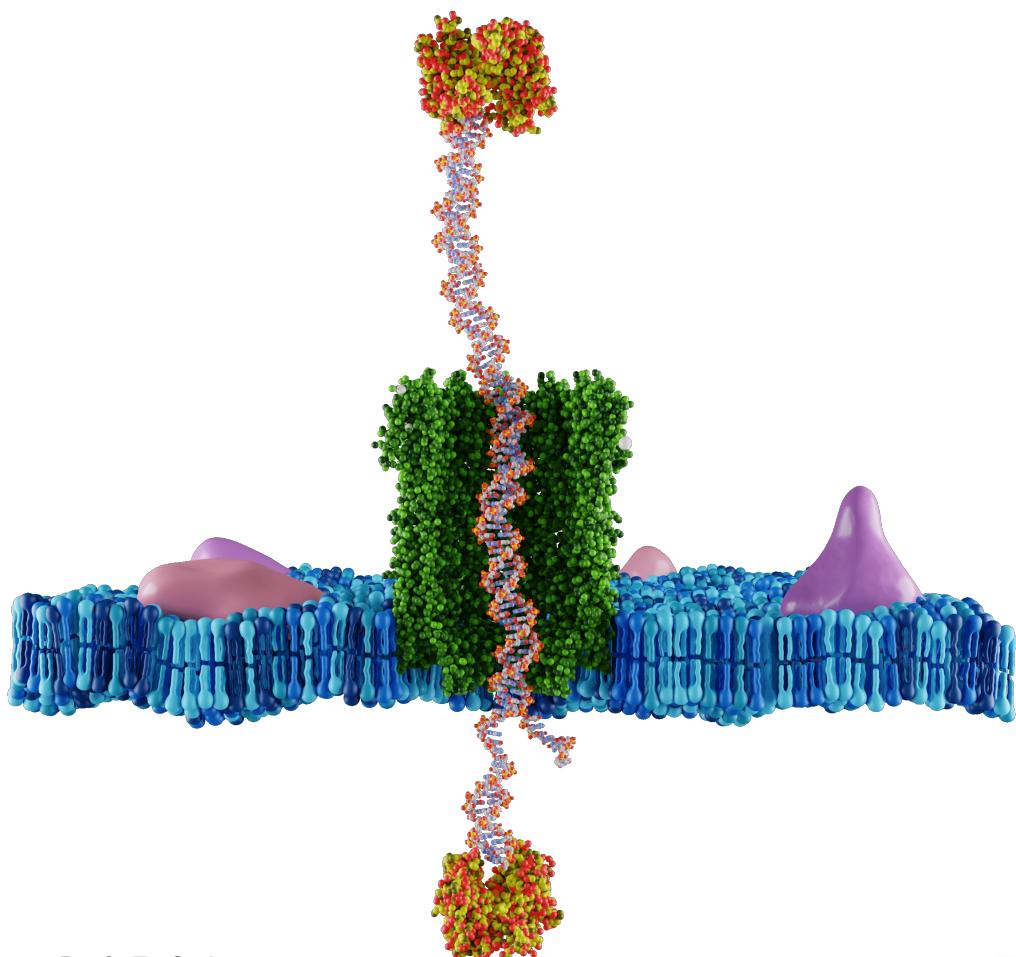


# Coarse-grained simulations of the DNA nanopiston

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# **Abstract**

abstract



# Vulgariserende Samenvatting

Summary in dutch.

asdf



# **Summary in Layman's Terms**

Summary in english.







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# CHAPTER

# 1

## Introduction

*...if we were to name the most powerful assumption of all, which leads one on and on in an attempt to understand life, it is that all things are made of atoms, and that everything that living things do can be understood in terms of the jigglings and wigglings of atoms.*

---

— Richard P. Feynman, *The Feynman Lectures on Physics*<sup>2</sup>

### 1.1 Deoxyribonucleic Acid

### 1.2 Polymer Physics

### 1.3 Computer Simulations

The theory of classical mechanics is regarded as the first major breakthrough in the field of physics. For every aspiring physicist this is still the starting point of their studies. Unfortunately getting to know these relatively simple laws of nature, leads to the inescapable realisation that these theories are expressed in mathematical formalism that are only analytically solvable in few idealised scenarios. Already when trying to apply these formulas to a problem consisting of just more than two bodies the equations become practically unsolvable, this fact has vexed physicists and mathematician for many years.

Although it is often times not possible to find an exact solution to many equations encountered in physics, finding reasonable approximations to their solution is. Various types of techniques are used to find these approximations, ranging from taylor expansions to simulations. Before the invention of computers these simulations where performed mechanical and consisted of large amounts of

... I took a number of rubber balls and stuck them together with rods of a selection of different lengths ranging from 2.75 to 4 inch. I tried to do this in the first place as casually as possible, working in my own office being interrupted every five minutes or so and not remembering what I had done before the interruption. However, . . .

mechanical simulations -> bernal

computer simulations -> los alamos, fast adaptation by industry.

Biophysics / condensed matter physics, many body systems, bridge between the microscopic and macroscopic

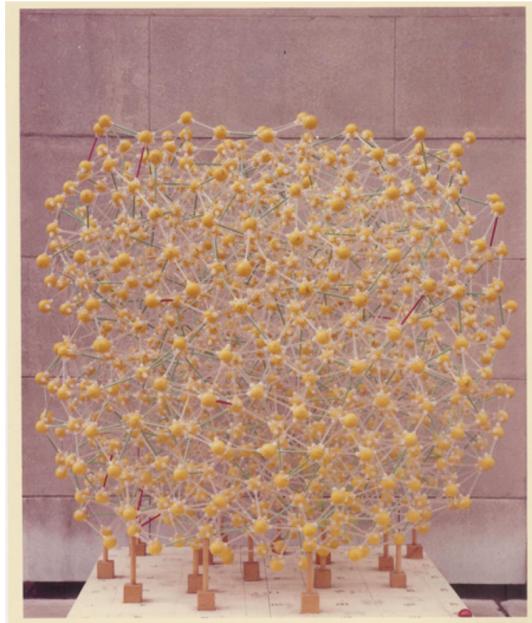


Figure 1.1: Example of an expanded model of a simple liquid (J L Finney, Ph.D thesis)

The central idea is that the trajectories of a system of  $N$  particles can be generated over time by numerically integrating the classical equations of motion.

- Iterating Newton laws of physics

$$m_i \ddot{\mathbf{r}}_i = \mathbf{f}_i, \quad \mathbf{f}_i = -\frac{\partial}{\partial \mathbf{r}_i} \mathcal{U}$$

- thermostats

- Integrations schemes
- Rare event sampling

---

**Algorithm 1:** The Velocity Verlet algorithm

---

```

Input : Configuration of the system at  $t = 0$ 
Output : Configuration of the system at  $t = t_f$ 
1 newList = [ ]
2 /* For odd elements in the list, we add 1, and for even
   elements, we add 2. After the loop, all elements are even.
   */
3 for  $i \leftarrow 0$  to  $n - 1$  do
4   if isOddNumber( $a_i$ ) then
5     | newList.append( $a_i + 1$ )           // Some thought-provoking comment.
6   else
7     | // Another comment
8     | newList.append( $a_i + 2$ )
9   end if
10 end for
11 return newList
```

---

- understanding many body - newtons algorithm - insight in the dynamics -> simulate trajectories - recent developments in techniques to simulate trajectories of rare event -increased computational power

### 1.3.2 Coarse Grained modelling



# CHAPTER 2

## nano pore

asldfkasdflj



# CHAPTER 3

## Methods

*Given for one instant an intelligence which could comprehend all the forces by which nature is animated and the respective positions of the beings which compose it, if moreover this intelligence were vast enough to submit these data to analysis, it would embrace in the same formula both the movements of the largest bodies in the universe and those of the lightest atom; to it nothing would be uncertain, and the future as the past would be present to its eyes.*

---

— Pierre-Simon Laplace

### 3.1 Figures

An example is Figure 3.1

### 3. METHODS

---

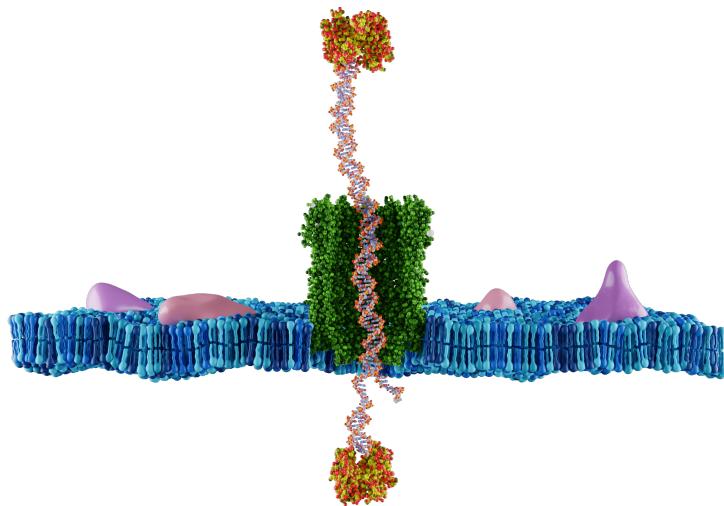


Figure 3.1: A landslide.

## 3.2 Tables

An example is Table 3.1

Model	Accuracy
regression	90%
random forests	95%

Table 3.1: A random table.

### 3.3 Equations

Equations can be inserted in the text itself, working in the *mathmode*(put text between \$-signs, for example  $Y_i = \frac{1}{x}$ ). Or put them in the text as a numbered floating element (e.g. Equation (3.1)).

$$y = \frac{1}{x} \quad (3.1)$$

$$y = \int_a^b x^2 dx \quad (3.2)$$

$$y = \sum_{i=1}^n x_i^2 \quad (3.3)$$

You can align the equations:

$$y = \frac{1}{x} \quad (3.4)$$

$$y = \int_a^b x^2 dx \quad (3.5)$$

$$y = \sum_{i=1}^n x_i^2 \quad (3.6)$$



# CHAPTER 4

## Rotaxane

### 4.1 Mixed Rotaxane

#### 4.1.1 Diffusion approximation

Studying the dynamics of the mixed rotaxane highlighted the importance of entropic interactions between the nano pore and the DNA strand. Here we observed that a fully double stranded DNA polymer represented a special case. The uniformity of the  $\mathcal{X}$  histogram corresponding to this 0 nt mixed rotaxane suggests a free diffusive motion of the rotaxane in a bounded one-dimensional domain. This isotropic behaviour was previously also observed in the bead-spring simulations by Bayoumi et al.<sup>1</sup>

$$\langle \Delta x^2 \rangle \simeq 2nDt.$$

$$\frac{\partial \psi}{\partial t} = D \frac{\partial^2 \psi}{\partial x^2}, P(x, t) = f(x)g(t)$$

Reflecting boundary conditions  $j = -D \frac{\partial \psi}{\partial x} = 0$ . Current vanishes at the boundaries

$$t : \quad \dot{g} = -\alpha g(t) \Rightarrow g(t) = e^{-\alpha t}$$

$$x : D\ddot{f} = -\alpha f(x) \Rightarrow f(x) = A \sin(Kx) + B \cos(Kx) \\ = B \cos\left(\frac{\pi n x}{L}\right)$$

$$\frac{\alpha}{D} = \frac{\pi^2 n^2}{L^2}$$

The general solution is given by the linear combination,

$$\psi(x, t) = \sum_{n=0}^{+\infty} C_n \cos\left(\frac{\pi n x}{L}\right) e^{-\frac{D\pi^2 n^2}{L^2} t} \\ = \frac{1}{L} \left\{ 1 + \sum_{n=1}^{+\infty} \cos\left(\frac{\pi n x_0}{L}\right) \cos\left(\frac{\pi n x}{L}\right) e^{-\frac{D\pi^2 n^2}{L^2} t} \right\}$$

$$\langle \Delta x^2 \rangle = \langle (x - x_0)^2 \rangle \\ = \frac{L^2}{6} \left( 1 - \frac{96}{\pi^4} \sum_{n=0}^{+\infty} \frac{1}{(2k+1)^4} e^{-\frac{D(2k+1)^2 \pi^2}{L^2} t} \right)$$

As expected, the mean squared distances saturates to  $\langle \Delta x^2 \rangle = L^2/6$  in the long-time limit  $t \gg L^2/D$ .

# CHAPTER 5

## hybrydisation

asldfkasdflj



# 6

CHAPTER

## Conclusions and Perspectives

6.0.1 asdf



## APPENDIX

A

# First appendix

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# Acknowledgements

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