}_{\bar{g}_{ac}} \cdot \bar{q}_c$$

$$\bar{g}_{ac} = \bar{g}_{ab} \cdot \bar{g}_{bc} = \begin{bmatrix} R_{ab}R_{bc} & R_{ab}p_{bc} + p_{ab} \\ 0 & 1 \end{bmatrix}$$

$$g_{ab} = (p_{ab}, R_{ab})$$



$$\bar{g}_{ab} = \begin{bmatrix} R_{ab} & p_{ab} \\ 0 & 1 \end{bmatrix}$$

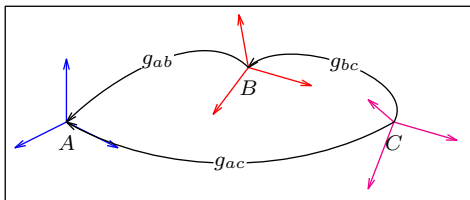


Figure 2.12

Special Euclidean Group

$$SE(3) = \left\{ \begin{bmatrix} R & p \\ 0 & 1 \end{bmatrix} \in \mathbb{R}^{4 \times 4} \mid p \in \mathbb{R}^3, R \in SO(3) \right\}$$

Property 4: $SE(3)$ forms a group.

Proof :

- ① $g_1 \cdot g_2 \in SE(3)$
- ② $e = I_4$
- ③ $(\bar{g})^{-1} = \begin{bmatrix} R^T & -R^T p \\ 0 & 1 \end{bmatrix}$
- ④ Associativity: Follows from property of matrix multiplication □

$$\bar{v} = s - r = \begin{bmatrix} v_1 \\ v_2 \\ v_3 \\ 0 \end{bmatrix}, \bar{g}_* \bar{v} = \bar{g}s - \bar{g}r = \begin{bmatrix} R & p \\ 0 & 1 \end{bmatrix} \begin{bmatrix} v_1 \\ v_2 \\ v_3 \\ 0 \end{bmatrix} = \begin{bmatrix} Rv \\ 0 \end{bmatrix}$$

The bar will be dropped to simplify notations

Property 5: An element of $SE(3)$ is a rigid transformation.

Exponential coordinates of $SE(3)$

For rotational motion:

$$\dot{p}(t) = \omega \times (p(t) - q)$$

$$\begin{bmatrix} \dot{p} \\ 0 \end{bmatrix} = \begin{bmatrix} \hat{\omega} & -\omega \times q \\ 0 & 0 \end{bmatrix} \begin{bmatrix} p \\ 1 \end{bmatrix}$$

$$\text{or } \dot{\bar{p}} = \hat{\xi} \cdot \bar{p} \Rightarrow \bar{p}(t) = e^{\hat{\xi}t} \bar{p}(0)$$

$$\text{where } e^{\hat{\xi}t} = I + \hat{\xi}t + \frac{(\hat{\xi}t)^2}{2!} + \dots$$

For translational motion:

$$\dot{p}(t) = v$$

$$\begin{bmatrix} \dot{p}(t) \\ 0 \end{bmatrix} = \begin{bmatrix} 0 & v \\ 0 & 0 \end{bmatrix} \begin{bmatrix} p \\ 1 \end{bmatrix}$$

$$\dot{\bar{p}}(t) = \hat{\xi} \cdot \bar{p}(t) \Rightarrow \bar{p}(t) = e^{\hat{\xi}t} \bar{p}(0)$$

$$\hat{\xi} = \begin{bmatrix} 0 & v \\ 0 & 0 \end{bmatrix}$$

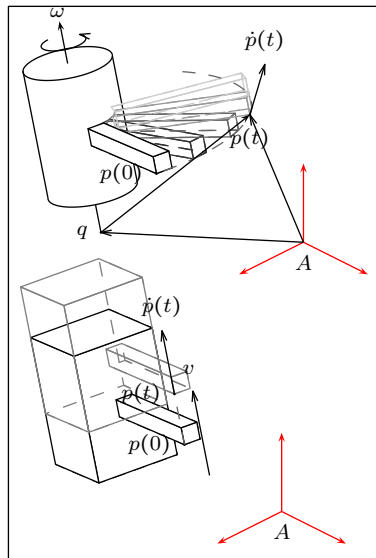


Figure 2.13

Exponential coordinates of $SE(3)$

Definition:

$$se(3) = \left\{ \begin{bmatrix} \hat{\omega} & v \\ 0 & 0 \end{bmatrix} \in \mathbb{R}^{4 \times 4} \mid v, \omega \in \mathbb{R}^3 \right\}$$

is called the twist space. There exists a 1-1 correspondence between $se(3)$ and \mathbb{R}^6 , defined by $\wedge : \mathbb{R}^6 \mapsto se(3)$

$$\xi := \begin{bmatrix} v \\ \omega \end{bmatrix} \mapsto \hat{\xi} = \begin{bmatrix} \hat{\omega} & v \\ 0 & 0 \end{bmatrix}$$

Property 6: $\exp : se(3) \mapsto SE(3), \hat{\xi}\theta \mapsto e^{\hat{\xi}\theta}$

Proof :

$$\text{Let } \hat{\xi} = \begin{bmatrix} \hat{\omega} & v \\ 0 & 0 \end{bmatrix}$$

- If $\omega = 0$, then $\hat{\xi}^2 = \hat{\xi}^3 = \dots = 0$, $e^{\hat{\xi}\theta} = \begin{bmatrix} I & v\theta \\ 0 & 1 \end{bmatrix} \in SE(3)$

(continues next slide)

Exponential coordinates of $SE(3)$

- If ω is not 0, assume $\|\omega\| = 1$.

Define:

$$g_0 = \begin{bmatrix} I & \omega \times v \\ 0 & 1 \end{bmatrix}, \hat{\xi}' = g_0^{-1} \cdot \hat{\xi} \cdot g_0 = \begin{bmatrix} \hat{\omega} & h\omega \\ 0 & 0 \end{bmatrix}$$

where $h = \omega^T \cdot v$.

$$e^{\hat{\xi}\theta} = e^{g_0 \cdot \hat{\xi}' \cdot g_0^{-1}} = g_0 \cdot e^{\hat{\xi}'\theta} \cdot g_0^{-1}$$

and as

$$\hat{\xi}'^2 = \begin{bmatrix} \hat{\omega}^2 & 0 \\ 0 & 0 \end{bmatrix}, \hat{\xi}'^3 = \begin{bmatrix} \hat{\omega}^3 & 0 \\ 0 & 0 \end{bmatrix}$$

we have

$$e^{\hat{\xi}'\theta} = \begin{bmatrix} e^{\hat{\omega}\theta} & h\omega\theta \\ 0 & 1 \end{bmatrix} \Rightarrow e^{\hat{\xi}\theta} = \begin{bmatrix} e^{\hat{\omega}\theta} & (I - e^{\hat{\omega}\theta})\hat{\omega}v + \omega\omega^T v\theta \\ 0 & 1 \end{bmatrix}$$



Exponential coordinates of $SE(3)$

$$p(\theta) = e^{\hat{\xi}\theta} \cdot p(0) \Rightarrow g_{ab}(\theta) = e^{\hat{\xi}\theta}$$

If there is offset,

$$g_{ab}(\theta) = e^{\hat{\xi}\theta} g_{ab}(0) \text{ (Why?)}$$

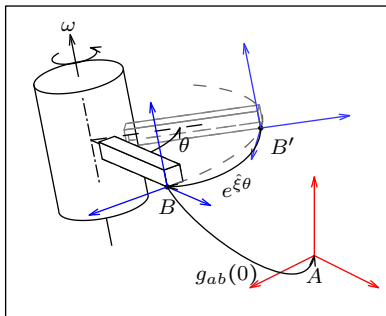


Figure 2.14

Exponential coordinates of $SE(3)$

Property 7: $\exp : se(3) \mapsto SE(3)$ is onto.

Proof :

Let $g = (p, R), R \in SO(3), p \in \mathbb{R}^3$

Case 1: ($R = I$) Let

$$\hat{\xi} = \begin{bmatrix} 0 & \frac{p}{\|p\|} \\ 0 & 0 \end{bmatrix}, \theta = \|p\| \Rightarrow e^{\hat{\xi}\theta} = g = \begin{bmatrix} I & p \\ 0 & 1 \end{bmatrix}$$

Case 2: ($R \neq I$)

$$e^{\hat{\xi}\theta} = \begin{bmatrix} e^{\hat{\omega}\theta} & (I - e^{\hat{\omega}\theta})(\omega \times v)_1 + \omega\omega^T v\theta \\ 0 & 1 \end{bmatrix} = \begin{bmatrix} R & p \\ 0 & 1 \end{bmatrix}$$

$$\Rightarrow \begin{cases} e^{\hat{\omega}\theta} = R \\ (I - e^{\hat{\omega}\theta})(\omega \times v) + \omega\omega^T v\theta = p \end{cases}$$

Solve for $\omega\theta$ from previous section. Let $A = (I - e^{\hat{\omega}\theta})\hat{\omega} + \omega\omega^T\theta$, $Av = p$.

Claim:

$$A = (I - e^{\hat{\omega}\theta})\hat{\omega} + \omega\omega^T\theta := A_1 + A_2$$

$$\ker A_1 \cap \ker A_2 = \phi \Rightarrow v = A^{-1}p$$



$\xi\theta \in \mathbb{R}^6$: Exponential coordinates of $g \in SE(3)$

Screws, twists and screw motion

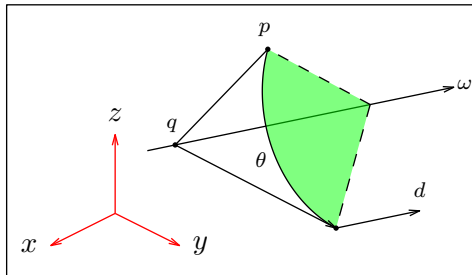


Figure 2.15

Screw attributes

Pitch: $h = \frac{d}{\theta} (\theta = 0, h = \infty), d = h \cdot \theta$
 Axis: $l = \{q + \lambda\omega \mid \lambda \in \mathbb{R}\}$
 Magnitude: $M = \theta$

Definition:

A **screw** S consists of an axis l , pitch h , and magnitude M . A **screw motion** is a rotation by $\theta = M$ about l , followed by translation by $h\theta$, parallel to l . If $h = \infty$, then, translation about v by $\theta = M$

Screws, twists and screw motion

Corresponding $g \in SE(3)$:

$$g \cdot p = q + e^{\hat{\omega}\theta}(p - q) + h\theta\omega$$

$$g \cdot \begin{bmatrix} p \\ 1 \end{bmatrix} = \begin{bmatrix} e^{\hat{\omega}\theta} & (I - e^{\hat{\omega}\theta})q + h\theta\omega \\ 0 & 1 \end{bmatrix} \begin{bmatrix} p \\ 1 \end{bmatrix} \Rightarrow$$

$$g = \begin{bmatrix} e^{\hat{\omega}\theta} & (I - e^{\hat{\omega}\theta})q + h\theta\omega \\ 0 & 1 \end{bmatrix}$$

On the other hand...

$$e^{\hat{\xi}\theta} = \begin{bmatrix} e^{\hat{\omega}\theta} & (I - e^{\hat{\omega}\theta})\omega \times v + \omega\omega^T v\theta \\ 0 & 1 \end{bmatrix}$$

If we let $v = -\omega \times q + h\omega$, then

$$(I - e^{\hat{\omega}\theta})(-\hat{\omega}^2 q) = (I - e^{\hat{\omega}\theta})(-\omega\omega^T q + q) = (I - e^{\hat{\omega}\theta})q$$

Thus, $e^{\hat{\xi}\theta} = g$

For pure rotation ($h = 0$): $\xi = (-\omega \times q, \omega)$

For pure translation: $g = \begin{bmatrix} I & v\theta \\ 0 & 1 \end{bmatrix}$, $\Rightarrow \xi = (v, 0)$, and $e^{\hat{\xi}\theta} = g$

Screw associated with a twist

$$\xi = (v, \omega) \in \mathbb{R}^6$$

- ① Pitch:
$$h = \begin{cases} \frac{\omega^T v}{\|\omega\|^2}, & \text{if } \omega \neq 0 \\ \infty, & \text{if } \omega = 0 \end{cases}$$
- ② Axis:
$$l = \begin{cases} \frac{\omega \times v}{\|\omega\|^2} + \lambda \omega, & \lambda \in \mathbb{R}, \text{ if } \omega \neq 0 \\ 0 + \lambda v & \lambda \in \mathbb{R}, \text{ if } \omega = 0 \end{cases}$$
- ③ Magnitude:
$$M = \begin{cases} \|\omega\|, & \text{if } \omega \neq 0 \\ \|v\|, & \text{if } \omega = 0 \end{cases}$$

Special cases:

- ① $h = \infty$, Pure translation (prismatic joint)
- ② $h = 0$, Pure rotation (revolute joint)

Screw associated with a twist

Screw	Twist: $\hat{\xi}\theta$
Case 1: Pitch: $h = \infty$ Axis: $l = \{q + \lambda v \mid \ v\ = 1, \lambda \in \mathbb{R}\}$ Magnitude: M	$\theta = M,$ $\hat{\xi} = \begin{bmatrix} 0 & v \\ 0 & 0 \end{bmatrix}$
Case 2: Pitch: $h \neq \infty$ Axis: $l = \{q + \lambda \omega \mid \ \omega\ = 1, \lambda \in \mathbb{R}\}$ Magnitude: M	$\theta = M,$ $\hat{\xi} = \begin{bmatrix} \hat{\omega} & -\hat{\omega}q + h\omega \\ 0 & 0 \end{bmatrix}$

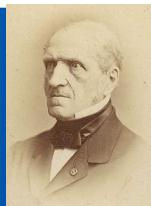
Definition: Screw Motion

Rotation about an axis by $\theta = M$, followed by translation about the same axis by $h\theta$

Chasles Theorem

Theorem 2 (Chasles):

Every rigid body motion can be realized by a rotation about an axis combined with a translation parallel to that axis.



1793–1880

Proof :

For $\hat{\xi} \in se(3)$:

$$\hat{\xi} = \hat{\xi}_1 + \hat{\xi}_2 = \begin{bmatrix} \hat{\omega} & -\omega \times q \\ 0 & 0 \end{bmatrix} + \begin{bmatrix} 0 & h\omega \\ 0 & 0 \end{bmatrix}$$

$$[\hat{\xi}_1, \hat{\xi}_2] = 0 \Rightarrow e^{\hat{\xi}\theta} = e^{\hat{\xi}_1\theta} e^{\hat{\xi}_2\theta}$$



† End of Section †

Velocity of a Rigid Body

Denote spatial angular velocity by:

$$\hat{\omega}_{ab}^s = \dot{R}_{ab} R_{ab}^T, \omega_{ab} \in \mathbb{R}^3$$

Then

$$V^a = \hat{\omega}_{ab}^s \cdot q_a = \omega_{ab}^s \times q_a$$

Body angular velocity:

$$\hat{\omega}_{ab}^b = R_{ab}^T \cdot \dot{R}_{ab}, v^b \triangleq R_{ab}^T \cdot v^a = \omega_{ab}^b \times q_b$$

Relation between body and spatial angular velocity:

$$\omega_{ab}^b = R_{ab}^T \cdot \omega_{ab}^s \text{ or } \hat{\omega}_{ab}^b = R_{ab}^T \hat{\omega}_{ab}^s R_{ab}$$

Velocity of a Rigid Body

□ Generalized Velocity:

$$g_{ab} = \begin{bmatrix} R_{ab}(t) & p_{ab}(t) \\ 0 & 1 \end{bmatrix}, q_a(t) = g_{ab}(t)q_b$$

$$\frac{d}{dt}q_a(t) = \dot{g}_{ab}(t)q_b = \dot{g}_{ab} \cdot g_{ab}^{-1} \cdot g_{ab} \cdot q_b = \hat{V}_{ab}^s \cdot q_a$$

$$\begin{aligned} \hat{V}_{ab}^s &= \dot{g}_{ab} \cdot g_{ab}^{-1} = \begin{bmatrix} \dot{R}_{ab} & \dot{p}_{ab} \\ 0 & 0 \end{bmatrix} \begin{bmatrix} R_{ab}^T & -R_{ab}^T p_{ab} \\ 0 & 1 \end{bmatrix} \\ &= \begin{bmatrix} \dot{R}_{ab} R_{ab}^T & -\dot{R}_{ab} R_{ab}^T p_{ab} + \dot{p}_{ab} \\ 0 & 0 \end{bmatrix} \\ &= \begin{bmatrix} \hat{\omega}_{ab}^s & -\omega_{ab}^s \times p_{ab} + \dot{p}_{ab} \\ 0 & 0 \end{bmatrix} \triangleq \begin{bmatrix} \hat{\omega}_{ab}^s & v_{ab}^s \\ 0 & 0 \end{bmatrix} \end{aligned}$$

2.4 Velocity of a Rigid Body

□ (Generalized) Spatial Velocity:

$$V_{ab}^s = \begin{bmatrix} v_{ab}^s \\ \omega_{ab}^s \end{bmatrix} = \begin{bmatrix} -\omega_{ab}^s \times p_{ab} + \dot{p}_{ab} \\ (R_{ab} R_{ab}^T)^\vee \end{bmatrix}$$

$$v_{q_a} = \omega_{ab}^s \times q_a + v_{ab}^s$$

Note: $v_{q_b} = g_{ab}^{-1} \cdot v_{q_a} = g_{ab}^{-1} \cdot \dot{g}_{ab} \cdot q_b = \hat{V}_{ab}^b \cdot q_b$

□ (Generalized) Body Velocity:

$$\hat{V}_{ab}^b = g_{ab}^{-1} \dot{g}_{ab} = \begin{bmatrix} R_{ab}^T \dot{R}_{ab} & R_{ab}^T \dot{p}_{ab} \\ 0 & 0 \end{bmatrix} \triangleq \begin{bmatrix} \hat{\omega}_{ab}^b & v_{ab}^b \\ 0 & 0 \end{bmatrix}$$

$$V_{ab}^b = \begin{bmatrix} v_{ab}^b \\ \omega_{ab}^b \end{bmatrix} = \begin{bmatrix} R_{ab}^T \dot{p}_{ab} \\ (R_{ab}^T \dot{R}_{ab})^\vee \end{bmatrix}$$

Relation between body and spatial velocity

$$\begin{aligned}
 \hat{V}_{ab}^s &= \dot{g}_{ab} \cdot g_{ab}^{-1} = g_{ab} \cdot g_{ab}^{-1} \cdot \dot{g}_{ab} \cdot g_{ab}^{-1} = g_{ab} \cdot \hat{V}_{ab}^b \cdot g_{ab}^{-1} \\
 &= \begin{bmatrix} R_{ab} & p_{ab} \\ 0 & 1 \end{bmatrix} \begin{bmatrix} \hat{\omega}_{ab}^b & v_{ab}^b \\ 0 & 0 \end{bmatrix} \begin{bmatrix} R_{ab}^T & -R_{ab}^T p_{ab} \\ 0 & 1 \end{bmatrix} \\
 &= \begin{bmatrix} R_{ab} & p_{ab} \\ 0 & 1 \end{bmatrix} \begin{bmatrix} \hat{\omega}_{ab}^b R_{ab}^T & -\hat{\omega}_{ab}^b R_{ab}^T p_{ab} + v_{ab}^b \\ 0 & 0 \end{bmatrix} \\
 &= \begin{bmatrix} R_{ab} \hat{\omega}_{ab}^b R_{ab}^T & -R_{ab} \hat{\omega}_{ab}^b R_{ab}^T p_{ab} + R_{ab} v_{ab}^b \\ 0 & 0 \end{bmatrix} \\
 V_{ab}^s &= \begin{bmatrix} v_{ab}^s \\ \omega_{ab}^s \end{bmatrix} = \underbrace{\begin{bmatrix} R_{ab} & \hat{p}_{ab} R_{ab} \\ 0 & R_{ab} \end{bmatrix}}_{\text{Ad}_g} V_{ab}^b
 \end{aligned}$$

$$\text{Ad}_g = \begin{bmatrix} R & \hat{p}R \\ 0 & R \end{bmatrix} \in \mathbb{R}^{6 \times 6}, \text{ for } g = (p, R)$$

Properties of Adjoint mapping

$$\begin{aligned}
 g^{-1} &= \begin{bmatrix} R^T & -R^T p \\ 0 & 1 \end{bmatrix} \Rightarrow \\
 \text{Ad}_{g^{-1}} &= \begin{bmatrix} R^T & (-R^T p)^\wedge R^T \\ 0 & R^T \end{bmatrix} \\
 &= \begin{bmatrix} R^T & -R^T \hat{p} \\ 0 & R^T \end{bmatrix} = (\text{Ad}_g)^{-1}
 \end{aligned}$$

$$\text{and } \text{Ad}_{g_1 \cdot g_2} = \text{Ad}_{g_1} \cdot \text{Ad}_{g_2}$$

The map $\text{Ad} : SE(3) \mapsto GL(\mathbb{R}^6)$, $\text{Ad}(g) = \text{Ad}_g$ is a group homomorphism

Matrix Rep	Vector Rep
$\hat{\xi} \in se(3)$	$\xi \in \mathbb{R}^6$
$g \cdot \hat{\xi} \cdot g^{-1} \in se(3)$	$\text{Ad}_g \xi \in \mathbb{R}^6$

Velocity of Screw Motion

$$g_{ab}(\theta) = e^{\hat{\xi}\theta(t)} g_{ab}(0), \frac{d}{dt} e^{\hat{\xi}\theta(t)} = \hat{\xi}\dot{\theta}(t) e^{\hat{\xi}\theta(t)} = \dot{\theta}(t) e^{\hat{\xi}\theta(t)} \hat{\xi}$$

$$\begin{aligned} \hat{V}_{ab}^s &= \dot{g}_{ab} \cdot g_{ab}^{-1} = (\hat{\xi}\dot{\theta} e^{\hat{\xi}\theta(t)} g_{ab}(0)) \cdot (g_{ab}^{-1}(0) e^{-\hat{\xi}\theta(t)}) \\ &= \hat{\xi}\dot{\theta} \Rightarrow V_{ab}^s = \xi\dot{\theta} \end{aligned}$$

$$\begin{aligned} \hat{V}_{ab}^b &= g_{ab}^{-1} \cdot \dot{g}_{ab} = g_{ab}^{-1}(0) e^{-\hat{\xi}\theta} \cdot e^{\hat{\xi}\theta} \hat{\xi}\dot{\theta} g_{ab}(0) \\ &= g_{ab}^{-1}(0) \hat{\xi}\dot{\theta} g_{ab}(0) = (\text{Ad}_{g_{ab}^{-1}(0)} \xi)^\wedge \dot{\theta} \Rightarrow V_{ab}^b = \text{Ad}_{g_{ab}^{-1}(0)} \xi \dot{\theta} \end{aligned}$$

Metric Property of $se(3)$

Let $g_i(t) \in SE(3)$, $i = 1, 2$, be representations of the same motion, obtained using coordinate frame A and B. Then,

$$g_2(t) = g_0 \cdot g_1(t) \cdot g_0^{-1} \Rightarrow V_2^s = \text{Ad}_{g_0} \cdot V_1^s$$

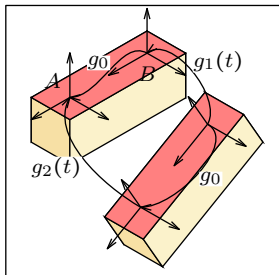


Figure 2.2

(Continues next slide)

Metric Property of $se(3)$

$$\|V_2^s\|^2 = (\text{Ad}_{g_0} \cdot V_1^s)^T (\text{Ad}_{g_0} \cdot V_1^s) = (V_1^s)^T \text{Ad}_{g_0}^T \cdot \text{Ad}_{g_0} \cdot V_1^s$$

$$\begin{aligned} \text{Ad}_{g_0}^T \cdot \text{Ad}_{g_0} &= \begin{bmatrix} R_0^T & 0 \\ -R_0^T \hat{p}_0 & R_0^T \end{bmatrix} \begin{bmatrix} R_0 & \hat{p}_0 R_0 \\ 0 & R_0 \end{bmatrix} \\ &= \begin{bmatrix} I & R_0^T \hat{p}_0 R_0 \\ -R_0^T \hat{p}_0 R_0 & I - R_0^T \hat{p}_0^2 R_0 \end{bmatrix} \end{aligned}$$

In general, $\|V_2^s\| \neq \|V_1^s\|$, or there exists no bi-invariant metric on $se(3)$.

Coordinate Transformation

$$g_{ac}(t) = g_{ab}(t) \cdot g_{bc}(t)$$

$$\hat{V}_{ac}^s = \dot{g}_{ac} \cdot g_{ac}^{-1}$$

$$= (\dot{g}_{ab} \cdot g_{bc} + g_{ab} \cdot \dot{g}_{bc})(g_{bc}^{-1} \cdot g_{ab}^{-1})$$

$$= \dot{g}_{ab} \cdot g_{ab}^{-1} + g_{ab} \cdot \dot{g}_{bc} \cdot g_{bc}^{-1} \cdot g_{ab}^{-1} = \hat{V}_{ab}^s + g_{ab} \hat{V}_{bc}^s g_{ab}^{-1}$$

$$\Rightarrow V_{ac}^s = V_{ab}^s + Ad_{g_{ab}} V_{bc}^s$$

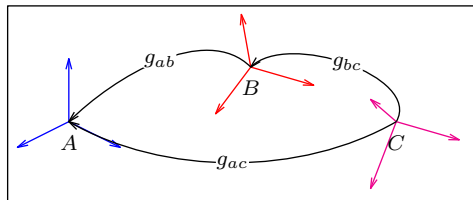


Figure 2.12

Similarly: $V_{ac}^b = Ad_{g_{bc}^{-1}} V_{ab}^b + V_{bc}^b$

Note: $V_{bc}^s = 0 \Rightarrow V_{ac}^s = V_{ab}^s, V_{ab}^b = 0 \Rightarrow V_{ac}^b = V_{bc}^b$

Example

$$g_{ab}(\theta_1) = \begin{bmatrix} c_{\theta_1} & -s_{\theta_1} & 0 & 0 \\ s_{\theta_1} & c_{\theta_1} & 0 & 0 \\ 0 & 0 & 1 & l_0 \\ 0 & 0 & 0 & 1 \end{bmatrix}, V_{ab}^s = \begin{bmatrix} 0 \\ 0 \\ 0 \\ 1 \end{bmatrix} \dot{\theta}_1$$

$$g_{bc}(\theta_2) = \begin{bmatrix} c_{\theta_2} & -s_{\theta_2} & 0 & 0 \\ s_{\theta_2} & c_{\theta_2} & 0 & l_1 \\ 0 & 0 & 1 & l_2 \\ 0 & 0 & 0 & 1 \end{bmatrix}, V_{bc}^s = \begin{bmatrix} l_1 \\ 0 \\ 0 \\ 1 \end{bmatrix} \dot{\theta}_2$$

$$V_{ac}^s = V_{ab}^s + Ad_{g_{ab}} \cdot V_{bc}^s = \begin{bmatrix} 0 \\ 0 \\ 0 \\ 1 \end{bmatrix} \dot{\theta}_1 + \begin{bmatrix} l_1 c_{\theta_1} \\ l_1 s_{\theta_1} \\ 0 \\ 1 \end{bmatrix} \dot{\theta}_2$$

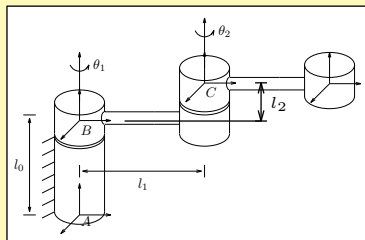


Figure 2.16

† End of Section †

Wrenches & Reciprocal Screws

Let

$$F_c = \begin{bmatrix} f_c \\ \tau_c \end{bmatrix} \in \mathbb{R}^6, f_c, \tau_c \in \mathbb{R}^3$$

be force

or moment applied at the origin of C

Generalized power:

$$\delta W = F_c \cdot V_{ac}^b = \langle f_c, v_{ac}^b \rangle + \langle \tau_c, \omega_{ac}^b \rangle$$

Work:

$$W = \int_{t_1}^{t_2} V_{ac}^b \cdot F_c dt$$

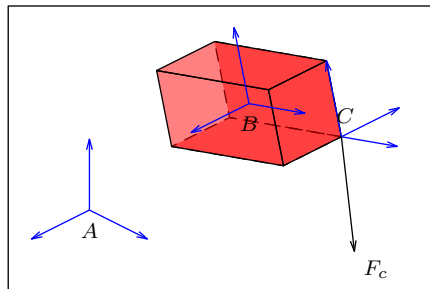


Figure 2.17

$$\begin{aligned} V_{ab}^b \cdot F_b &= (\text{Ad}_{g_{bc}} \cdot V_{ac}^b)^T \cdot F_b \\ &= (V_{ac}^b)^T \text{Ad}_{g_{bc}}^T \cdot F_b = (V_{ac}^b)^T \cdot F_c, \forall V_{ac}^b \\ \Rightarrow F_c &= \text{Ad}_{g_{bc}}^T \cdot F_b \end{aligned}$$

(see next page)

Wrenches & Reciprocal Screws

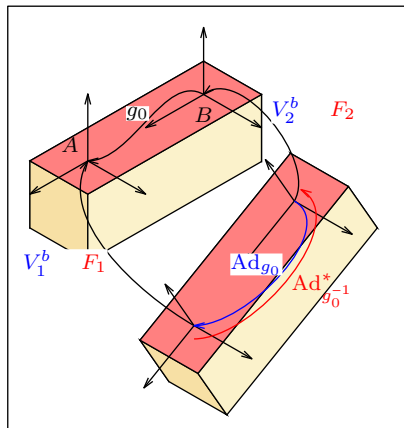


Figure 2.18

$$\begin{aligned}
 V_2^s &= \text{Ad}_{g_0^{-1}} \cdot V_1^s \\
 (V_2^b &= \text{Ad}_{g_0^{-1}} \cdot V_1^b) \\
 \Rightarrow V_1^b &= \text{Ad}_{g_0} \cdot V_2^b \\
 F_2 &= \text{Ad}_{g_0}^* F_1
 \end{aligned}$$

Screw coordinates for a wrench

Generate a wrench associated with S :

- ($h \neq \infty$): force of mag. M along l , and torque of mag. hM about l .
- ($h = \infty$): pure torque of mag. M about l

$$F = \begin{cases} M \begin{bmatrix} \omega \\ -\omega \times q + h\omega \end{bmatrix} & h \neq \infty \\ M \begin{bmatrix} 0 \\ \omega \end{bmatrix} & h = \infty \end{cases}$$

F : wrench along the screw S .

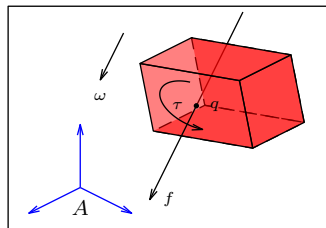


Figure 2.19

(see next page)

Screw coordinates for a wrench (Continued)

1 Pitch:

$$h = \begin{cases} \frac{f^T \tau}{\|f\|^2} & \text{if } f \neq 0 \\ \infty & \text{if } f = 0 \end{cases}$$

2 Axis:

$$l = \begin{cases} \frac{f \times \tau}{\|f\|^2} + \lambda f, \lambda \in \mathbb{R} & \text{if } f \neq 0 \\ 0 + \lambda \tau, \lambda \in \mathbb{R} & \text{if } f = 0 \end{cases}$$

3 Magnitude:

$$M = \begin{cases} \|f\| & \text{if } f \neq 0 \\ \|\tau\| & \text{if } f = 0 \end{cases}$$

Poinsot Theorem

Theorem 3 (Poinsot):

Every collection of wrenches applied to a rigid body is equivalent to a force applied along a fixed axis plus a torque about the axis.



1777-1859

□ Multi-fingered grasp:

$$F_O = \sum_{i=1}^k \text{Ad}_{g_{oc_i}^{-1}}^T \cdot F_{c_i}$$

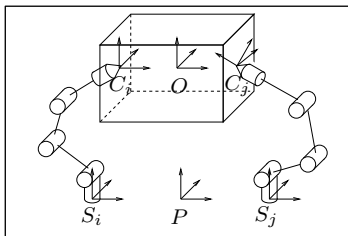


Figure 2.20

Reciprocal screws

$$V = \begin{bmatrix} v \\ \omega \end{bmatrix}, F = \begin{bmatrix} f \\ \tau \end{bmatrix}$$

$$F \cdot V = f^T \cdot v + \tau^T \cdot \omega$$

$$\downarrow \downarrow$$

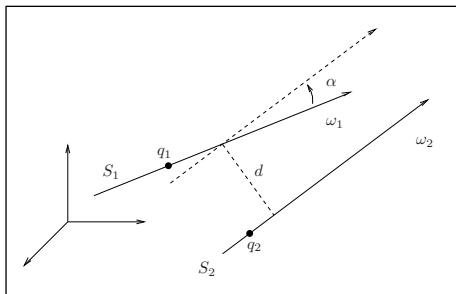
$$S_2 \ S_1$$


Figure 2.21

$$\alpha = \text{atan2}((\omega_1 \times \omega_2) \cdot n, \omega_1 \cdot \omega_2)$$

$$\begin{aligned} S_1 \odot S_2 &= M_1 M_2 ((h_1 + h_2) \cos \alpha - d \sin \alpha) \\ &= 0 \text{ if reciprocal} \end{aligned}$$

(continues next slide)

Reciprocal screws

Given $V = M_1 \begin{bmatrix} q_1 \times \omega_1 + h_1 \omega_1 \\ \omega_1 \end{bmatrix}$, $F = M_2 \begin{bmatrix} \omega_2 \\ q_2 \times \omega_2 + h_2 \omega_2 \end{bmatrix}$,

Let $q_2 = q_1 + dn$, then

$$\begin{aligned} V \cdot F &= M_1 M_2 (\omega_2 \cdot (q_1 \times \omega_1 + h_1 \omega_1) + \omega_1 \cdot (q_2 \times \omega_2 + h_2 \omega_2)) \\ &= M_1 M_2 (\omega_2 \cdot (q_1 \times \omega_1) + h_1 \omega_1 \cdot \omega_2 \\ &\quad + \omega_1 \cdot ((q_1 + dn) \times \omega_2) + h_2 \omega_1 \cdot \omega_2) \\ &= M_1 M_2 ((h_1 + h_2) \cos \alpha - d \sin \alpha) \end{aligned}$$

Example: basic joints

- Revolute joint: $\xi = \begin{bmatrix} -\omega \times q \\ \omega \end{bmatrix}$

$$\xi^\perp = \text{span} \left\{ \begin{bmatrix} \omega_i \\ q \times \omega_i \end{bmatrix}, \begin{bmatrix} 0 \\ v_j \end{bmatrix} \mid \omega_i \in S^2, i = 1, 2, 3 \right\}: 5\text{-system}$$

$$v_j \cdot \omega = 0, j = 1, 2$$

- Prismatic joint: $\xi = \begin{bmatrix} v \\ 0 \end{bmatrix}$

$$\xi^\perp = \text{span} \left\{ \begin{bmatrix} \omega_i \\ q \times \omega_i \end{bmatrix}, \begin{bmatrix} 0 \\ v_j \end{bmatrix} \mid \omega_i \cdot v = 0, i = 1, 2 \right\}: 5\text{-system}$$

$$v_j \in S^2, j = 1, 2, 3$$

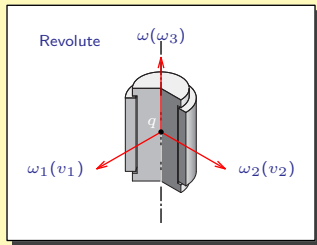


Figure 2.22

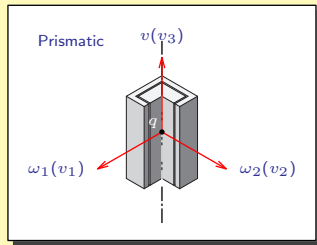


Figure 2.23

Basic joints (continued)

- Spherical joint: $\xi = \text{span} \left\{ \begin{bmatrix} -\omega_i \times q \\ \omega_i \end{bmatrix} \mid \omega_i \in S^2, i = 1, 2, 3 \right\}$

$$\xi^\perp = \text{span} \left\{ \begin{bmatrix} \omega_i \\ q \times \omega_i \end{bmatrix} \mid \omega_i \in S^2, i = 1, 2, 3 \right\}: 3\text{-system}$$

- Universal joint: $\xi = \text{span} \left\{ \begin{bmatrix} q \times x \\ x \end{bmatrix}, \begin{bmatrix} q \times y \\ y \end{bmatrix} \right\}$

$$\xi^\perp = \text{span} \left\{ \begin{bmatrix} \omega_i \\ q \times \omega_i \end{bmatrix}, \begin{bmatrix} 0 \\ z \end{bmatrix} \mid \omega_i \in S^2, i = 1, 2, 3 \right\}: 4\text{-system}$$

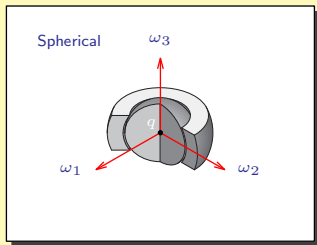


Figure 2.24

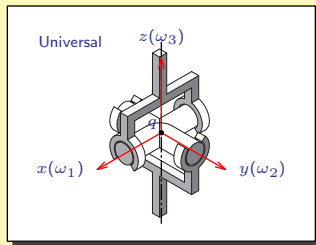


Figure 2.25

Kinematic chains

- Universal-Spherical Dyad:

$$\xi = \text{span} \left\{ \begin{bmatrix} q_1 \times x \\ x \end{bmatrix}, \begin{bmatrix} q_1 \times y \\ y \end{bmatrix}, \begin{bmatrix} q_2 \times \omega_i \\ \omega_i \end{bmatrix} \middle| \omega_i \in S^2, i = 1, 2, 3 \right\}$$

$$\xi^\perp = \text{span} \left\{ \begin{bmatrix} v \\ q_1 \times v \end{bmatrix} \middle| v = \frac{q_2 - q_1}{\|q_2 - q_1\|} \right\}$$

- Revolute-Spherical Dyad: zero pitch screws passing through the center of the sphere, lie on a plane containing the axis of the revolute joint: 2-system

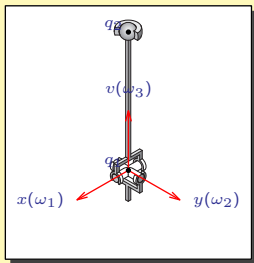


Figure 2.26

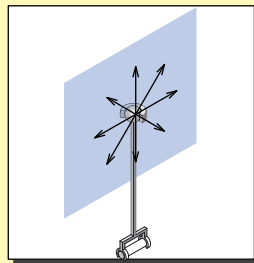


Figure 2.27

† End of Section †

References

□ Reference:

- [1] Murray, R.M. and Li, Z.X. and Sastry, S.S., **A mathematical introduction to robotic manipulation**. CRC Press, 1994.
- [2] Ball, R.S., **A treatise on the theory of screws**. University Press, 1900.
- [3] Bottema, O. and Roth, B. , **Theoretical kinematics**. Dover Publications, 1990.
- [4] Craig, J.J., **Introduction to robotics : mechanics and control**, 3rd ed. Prentice Hall, 2004.
- [5] Fu, K.S. and Gonzalez, R.C. and Lee, C.S.G., **Robotics : control, sensing, vision, and intelligence**. CAD/CAM, robotics, and computer vision. McGraw-Hill, 1987.
- [6] Hunt, K.H., **Kinematic geometry of mechanisms**. 1978, Oxford, New York: Clarendon Press, 1978.
- [7] Paul, R.P., **Robot manipulators : mathematics, programming, and control**. The MIT Press series in artificial intelligence. MIT Press, 1981.
- [8] Park, F. C., **A first course in robot mechanics**. Available online, 2006.
- [9] Tsai, L.-W., **Robot analysis : the mechanics of serial and parallel manipulators**. Wiley, 1999.
- [10] Spong, M.W. and Hutchinson, S. and Vidyasagar, M. , **Robot modeling and control**. John Wiley & Sons, 2006.
- [11] Selig, J., **Geometric Fundamentals of Robotics**. Springer, 2008.