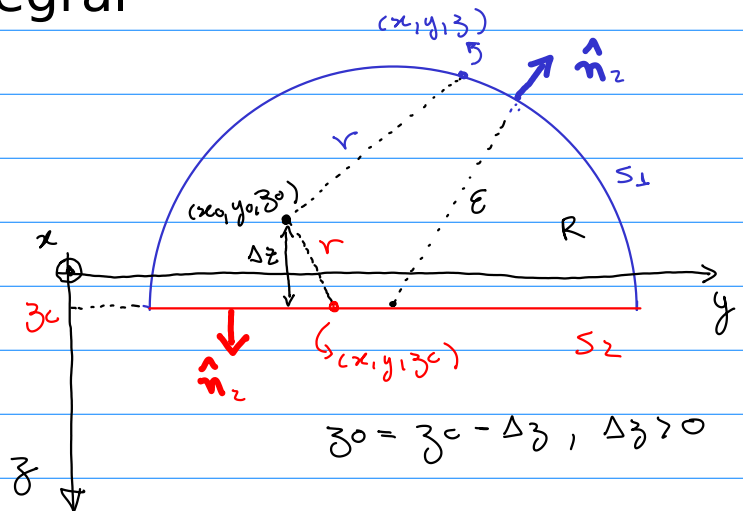


# Upward continuation integral

Green's third identity (Kellogg, 1967, p. 219)

$$U_0 = -\frac{1}{4\pi} \iiint_R \frac{1}{r} \nabla^2 U \, d\sigma + \frac{1}{4\pi} \iint_S \frac{1}{r} \partial_n U \, dS - \frac{1}{4\pi} \iint_S U \partial_n \frac{1}{r} \, dS$$

Split the surface  $S$  into the surfaces  $S_1$  and  $S_2$  and consider that  $U$  is harmonic in  $R$ .



$$U_0 = \frac{1}{4\pi} \iint_{S_1} \frac{1}{r} \partial_n U \, dS_1 + \frac{1}{4\pi} \iint_{S_2} \frac{1}{r} \partial_n U \, dS_2 - \frac{1}{4\pi} \iint_{S_1} U \partial_n \frac{1}{r} \, dS_1 - \frac{1}{4\pi} \iint_{S_2} U \partial_n \frac{1}{r} \, dS_2$$

$\partial_n U = \nabla U^T \hat{n}_2$   
 $\partial_n \frac{1}{r} = \nabla \frac{1}{r}^T \hat{n}_2$

We consider that  $U$  and its derivatives are regular at infinite (Kellogg, 1967, p. 217)

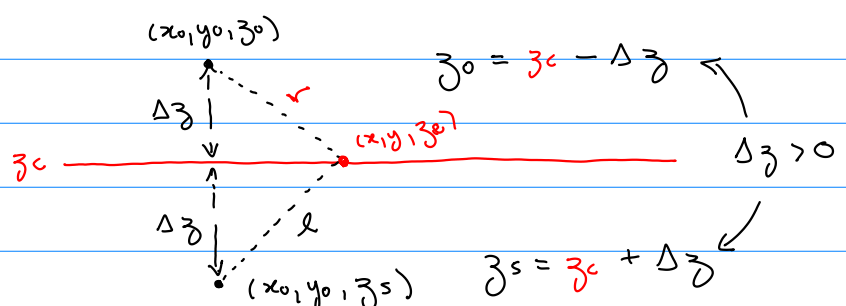
By letting  $\varepsilon \rightarrow \infty$ , the integrals on  $S_1$  vanish and we obtain:

$$U_0 = \frac{1}{4\pi} \iint_{-\infty}^{\infty} \iint_{-\infty}^{\infty} \frac{1}{r} \partial_z U - U \partial_z \frac{1}{r} \, dx dy$$

$\hat{n}_2 = \hat{z}$   
 $\partial_n \square = \nabla \square^T \hat{z} = \partial_z \square$   
 $dS_2 = dx dy$

$$\frac{1}{r} = \frac{1}{[(x_0 - x)^2 + (y_0 - y)^2 + (z_0 - z_c)^2]^{1/2}}$$

$$\frac{1}{r} = \frac{1}{[(x_0 - x)^2 + (y_0 - y)^2 + (z_s - z_c)^2]^{1/2}}$$

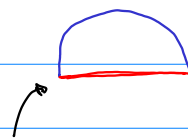


Ex: Show that  $1/r$  is harmonic in  $R$ .

Applying the Green's second identity with  $U$  and  $\frac{1}{r}$ , we obtain:

$$\iiint_R \frac{1}{r} \nabla^2 U - U \nabla^2 \frac{1}{r} dV = \iint_S \frac{1}{r} \partial_n U - U \partial_n \frac{1}{r} dS$$

$= 0 \text{ ok?} \quad \quad = 0 \text{ ok?}$



$$\iint_S \frac{1}{r} \partial_n U dS - \iint_S U \partial_n \frac{1}{r} dS = 0$$

We consider that  $U$  and its derivatives are regular at infinite (Kellogg, 1967, p. 217)

By letting  $\epsilon \rightarrow \infty$ , the integrals on  $S_\epsilon$  vanish and we obtain:

$$\int_{-\infty}^{\infty} \int_{-\infty}^{\infty} \frac{1}{r} \partial_z U dx dy - \int_{-\infty}^{\infty} \int_{-\infty}^{\infty} U \partial_z \frac{1}{r} dx dy = 0$$

Now, multiply this equation by  $\frac{1}{4\pi}$  and subtract or add the result from the previous equation for  $U_0$ :

$$U_0 = \frac{1}{4\pi} \int_{-\infty}^{\infty} \int_{-\infty}^{\infty} \left( \frac{1}{r} \mp \frac{1}{r} \right) \partial_z U dx dy - \int_{-\infty}^{\infty} \int_{-\infty}^{\infty} U \left( \partial_z \frac{1}{r} \mp \partial_z \frac{1}{r} \right) dx dy$$

Ex: Show that the  $\frac{1}{r} = \frac{1}{r}$  and  $\partial_z \frac{1}{r} = -\partial_z \frac{1}{r}$  for points on the surface  $z = z_0$

Case 1) Result obtained by subtracting

$$U_0 = \frac{1}{4\pi} \int_{-\infty}^{\infty} \int_{-\infty}^{\infty} \left( \frac{1}{r} - \frac{1}{r} \right) \partial_z U dx dy - \int_{-\infty}^{\infty} \int_{-\infty}^{\infty} U \left( \partial_z \frac{1}{r} - \partial_z \frac{1}{r} \right) dx dy$$

$$U_0 = -\frac{1}{4\pi} \int_{-\infty}^{\infty} \int_{-\infty}^{\infty} U \left( 2 \partial_z \frac{1}{r} \right) dx dy, \quad \partial_z \frac{1}{r} = -\frac{z_0 - z_0}{r^3}$$

$$U(x_0, y_0, z_0) = \frac{z_0 - z_0}{2\pi} \int_{-\infty}^{\infty} \int_{-\infty}^{\infty} \frac{U(x, y, z_0)}{[(x_0 - x)^2 + (y_0 - y)^2 + (z_0 - z_0)^2]^{3/2}} dx dy$$

upward continuation integral

(Skeels, 1947; Henderson and Zietz, 1949; Henderson, 1960; Roy, 1962; Bhattacharyya, 1967; Henderson, 1970; Blakely, 1996, p. 40)

The upward continuation integral states that the values of a harmonic function  $U(x_0, y_0, z_0)$  at any point  $(x_0, y_0, z_0)$ ,  $z_0 < z_c$ , can be exactly reproduced by the convolution of its values  $U(x, y, z_c)$  and the vertical derivative of the function  $1/r$ , both evaluated on the horizontal plane  $z = z_c$ . This equation also shows that any spatial derivative of the harmonic function  $U(x_0, y_0, z_0)$  can be obtained by properly differentiating the integrand. Then, by assuming the knowledge of the harmonic function on the horizontal plane  $z = z_c$ , it is possible to compute not only  $U(x_0, y_0, z_0)$ , but also any of its spatial derivatives at any point

Case 2) Result obtained by adding

$$U_0 = \frac{1}{4\pi} \int_{-\infty}^{\infty} \int_{-\infty}^{\infty} \left( \frac{1}{r} + \frac{1}{r} \right) \partial_z U \, dx \, dy - \int_{-\infty}^{\infty} \int_{-\infty}^{\infty} U \left( \cancel{\partial_z \frac{1}{r}} + \partial_z \frac{1}{r} \right) dx \, dy$$

= 0

$$U_0 = \frac{1}{4\pi} \int_{-\infty}^{\infty} \int_{-\infty}^{\infty} \left( 2 \frac{1}{r} \right) \partial_z U \, dx \, dy$$

$$U(x_0, y_0, z_0) = \frac{1}{2\pi} \int_{-\infty}^{\infty} \int_{-\infty}^{\infty} \frac{\partial_z U(x, y, z_c)}{[(x_0 - x)^2 + (y_0 - y)^2 + (z_0 - z_c)^2]^{\frac{3}{2}}} dx \, dy$$

(Roy, 1962)

\* Skeels, D. C., 1947, Ambiguity in gravity interpretation: GEOPHYSICS, 12, 43-56. doi: 10.1190/1.1437295

\* Henderson, R. G., and I. Zietz, 1949, The upward continuation of anomalies in total magnetic intensity fields: GEOPHYSICS, 14, 517-534. doi: 10.1190/1.1437560

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\* Roy, A., 1962, Ambiguity in geophysical interpretation: GEOPHYSICS, 27, 90-99. doi: 10.1190/1.1438985

\* Bhattacharyya, B. K., 1967, Some general properties of potential fields in space and frequency domain: a review: GeosExploration, 5, 127-143. doi: 10.1016/0016-7142(67)90021-X

\* Henderson, R. G., 1970, On the validity of the use of the upward continuation integral for total magnetic intensity data: GEOPHYSICS, 35, 916-919. doi: 10.1190/1.1440137