## Datreat theory: ring2

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theory id-name: ring2
Dynamic scattering function for a polymer ring melt as developed in the PRL of S. Goossen
Parameters (17):
 1: ampli
                     prefactor
  2: diff
              > ABS: limiting diffusion (NMR value) in [A**2/ns]
  3: r021
              > ABS: reference mean squared displacement at first transition point to medium timed diff.
  4: alpha
              > ABS: sub-diffusion exponent of SHORT time diffusion
 5: r022
              > ABS: reference mean squared displacement at second transition point to D0 in [A**2]
  6: beta
              > ABS: sub-diffusion exponent of medium time diffusion
  7: a_cross > ABS: transition exponent between short and long time diffusion (sharper kink for larger a)
              > INT: number of segments in one ring
              > ABS: effective segment length
 10: nue
              > ABS: chain statistics exponent (nu=0.5 => random walk, Gaussian)
 11: wl4
              > ABS: Rouse rate in [A**4/ns]
              > ABS: transition mode number between simple ring-Rouse and large p modification
12: pmin
             > ABS: sharpness of transition
> ABS: prefactor f(p) limit for small p values (default 1)
> ABS: prefactor f(p) limit for large p values (default F=0.9??) transitin width is pwidth
13: pwidth
 14: f0
15: finf
            > ABS: large p tau(p) = tauinf/p**pexinf
> ABS: large p tau(p) = tauinf/p**pexinf
16: tauinf
17: pexinf
INPUT: Parameters that are extracted from the actual considered data records:
..there may be default assumptions, but better make sure that these parameters are set properly!
             > a
OUTPUT: Parameters that are computed and added to the records parameters as information:
          > predicted ring radius of gyration
cite: S. Goossen et al., PRL 2014, 113, 168302 !
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Modified polymer ring structure factor with sublinear diffusion with 3 regimes.

## Diffusion factor:

$$\exp\left(-\langle r^2(t)\rangle q^2/6\right)$$

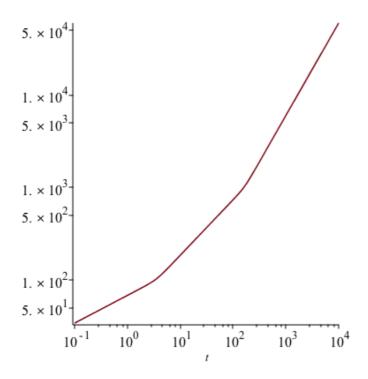
with

$$\langle r^{2}(t) \rangle = \left[ \left( e^{\frac{1}{\beta} \left( (-\alpha + \beta) \ln \left( \frac{r_{I}^{2}}{r_{2}^{2}} \right) + \beta \alpha \left( -\ln \left( \frac{r_{2}^{2}}{6D_{0}} \right) \right) \right)} r_{2}^{2} t^{\alpha} \right]^{a} + \left( e^{\left( -\ln \left( \frac{r_{2}^{2}}{6D_{0}} \right) \right) \beta} r_{2}^{2} t^{\beta} \right)^{a} + \left( 6D_{0} t \right)^{a} \right]^{1/a}$$

describing sublinear-linear center-of-mass diffusion with a normal long time diffusion constant  $D_0$  and a sublinear diffusion with exponent  $\alpha$  until the first kink at the msd level  $r_1^2$  and with an exponent  $\beta$  until the second kink at the msd level  $r_2^2$ . The sharpness of the transitions is controlled by a, the larger the value the sharper is the transition.

Intended use assumes that  $0 < \alpha < \beta < 1$  and  $r_1^2 < r_2^2$ . This is NOT automatic !

The figure shows  $\langle r^2(t) \rangle$  for  $r_1^2=100 {\rm \mathring{A}}^2$  ;  $r_2^2=1000 {\rm \mathring{A}}^2$  ;  $\alpha=0.3$  and  $\beta=0.6$  (a=20).



Ring structure factor (without diffusion):

$$S(q,t) = \frac{1}{N} \sum_{i,j}^{N} \exp \left[ \frac{(ql)^2}{6} \left( |i-j| \{N - |i-j| \}/N \right)^{2\nu} - B_{i,j}(t) \right]$$

with

$$B_{i,j}(t) = A \sum_{p \text{ even}} \frac{F(p)}{p^2} \cos(p\pi[i-j]/N) [1 - \exp(-t\Gamma(p))]$$

with the mode p dependent amplitude F(p) and rate  $\Gamma(p)$  parameters. Here

$$A = 2N^{2\nu} (lq)^2 / (3\pi^2)$$

and

$$F(p) = F_0[1 - T_f(p)] + F_\infty T_f(p)$$

and

$$\Gamma(p) = [1 - T_f(p)]p^2/\tau_R + T_f(p)p^\mu/\tau_\infty$$

the transition function is

$$T_f(p) = \{1 + \exp([p - p_{min}]/p_{width})\}^{-1}$$

the Rouse time is computed as

$$au_R = N^2/(W\pi^2)$$
 with  $W = W l^4/l^4$ 

## **Problems (scaling behavior)**

The above formulation (and also the following) exhibits only works approximately if we use a (N, I) scaling:  $N \to \eta N$  and  $l \to l/\eta^{\nu}$ !.

In addition I guess the following modification (motivated by the old expression for Zimm in the norbornene dynamics paper, sources for that were Doi-Edwards and papers of Harnau and Winkler (I would have to rework the matter but changes made in analogy to that are shown below) should be made.

$$S(q,t) = \frac{1}{N} \sum_{i,j}^{N} \exp \left[ \frac{(ql)^2}{6} \left( |i-j| \{N-|i-j| \}/N \right)^{2\nu} - B_{i,j}(t) \right]$$

with

$$B_{i,j}(t) = A \sum_{p,even} \frac{F(p)}{p^{1+2\nu}} \cos(p\pi[i-j]/N) [1 - \exp(-t\Gamma(p))]$$

with the mode p dependent amplitude F(p) and rate  $\Gamma(p)$  parameters. Here

$$A = 2N^{2\nu} (lq)^2 / (3\pi^2)$$

and

$$F(p) = F_0[1 - T_f(p)] + F_{\infty}T_f(p)$$

and

$$\Gamma(p) = [1 - T_f(p)] p^{4\nu} / \tau_R + T_f(p) p^{\mu} / \tau_\infty$$

the transition function is

$$T_f(p) = \{1 + \exp([p - p_{min}]/p_{width})\}^{-1}$$

the Rouse time is computed as

$$au_R = N^2/(W\pi^2)$$
 with  $W = W l^4/l^4$ 

Scaling works approximately with  $N \to \eta N$  and  $l \to l/\eta^{\nu}$ !

This formulation (still missing possible modifications for the  $p^4$ -terms) is implemented as theory  ${\bf ring2a}$ .

## Parameter correspondences

Observe units, if not other specified: Angstroems and nano-seconds

$$\begin{aligned}
&\text{diff} \to D_0 \\
&\text{r021} \to r_1^2 \\
&\text{alpha} \to \alpha \\
&\text{r022} \to r_2^2 \\
&\text{beta} \to \beta \\
&\text{a\_cross} \to a \\
&\text{n} \to N
\end{aligned}$$

 $\begin{array}{l} 1 \rightarrow l \\ \text{nue} \rightarrow \nu \\ \text{wl4} \rightarrow Wl^4 \\ \text{pmin} \rightarrow p_{min} \\ \text{pwidth} \rightarrow p_{width} \\ \text{f0} \rightarrow F_0 \\ \text{finf} \rightarrow F_\infty \\ \text{tauinf} \rightarrow \tau_\infty \\ \text{pexinf} \rightarrow \mu \end{array}$