Chapter 4 Manipulator Dynamics

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Coordinate Invariant Algorithms

Lagrange's Equations with

References

Lecture Notes for A Geometrical Introduction to Robotics and Manipulation

Richard Murray and Zexiang Li and Shankar S. Sastry CRC Press

Zexiang Li¹ and Yuanqing Wu¹

¹ECE, Hong Kong University of Science & Technology

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Definition: Dynamics

Physical laws governing the motions of bodies and aggregates of bodies.

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Definition: Dynamics

Physical laws governing the motions of bodies and aggregates of bodies.

□ A short history:



"Everything happens for a reason."

Aristotle (384 BC - 322 BC)

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Definition: Dynamics

Physical laws governing the motions of bodies and aggregates of bodies.

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Experiments with cannon balls from the tower of Pisa.

G. Galilei (1564 - 1642)

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G. Galilei (1564 - 1642)



Laws of motion.

I. Newton (1642 - 1726)

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Laws of motion from particles to rigid bodies.

L. Euler (1707 - 1783)

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Laws of motion from particles to rigid bodies.

L. Euler (1707 - 1783)



Calculus of Variation and the Principles of least action.

J. Lagrange (1736 - 1813)

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Laws of motion from particles to rigid bodies.

L. Euler (1707 - 1783)



Calculus of Variation and the Principles of least action.

J. Lagrange (1736 - 1813)



Quaternions and Hamilton's Principle.

W. Hamilton (1805 - 1865)

† End of Section †



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Lagrangian Equation:

4.2 Lagrange Equations

□ A simple example:

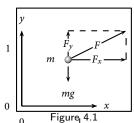
Newton's Equation:

 $m\ddot{x} = F_x$ $m\ddot{y} = F_y - mg$

Momentum: $P_x = m\dot{x}$

$$P_y = m\dot{y}$$

$$\frac{d}{dt}P_x = F_x, \frac{d}{dt}P_y = F_y - mg$$



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□ A simple example:

Newton's Equation:

 $m\ddot{x} = F_x$ $m\ddot{y} = F_y - mg$

Momentum: $P_x = m\dot{x}$

$$\begin{aligned} P_y &= m\dot{y} \\ \frac{\mathrm{d}}{\mathrm{d}t}P_x &= F_x, \frac{\mathrm{d}}{\mathrm{d}t}P_y &= F_y - mg \end{aligned}$$

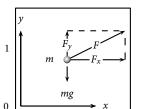


Figure 4.1

Lagrangian Equation:

$$\frac{\mathrm{d}}{\mathrm{d}t} \frac{\partial L}{\partial \dot{x}} - \frac{\partial L}{\partial x} = F_x$$

$$\frac{\mathrm{d}}{\mathrm{d}t} \frac{\partial L}{\partial \dot{y}} - \frac{\partial L}{\partial y} = F_y$$

Lagrangian function:

$$L=T-V, P_x=\frac{\partial L}{\partial \dot{x}}, P_y=\frac{\partial L}{\partial \dot{y}}$$

Kinetic energy:

$$T=\frac{1}{2}m(\dot{x}^2+\dot{y}^2)$$

Potential energy:

$$V = mgy$$

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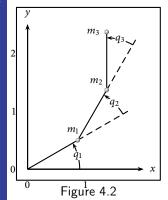
Coordinate Invariant Algorithms

Equations with

□ Generalization to multibody systems:

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Equations



 $q_i, i = 1, \dots, n$: generalized coordinates Kinetic energy:

$$T = T(q, \dot{q})$$

Potential energy:

$$V = V(q)$$

Lagrangian:

$$L(q,\dot{q}) = T(q,\dot{q}) - V(q)$$

 τ_i , i = 1, ..., n: external force on q_i Lagrangian Equation:

$$\frac{\mathrm{d}}{\mathrm{d}t}\frac{\partial L}{\partial \dot{q}_i} - \frac{\partial L}{\partial q_i} = \tau_i, i = 1, \dots, n$$

⋄ Example: Pendulum equation

Generalized coordinate:

 $\theta \in S^1$

Kinematics:

Lagrange's Equations

$$x = l \sin \theta, y = -l \cos \theta$$

$$\dot{x} = l\cos\theta \cdot \dot{\theta}, \dot{y} = l\sin\theta \cdot \dot{\theta}$$

Kinetic energy:

$$T(\theta, \dot{\theta}) = \frac{1}{2}m(\dot{x}^2 + \dot{y}^2) = \frac{1}{2}ml^2\dot{\theta}^2$$

Potential energy:

$$V = mgl(1 - \cos\theta)$$

Lagrangian function:

$$L = T - V = \frac{1}{2}ml^2\dot{\theta} - mgl(1 - \cos\theta), \Rightarrow \frac{\partial L}{\partial \dot{\theta}} = ml^2\dot{\theta}, \frac{\partial L}{\partial \theta} = -mgl\sin\theta$$

Equation of motion:

$$\frac{\mathrm{d}}{\mathrm{d}t}\frac{\partial L}{\partial \dot{\theta}} - \frac{\partial L}{\partial \theta} = \tau \Rightarrow ml^2\ddot{\theta} + mgl\sin\theta = \tau$$

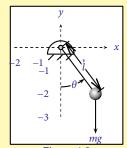


Figure 4.3

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♦ Example: A spherical pendulum

Generalized coordinate:

$$r(\theta, \phi) = \begin{bmatrix} l\sin\theta\cos\phi \\ l\sin\theta\sin\phi \\ -l\cos\theta \end{bmatrix}$$

grange's Kinetic energy:

$$T = \frac{1}{2}m\|\dot{r}\|^2 = \frac{1}{2}ml^2(\dot{\theta}^2 + (1-\cos^2\theta)\dot{\phi}^2)$$

oordinate Potential energy:

$$V = -mgl\cos\theta$$

Lagrangian function:

$$L(q, \dot{q}) = \frac{1}{2}ml^2(\dot{\theta} + (1 - \cos^2\theta)\dot{\phi}^2) + mgl\cos\theta$$

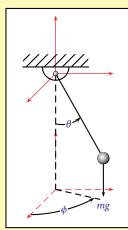


Figure 4.4

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$$\begin{cases} \frac{\mathrm{d}}{\mathrm{d}t} \frac{\partial L}{\partial \dot{\theta}} = \frac{\mathrm{d}}{\mathrm{d}t} (ml^2 \dot{\theta}) = ml^2 \ddot{\theta}, \\ \frac{\partial L}{\partial \theta} = ml^2 \sin \theta \cos \theta \dot{\phi}^2 - mgl \sin \theta \\ \frac{\mathrm{d}}{\mathrm{d}t} \frac{\partial L}{\partial \dot{\phi}} = \frac{\mathrm{d}}{\mathrm{d}t} (ml^2 \sin^2 \theta \dot{\phi}) = ml^2 \sin^2 \theta \ddot{\phi} + 2ml^2 \sin \theta \cos \theta \dot{\theta} \dot{\phi}, \\ \frac{\partial L}{\partial \phi} = 0 \end{cases}$$

 $\begin{bmatrix} ml^2 & 0 \\ 0 & ml^2s_\theta^2 \end{bmatrix} \begin{bmatrix} \ddot{\theta} \\ \ddot{\phi} \end{bmatrix} + \begin{bmatrix} -ml^2s_\theta c_\theta \dot{\phi}^2 \\ 2ml^2s_\theta c_\theta \dot{\theta} \dot{\phi} \end{bmatrix} + \begin{bmatrix} mgls_\theta \\ 0 \end{bmatrix} = \begin{bmatrix} 0 \\ 0 \end{bmatrix}$

·

□ Kinetic energy of a rigid body:

Figure 4.5

Lagrange's Equations

Volume occupied by the body:
$$V$$
 Mass density:
$$\rho(r)$$
 Mass:
$$m = \int_{V} \rho(r) \mathrm{d}V$$
 Mass center:
$$\overline{r} \triangleq \frac{1}{m} \int_{V} \rho(r) r \mathrm{d}V$$
 Relative to frame at the mass center:
$$\overline{r} = 0$$

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Kinetic energy (In A-frame):

 $T = \frac{1}{2} \int_{V} \rho(r) \|\dot{p} + \dot{R}r\|^{2} dV = \frac{1}{2} \int_{V} \rho(r) (\|\dot{p}\|^{2} + 2\dot{p}^{T}\dot{R}r + \|\dot{R}r\|^{2}) dV$ $= \frac{1}{2} m \|\dot{p}\|^{2} + \dot{p}^{T}\dot{R} \int_{V} \rho(r) r dV + \frac{1}{2} \int_{V} \rho(r) \|\dot{R}r\|^{2} dV$

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Kinetic energy (In A-frame):

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$$T = \frac{1}{2} \int_{V} \rho(r) \|\dot{p} + \dot{R}r\|^{2} dV = \frac{1}{2} \int_{V} \rho(r) (\|\dot{p}\|^{2} + 2\dot{p}^{T}\dot{R}r + \|\dot{R}r\|^{2}) dV$$

$$= \frac{1}{2} m \|\dot{p}\|^{2} + \dot{p}^{T}\dot{R} \int_{V} \rho(r) r dV + \frac{1}{2} \int_{V} \rho(r) \|\dot{R}r\|^{2} dV$$

$$= \frac{1}{2} \int_{V} \rho(r) \|\dot{R}r\|^{2} dV$$

$$= \frac{1}{2} \int_{V} \rho(r) \|R^{T}\dot{R}r\|^{2} dV = \frac{1}{2} \int_{V} \rho(r) \|\hat{\omega}r\|^{2} dV = \frac{1}{2} \int_{V} \rho(r) \|\hat{r}\omega\|^{2} dV$$

$$= \frac{1}{2} \int_{V} \rho(r) (-\omega^{T}\hat{r}^{2}\omega) dV = \frac{1}{2} \omega^{T} \left(-\int_{V} \rho(r) \hat{r}^{2} dV\right) \omega \triangleq \frac{1}{2} \omega^{T} \mathcal{I}\omega$$

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where

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Opnamics $\mathcal{I} = -\int \rho(r)\hat{r}^2 dV \triangleq \begin{bmatrix} \mathcal{I}_{xx} & \mathcal{I}_{xy} & \mathcal{I}_{xz} \\ \mathcal{I}_{xy} & \mathcal{I}_{yy} & \mathcal{I}_{yz} \\ \mathcal{I}_{xy} & \mathcal{I}_{yz} & \mathcal{I}_{zz} \end{bmatrix}$

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is the *Inertia tensor* with

he Inertia tensor with
$$\mathcal{I}_{xx} = \int \rho(r)(y^2 + z^2) dx dy dz, \quad \mathcal{I}_{xy} = -\int \rho(r) xy dx dy dz$$

$$T = \frac{1}{2} m \|\dot{p}\|^2 + \frac{1}{2} (\omega^b)^T \mathcal{I} \omega^b = \frac{1}{2} m \|R^T \dot{p}\|^2 + \frac{1}{2} (\omega^b)^T \mathcal{I} \omega^b$$

$$\begin{aligned}
T &= \frac{1}{2}m\|p\|^{2} + \frac{1}{2}(\omega^{b})^{2}\mathcal{L}\omega^{b} = \frac{1}{2}m\|R^{2}p\|^{2} + \frac{1}{2}(\omega^{b})^{2} \\
&= \frac{1}{2}(V^{b})^{T} \underbrace{\begin{bmatrix} mI & 0 \\ 0 & \mathcal{L} \end{bmatrix}}_{M^{b}} V^{b}
\end{aligned}$$

 M^b : Generalized inertia matrix in B-frame.

\diamond Example: M^b for a rectangular object

Lagrange's Equations

$$\rho = \frac{m}{l\omega h}$$

$$\mathcal{I}_{xx} = \int_{V} \rho(y^{2} + z^{2}) dx dy dz$$

$$= \rho \int_{-\frac{h}{2}}^{\frac{h}{2}} \int_{-\frac{\omega}{2}}^{\frac{\omega}{2}} \int_{-\frac{1}{2}}^{\frac{1}{2}} (y^{2} + z^{2}) dx dy dz$$

$$= \rho \left(\frac{1}{12} (l\omega^{3} h + l\omega h^{3}) \right) = \frac{m}{12} (\omega^{2} + h^{2})$$

$$\mathcal{I}_{xy} = -\int_{V} \rho xy dV = -\rho \int_{-\frac{h}{2}}^{\frac{h}{2}} \int_{-\frac{\omega}{2}}^{\frac{\omega}{2}} \int_{-\frac{1}{2}}^{\frac{1}{2}} xy dx dy dz$$

$$= -\rho \int_{-\frac{h}{2}}^{\frac{h}{2}} \int_{-\frac{\omega}{2}}^{\frac{\omega}{2}} \frac{y}{2} x^{2} \Big|_{-\frac{1}{2}}^{\frac{1}{2}} dy dz = 0$$

Figure 4.6

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$$\mathcal{I} = \begin{bmatrix} \frac{m}{12}(w^2 + h^2) & 0 & 0\\ 0 & \frac{m}{12}(l^2 + h^2) & 0\\ 0 & 0 & \frac{m}{12}(w^2 + l^2) \end{bmatrix}, M^b = \begin{bmatrix} mI_{3\times3} & 0\\ 0 & \mathcal{I} \end{bmatrix}$$

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$$\mathcal{I} = \begin{bmatrix} \frac{m}{12} (w^2 + h^2) & 0 & 0\\ 0 & \frac{m}{12} (l^2 + h^2) & 0\\ 0 & 0 & \frac{m}{12} (w^2 + l^2) \end{bmatrix}, M^b = \begin{bmatrix} mI_{3\times3} & 0\\ 0 & \mathcal{I} \end{bmatrix}$$

\square M^b under change of frames:

$$\begin{split} \hat{V}_1 &= g_1^{-1} \cdot \dot{g}_1, \quad T &= \frac{1}{2} V_1^T M_1^b V_1 \\ V_1 &= \mathrm{Ad}_{g_0} V_2 \\ T &= \frac{1}{2} (\mathrm{Ad}_{g_0} V_2)^T M_1^b (\mathrm{Ad}_{g_0} V_2) \\ &= \frac{1}{2} V_2^T \mathrm{Ad}_{g_0}^T M_1^b \mathrm{Ad}_{g_0} V_2 \triangleq \frac{1}{2} V_2^T M_2^b V_2 \end{split}$$

$$M_2^b = \mathrm{Ad}_{g_0}^T M_1^b \mathrm{Ad}_{g_0}$$

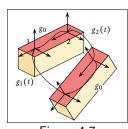


Figure 4.7

⋄ Example: Dynamics of a 2-dof planar robot

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$$\mathcal{I}_i = \left[\begin{array}{ccc} \mathcal{I}_{xx_i} & 0 & 0 \\ 0 & \mathcal{I}_{yy_i} & 0 \\ 0 & 0 & \mathcal{I}_{zz_i} \end{array} \right], i = 1, 2$$

$$T(\theta, \dot{\theta}) = \frac{1}{2} m_1 \|v_1\|^2 + \frac{1}{2} \omega_1^T \mathcal{I}_1 \omega_1$$
$$+ \frac{1}{2} m_2 \|v_2\|^2 + \frac{1}{2} \omega_2^T \mathcal{I}_2 \omega_2$$

$$\omega_1 = \begin{bmatrix} 0 \\ 0 \\ \dot{\theta}_1 \end{bmatrix} \qquad \omega_2 = \begin{bmatrix} 0 \\ 0 \\ \dot{\theta}_1 + \dot{\theta}_2 \end{bmatrix}$$

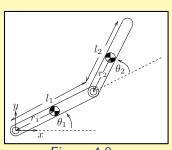


Figure 4.8

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 $P_i = \begin{bmatrix} \frac{x_i}{y_i} \\ 0 \end{bmatrix}$: Mass center

 γ_i : Distance from joint i to mass center

Change of Coordinates:

$$\begin{cases} \overline{x}_1 = r_1 c_1 \\ \overline{y}_1 = r_1 s_1 \end{cases} \Rightarrow \begin{cases} \dot{\overline{x}}_1 = -r_1 s_1 \dot{\theta}_1 \\ \dot{\overline{y}}_1 = r_1 c_1 \dot{\theta}_1 \end{cases}$$

$$\begin{cases} \overline{x}_2 = l_1 c_1 + r_2 c_{12} \\ \overline{y}_2 = l_1 s_1 + r_2 s_{12} \end{cases} \Rightarrow \begin{cases} \dot{\overline{x}}_2 = -(l_1 s_1 + r_2 s_{12}) \dot{\theta}_1 - r_2 s_{12} \dot{\theta}_2 \\ \dot{\overline{y}}_1 = (l_1 c_1 + r_2 c_{12}) \dot{\theta}_1 + r_2 c_{12} \dot{\theta}_2 \end{cases}$$

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4.2 Lagrange's Equations

Kinetic energy:

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with Constraints $T(\theta, \dot{\theta}) = \frac{1}{2} m_1 (\dot{x}_1^2 + \dot{y}_1^2) + \frac{1}{2} \mathcal{I}_{z_1} \dot{\theta}_1^2 + \frac{1}{2} m_2 (\dot{x}_2^2 + \dot{y}_2^2) + \frac{1}{2} \mathcal{I}_{z_2} (\dot{\theta}_1^2 + \dot{\theta}_2^2)$ $= \frac{1}{2} [\dot{\theta}_1 \ \dot{\theta}_2] \underbrace{\begin{bmatrix} \alpha + 2\beta c_2 & \delta + \beta c_2 \\ \delta + \beta c_2 & \delta \end{bmatrix}}_{M(\theta)} \underbrace{\begin{bmatrix} \dot{\theta}_1 \\ \dot{\theta}_2 \end{bmatrix}}_{(\phi)}$

$$\beta = m_2 l_1 r_2, \delta = \mathcal{I}_{z_2} + m_2 r_2^2, L = T$$

Equation of motion:

$$M(\theta) \begin{bmatrix} \ddot{\theta}_1 \\ \ddot{\theta}_2 \end{bmatrix} + \begin{bmatrix} -\beta s_2 \dot{\theta}_2 & -\beta s_2 (\dot{\theta}_1 + \dot{\theta}_2) \\ \beta s_2 \dot{\theta}_1 & 0 \end{bmatrix} \begin{bmatrix} \dot{\theta}_1 \\ \dot{\theta}_2 \end{bmatrix} = \begin{bmatrix} \tau_1 \\ \tau_2 \end{bmatrix}$$

 $\alpha = \mathcal{I}_{z_1} + \mathcal{I}_{z_2} + m_1 r_1^2 + m_2 (l_1^2 + r_2^2),$

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Dynamics of Open-chain Manipulators

□ Dynamics of an open-chain manipulator:

Definition:

 L_i : frame at mass center of link i, $g_{sl_i}(\theta) = e^{\hat{\xi}_1\theta_1} \cdots e^{\hat{\xi}_i\theta_i} g_{sl_i}(o)$

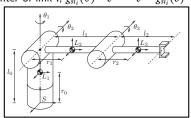


Figure 4.9

$$V_{sl_i}^b = J_{sl_i}^b(\theta)\dot{\theta} = \begin{bmatrix} \xi_1^{\dagger} & \xi_2^{\dagger} & \cdots & \xi_i^{\dagger} & 0 & \cdots & 0 \end{bmatrix} \begin{bmatrix} \dot{\theta}_1 \\ \vdots \\ \dot{\theta}_i \\ \dot{\theta}_{i+1} \\ \vdots \\ \dot{\theta} \end{bmatrix} = J_i(\theta)\dot{\theta}$$

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 $\xi_{j}^{\dagger} = \operatorname{Ad}^{-1}(e^{\hat{\xi}_{j+1}\theta_{j+1}} \cdots e^{\hat{\xi}_{i}\theta_{i}} g_{sl_{i}}(0)) \xi_{j}, \ j \leq i$ Chapter
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Dynamics $T_{i}(\theta, \dot{\theta}) = \frac{1}{2} (V_{sl_{i}}^{b})^{T} M_{i}^{b} V_{sl_{i}}^{b} = \frac{1}{2} \dot{\theta}^{T} J_{i}^{T}(\theta) M_{i}^{b} J_{i}(\theta) \dot{\theta}$ Introduction $T(\theta) = \sum_{i=1}^{n} T_{i}(\theta, \dot{\theta}) = \frac{1}{2} \dot{\theta}^{T} M(\theta) \dot{\theta},$ Lagrange's
Equations $M(\theta) = \sum_{i} J_{i}^{T}(\theta) M_{i}^{b} J_{i}(\theta) = \frac{1}{2} \sum_{i=1}^{n} M_{ij}(\theta) \dot{\theta}_{i} \dot{\theta}_{j}$ Dynamics of

 $h_i(\theta)$: Height of L_i , $V_i(\theta) = m_i g h_i(\theta)$, $V(\theta) = \sum_{i=1} m_i g h_i(\theta)$

Lagrange's Equation:

Open-chain Manipulators

$$\frac{\mathrm{d}}{\mathrm{d}t} \frac{\partial L}{\partial \dot{\theta}_i} - \frac{\partial L}{\partial \theta_i} = \tau_i, \quad i = 1, \dots, n,$$

$$\frac{\mathrm{d}}{\mathrm{d}t} \frac{\partial L}{\partial \dot{\theta}_i} = \frac{\mathrm{d}}{\mathrm{d}t} \left(\sum_{j=1}^n M_{ij} \dot{\theta}_j \right) = \sum_{j=1}^n M_{ij} \ddot{\theta}_j + \dot{M}_{ij} \dot{\theta}_j$$

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Lagrange's Equations with

$$\begin{split} \frac{\partial L}{\partial \theta_i} &= \frac{1}{2} \sum_{j,k=1}^n \frac{\partial M_{kj}}{\partial \theta_i} \dot{\theta}_k \dot{\theta}_j - \frac{\partial V}{\partial \theta_i}, \qquad \dot{M}_{ij} = \sum_k \frac{\partial M_{ij}}{\partial \theta_k} \dot{\theta}_k \\ \\ &\Rightarrow \sum_{j=1}^n M_{ij} \ddot{\theta}_j + \sum_{j,k=1}^n \left(\frac{\partial M_{ij}}{\partial \theta_k} \dot{\theta}_j \dot{\theta}_k - \frac{1}{2} \frac{\partial M_{kj}}{\partial \theta_i} \dot{\theta}_k \dot{\theta}_j \right) + \frac{\partial V}{\partial \theta_i} = \tau_i \\ \\ &\Rightarrow \sum_{j=1}^n M_{ij} \ddot{\theta}_j + \sum_{j,k=1}^n \Gamma_{ijk} \dot{\theta}_k \dot{\theta}_j + \frac{\partial V}{\partial \theta_i} = \tau_i \\ \\ &\Gamma_{ijk} \triangleq \frac{1}{2} \left(\frac{\partial M_{ij}}{\partial \theta_k} + \frac{\partial M_{ik}}{\partial \theta_j} - \frac{\partial M_{kj}}{\partial \theta_i} \right) \\ \\ &\dot{\theta}_i \cdot \dot{\theta}_j, i \neq j \colon \text{ Coriolis term,} \qquad \dot{\theta}_i^2 \colon \text{ Centrifugal term} \end{split}$$

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$$\begin{split} \frac{\partial L}{\partial \theta_i} &= \frac{1}{2} \sum_{j,k=1}^n \frac{\partial M_{kj}}{\partial \theta_i} \dot{\theta}_k \dot{\theta}_j - \frac{\partial V}{\partial \theta_i}, \qquad \dot{M}_{ij} = \sum_k \frac{\partial M_{ij}}{\partial \theta_k} \dot{\theta}_k \\ &\Rightarrow \sum_{j=1}^n M_{ij} \ddot{\theta}_j + \sum_{j,k=1}^n \left(\frac{\partial M_{ij}}{\partial \theta_k} \dot{\theta}_j \dot{\theta}_k - \frac{1}{2} \frac{\partial M_{kj}}{\partial \theta_i} \dot{\theta}_k \dot{\theta}_j \right) + \frac{\partial V}{\partial \theta_i} = \tau_i \\ &\Rightarrow \sum_{j=1}^n M_{ij} \ddot{\theta}_j + \sum_{j,k=1}^n \Gamma_{ijk} \dot{\theta}_k \dot{\theta}_j + \frac{\partial V}{\partial \theta_i} = \tau_i \\ &\Gamma_{ijk} \triangleq \frac{1}{2} \left(\frac{\partial M_{ij}}{\partial \theta_k} + \frac{\partial M_{ik}}{\partial \theta_j} - \frac{\partial M_{kj}}{\partial \theta_i} \right) \\ &\dot{\theta}_i \cdot \dot{\theta}_j, i \neq j \colon \text{ Coriolis term,} \qquad \dot{\theta}_i^2 \colon \text{ Centrifugal term} \end{split}$$

Define:
$$C_{ij}(\theta, \dot{\theta}) = \sum_{k=1}^{n} \Gamma_{ijk} \dot{\theta}_k = \frac{1}{2} \sum_{k=1}^{n} \left(\frac{\partial M_{ij}}{\partial \theta_k} + \frac{\partial M_{ik}}{\partial \theta_j} - \frac{\partial M_{kj}}{\partial \theta_i} \right) \dot{\theta}_k$$

$$\Rightarrow \boxed{M(\theta) \ddot{\theta} + C(\theta, \dot{\theta}) \dot{\theta} + N(\theta) = \tau}$$

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Property 1:

- $\dot{M} 2C \in \mathbb{R}^{n \times n}$ is skew symmetric

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Property 1:

$$\mathbf{I} M(\theta) = M^{T}(\theta), \dot{\theta}^{T} M(\theta) \dot{\theta} \ge 0, \dot{\theta}^{T} M(\theta) \dot{\theta} = 0 \Leftrightarrow \dot{\theta} = 0$$

 $\dot{M} - 2C \in \mathbb{R}^{n \times n}$ is skew symmetric

Proof:

$$\begin{split} (\dot{M} - 2C)_{ij} &= \dot{M}_{ij} - 2C_{ij}(\theta) \\ &= \sum_{k=1}^{n} \frac{\partial M_{ij}}{\partial \theta_{k}} \dot{\theta}_{k} - \frac{\partial M_{ij}}{\partial \theta_{k}} \dot{\theta}_{k} - \frac{\partial M_{ik}}{\partial \theta_{j}} \dot{\theta}_{k} + \frac{\partial M_{kj}}{\partial \theta_{i}} \dot{\theta}_{k} \\ &= \sum_{k=1}^{n} \frac{\partial M_{kj}}{\partial \theta_{i}} \dot{\theta}_{k} - \frac{\partial M_{ik}}{\partial \theta_{j}} \dot{\theta}_{k} \end{split}$$

Switching i and j shows $(\dot{M} - 2C)^T = -(\dot{M} - 2C)$

⋄ Example: Planar 2-DoF Robot (continued)

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$$\begin{split} m_{11}(\theta) &= \alpha + 2\beta \cos \theta_2, m_{22} = \delta \\ m_{12}(\theta) &= m_{21}(\theta) = \delta + \beta \cos \theta_2 \\ c_{11}(\theta, \dot{\theta}) &= -\beta \sin \theta_2 \cdot \dot{\theta}_2, c_{12}(\theta, \dot{\theta}) = -\beta \sin \theta_2 (\dot{\theta}_1 + \dot{\theta}_2) \\ c_{21}(\theta, \dot{\theta}) &= \beta \sin \theta_2 \cdot \dot{\theta}_1, c_{22}(\theta, \dot{\theta}) = 0 \end{split}$$

$$\begin{split} &\Gamma_{111} = \frac{1}{2} \left(\frac{\partial M_{11}}{\partial \theta_1} + \frac{\partial M_{11}}{\partial \theta_1} - \frac{\partial M_{11}}{\partial \theta_1} \right) = \frac{1}{2} \frac{\partial M_{11}}{\partial \theta_1} = 0 \\ &\Gamma_{112} = \frac{1}{2} \left(\frac{\partial M_{11}}{\partial \theta_2} + \frac{\partial M_{12}}{\partial \theta_1} - \frac{\partial M_{21}}{\partial \theta_1} \right) = \frac{1}{2} \frac{\partial M_{11}}{\partial \theta_2} = -\beta \sin \theta_2 \\ &\Gamma_{121} = \frac{1}{2} \left(\frac{\partial M_{12}}{\partial \theta_1} + \frac{\partial M_{11}}{\partial \theta_2} - \frac{\partial M_{12}}{\partial \theta_1} \right) = \frac{1}{2} \frac{\partial M_{11}}{\partial \theta_2} = -\beta \sin \theta_2 \\ &\Gamma_{122} = \frac{1}{2} \left(\frac{\partial M_{12}}{\partial \theta_2} + \frac{\partial M_{12}}{\partial \theta_2} - \frac{\partial M_{22}}{\partial \theta_1} \right) = \frac{\partial M_{12}}{\partial \theta_2} - \frac{1}{2} \frac{\partial M_{22}}{\partial \theta_1} = -\beta \sin \theta_2 \end{split}$$

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$$\begin{split} &\Gamma_{211} = \frac{1}{2} \left(\frac{\partial M_{21}}{\partial \theta_1} + \frac{\partial M_{21}}{\partial \theta_1} - \frac{\partial M_{11}}{\partial \theta_2} \right) = \frac{\partial M_{21}}{\partial \theta_1} - \frac{1}{2} \frac{\partial M_{11}}{\partial \theta_2} = \beta \sin \theta_2 \\ &\Gamma_{212} = \frac{1}{2} \left(\frac{\partial M_{21}}{\partial \theta_2} + \frac{\partial M_{22}}{\partial \theta_1} - \frac{\partial M_{21}}{\partial \theta_2} \right) = \frac{1}{2} \frac{\partial M_{22}}{\partial \theta_1} = 0 \\ &\Gamma_{221} = \frac{1}{2} \left(\frac{\partial M_{22}}{\partial \theta_1} + \frac{\partial M_{21}}{\partial \theta_2} - \frac{\partial M_{12}}{\partial \theta_2} \right) = \frac{1}{2} \frac{\partial M_{22}}{\partial \theta_1} = 0 \\ &\Gamma_{222} = \frac{1}{2} \left(\frac{\partial M_{22}}{\partial \theta_2} + \frac{\partial M_{22}}{\partial \theta_2} - \frac{\partial M_{22}}{\partial \theta_2} \right) = \frac{1}{2} \frac{\partial M_{22}}{\partial \theta_2} = 0 \\ &\dot{M} - 2C = \begin{bmatrix} -2\beta \sin \theta_2 \cdot \dot{\theta}_2 & -\beta \sin \theta_2 \cdot \dot{\theta}_2 \\ -\beta \sin \theta_2 \cdot \dot{\theta}_2 & 0 \end{bmatrix} \\ &- \begin{bmatrix} -2\beta \sin \theta_2 \cdot \dot{\theta}_2 & -2\beta \sin \theta_2 (\dot{\theta}_1 + \dot{\theta}_2) \\ 2\beta \sin \theta_2 \cdot \dot{\theta}_1 & 0 \end{bmatrix} \\ &= \begin{bmatrix} 0 & \beta \sin \theta_2 (2\dot{\theta}_1 + \dot{\theta}_2) \\ -\beta \sin \theta_2 (2\dot{\theta}_1 + \dot{\theta}_2) \end{bmatrix} \\ &\in \text{skew-symmetric} \end{split}$$

⋄ Example: Dynamics of a 3-dof robot

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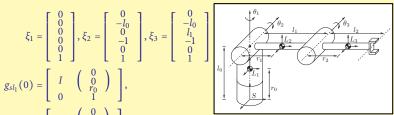


Figure 4.9

$$g_{sl_2}(0) = \begin{bmatrix} I & \begin{pmatrix} 0 \\ r_1 \\ l_0 \end{bmatrix},$$

$$g_{sl_3}(0) = \left[\begin{array}{cc} I & \begin{pmatrix} 0 \\ l_1 + r_2 \\ 0 & 1 \end{array} \right) \right]$$

$$M_i = \begin{bmatrix} \begin{array}{c|ccc} m_i & 0 & 0 & & & & \\ 0 & m_i & 0 & & & & \\ 0 & 0 & m_i & & & & \\ \hline & & & & \mathcal{I}_{x_i} & 0 & 0 \\ & & & & 0 & \mathcal{I}_{y_i} & 0 \\ & & & & 0 & 0 & \mathcal{I}_{z_i} \end{array}$$

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 m_i : The mass of the object

 \mathcal{I}_{x_i} : The moment of inertia about the x axis

$$\Gamma_{112} = (\mathcal{I}_{y_2} - \mathcal{I}_{z_2} - m_2 r_1^2) c_2 s_2 + (\mathcal{I}_{y_2} - \mathcal{I}_{z_3}) c_{23} s_{23} - t (l_1 s_2 + r_2 s_{23})$$

$$\Gamma_{113} = (\mathcal{I}_{y_3} - \mathcal{I}_{z_3})c_{23}s_{23} - tr_2s_{23}$$

$$\Gamma_{121} = (\mathcal{I}_{y_2} - \mathcal{I}_{z_2} - m_2 r_1^2) c_2 s_2 + (\mathcal{I}_{y_3} - \mathcal{I}_{z_3}) c_{23} s_{23} - t (l_1 s_2 + r_2 s_{23})$$

$$\Gamma_{131} = (\mathcal{I}_{y_3} - \mathcal{I}_{z_3})c_{23}s_{23} - tr_2s_{23}$$

$$\Gamma_{211} = (\mathcal{I}_{z_2} - \mathcal{I}_{y_2} + m_2 r_1^2) c_2 s_2 + (\mathcal{I}_{z_3} - \mathcal{I}_{y_3}) c_{23} s_{23} + t (l_1 s_2 + r_2 s_{23})$$

$$\Gamma_{223} = -l_1 m_3 r_2 s_3$$

$$\Gamma_{232} = -l_1 m_3 r_2 s_3$$

$$\Gamma_{233} = -l_1 m_3 r_2 s_3$$

$$\Gamma_{311} = (\mathcal{I}_{z_3} - \mathcal{I}_{y_3})c_{23}s_{23} + tr_2s_{23}$$

$$\Gamma_{322}=l_1m_3r_2s_3$$

where $t = m_3(l_1c_2 + r_2c_{23})$.

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$$\begin{split} M(\theta) &= \begin{bmatrix} M_{11} & M_{12} & M_{13} \\ M_{21} & M_{22} & M_{23} \\ M_{31} & M_{32} & M_{33} \end{bmatrix} = J_1^T M_1 J_1 + J_2^T M_2 J_2 + J_3^T M_3 J_3 \\ M_{11} &= I_{y2} s_2^2 + I_{y3} s_{23}^2 + I_{z1} + I_{z2} c_2^2 + I_{z3} c_{23}^2 + m_2 r_1^2 c_2^2 + m_3 \left(l_1 c_2 + r_2 c_{23} \right)^2 \\ M_{12} &= M_{13} = M_{21} = M_{31} = 0 \\ M_{22} &= I_{x2} + I_{x3} + m_2 l_1^2 + M_2 r_1^2 + m_3 r_2^2 + 2 m_3 l_1 r_2 c_3 \\ M_{23} &= I_{x3} + m_3 r_2^2 + m_3 l_1 r_2 c_3 \\ M_{32} &= I_{x3} + m_3 r_2^2 + m_3 l_1 r_2 c_3 \\ M_{33} &= I_{x_3} + m_3 r_2^2 \end{bmatrix} \\ C_{ij}(\theta, \dot{\theta}) &= \sum_{k=1}^{n} \Gamma_{ijk} \dot{\theta}_k = \frac{1}{2} \sum_{k=1}^{n} \left(\frac{\partial M_{ij}}{\partial \theta_k} + \frac{\partial M_{ik}}{\partial \theta_j} - \frac{\partial M_{kj}}{\partial \theta_i} \right) \dot{\theta}_k \\ & \diamond \end{split}$$

☐ Additional Properties of the dynamics in terms of POE:

Define:

$$A_{ij} = \begin{cases} A \mathbf{d}_{e^{\hat{\xi}_{j+1}\theta_{j+1}} \dots e^{\hat{\xi}_{i}\theta_{i}}}^{-1} & i > j \\ I & i = j \\ 0 & i < j \end{cases}$$

$$J_i(\theta) = \operatorname{Ad}_{g_{sl_i}^{-1}(0)} \big[A_{i1} \xi_1 \cdots A_{ii} \xi_i \ 0 \cdots 0 \big]$$

$$M_i' = \operatorname{Ad}_{g_{sl_i(0)}}^T M_i \operatorname{Ad}_{g_{sl_i(0)}^{-1}}$$
 (intertia of i^{th} link in S)

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☐ Additional Properties of the dynamics in terms of POE:

Define:

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$$\begin{split} A_{ij} &= \begin{cases} \mathrm{Ad}_{e^{\hat{\xi}_{j+1}\theta_{j+1}} \dots e^{\hat{\xi}_{i}\theta_{i}} & i > j \\ I & i = j \\ 0 & i < j \end{cases} \\ J_{i}(\theta) &= \mathrm{Ad}_{g_{sl_{i}}^{-1}(0)} [A_{i1}\xi_{1} \dots A_{ii}\xi_{i} \ 0 \dots 0] \\ M'_{i} &= \mathrm{Ad}_{g_{sl_{i}}^{-1}(0)}^{T} M_{i} \mathrm{Ad}_{g_{sl_{i}}^{-1}(0)} \ (\ \text{intertia of} \ i^{th} \ \text{link in} \ S) \end{cases} \end{split}$$

Property 2:

$$M_{ij}(\theta) = \sum_{l=\max(i,j)}^{n} \xi_{i}^{T} A_{li}^{T} M_{l}' A_{lj} \xi_{j}, C_{ij}(\theta, \dot{\theta}) = \frac{1}{2} \sum_{k=1}^{n} \left(\frac{\partial M_{ij}}{\partial \theta_{k}} + \frac{\partial M_{ik}}{\partial \theta_{j}} - \frac{\partial M_{kj}}{\partial \theta_{i}} \right) \dot{\theta}_{k}$$

where

$$\frac{\partial M_{ij}}{\partial \theta_k} = \sum_{l=\max(i,j)}^{n} \left(\left[A_{k-1,i} \xi_i, \xi_k \right]^T A_{lk}^T M_l A_{lj} \xi_j + \xi_i^T A_{li}^T M_l A_{lk} \left[A_{k-1,j} \xi_j, \xi_k \right] \right)$$

† End of Section †

Section 4.4

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- 3 Dynamics of Open-chain Manipulators
- 4 Coordinate Invariant Algorithms
- 5 Lagrange's Equations with Constraints
 - General formulation
 - Unified geometric approach (Liu G.F., Li Z.X.)
 - Gauge invariant formulation (Aghili, F.)
- 6 References

Newton-Euler equations in spatial frame:

Newton's Equation:

$$f^s = \frac{d}{dt}(m\dot{p}) = m\ddot{p}$$

Spatial angular momentum:

$$\mathcal{I}^{s} \cdot \omega^{s} = R(\mathcal{I} \cdot \omega^{b}) = \underbrace{R \cdot \mathcal{I} \cdot R^{T}}_{\mathcal{I}^{s}} \cdot \omega^{s}$$

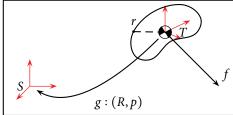


Figure 4.10 (Continues next slide)

Coordinate Invariant Algorithms

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$$\tau^{s} = \frac{\mathrm{d}}{\mathrm{d}t} (\mathcal{I}^{s} \omega^{s}) = \frac{\mathrm{d}}{\mathrm{d}t} (R \mathcal{I} R^{T} \omega^{s}) = \mathcal{I}^{s} \dot{\omega}^{s} + \dot{R} \mathcal{I} R^{T} \omega^{s} + R \mathcal{I} \dot{R}^{T} \omega^{s}$$
$$= \mathcal{I}^{s} \dot{\omega}^{s} + \underbrace{\dot{R} R^{T}}_{\hat{\omega}^{s}} \mathcal{I}^{s} \omega^{s} - R \mathcal{I} R^{T} \hat{\omega}^{s} \omega^{s} = \mathcal{I}^{s} \dot{\omega}^{s} + \omega^{s} \times (\mathcal{I}^{s} \omega^{s})$$

□ Transform equations to body frame:

$$\frac{\mathrm{d}}{\mathrm{d}t} (m\dot{p}) = \frac{\mathrm{d}}{\mathrm{d}t} (mRv^b) = m\dot{R}v^b + mR\dot{v}^b, R^T f^s = mR^T \dot{R}v^b + m\dot{v}^b$$

$$\Rightarrow f^b = m\omega^b \times v^b + m\dot{v}^b,$$

$$\tau^b = R^T \tau^s = R^T \frac{\mathrm{d}}{\mathrm{d}t} (R\mathcal{I}\omega^b) = \mathcal{I}\dot{\omega}^b + \omega^b \times \mathcal{I}\omega^b$$

$$\Rightarrow \underbrace{\begin{bmatrix} mI & 0 \\ 0 & \mathcal{I} \end{bmatrix}}_{\dot{\omega}b} \underbrace{\begin{bmatrix} \dot{v}^b \\ \dot{\omega}^b \end{bmatrix}}_{\dot{\omega}b} + \underbrace{\begin{bmatrix} \omega^b \times mv^b \\ \omega^b \times \mathcal{I}\omega^b \end{bmatrix}}_{b} = \begin{bmatrix} f^b \\ \tau^b \end{bmatrix} = F^b$$
(*)

Define:

$$[,]: se(3) \times se(3) \mapsto se(3), [\hat{\xi}_1, \hat{\xi}_2] \triangleq \hat{\xi}_1 \hat{\xi}_2 - \hat{\xi}_2 \hat{\xi}_1$$

if

$$\begin{bmatrix} \hat{\xi}_1, \hat{\xi}_2 \end{bmatrix} = \begin{bmatrix} \hat{\omega}_i & v_i \\ 0 & 0 \end{bmatrix}, i = 1, 2$$

then

$$\hat{\xi} = \left[\begin{array}{cc} (\omega_1 \times \omega_2)^{\wedge} & \hat{\omega}_1 v_2 - \hat{\omega}_2 v_1 \\ 0 & 0 \end{array} \right] = ad_{\xi_i} \cdot \xi_2$$

where

$$ad_{\xi_1} = \left[\begin{array}{cc} \hat{\omega}_1 & \hat{\nu}_1 \\ 0 & \hat{\omega}_1 \end{array} \right]$$

It is straightforward computation to see that:

$$(*) \Leftrightarrow M^b \dot{V}^b - \operatorname{ad}_{V^b}^T M^b V^b = F^b$$

Property 3:

$$\mathrm{Ad}_g\big[\hat{\xi}_1,\hat{\xi}_2\big] = \big[\mathrm{Ad}_g\hat{\xi}_1,\mathrm{Ad}_g\hat{\xi}_2\big] \Rightarrow \mathrm{Ad}_g\mathrm{ad}_{\hat{\xi}_1} = \mathrm{ad}_{\mathrm{Ad}_g\hat{\xi}_1}\mathrm{Ad}_g$$

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☐ Coordinate invariance of Newton-Euler equations:

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$$M_{1}\dot{V}_{1} - \operatorname{ad}_{V_{1}}^{T}M_{1}V_{1} = F_{1}$$

$$\begin{cases} V_{1} = \operatorname{Ad}_{g_{0}}V_{2}, \\ F_{1} = \operatorname{Ad}_{g_{0}^{-1}}^{T}F_{2} \end{cases} \Rightarrow$$

$$F_{1} = (\operatorname{Ad}_{g_{0}}^{T})^{-1}F_{2} = (\operatorname{Ad}_{g_{0}}^{-1})^{T}F_{2}$$

$$M_{1} = \operatorname{Ad}_{g_{0}}^{-T}M_{2}\operatorname{Ad}_{g_{0}}^{-1}$$

$$\Rightarrow \operatorname{Ad}_{g_{0}}^{-T}M_{2}\operatorname{Ad}_{g_{0}}^{-1}\operatorname{Ad}_{g_{0}}V_{2} - \operatorname{ad}_{(\operatorname{Ad}_{g_{0}},V_{2})}^{T}\operatorname{Ad}_{g_{0}}^{-T}M_{2}\operatorname{Ad}_{g_{0}}^{-1}\operatorname{Ad}_{g_{0}}V_{1} = \operatorname{Ad}_{g_{0}}^{-T}F_{2}$$

Since $ad_{Ad_{g_0}V} = Ad_{g_0}ad_VAd_{g_0}^{-1}$ by Property 3, we have

$$M_2 \dot{V}_2 - \text{ad}_{V_2}^T M_2 V_2 = F_2$$

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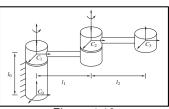


Figure 4.12

 C_i : Frame fixed to link i, located along the i^{th} axis

 F_i : Generalized force link i-1 exerting on link i, expressed in C_i

 τ_i : Joint torque of link i

 $g_{i-1,i}$: Transformation of C_i relative to C_{i-1}

$$g_{i-1,i}(\theta_i) = e^{\hat{\xi}_i'\theta_i} \cdot g_{i-1,i}(0) = g_{i-1,i}(0)e^{\hat{\xi}_i\theta_i}$$

 $\xi_i = \operatorname{Ad}_{g_{i-1,i}^{-1}(0)} \cdot \xi_i'$: i^{th} axis in C_i frame.

$$\xi_i = \begin{cases} \begin{bmatrix} 0 \\ z_i \end{bmatrix} : & \text{Revolute joint.} \\ \begin{bmatrix} z_i \\ 0 \end{bmatrix} : & \text{Prismatic joint.} \end{cases}$$

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$$\Rightarrow g_{i-1,i}^{-1} \cdot \dot{g}_{i-1,i} = \hat{\xi}_i \cdot \dot{\theta}_i$$

$$M_i$$
: Moment of inertia in C_i

$$M_i = \begin{bmatrix} m_i I & -m_i \hat{r}_i \\ m_i \hat{r}_i & \mathcal{I}_i - m_i \hat{r}_i^2 \end{bmatrix} \qquad \begin{array}{c} m_i : & \text{Mass of link } i \\ \mathcal{I}_i : & \text{inertia tensor} \end{array}$$

$$g_i = g_{i-1}g_{i-1,i}$$

$$\hat{V}_i = g_i^{-1} \cdot \dot{g}_i = g_{i-1,i}^{-1} \hat{V}_{i-1} g_{i-1,i} + \hat{\xi}_i \dot{\theta}_i$$

$$V_i = Ad_{g_{i-1,i}^{-1}}V_{i-1} + \xi_i\dot{\theta}_i$$

$$\dot{\hat{V}}_{i} = \dot{g}_{i-1,i}^{-1} \hat{V}_{i-1} g_{i-1,i} + g_{i-1,i}^{-1} \hat{V}_{i-1} \dot{g}_{i-1,i} + g_{i-1,i}^{-1} \dot{\hat{V}}_{i-1} g_{i-1,i} + \hat{\xi}_{i} \ddot{\theta}_{i}
= -\dot{g}_{i-1,i}^{-1} g_{i-1,i} g_{i-1,i}^{-1} \hat{V}_{i-1} g_{i-1,i} + g_{i-1,i}^{-1} \hat{V}_{i-1} g_{i-1,i} g_{i-1,i}^{-1} \dot{g}_{i-1,i}
+ g_{i-1,i}^{-1} \dot{\hat{V}}_{i-1} g_{i-1,i} + \hat{\xi}_{i} \ddot{\theta}_{i}$$

$$\begin{split} &= -\hat{\xi}_{i}\dot{\theta}_{i} \left(\mathrm{Ad}_{g_{i-1,i}^{-1}} V_{i-1} \right)^{\wedge} + \left(\mathrm{Ad}_{g_{i-1,i}^{-1}} V_{i-1} \right)^{\wedge} \hat{\xi}_{i}\dot{\theta}_{i} + \left(\mathrm{Ad}_{g_{i-1,i}^{-1}} \dot{V}_{i-1} \right)^{\wedge} + \hat{\xi}_{i}\ddot{\theta}_{i} \\ &\Rightarrow \dot{V}_{i} = \xi_{i}\ddot{\theta}_{i} + \mathrm{Ad}_{g_{i-1,i}^{-1}} \dot{V}_{i-1} - \mathrm{ad}_{\xi_{i}\dot{\theta}_{i}} \left(\mathrm{Ad}_{g_{i-1,i}^{-1}} V_{i-1} \right) \end{split}$$

□ Forward Recursion (kinematics):

$$\begin{cases} \text{init.}: V_0 = 0, \dot{V}_0 = \left[\begin{array}{c} g \\ 0 \end{array} \right] \text{(gravity vector)} \\ \\ g_{i-1,i} = g_{i-1,i}(0)e^{\dot{\xi}_i\theta_i} \\ \\ V_i = \operatorname{Ad}_{g_{i-1,i}^{-1}}V_{i-1} + \xi_i\dot{\theta}_i \\ \\ \dot{V}_i = \xi_i\ddot{\theta}_i + \operatorname{Ad}_{g_{i-1,i}^{-1}}\dot{V}_{i-1} - \operatorname{ad}_{\xi_i\dot{\theta}_i} (\operatorname{Ad}_{g_{i-1,i}^{-1}}V_{i-1}) \end{cases}$$

□ Backward Recursion (inverse dynamics):

 F_{n+1} : End-effector wrench, $g_{n,n+1}$: transform from tool frame to C_n

$$F_{i} = \operatorname{Ad}_{g_{i,i+1}}^{T} \cdot F_{i+1} + M_{i}\dot{V}_{i} - \operatorname{ad}_{V_{i}}^{T} \cdot M_{i}V_{i}$$

$$\tau_{i} = \xi_{i}^{T} \cdot F_{i}$$

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$$\begin{aligned} V_1 &= \mathrm{Ad}_{g_{0,1}^{-1}} \cdot V_0 + \xi_1 \dot{\theta}_1 \\ F_n &= \mathrm{Ad}_{g_{n,n+1}^{-1}}^T \cdot F_{n+1} + M_n \dot{V}_n - \mathrm{ad}_{V_n}^T \cdot (M_n V_n) \end{aligned}$$

Define:

$$\begin{split} V &= \left[\begin{array}{c} V_1 \\ \vdots \\ V_n \end{array} \right] \in \mathbb{R}^{6n \times 1}, \dot{\theta} = \left[\begin{array}{c} \dot{\theta}_1 \\ \vdots \\ \dot{\theta}_n \end{array} \right] \in \mathbb{R}^n, \xi = \left[\begin{array}{c} \xi_1 & 0 & 0 \\ 0 & \ddots & 0 \\ 0 & 0 & \xi_n \end{array} \right] \in \mathbb{R}^{6n \times n} \\ F &= \left[\begin{array}{c} F_1 \\ \vdots \\ F_n \end{array} \right] \in \mathbb{R}^{6n \times 1}, \tau = \left[\begin{array}{c} \tau_1 \\ \vdots \\ \dot{\tau}_n \end{array} \right] \in \mathbb{R}^n, P_0 = \left[\begin{array}{c} \operatorname{Ad}_{g_{0,1}^{-1}} \\ 0 \\ \vdots \\ \dot{0} \end{array} \right] \in \mathbb{R}^{6n \times 6} \\ P_t &= \left[0 \cdots 0 \quad \operatorname{Ad}_{g_{n+1}^{-1}} \right] \in \mathbb{R}^{6 \times 6n} \end{split}$$

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$$\begin{split} V_1 &= \mathrm{Ad}_{g_{0,1}^{-1}} \cdot V_0 + \xi_1 \dot{\theta}_1 \\ V_2 &- \mathrm{Ad}_{g_{1,2}^{-1}} V_1 = \xi_2 \dot{\theta}_2 \\ &\vdots \\ V_n &- \mathrm{Ad}_{g_{n-1,n}^{-1}} V_{n-1} = \xi_n \dot{\theta}_n \\ &\Rightarrow \underbrace{ \begin{bmatrix} I & 0 & \cdots & 0 \\ -\mathrm{Ad}_{g_{1,2}^{-1}} & I & \ddots & \vdots \\ 0 & \ddots & \ddots & 0 \\ 0 & 0 & -\mathrm{Ad}_{g_{n-1,n}^{-1}} & I \end{bmatrix} \underbrace{ \begin{bmatrix} V_1 \\ V_2 \\ \vdots \\ V_n \end{bmatrix}}_{V} \\ &= \underbrace{ \begin{bmatrix} \mathrm{Ad}_{g_{0,1}^{-1}} \\ 0 \\ \vdots \\ 0 \end{bmatrix}}_{P_0} V_0 + \underbrace{ \begin{bmatrix} \xi_1 & \xi_2 & \ddots & \vdots \\ \xi_n & \vdots \\ \vdots \\ \vdots \\ \theta_n \end{bmatrix}}_{\xi} \underbrace{ \begin{bmatrix} \dot{\theta}_1 \\ \dot{\theta}_2 \\ \vdots \\ \dot{\theta}_n \end{bmatrix}}_{\dot{\theta}} \end{split}$$

Thus
$$V = GP_0V_0 + G\xi\dot{\theta}$$

where

$$G = \left[\begin{array}{ccccc} I & 0 & 0 & \cdots & 0 \\ \mathrm{Ad}_{g_{1,2}^{-1}} & I & 0 & \cdots & 0 \\ \mathrm{Ad}_{g_{1,3}^{-1}} & \mathrm{Ad}_{g_{2,3}^{-1}} & I & \ddots & \vdots \\ \vdots & \vdots & \ddots & I & 0 \\ \mathrm{Ad}_{g_{1,n}^{-1}} & \mathrm{Ad}_{g_{2,n}^{-1}} & \cdots & \mathrm{Ad}_{g_{n-1,n}^{-1}} & I \end{array} \right] \in \mathbb{R}^{6n \times 6n}$$

$$\begin{split} \dot{V}_1 &= \xi_1 \ddot{\theta}_1 + \mathrm{Ad}_{g_{0,1}^{-1}} \dot{V}_0 - \mathrm{ad}_{\xi_1 \dot{\theta}_1} \big(\mathrm{Ad}_{g_{0,1}^{-1}} V_0 \big) \\ \dot{V}_2 - \mathrm{Ad}_{g_{1,2}^{-1}} \dot{V}_1 &= \xi_2 \ddot{\theta}_2 - \mathrm{ad}_{\xi_2 \dot{\theta}_2} \big(\mathrm{Ad}_{g_{1,2}^{-1}} V_1 \big) \\ \dot{V}_n - \mathrm{Ad}_{g_{n-1,n}^{-1}} \dot{V}_{n-1} &= \xi_n \ddot{\theta}_n - \mathrm{ad}_{\xi_n \dot{\theta}_n} \big(\mathrm{Ad}_{g_{n-1,n}^{-1}} V_{n-1} \big) \end{split}$$

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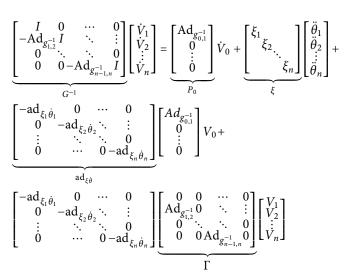
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Thus

$$\dot{V} = G \cdot \xi \ddot{\theta} + G \cdot P_0 \dot{V}_0 + G \cdot \operatorname{ad}_{\xi \dot{\theta}} P_0 V_0 + G \cdot \operatorname{ad}_{\xi \dot{\theta}} \Gamma V$$

Finally the backward recursion:

$$\begin{split} F_n &= \operatorname{Ad}_{g_{n,n+1}^{-1}}^T F_{n+1} + M_n \dot{V}_n - \operatorname{ad}_{V_n}^T \cdot \left(M_n V_n \right) \\ F_{n-1} &= \operatorname{Ad}_{g_{n-1,n}^{-1}}^T F_n + M_{n-1} \dot{V}_{n-1} - \operatorname{ad}_{V_{n-1}}^T \cdot \left(M_{n-1} V_{n-1} \right) \\ &\vdots \\ F_1 &= \operatorname{Ad}_{g_{1,2}^{-1}}^T F_2 + M_1 \dot{V}_1 - \operatorname{ad}_{V_1}^T \left(M_1 V_1 \right) \end{split}$$

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$$\Rightarrow \begin{bmatrix} I - \operatorname{Ad}_{g_{1,2}^{-1}}^{T} 0 & 0 \\ 0 & I & \ddots & 0 \\ \vdots & \ddots & \ddots - \operatorname{Ad}_{g_{n-1,n}^{-1}}^{T} \end{bmatrix} \begin{bmatrix} F_1 \\ F_2 \\ \vdots \\ F_n \end{bmatrix} = \underbrace{\begin{bmatrix} M_1 \\ M_2 \\ \vdots \\ M_n \end{bmatrix}}_{M} \begin{bmatrix} \dot{V}_1 \\ \dot{V}_2 \\ \vdots \\ \dot{V}_n \end{bmatrix} + \underbrace{\begin{bmatrix} 0 \\ \vdots \\ 0 \\ \operatorname{Ad}_{g_{n,n+1}^{-1}}^{T} \end{bmatrix}}_{P_t^T} F_{n+1} + \underbrace{\begin{bmatrix} 0 \\ \vdots \\ 0 \\ \operatorname{Ad}_{g_{n,n+1}^{-1}}^{T} \end{bmatrix}}_{P_t^T} F_{n+1} + \underbrace{\begin{bmatrix} 0 \\ \vdots \\ 0 \\ \operatorname{Ad}_{g_{n,n+1}^{-1}}^{T} \end{bmatrix}}_{P_t^T} F_{n+1} + \underbrace{\begin{bmatrix} 0 \\ \vdots \\ 0 \\ \operatorname{Ad}_{g_{n,n+1}^{-1}}^{T} \end{bmatrix}}_{P_t^T} F_{n+1} + \underbrace{\begin{bmatrix} 0 \\ \vdots \\ 0 \\ \operatorname{Ad}_{g_{n,n+1}^{-1}}^{T} \end{bmatrix}}_{P_t^T} F_{n+1} + \underbrace{\begin{bmatrix} 0 \\ \vdots \\ 0 \\ \operatorname{Ad}_{g_{n,n+1}^{-1}}^{T} \end{bmatrix}}_{P_t^T} F_{n+1} + \underbrace{\begin{bmatrix} 0 \\ \vdots \\ 0 \\ \operatorname{Ad}_{g_{n,n+1}^{-1}}^{T} \end{bmatrix}}_{P_t^T} F_{n+1} + \underbrace{\begin{bmatrix} 0 \\ \vdots \\ 0 \\ \operatorname{Ad}_{g_{n,n+1}^{-1}}^{T} \end{bmatrix}}_{P_t^T} F_{n+1} + \underbrace{\begin{bmatrix} 0 \\ \vdots \\ 0 \\ \operatorname{Ad}_{g_{n,n+1}^{-1}}^{T} \end{bmatrix}}_{P_t^T} F_{n+1} + \underbrace{\begin{bmatrix} 0 \\ \vdots \\ 0 \\ \operatorname{Ad}_{g_{n,n+1}^{-1}}^{T} \end{bmatrix}}_{P_t^T} F_{n+1} + \underbrace{\begin{bmatrix} 0 \\ \vdots \\ 0 \\ \operatorname{Ad}_{g_{n,n+1}^{-1}}^{T} \end{bmatrix}}_{P_t^T} F_{n+1} + \underbrace{\begin{bmatrix} 0 \\ \vdots \\ 0 \\ \operatorname{Ad}_{g_{n,n+1}^{-1}}^{T} \end{bmatrix}}_{P_t^T} F_{n+1} + \underbrace{\begin{bmatrix} 0 \\ \vdots \\ 0 \\ \operatorname{Ad}_{g_{n,n+1}^{-1}}^{T} \end{bmatrix}}_{P_t^T} F_{n+1} + \underbrace{\begin{bmatrix} 0 \\ \vdots \\ 0 \\ \operatorname{Ad}_{g_{n,n+1}^{-1}}^{T} \end{bmatrix}}_{P_t^T} F_{n+1} + \underbrace{\begin{bmatrix} 0 \\ \vdots \\ 0 \\ \operatorname{Ad}_{g_{n,n+1}^{-1}}^{T} \end{bmatrix}}_{P_t^T} F_{n+1} + \underbrace{\begin{bmatrix} 0 \\ \vdots \\ 0 \\ \operatorname{Ad}_{g_{n,n+1}^{-1}}^{T} \end{bmatrix}}_{P_t^T} F_{n+1} + \underbrace{\begin{bmatrix} 0 \\ \vdots \\ 0 \\ \operatorname{Ad}_{g_{n,n+1}^{-1}}^{T} \end{bmatrix}}_{P_t^T} F_{n+1} + \underbrace{\begin{bmatrix} 0 \\ \vdots \\ 0 \\ \vdots \\ 0 \\ \operatorname{Ad}_{g_{n,n+1}^{-1}}^{T} \end{bmatrix}}_{P_t^T} F_{n+1} + \underbrace{\begin{bmatrix} 0 \\ \vdots \\ 0 \\ \vdots \\ 0 \\ \vdots \\ 0 \\ \operatorname{Ad}_{g_{n,n+1}^{-1}}^{T} \end{bmatrix}}_{P_t^T} F_{n+1} + \underbrace{\begin{bmatrix} 0 \\ \vdots \\ 0$$

$$\underbrace{\begin{bmatrix} -\operatorname{ad}_{V_1}^T \\ -\operatorname{ad}_{V_n}^T \end{bmatrix}}_{\operatorname{ad}_{V}^T} \begin{bmatrix} M_1 \\ \cdot \\ M_n \end{bmatrix} \begin{bmatrix} V_1 \\ \vdots \\ V_n \end{bmatrix}$$

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$$\Rightarrow \begin{bmatrix} I - \operatorname{Ad}_{g_{1,2}^{-1}}^{T} 0 & 0 \\ 0 & I & \ddots & 0 \\ \vdots & \ddots & \ddots - \operatorname{Ad}_{g_{n-1,n}^{-1}}^{T} \end{bmatrix} \begin{bmatrix} F_1 \\ F_2 \\ \vdots \\ F_n \end{bmatrix} = \underbrace{\begin{bmatrix} M_1 \\ M_2 \\ \vdots \\ M_n \end{bmatrix}}_{M} \begin{bmatrix} \dot{V}_1 \\ \dot{V}_2 \\ \vdots \\ \dot{V}_n \end{bmatrix} + \underbrace{\begin{bmatrix} 0 \\ \vdots \\ 0 \\ \operatorname{Ad}_{g_{n,n+1}^{-1}}^{T} \end{bmatrix}}_{P_t^T} F_{n+1} + \underbrace{\begin{bmatrix} 0 \\ \vdots \\ 0 \\ \operatorname{Ad}_{g_{n,n+1}^{-1}}^{T} \end{bmatrix}}_{P_t^T} F_{n+1} + \underbrace{\begin{bmatrix} 0 \\ \vdots \\ 0 \\ \operatorname{Ad}_{g_{n,n+1}^{-1}}^{T} \end{bmatrix}}_{P_t^T} F_{n+1} + \underbrace{\begin{bmatrix} 0 \\ \vdots \\ 0 \\ \operatorname{Ad}_{g_{n,n+1}^{-1}}^{T} \end{bmatrix}}_{P_t^T} F_{n+1} + \underbrace{\begin{bmatrix} 0 \\ \vdots \\ 0 \\ \operatorname{Ad}_{g_{n,n+1}^{-1}}^{T} \end{bmatrix}}_{P_t^T} F_{n+1} + \underbrace{\begin{bmatrix} 0 \\ \vdots \\ 0 \\ \operatorname{Ad}_{g_{n,n+1}^{-1}}^{T} \end{bmatrix}}_{P_t^T} F_{n+1} + \underbrace{\begin{bmatrix} 0 \\ \vdots \\ 0 \\ \operatorname{Ad}_{g_{n,n+1}^{-1}}^{T} \end{bmatrix}}_{P_t^T} F_{n+1} + \underbrace{\begin{bmatrix} 0 \\ \vdots \\ 0 \\ \operatorname{Ad}_{g_{n,n+1}^{-1}}^{T} \end{bmatrix}}_{P_t^T} F_{n+1} + \underbrace{\begin{bmatrix} 0 \\ \vdots \\ 0 \\ \operatorname{Ad}_{g_{n,n+1}^{-1}}^{T} \end{bmatrix}}_{P_t^T} F_{n+1} + \underbrace{\begin{bmatrix} 0 \\ \vdots \\ 0 \\ \operatorname{Ad}_{g_{n,n+1}^{-1}}^{T} \end{bmatrix}}_{P_t^T} F_{n+1} + \underbrace{\begin{bmatrix} 0 \\ \vdots \\ 0 \\ \operatorname{Ad}_{g_{n,n+1}^{-1}}^{T} \end{bmatrix}}_{P_t^T} F_{n+1} + \underbrace{\begin{bmatrix} 0 \\ \vdots \\ 0 \\ \operatorname{Ad}_{g_{n,n+1}^{-1}}^{T} \end{bmatrix}}_{P_t^T} F_{n+1} + \underbrace{\begin{bmatrix} 0 \\ \vdots \\ 0 \\ \operatorname{Ad}_{g_{n,n+1}^{-1}}^{T} \end{bmatrix}}_{P_t^T} F_{n+1} + \underbrace{\begin{bmatrix} 0 \\ \vdots \\ 0 \\ \operatorname{Ad}_{g_{n,n+1}^{-1}}^{T} \end{bmatrix}}_{P_t^T} F_{n+1} + \underbrace{\begin{bmatrix} 0 \\ \vdots \\ 0 \\ \operatorname{Ad}_{g_{n,n+1}^{-1}}^{T} \end{bmatrix}}_{P_t^T} F_{n+1} + \underbrace{\begin{bmatrix} 0 \\ \vdots \\ 0 \\ \operatorname{Ad}_{g_{n,n+1}^{-1}}^{T} \end{bmatrix}}_{P_t^T} F_{n+1} + \underbrace{\begin{bmatrix} 0 \\ \vdots \\ 0 \\ \operatorname{Ad}_{g_{n,n+1}^{-1}}^{T} \end{bmatrix}}_{P_t^T} F_{n+1} + \underbrace{\begin{bmatrix} 0 \\ \vdots \\ 0 \\ \vdots \\ 0 \\ \operatorname{Ad}_{g_{n,n+1}^{-1}}^{T} \end{bmatrix}}_{P_t^T} F_{n+1} + \underbrace{\begin{bmatrix} 0 \\ \vdots \\ 0 \\ \vdots \\ 0 \\ \vdots \\ 0 \\ \operatorname{Ad}_{g_{n,n+1}^{-1}}^{T} \end{bmatrix}}_{P_t^T} F_{n+1} + \underbrace{\begin{bmatrix} 0 \\ \vdots \\ 0$$

$$\underbrace{\begin{bmatrix} -\operatorname{ad}_{V_1}^T \\ -\operatorname{ad}_{V_n}^T \end{bmatrix}}_{\operatorname{ad}_{V_n}^T} \begin{bmatrix} M_1 \\ \cdot \\ M_n \end{bmatrix} \begin{bmatrix} V_1 \\ \vdots \\ V_n \end{bmatrix}$$

$$F = G^{T}M\dot{V} + G^{T}P_{t}^{T}\underbrace{F_{n+1}}_{F_{t}} + G^{T} \cdot \operatorname{ad}_{V}^{T}MV$$

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$$\begin{split} \tau &= \xi^T \cdot F \\ \tau &= \xi^T G^T M \dot{V} + \xi^T G^T P_t^T F_t + \xi^T G^T \cdot \operatorname{ad}_V^T M V \\ &\Rightarrow \begin{cases} &= \xi^T G^T M (G \xi \ddot{\theta} + G P_0 \dot{V}_0 + G \cdot \operatorname{ad}_{\xi \dot{\theta}} P_0 V_0 + G \cdot \operatorname{ad}_{\xi \dot{\theta}} \Gamma V) + \\ &\xi^T G^T P_t^T F_t + \xi^T G^T \cdot \operatorname{ad}_V^T M V \\ &= \xi^T G^T M G \xi \ddot{\theta} + \xi^T G^T M G P_0 \dot{V}_0 + \xi^T G^T M G \cdot \operatorname{ad}_{\xi \dot{\theta}} \Gamma V + \\ &\xi^T G^T P_t^T F_t + \xi^T G^T \cdot \operatorname{ad}_V^T M V \end{split}$$

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$$\tau = \xi^{T} \cdot F$$

$$\tau = \xi^{T} G^{T} M \dot{V} + \xi^{T} G^{T} P_{t}^{T} F_{t} + \xi^{T} G^{T} \cdot \operatorname{ad}_{V}^{T} M V$$

$$\Rightarrow \frac{\xi^{T} G^{T} M (G \xi \ddot{\theta} + G P_{0} \dot{V}_{0} + G \cdot \operatorname{ad}_{\xi \dot{\theta}} P_{0} V_{0} + G \cdot \operatorname{ad}_{\xi \dot{\theta}} \Gamma V) + G \cdot \operatorname{ad}_{\xi \dot{\theta}} \Gamma V}{\xi^{T} G^{T} P_{t}^{T} F_{t} + \xi^{T} G^{T} \cdot \operatorname{ad}_{V}^{T} M V}$$

$$= \xi^{T} G^{T} M G \xi \ddot{\theta} + \xi^{T} G^{T} M G P_{0} \dot{V}_{0} + \xi^{T} G^{T} M G \cdot \operatorname{ad}_{\xi \dot{\theta}} \Gamma V + \xi^{T} G^{T} P_{t}^{T} F_{t} + \xi^{T} G^{T} \cdot \operatorname{ad}_{V}^{T} M V$$

Finally we get (recall that $V_0 = 0$, we have $V = G\xi\dot{\theta}$):

$$\underbrace{\xi^{T}G^{T}MG\xi \ddot{\theta} + \xi^{T}G^{T}(MG \cdot \operatorname{ad}_{\xi\dot{\theta}}\Gamma + \operatorname{ad}_{V}^{T}M)G\xi \dot{\theta} + \xi^{T}G^{T}MGP_{0}\dot{V}_{0}}_{\phi(\theta)} + \xi^{T}G^{T}P_{t}^{T}F_{t} = \tau$$

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$$M(\theta)\ddot{\theta} + C(\theta,\dot{\theta}) + \phi(\theta) + J_t^T(\theta)F_t = \tau$$

$$M(\theta) = \xi^T G^T M G \xi$$

$$C(\theta, \dot{\theta}) = \xi^T G^T (M G \operatorname{ad}_{\xi \dot{\theta}} \Gamma + \operatorname{ad}_V^T M) G \xi \dot{\theta}$$

$$\phi(\theta) = \xi^T G^T M G P_0 \dot{V}_0$$

$$I_t = P_t G \xi$$

$$M = \left[\begin{array}{ccc} M_1 & & 0 \\ 0 & \ddots & M_n \end{array} \right]$$

Property 4:

$$\Gamma^{n} = 0,$$

$$G = (I - \Gamma)^{-1} = I + \Gamma + \Gamma^{2} + \dots + \Gamma^{n-1}$$

$$I + \Gamma G = G$$

\square Square root factorization of M:

Definition: Articulated body inertia algorithm (Featherstone)

$$\begin{split} \hat{M}_n &= M_n, b_n = -\mathrm{ad}_{V_n}^T \big(M_n V_n \big) \\ F_i &= \hat{M}_i \dot{V}_i + b_i, i = n, \dots, 1 \\ b_i &= b_i \big(V_i, V_{i+1}, \xi_{i+1}, \hat{M}_{i+1}, \tau_{i+1} \big) \text{ (bias force)} \end{split}$$

$$\begin{cases} F_{i} = \operatorname{Ad}_{g_{i,i+1}^{-1}}^{T} F_{i+1} + M_{i} \dot{V}_{i} - \operatorname{ad}_{V_{i}}^{T} M_{i} V_{i} \\ F_{i+1} = \hat{M}_{i+1} \dot{V}_{i+1} + b_{i+1} \\ \dot{V}_{i+1} = \operatorname{Ad}_{g_{i,i+1}^{-1}} \dot{V}_{i} + \xi_{i+1} \ddot{\theta}_{i+1} - \operatorname{ad}_{\xi_{i+1} \dot{\theta}_{i+1}} \operatorname{Ad}_{g_{i,i+1}^{-1}} V_{i} \end{cases} \Rightarrow$$

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\Box Square root factorization of M:

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 \Rightarrow

$$F_{i+1} = \hat{M}_{i+1} \left(\operatorname{Ad}_{g_{i,i+1}^{-1}} \dot{V}_{i} + \xi_{i+1} \ddot{\theta}_{i+1} - \operatorname{ad}_{\xi_{i+1} \dot{\theta}_{i+1}} \operatorname{Ad}_{g_{i,i+1}^{-1}} V_{i} \right) + b_{i+1} \Rightarrow$$

$$\tau_{i+1} = \xi_{i+1}^{T} F_{i+1} = \xi_{i+1}^{T} \left(\hat{M}_{i+1} \left(\operatorname{Ad}_{g_{i,i+1}^{-1}} \dot{V}_{i} + \xi_{i+1} \ddot{\theta}_{i+1} - \operatorname{ad}_{\xi_{i+1} \dot{\theta}_{i+1}} \operatorname{Ad}_{g_{i,i+1}^{-1}} V \right) + b_{i+1} \right) \Rightarrow$$

$$\ddot{\theta}_{i+1} = \frac{\tau_{i+1} + \xi_{i+1}^{T} \left(\hat{M}_{i+1} \left(-\operatorname{Ad}_{g_{i,i+1}^{-1}} \dot{V}_{i} + \operatorname{ad}_{\xi_{i+1} \dot{\theta}_{i+1}} \operatorname{Ad}_{g_{i,i+1}^{-1}} V_{i} \right) - b_{i+1} \right)}{\xi_{i+1}^{T} \hat{M}_{i+1} \xi_{i+1}}$$

$$\Rightarrow$$

$$F_{i} = \operatorname{Ad}_{g_{i,i+1}^{-1}} \left(\hat{M}_{i+1} \left(\operatorname{Ad}_{g_{i,i+1}^{-1}} \dot{V}_{i} + \operatorname{ad}_{\xi_{i+1} \dot{\theta}_{i+1}} \operatorname{Ad}_{g_{i,i+1}^{-1}} V_{i} \right) - b_{i+1} \right)}{\xi_{i+1}^{T} \hat{M}_{i+1} \xi_{i+1}} \right)$$

$$- \operatorname{ad}_{\xi_{i+1} \dot{\theta}_{i+1}} \operatorname{Ad}_{g_{i,i+1}^{-1}} V_{i} \right) + b_{i+1} \right) + M_{i} \dot{V}_{i} - \operatorname{ad}_{V_{i}}^{T} M_{i} V_{i}$$

Square root factorization of *M*:

 $\hat{M}_{i} = M_{i} + Ad_{g_{i,i+1}^{-1}}^{T} \hat{M}_{i+1} Ad_{g_{i,i+1}^{-1}} - \frac{Ad_{g_{i,i+1}^{-1}}^{T} \hat{M}_{i+1} \xi_{i+1} \xi_{i+1}^{T} \hat{M}_{i+1} Ad_{g_{i,i+1}^{-1}}}{\xi_{i+1}^{T} \hat{M}_{i+1} \xi_{i+1}}$

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 $b_{i} = \operatorname{Ad}_{g_{i+1}^{-1}}^{T} b_{i+1} - \operatorname{ad}_{V_{i}}^{T} M_{i} V_{i} - \operatorname{Ad}_{g_{i+1}^{-1}}^{T} \hat{M}_{i+1} \operatorname{ad}_{\xi_{i+1} \dot{\theta}_{i+1}} \operatorname{Ad}_{g_{i+1}^{-1}} V_{i} +$ $\operatorname{Ad}_{\sigma^{-1}}^{T} \hat{M}_{i+1} \xi_{i+1} (\tau_{i+1} - \xi_{i+1}^{T} b_{i+1} + \xi_{i+1}^{T} \hat{M}_{i+1} \operatorname{ad}_{\xi_{i+1} \dot{\theta}_{i+1}} \operatorname{Ad}_{\sigma^{-1}_{i+1}} V_{i})$ $\xi_{i+1}^T \hat{M}_{i+1} \xi_{i+1}$ $\Rightarrow \hat{M} = M + \Gamma^T \hat{M} \Gamma - \Gamma^T \hat{M} \xi (\xi^T \hat{M} \xi)^{-1} \xi^T \hat{M} \Gamma$ $\Rightarrow M = \hat{M} - \Gamma^T \hat{M} \Gamma + \Gamma^T \hat{M} \xi (\xi^T \hat{M} \xi)^{-1} \xi^T \hat{M} \Gamma$ $\Rightarrow M(\theta) = \xi^T G^T (\hat{M} - \Gamma^T \hat{M} \Gamma + \Gamma^T \hat{M} \xi (\xi^T \hat{M} \xi)^{-1} \xi^T \hat{M} \Gamma) G \xi$ $\Rightarrow M(\theta) = \xi^T G^T \hat{M}^T \xi (\xi^T \hat{M} \xi)^{-1} \xi^T \hat{M} G \xi \triangleq W W^T$ $W = \xi^T G^T \hat{M}^T \xi (\xi^T \hat{M} \xi)^{-\frac{1}{2}}$

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 - 4 Coordinate Invariant Algorithms
 - 5 Lagrange's Equations with Constraints
 - General formulation
 - Unified geometric approach (Liu G.F., Li Z.X.)
 - Gauge invariant formulation (Aghili, F.)
 - 6 References

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Definition: Holonomic constraints

Let $q = (q_1, \ldots, q_n) \in E$, Ambient Space. A holonomic constraint is a set of constraint equations:

$$h_i(q)=0, i=1,\ldots,k$$

If $\{dh_i(q)\}_{i=1}^k$ is linearly independent, then $Q = h^{-1}(0)$ is a manifold of dim $m \triangleq n - k$.

$$T_qQ: \{V \in T_qE | \mathrm{d}h_i \cdot V = 0, \forall i = 1, \dots, k\} \subset T_qE:$$

subspace of permissible velocities.

$$T_qQ: \{V \in T_qE | \mathrm{d}h_i \cdot V = 0, \forall i = 1, \dots, k\} \subset T_qE$$

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Lagrange's Equations with Constraints

 $T_a^* Q^{\perp} : \{ f \in T_q^* E | \langle f, \nu \rangle = 0, \forall V \in T_q Q \}$ $= \operatorname{span}\{\mathrm{d}h_1,\ldots,\mathrm{d}h_k\} \subset T_a^*E:$

the subspace of constraint forces (Annihilator of T_aQ).

$$\dim(T_qQ)=m,\dim(T_q^*Q^\perp)=n-m=k$$

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1 General formulation

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$$T_q^* Q^{\perp} : \{ f \in T_q^* E | \langle f, v \rangle = 0, \forall V \in T_q Q \}$$
$$= \operatorname{span} \{ \operatorname{d} h_1, \dots, \operatorname{d} h_k \} \subset T_q^* E :$$

the subspace of constraint forces (Annihilator of T_qQ).

$$\dim(T_qQ)=m,\dim(T_q^*Q^\perp)=n-m=k$$

Definition: Constraint forces

$$\Gamma = \frac{\partial h}{\partial q}^T \cdot \lambda$$

 $\lambda \in \mathbb{R}^k$: the vector of relative magnitudes of constraint forces.

Definition: Pfaffian Constraints

A Pfaffian constraint has the form:

$$A(q)\dot{q}=0, A(q)\in\mathbb{R}^{k\times n}$$

Given a Pfaffian constraint,

$$\Delta_q = \{ v_q \in T_q E | A(q) \cdot v_q = 0 \} \subset T_q E :$$

distribution of permissible velocities.

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Definition: Pfaffian Constraints

A Pfaffian constraint has the form:

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Given a Pfaffian constraint,

$$\Delta_q = \{ v_q \in T_q E | A(q) \cdot v_q = 0 \} \subset T_q E :$$

distribution of permissible velocities.

Definition:

 $A(q)\dot{q}=0$ is holonomic (or integrable) iff Δ_a is an involutive distribution, iff

$$\exists h_i : E \mapsto \mathbb{R}, i = 1, \dots, k \text{ s.t. } \Delta_q = T_q Q, Q = h^{-1}(0)$$

Lagrange's

Equations with Constraints

Constraint forces:

 $\Gamma = A^T(q) \cdot \lambda, \lambda \in \mathbb{R}^k$

Lagrange's Equations with Constraints

1 General formulation

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§ Constraint forces:

$$\Gamma = A^T(q) \cdot \lambda, \lambda \in \mathbb{R}^k$$

§ Kinetic energy:

$$T(q,\dot{q}) = \frac{1}{2}\dot{q}^T \cdot M(q) \cdot \dot{q}$$

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§ Constraint forces:

$$\Gamma = A^T(q) \cdot \lambda, \lambda \in \mathbb{R}^k$$

Kinetic energy:

$$T(q,\dot{q}) = \frac{1}{2}\dot{q}^T \cdot M(q) \cdot \dot{q}$$

Potential energy:

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§ Constraint forces:

$$\Gamma = A^T(q) \cdot \lambda, \lambda \in \mathbb{R}^k$$

§ Kinetic energy:

$$T(q,\dot{q}) = \frac{1}{2}\dot{q}^T \cdot M(q) \cdot \dot{q}$$

§ Potential energy:

§ Lagrangian:

$$L(q, \dot{q}) = T(q, \dot{q}) - V(q)$$

1 General formulation

□ Lagrange's equations with constraints:

$$M(q)\ddot{q} + C(q,\dot{q}) + N(q) + A^{T}(q)\lambda = F$$

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1 General formulation

☐ Lagrange's equations with constraints:

$$M(q)\ddot{q} + C(q,\dot{q}) + N(q) + A^{T}(q)\lambda = F$$

□ Explicit solution for constraint forces:

$$A(q)\ddot{q} + \dot{A}(q)\dot{q} = 0$$

$$(AM^{-1}A^{T})\lambda = AM^{-1}(F - C - N) + \dot{A}\dot{q}$$

$$\lambda = (AM^{-1}A^{T})^{-1}(AM^{-1}(F - C - N) + \dot{A}\dot{q})$$

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♦ Example:

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$$x^2 + y^2 = l^2$$

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$$\underbrace{\begin{bmatrix} x & y \end{bmatrix}}_{A(q)} \begin{bmatrix} \dot{x} \\ \dot{y} \end{bmatrix} = 0$$
grange's

$$L(q, \dot{q}) = \frac{1}{2}m(\dot{x}^2 + \dot{y}^2) - mgy$$

$$\begin{bmatrix} m & 0 \\ 0 & m \end{bmatrix} \begin{bmatrix} \ddot{x} \\ \ddot{y} \end{bmatrix} + \begin{bmatrix} 0 \\ mg \end{bmatrix} + \begin{bmatrix} x \\ y \end{bmatrix} \lambda = 0$$

$$\lambda = (AM^{-1}A^{T})^{-1}(AM^{-1}(F - C - N) + \dot{A}\dot{q})$$

$$=-\frac{m}{l^2}(gy+\dot{x}^2+\dot{y}^2)$$

 $\begin{array}{c|c}
y \\
-2 & -1 \\
-2 & \theta \\
\end{array}$ $\begin{array}{c|c}
 & x \\
 & (x,y) \\
\end{array}$ $\begin{array}{c|c}
 & mg \\
\end{array}$

Figure 4.13

Tension =
$$\|\begin{bmatrix} x \\ y \end{bmatrix} \lambda\| = \frac{mg}{l}y + \frac{m}{l}(\dot{x}^2 + \dot{y}^2)$$

Lagrange's Equations with

Constraints

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□ Lagrange-d'Alembert formulation:

Given the Pfaffian constraint $A(q)\dot{q} = 0$ and virtual displacement $\delta a \in \mathbb{R}^k$, we have:

Theorem 1 (D'alembert Principle): Forces of constraints do no virtual work!

 $(A^{T}(q)\lambda) \cdot \delta q = 0$ for $A(q)\delta q = 0$



1717-1783

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1 General formulation

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□ Lagrange-d'Alembert formulation:

Given the Pfaffian constraint $A(q)\dot{q}=0$ and virtual displacement $\delta q\in\mathbb{R}^k$, we have:

Theorem 1 (D'alembert Principle): Forces of constraints do no virtual work!

$$(A^{T}(q)\lambda) \cdot \delta q = 0$$
 for $A(q)\delta q = 0$



1717-1783

Theorem 2 (Lagrange-d'Alembert Equation):

$$\left(\frac{\mathrm{d}}{\mathrm{d}t}\frac{\partial L}{\partial \dot{q}} - \frac{\partial L}{\partial q} - \tau\right) \cdot \delta q = 0, A(q)\delta q = 0$$



1736-1813

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1 General formulation

Let $A(q) = [A_1(q) \ A_2(q)]$, and $A_2(q) \in \mathbb{R}^{k \times k}$ is invertible, then $\delta q_1 \in \mathbb{R}^{n-k}$ are free variables:

$$\delta q_{2} = -A_{2}^{-1}(q)A_{1}(q)\delta q_{1}$$

$$\Rightarrow \left(\frac{\mathrm{d}}{\mathrm{d}t}\frac{\partial L}{\partial \dot{q}} - \frac{\partial L}{\partial q} - \tau\right) \cdot \delta q$$

$$= \left(\frac{\mathrm{d}}{\mathrm{d}t}\frac{\partial L}{\partial \dot{q}_{1}} - \frac{\partial L}{\partial q_{1}} - \tau_{1}\right) \cdot \delta q_{1} + \left(\frac{\mathrm{d}}{\mathrm{d}t}\frac{\partial L}{\partial \dot{q}_{2}} - \frac{\partial L}{\partial q_{2}} - \tau_{2}\right) \cdot \delta q_{2}$$

 $= \left(\frac{\mathrm{d}}{\mathrm{d}t}\frac{\partial L}{\partial \dot{q}_{1}} - \frac{\partial L}{\partial q_{1}} - \tau_{1}\right) \cdot \delta q_{1} + \left(\frac{\mathrm{d}}{\mathrm{d}t}\frac{\partial L}{\partial \dot{q}_{2}} - \frac{\partial L}{\partial q_{2}} - \tau_{2}\right) \cdot \left(-A_{2}^{-1}A_{1}\right)\delta q_{1}$

(Continues next slide)

Lagrange's Equations with Constraints

1 General formulation

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As $\delta q_1 \in \mathbb{R}^{n-k}$ is free,

$$\left(\frac{\mathrm{d}}{\mathrm{d}t}\frac{\partial L}{\partial \dot{q}_1} - \frac{\partial L}{\partial q_1} - \tau_1\right) - A_1^T A_2^{-T} \left(\frac{\mathrm{d}}{\mathrm{d}t}\frac{\partial L}{\partial \dot{q}_2} - \frac{\partial L}{\partial q_2} - \tau_2\right) = 0$$

Lagrange-d'Alembert equation

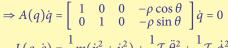
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Pfaffian constraint:

⋄ Example: Dynamics of a rolling disk

Lagrange's Equations with Constraints

 $\begin{cases} \dot{x} - \rho \cos \theta \dot{\phi} = 0 \\ \dot{y} - \rho \sin \theta \dot{\phi} = 0 \end{cases}$



$$L(q, \dot{q}) = \frac{1}{2}m(\dot{x}^2 + \dot{y}^2) + \frac{1}{2}\mathcal{I}_1\ddot{\theta}^2 + \frac{1}{2}\mathcal{I}_2\dot{\phi}^2$$

Lagrange-d'Alembert equation:

$$\left(\begin{bmatrix} m & m & & \\ & m & & \\ & & \mathcal{I}_1 & \\ & & & \mathcal{I}_2 \end{bmatrix} \ddot{q} - \begin{bmatrix} 0 & 0 \\ 0 & \\ \tau_{\theta} \\ \tau_{\phi} \end{bmatrix} \right) \cdot \delta q = 0$$

where $A(q) \cdot \delta q = 0$.

(Continues next slide)

Figure 4.14

1 General formulation

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As
$$\begin{cases} \delta x = \rho \cos \theta \cdot \delta \phi \\ \delta y = \rho \sin \theta \cdot \delta \phi \end{cases}$$

the equation of motion is:

As

$$\begin{cases} \ddot{x} = \rho \cos \theta \cdot \ddot{\phi} - \rho \sin \theta \cdot \dot{\theta} \dot{\phi} \\ \ddot{y} = \rho \sin \theta \cdot \ddot{\phi} + \rho \cos \theta \cdot \dot{\theta} \dot{\phi} \end{cases}$$
$$\Rightarrow \begin{bmatrix} \mathcal{I}_{1} & 0 \\ 0 & \mathcal{I}_{2} + m\rho^{2} \end{bmatrix} \begin{bmatrix} \ddot{\theta} \\ \ddot{\phi} \end{bmatrix} = \begin{bmatrix} \tau_{\theta} \\ \tau_{\phi} \end{bmatrix}$$

Solve for $(\theta(t), \phi(t))$, and then solve for (x(t), y(t)) from:

$$\begin{cases} \dot{x} = \rho \cos \theta \cdot \dot{\phi} \\ \dot{y} = \rho \sin \theta \cdot \dot{\phi} \end{cases} \Leftarrow \text{1st order differential equation}$$

(Continues next slide)

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□ Nature of nonholonomic constraints:

Consider $q = (r, s) \in \mathbb{R}^2 \times \mathbb{R}$ with Pfaffian constraint,

$$\dot{s} + a^T(r)\dot{r} = 0, a(r) \in \mathbb{R}^2$$

Lagrangian:

$$L = L(r, \dot{r}, \dot{s})$$

and constrained Lagrangian:

$$L_c(r,\dot{r}) = L(r,\dot{r},-a^T(r)\dot{r})$$

⇒ Lagrange's equation:

$$\frac{\mathrm{d}}{\mathrm{d}t}\frac{\partial L_c}{\partial \dot{r}_i} - \frac{\partial L_c}{\partial r_i} = 0, i = 1, 2$$

$$\Rightarrow \frac{\mathrm{d}}{\mathrm{d}t} \left(\frac{\partial L}{\partial \dot{r}_i} - a_i(r) \frac{\partial L}{\partial \dot{s}} \right) - \left(\frac{\partial L}{\partial r_i} - \frac{\partial L}{\partial \dot{s}} \sum_j \frac{\partial a_j}{\partial r_i} \dot{r}_j \right) = 0$$

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$$\Rightarrow \left(\frac{\mathrm{d}}{\mathrm{d}t}\frac{\partial L}{\partial \dot{r}_{i}} - \frac{\partial L}{\partial r_{i}}\right) - a_{i}(r)\left(\frac{\mathrm{d}}{\mathrm{d}t}\frac{\partial L}{\partial \dot{s}} - \frac{\partial L}{\partial s}\right) = \underbrace{\frac{\partial L}{\partial \dot{s}}\left(\dot{a}_{i}(r) - \sum_{j}\frac{\partial a_{j}}{\partial r_{i}}\dot{r}_{j}\right)}_{\neq 0}(*)$$

If the constraint is holonomic, i.e.

$$a_i(r) = \frac{\partial h}{\partial r_i}$$
 for some $h: E \mapsto \mathbb{R}$

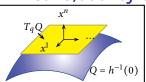
then RHS (right hand side) of (*) equals

$$\frac{\partial L}{\partial \dot{s}} \left(\sum_{j} \frac{\partial^{2} h}{\partial r_{i} \partial r_{j}} \dot{r}_{j} - \sum_{j} \frac{\partial^{2} h}{\partial r_{j} \partial r_{i}} \dot{r}_{j} \right) = 0$$

2 Unified geometric approach (Liu G.F., Li Z.X.)

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Metric, duality and orthogonality on T_aE :



$$\mathcal{K} \triangleq \frac{1}{2} \ll \dot{q}, \dot{q} \gg_{M}$$
$$= \frac{1}{2} \dot{q}^{T} M(q) \dot{q}$$

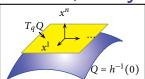
Figure 4.15

$$\begin{split} T_q Q^\perp &= \big\{ V_1 \in T_q E \big| \ll V_1, V_2 \gg_M \triangleq V_1^T M V_2 = 0, \, \forall \, V_2 \in T_q Q \big\} \\ T_q^* Q^\perp &= \big\{ f \in T_q^* E \big| \big\langle f, V \big\rangle = 0, \, \forall \, V \in T_q Q \big\} : \quad \text{constraint forces} \\ T_q E &= T_q Q \oplus T_q Q^\perp, \, T_q^* E = T_q^* Q \oplus T_q^* Q^\perp \end{split}$$

Lagrange's Equations with Constraints

2 Unified geometric approach (Liu G.F., Li Z.X.)

□ Metric, duality and orthogonality on T_qE :



$$\mathcal{K} \triangleq \frac{1}{2} \ll \dot{q}, \dot{q} \gg_{M}$$
$$= \frac{1}{2} \dot{q}^{T} M(q) \dot{q}$$

Figure 4.15

$$\begin{split} T_q Q^\perp &= \left\{ V_1 \in T_q E | \ll V_1, V_2 \gg_M \triangleq V_1^T M V_2 = 0, \forall \, V_2 \in T_q Q \right\} \\ T_q^* Q^\perp &= \left\{ f \in T_q^* E | \left\langle f, V \right\rangle = 0, \forall \, V \in T_q Q \right\} : \quad \text{constraint forces} \\ T_q E &= T_q Q \oplus T_q Q^\perp, T_q^* E = T_q^* Q \oplus T_q^* Q^\perp \end{split}$$

Definition:

$$M^{\flat}: T_q E \mapsto T_q^* E, \langle M^{\flat} V_1, V_2 \rangle = V_1^T M V_2 = \ll V_1, V_2 \gg_M$$

$$M^{\sharp}: T_q^* E \mapsto T_q E, M^{\sharp} = M^{\flat^{-1}}$$

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2 Unified geometric approach (Liu G.F., Li Z.X.)

 $T_qE \longrightarrow M^b \longrightarrow T_q^*E$ $T_qQ^1 \longrightarrow T_q^*Q$ $T_qQ \longrightarrow T_q^*Q$ $T_qQ \longrightarrow T_q^*Q$ $T_qQ \longrightarrow T_q^*Q$

Figure 4.16

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2 Unified geometric approach (Liu G.F., Li Z.X.)

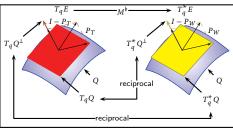


Figure 4.16

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Property 5:

Under the basis $\frac{\partial}{\partial q_i}$ and dq_i , $i=1,\ldots,n$ of T_qE and T_q^*E respectively, the matrix representation of M^{\flat} and M^{\dagger} is M and M^{-1} respectively.

2 Unified geometric approach (Liu G.F., Li Z.X.)

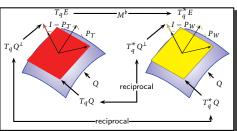


Figure 4.16

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Property 5:

Under the basis $\frac{\partial}{\partial q_i}$ and $\mathrm{d}q_i, i=1,\ldots,n$ of T_qE and T_q^*E respectively, the matrix representation of M^{\flat} and M^{\sharp} is M and M^{-1} respectively.

Property 6:

$$M^{\sharp}(T_q^*Q) = T_qQ$$
$$M^{\sharp}(T_q^*Q^{\perp}) = T_qQ^{\perp}$$

2 Unified geometric approach (Liu G.F., Li Z.X.)

Given

$$h: E \mapsto \mathbb{R}^k, m = n - k$$

$$h_* \triangleq T_q h: T_q E \mapsto T_{h(q)} \mathbb{R}^k$$

$$h^* \triangleq T_q^* h: T_{h(q)}^* \mathbb{R}^k \mapsto T_q^* E$$

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we have:

$$\ker h_* = T_q Q, h_* (T_q Q^{\perp}) = T_{h(q)} \mathbb{R}^k, h^* (T_{h(q)}^* \mathbb{R}^k) = T_q^* Q^{\perp}$$

$$T_{q}^{*}E \xleftarrow{h^{*}} T_{q}^{*}\mathbb{R}^{k}$$

$$M^{\sharp} \downarrow \qquad \qquad \downarrow M_{2}^{\sharp} \qquad \qquad M_{2}^{\sharp} = h_{*} \circ M^{\sharp} \circ h^{*}$$

$$T_{q}E \xrightarrow{l_{*}} T_{h(q)}\mathbb{R}^{k}$$

2 Unified geometric approach (Liu G.F., Li Z.X.)

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Lemma 1:

The map $(I - P_{\omega}) : T_a^* E \mapsto T_a^* Q^{\perp}$ given by

$$\left(I-P_{\omega}\right)=h^{*}\circ M_{2}^{\flat}\circ h_{*}\circ M^{\sharp}$$

is a well-defined projection map, with the property:

$$(I-P_{\omega})f_1=0, \forall f_1\in T_q^*Q$$

$$(I-P_{\omega})f_2=f_2, \forall f_2\in T_q^*Q^{\perp}$$

2 Unified geometric approach (Liu G.F., Li Z.X.)

Lemma 1:

The map $(I - P_{\omega}) : T_a^* E \mapsto T_a^* Q^{\perp}$ given by

$$(I - P_{\omega}) = h^* \circ M_2^{\flat} \circ h_* \circ M^{\sharp}$$

is a well-defined projection map, with the property:

$$(I - P_{\omega})f_1 = 0, \forall f_1 \in T_q^* Q$$

$$(I - P_{\omega})f_2 = f_2, \forall f_2 \in T_a^* Q^{\perp}$$

Lagrange's Equations with Constraints

Proof:

Given $f_1 \in T_q^*Q$, $M^{\sharp}(f_1) \in T_qQ = \ker h_*$, then $(I - P_{\omega})f_1 = 0$. For $f_2 \in T_q^*Q^{\perp}$, $\exists \lambda \in \mathbb{R}^{n-m}$ s.t. $f_2 = h^* \lambda$, and

$$(I-P_{\omega})f_2=h^*M_2^{\flat}h_*M^{\sharp}h^*\lambda=h^*\lambda=f_2$$

thus $P_{\omega}: T_a^* E \mapsto T_a^* Q$ is a well-defined projection map. Similarly,

$$P_T: T_qE \mapsto T_qQ, P_T = I - M^{\sharp}h^*M_2^bh_*$$

and

$$(I-P_T): T_qE \mapsto T_qQ^{\perp}$$

are projection maps.

2 Unified geometric approach (Liu G.F., Li Z.X.)

Lemma 2:

 $P_{\omega}M = MP_{T}$ $P_{\omega}h^{*} = h_{*}P_{T} = 0$

 $P_T = P_{\omega}^T$

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2 Unified geometric approach (Liu G.F., Li Z.X.)

Lemma 2:

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$$P_{\omega}M = MP_{T}$$

$$P_{\omega}h^{*} = h_{*}P_{T} = 0$$

$$P_{T} = P_{\omega}^{T}$$

For nonholonomic constraints:

$$h_* \leftarrow A(q)$$

 $h^* \leftarrow A^*(q)$
 $T_q Q \leftarrow \Delta_q$
 $T_a^* Q^{\perp} \leftarrow \operatorname{span}\{a_i(q), i = 1, \dots, k\}$

application in hybrid velocity/force control.

2 Unified geometric approach (Liu G.F., Li Z.X.)

□ Lagrange's equations of motion:

 $M(q)\ddot{q} + C(q, \dot{q}) + N = \tau + A^{T}(q)\lambda \Rightarrow$ $\lambda = (AM^{-1}A^{T})^{-1}AM^{-1}(M\ddot{q} + (C + N - \tau))$

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2 Unified geometric approach (Liu G.F., Li Z.X.)

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□ Lagrange's equations of motion:

$$M(q)\ddot{q} + C(q, \dot{q}) + N = \tau + A^{T}(q)\lambda \Rightarrow$$
$$\lambda = (AM^{-1}A^{T})^{-1}AM^{-1}(M\ddot{q} + (C + N - \tau))$$

Define $P_{\omega} = I - A^{T} (AM^{-1}A^{T})^{-1}AM^{-1}$, then:

$$P_{\omega} M \ddot{q} + P_{\omega} C + P_{\omega} N = P_{\omega} \tau$$

Denote $\tilde{C} = P_{\omega}C$, $\tilde{N} = P_{\omega}N$, $\tilde{\tau} = P_{\omega}\tau$, $P_{\omega}M\ddot{\theta} \triangleq \tilde{M}\ddot{\theta}$ is the intertia force in T_q^*Q .

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2 Unified geometric approach (Liu G.F., Li Z.X.)

□ Lagrange's equations of motion:

$$M(q)\ddot{q} + C(q, \dot{q}) + N = \tau + A^{T}(q)\lambda \Rightarrow$$

 $\lambda = (AM^{-1}A^{T})^{-1}AM^{-1}(M\ddot{q} + (C + N - \tau))$

Define $P_{\omega} = I - A^{T} (AM^{-1}A^{T})^{-1}AM^{-1}$, then:

$$P_{\omega}M\ddot{q}+P_{\omega}C+P_{\omega}N=P_{\omega}\tau$$

Denote $\tilde{C}=P_{\omega}C, \tilde{N}=P_{\omega}N, \tilde{\tau}=P_{\omega}\tau, P_{\omega}M\ddot{\theta}\triangleq \tilde{M}\ddot{\theta}$ is the intertia force in $T_{q}^{*}Q$.

Definition: Dynamics in T_q^*Q

$$\tilde{M}\ddot{\theta} + \tilde{C} + \tilde{N} = \tilde{\tau}$$

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2 Unified geometric approach (Liu G.F., Li Z.X.)

Similarly

$$(I - P_{\omega})(M\ddot{q} + C + N) = (I - P_{\omega})\tau + A^{T}\tau$$

Let

$$P_T = I - M^{-1}A^T (AM^{-1}A^T)^{-1}A$$

then since $P_{\omega}M = MP_T$, we have:

Definition: Dynamics in $T_q^*Q^{\perp}$

$$M(I-P_T)(\ddot{q}+M^{-1}C)=(I-P_\omega)(\tau-N)+A^T\lambda$$

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Similarly

$$(I - P_{\omega})(M\ddot{q} + C + N) = (I - P_{\omega})\tau + A^{T}\tau$$

Let

$$P_T = I - M^{-1}A^T (AM^{-1}A^T)^{-1}A$$

then since $P_{\omega}M = MP_T$, we have:

Definition: Dynamics in $T_q^*Q^{\perp}$

$$M(I-P_T)(\ddot{q}+M^{-1}C)=(I-P_{\omega})(\tau-N)+A^T\lambda$$

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□ **Geometric Interpretation:**

$$\nabla \leftrightarrow M$$

$$M\ddot{q} + C + N = \tau + A^T \lambda \iff M \nabla_{\dot{q}} \dot{q} = \tau - N + A^T \lambda$$

 $\tilde{\nabla} \leftrightarrow \text{ induced metric on } T_q Q$

 $S: TQ \otimes TQ \mapsto N(Q): 2^{nd}$ fundamental form

TQ: tangent vector field

N(Q): normal vector field

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4.5 Lagrange's Equations with Constraints

2 Unified geometric approach (Liu G.F., Li Z.X.)

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$$\nabla_X Y = \tilde{\nabla}_X Y + S(X, Y)$$

$$M\underbrace{(I - P_T)(\ddot{q} + M^{-1}C)}_{S(\dot{q}, \dot{q})} = (I - P_{\omega})(\tau - N) + A^T \lambda$$

 $MS(\dot{q}, \dot{q})$: centrifugal force due to curvature of Q in E

2 Unified geometric approach (Liu G.F., Li Z.X.)

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$$\nabla_X Y = \tilde{\nabla}_X Y + S(X, Y)$$

$$M\underbrace{(I - P_T)(\ddot{q} + M^{-1}C)}_{S(\dot{q}, \dot{q})} = (I - P_{\omega})(\tau - N) + A^T \lambda$$

 $MS(\dot{q}, \dot{q})$: centrifugal force due to curvature of Q in E

Definition: Hybrid position/force control

$$M\tilde{\nabla}_{\dot{q}}\dot{q} = \tilde{\tau} - \tilde{N}$$

$$MS(\dot{q}, \dot{q}) = (I - P_{\omega})(\tau - N) + A^{T}\lambda$$

2 Unified geometric approach (Liu G.F., Li Z.X.)

♦ Example: Dynamics of a Spherical Pendulum

Lagrange's Equations with Constraints

$$K = \frac{1}{2}m(\dot{x}^2 + \dot{y}^2 + \dot{z}^2) = \frac{1}{2}\dot{q}^T M \dot{q}$$

$$q = (x, y, z)^T, M = mI$$

$$h(q) = q^T q - r^2 = 0$$

 $A = (x, y, z), M_2^{\sharp} = AM^{\sharp}A^T = r^2/m$

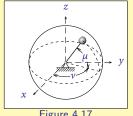


Figure 4.17

$$P_{\omega} = \frac{1}{r^{2}} \begin{bmatrix} y^{2} + z^{2} & -xy & -xz \\ -yz & x^{2} + z^{2} & -yz \\ -zx & -zy & x^{2} + y^{2} \end{bmatrix}$$
$$I - P_{\omega} = \frac{1}{r^{2}} \begin{bmatrix} x^{2} & xy & xz \\ yx & y^{2} & yz \\ zx & zy & z^{2} \end{bmatrix}$$

2 Unified geometric approach (Liu G.F., Li Z.X.)

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$$P_T = P_{\omega}^T = P_{\omega}$$

 (μ, ν) : Spherical coordinates

$$q = (r \cos \mu \cos \nu, r \cos \mu \sin \nu, r \sin \mu)^T$$

$$\begin{split} \tilde{\nabla}_{\dot{q}}\dot{q} &= P_T(\nabla_{\dot{q}}\dot{q}) \\ &= \begin{bmatrix} -r\sin\mu\cos\nu & -r\sin\nu \\ -r\sin\mu\sin\nu & r\cos\nu \end{bmatrix} \begin{bmatrix} \nu_1 \\ \nu_2 \end{bmatrix} \end{split}$$

$$S(\dot{q}, \dot{q}) = (I - P_T)(\nabla_{\dot{q}}\dot{q})$$

$$= (-\dot{\mu}^2 - \cos^2 \mu \dot{v}^2) \begin{bmatrix} r\cos \mu \cos v \\ r\cos \mu \sin v \\ r\sin \mu \end{bmatrix}$$

where

$$v_1 = \ddot{\mu} + \sin \mu \cos \mu \dot{v}^2$$

$$v_2 = \cos \mu \ddot{v} - 2 \sin \mu \dot{\mu} \dot{v}$$



2 Unified geometric approach (Liu G.F., Li Z.X.)

□ Control Algorithm:

holonomic constraints:

 \tilde{q} : coordinates of Q

$$q = \psi(\tilde{q}) \Rightarrow \dot{q} = J \cdot \dot{\tilde{q}}$$

$$q - \psi(q) \rightarrow q - \gamma \cdot q$$

$$\tau = MJ(\ddot{q}_d - K_v \dot{\tilde{e}} - K_p \tilde{e}) + C_1 + N + A^T(-\lambda_d + K_I \int (\lambda - \lambda_d))$$

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2 Unified geometric approach (Liu G.F., Li Z.X.)

□ Control Algorithm:

holonomic constraints:

 \tilde{q} : coordinates of Q

$$q = \psi(\tilde{q}) \Rightarrow \dot{q} = J \cdot \dot{\tilde{q}}$$

$$\tau = MJ(\ddot{\ddot{q}}_d - K_v \dot{\ddot{e}} - K_p \tilde{e}) + C_1 + N + A^T \left(-\lambda_d + K_I \int \left(\lambda - \lambda_d \right) \right)$$

2 nonholonomic constraints:

Let $J(q) \in \mathbb{R}^{n \times m}$ be s.t. AJ = 0. Write $\dot{q} = J \cdot u$ for some u

$$\tau = MJ(\dot{u}_d - K_p(u - u_d)) + M\dot{J}u + C + N + A^T(-\lambda_d + K_I \int (\lambda - \lambda_d))$$

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2 Unified geometric approach (Liu G.F., Li Z.X.)

♦ Example: 6-DoF manipulator on a sphere with frictionless point contact

■ Contact constraint:

$$v_z = 0 \Leftrightarrow \begin{bmatrix} 0 \ 0 \ 1 \ 0 \ 0 \end{bmatrix} \mathrm{Ad}_{g_{fl_f}^{-1}} V_{of} = 0$$

⇒ Holonomic constraint:

$$\eta = (\alpha_o^T, \alpha_f^T, \psi)$$
: Parametrization of Q

$$P_{\omega} = \text{diag}(1, 1, 0, 1, 1, 1)$$

Newton-Euler Equations of motion:

$$M\dot{V}_{of} - ad_{V_{of}}^T MV_{of} = F_m + G + A^T \lambda$$

$$V_{of} = \left[\begin{array}{ccc} R_{\psi}M_o & -M_f & 0 \\ 0 & 0 & 0 \\ R_{\psi}R_oK_oM_o & -R_oK_fM_f & 0 \\ -T_oM_o & -T_fM_f & 1 \\ \end{array} \right] \left[\begin{array}{c} \dot{\alpha}_o \\ \dot{\alpha}_f \\ \dot{\psi} \end{array} \right] \triangleq J\dot{\eta}$$

$$MJ\ddot{\eta} + C_1 = F_m + G + A^T \lambda \tag{*}$$

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2 Unified geometric approach (Liu G.F., Li Z.X.)

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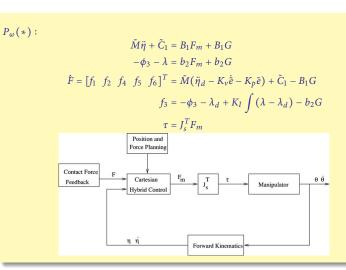
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2 Unified geometric approach (Liu G.F., Li Z.X.)

♦ Example: 6-DoF manipulator rolling on a sphere

$$\dot{A}_1(q)$$

$$f_c = A_1^T \lambda, \lambda \in \mathbb{R}^4$$

$$\begin{bmatrix} \omega_x \\ \omega_y \end{bmatrix} = -R_0 (K_f + R_{\psi} K_o R_{\psi}) M_f \dot{\alpha}_f$$

$$V_{of} = \mathrm{Ad}_{g_{fl_f}} \cdot V_{l_o l_f}$$

$$= \operatorname{Ad}_{g_{fi_f}} \left[\begin{array}{c} 0 \\ 0 \\ 0 \\ -R_o \left(K_f + R_{\psi} K_o R_{\psi} \right) M_f \end{array} \right] \dot{\alpha}_f$$

$$\triangleq J_f \dot{\alpha}_f$$

 $\operatorname{span}\{J_f\}$: Not involutive

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2 Unified geometric approach (Liu G.F., Li Z.X.)

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$$\begin{split} MJ_f \ddot{\alpha}_f + \left(M \dot{J}_f \dot{\alpha}_f - \operatorname{ad}_{J_f \dot{\alpha}_f}^T M J_f \dot{\alpha}_f \right) &= F_m + G + A^T \lambda \\ F_m &= MJ_f (\ddot{\alpha}_{fd} - K_p (\dot{\alpha}_f - \dot{\alpha}_{fd})) + \left(M \dot{J}_f \dot{\alpha}_f - \operatorname{ad}_{J_f \dot{\alpha}_f}^T M J_f \dot{\alpha}_f \right) \\ &+ A^T (-\lambda_d + \int \left(\lambda - \lambda_d \right) \right) - G \end{split}$$

2 Unified geometric approach (Liu G.F., Li Z.X.)

Lagrange's Equations with Constraints

$$MJ_{f}\ddot{\alpha}_{f} + (M\dot{J}_{f}\dot{\alpha}_{f} - \operatorname{ad}_{J_{f}\dot{\alpha}_{f}}^{T}MJ_{f}\dot{\alpha}_{f}) = F_{m} + G + A^{T}\lambda$$

$$F_{m} = MJ_{f}(\ddot{\alpha}_{fd} - K_{p}(\dot{\alpha}_{f} - \dot{\alpha}_{fd})) + (M\dot{J}_{f}\dot{\alpha}_{f} - \operatorname{ad}_{J_{f}\dot{\alpha}_{f}}^{T}MJ_{f}\dot{\alpha}_{f})$$

$$+ A^{T}(-\lambda_{d} + \int (\lambda - \lambda_{d})) - G$$

Example: Redundant parallel manipulator

$$\theta = (\theta_1, \dots, \theta_6) \in E$$

$$\theta_a = (\theta_1, \theta_3, \theta_5)$$

$$\theta_p = (\theta_2, \theta_4, \theta_6)$$

$$H(\theta) = \begin{bmatrix} x_a + lc_1 + lc_{12} - x_b - lc_3 - lc_{34} \\ y_a + ls_1 + ls_{12} - y_b - ls_3 - ls_{34} \\ x_a + lc_1 + lc_{12} - x_c - lc_5 - lc_{56} \\ y_a + ls_1 + ls_{12} - y_c - ls_5 - ls_{56} \end{bmatrix} = 0$$
where $c_{ii} = \cos(\theta_i + \theta_i)$, $s_{ii} = \sin(\theta_i + \theta_i)$.

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2 Unified geometric approach (Liu G.F., Li Z.X.)

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$$M_i(\theta) \in \mathbb{R}^{2 imes 2} : i^{th}$$
 chain

$$M(\theta) = \operatorname{diag}(M_1(\theta), \dots, M_3(\theta))$$

 $M(\theta)\ddot{\theta} + C + N = \tau + A^T\lambda$

If all joints are actuated, we can achieve:

Position control of end-effector

internal grasping force

As τ_2 , τ_4 , $\tau_6 = 0$,

$$ilde{ heta} \in \mathbb{R}^2$$
: local parametrization of $Q = H^{-1}(0)$

$$\theta = \psi(\tilde{\theta})$$
: embedding of Q in E

$$\dot{\theta} = J\dot{\tilde{\theta}}$$

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2 Unified geometric approach (Liu G.F., Li Z.X.)

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Given $P_{\omega}: T_{\theta}^* E \mapsto T_{\theta}^* Q$, the dynamics in $T_{\theta}^* Q$ is given by:

$$\begin{split} P_{\omega}MJ\ddot{\tilde{\theta}} + P_{\omega}(C_1 + N) &= P_{\omega}\tau \\ \tilde{\tau} &= (\tau_1, \tau_3, \tau_5) \\ \tilde{P}_{\omega} &= (P_1, P_3, P_5) \\ \hat{\tau} &= \hat{P}_{\omega}\tilde{\tau} = P_{\omega}\tau \in \mathbb{R}^6 \\ \hat{\tau} &= P_{\omega}MJ(\ddot{\tilde{\theta}}_d - K_v\dot{\tilde{e}} - K_p\tilde{e}) + P_{\omega}(C_1 + N) \end{split}$$



3 Gauge invariant formulation (Aghili, F.)

Gauge-invariant Formulation (Aghili):

♦ Square root factorization of inertia matrix:

 $M = WW^{T}$ (square root factorization)

$$\begin{cases} v \triangleq W^T \dot{q} \in \mathbb{R}^n \\ u \triangleq W^{-1} \tau \in \mathbb{R}^n \end{cases} \qquad T = \frac{1}{2} \dot{q}^T M \dot{q} = \frac{1}{2} v^T v$$

Lagrange's Equations with Constraints

3 Gauge invariant formulation (Aghili, F.)

Gauge-invariant Formulation (Aghili):

♦ Square root factorization of inertia matrix:

 $M = WW^{T}$ (square root factorization)

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 (square root factorization)

$$\begin{cases} v \triangleq W^T \dot{q} \in \mathbb{R}^n \\ u \triangleq W^{-1} \tau \in \mathbb{R}^n \end{cases} \qquad T = \frac{1}{2} \dot{q}^T M \dot{q} = \frac{1}{2} v^T v$$

♦ Lagrange's Equation:

$$\frac{\mathrm{d}}{\mathrm{d}t} \frac{\partial T}{\partial \dot{q}} - \frac{\partial T}{\partial q} = \tau \Rightarrow \frac{\mathrm{d}}{\mathrm{d}t} (Wv) - \frac{\partial v^T}{\partial q} v = \tau \Rightarrow$$

$$W\dot{v} + \dot{W}v - \frac{\partial v^T}{\partial q} v = \tau \Rightarrow \dot{v} + W^{-1} (\dot{W} - \frac{\partial v^T}{\partial q}) v = W^{-1} \tau = u$$

Define $\Gamma \triangleq W^{-1}(\dot{W} - \frac{\partial v^T}{\partial a})$, then:

$$\dot{v} + \Gamma v = u$$

Lagrange's

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3 Gauge invariant formulation (Aghili, F.)

□ Gauge-invariant Formulation (Aghili):

♦ Change of coordinates:

$$\bar{v} = V^T v, \bar{u} = V^T u, V \in U(n) \Rightarrow \frac{\mathrm{d}}{\mathrm{d}t} (V\bar{v}) + \Gamma(V\bar{v}) = V\bar{u} \Rightarrow V\dot{\bar{v}} + \dot{V}\bar{v} + \Gamma V\bar{v} = V\bar{u} \Rightarrow \dot{\bar{v}} + V^T (\Gamma + \dot{V}V^T)V\bar{v} = \bar{u}$$

$$VV + VV + 1VV = VU \rightarrow V + V + V + V + VV = U$$

$$VV^{T} = I \Rightarrow \dot{V}V^{T} + V\dot{V}^{T} = 0 \Rightarrow \dot{V}V^{T} = -(\dot{V}V^{T})^{T} =: -\Omega \Rightarrow$$

$$\bar{\Gamma} \triangleq V^T (\Gamma - \Omega) V, \Rightarrow \boxed{\dot{\bar{v}} + \bar{\Gamma} \bar{v} = \bar{u}}$$

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3 Gauge invariant formulation (Aghili, F.)

□ Gauge-invariant Formulation (Aghili):

♦ Change of coordinates:

$$\bar{v} = V^T v, \bar{u} = V^T u, V \in U(n) \Rightarrow \frac{\mathrm{d}}{\mathrm{d}t} (V\bar{v}) + \Gamma(V\bar{v}) = V\bar{u} \Rightarrow$$

$$V\dot{\bar{v}} + \dot{V}\bar{v} + \Gamma V\bar{v} = V\bar{u} \Rightarrow \dot{\bar{v}} + V^T (\Gamma + \dot{V}V^T) V\bar{v} = \bar{u}$$

$$VV^T = I \Rightarrow \dot{V}V^T + V\dot{V}^T = 0 \Rightarrow \dot{V}V^T = -(\dot{V}V^T)^T =: -\Omega \Rightarrow$$

$$\bar{\Gamma} \triangleq V^T (\Gamma - \Omega) V, \Rightarrow \boxed{\dot{v} + \bar{\Gamma}\bar{v} = \bar{u}}$$

♦ Pfaffian constraint:

$$A(q)\dot{q}=0, A(q)\in\mathbb{R}^{m\times n}, \Lambda\triangleq AW^{-T}\Rightarrow \Lambda\nu=0,$$

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3 Gauge invariant formulation (Aghili, F.)

□ Gauge-invariant Formulation (Aghili):

♦ Change of coordinates:

$$\bar{v} = V^T v, \, \bar{u} = V^T u, \, V \in U(n) \Rightarrow \frac{\mathrm{d}}{\mathrm{d}t} (V\bar{v}) + \Gamma(V\bar{v}) = V\bar{u} \Rightarrow$$

$$V\dot{\bar{v}} + \dot{V}\bar{v} + \Gamma V\bar{v} = V\bar{u} \Rightarrow \dot{\bar{v}} + V^T (\Gamma + \dot{V}V^T) V\bar{v} = \bar{u}$$

$$VV^T = I \Rightarrow \dot{V}V^T + V\dot{V}^T = 0 \Rightarrow \dot{V}V^T = -(\dot{V}V^T)^T =: -\Omega \Rightarrow$$

$$\bar{\Gamma} \triangleq V^T (\Gamma - \Omega) V, \Rightarrow \boxed{\dot{\bar{v}} + \bar{\Gamma}\bar{v} = \bar{u}}$$

♦ Pfaffian constraint:

$$A(q)\dot{q}=0, A(q)\in\mathbb{R}^{m\times n}, \Lambda\triangleq AW^{-T}\Rightarrow \Lambda\nu=0,$$

♦ Lagrange's equation with constraint:

$$\dot{\nu} + \Gamma \nu = u + \Lambda^T \lambda$$

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3 Gauge invariant formulation (Aghili, F.)

□ Gauge-invariant Formulation (Aghili):

 \diamond **SVD** of Λ :

$$\Lambda = U\Sigma V^T, \bar{v} = V^T v, \bar{u} = V^T u, \bar{\Lambda} \triangleq \Lambda V$$

where $\bar{\Lambda}\bar{\nu}=0$ and:

$$\sigma = \begin{bmatrix} S & 0 \\ 0 & 0 \end{bmatrix}, S = \operatorname{diag}(\sigma_1, \dots, \sigma_r), \sigma_1 \ge \dots \ge \sigma_r, r \le m$$

$$\begin{cases} U = \begin{bmatrix} U_1 & U_2 \end{bmatrix}, U_1 \in \mathbb{R}^{m \times r}, U_2 \in \mathbb{R}^{m \times (m-r)} \\ V = \begin{bmatrix} V_1 & V_2 \end{bmatrix}, V_1 \in \mathbb{R}^{n \times r}, V_2 \in \mathbb{R}^{n \times (n-r)} \end{cases} \Rightarrow$$

$$\bar{\Lambda} = \begin{bmatrix} \Lambda_r & 0_{m \times (n-r)} \end{bmatrix}, \Lambda_r \triangleq U_1 S, \bar{\Lambda} \bar{v} = 0 \Rightarrow$$

$$\bar{v} = \begin{bmatrix} 0_{r \times 1} \\ v_r \end{bmatrix}, v_r \triangleq V_2^T v = V_2^T W^T \dot{q}$$

(Continues next slide)

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3 Gauge invariant formulation (Aghili, F.)

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□ Gauge-invariant Formulation (Aghili):

$$\bar{u} = \begin{bmatrix} u_o \\ u_r \end{bmatrix}, \begin{cases} u_o = V_1^T W^{-1} \tau \\ u_r = V_2^T W^{-1} \tau \end{cases}$$

$$\bar{\Gamma}_{ij}\triangleq V_{i}^{T}\big(\Gamma-\Omega\big)V_{j}, i,j=1,2, \Gamma_{r}\triangleq\bar{\Gamma}_{22}, \Gamma_{o}\triangleq\bar{\Gamma}_{12}$$

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4.5 Lagrange's Equations with Constraints

3 Gauge invariant formulation (Aghili, F.)

☐ Gauge-invariant Formulation (Aghili):

$$\bar{u} = \begin{bmatrix} u_o \\ u_r \end{bmatrix}, \begin{cases} u_o = V_1^T W^{-1} \tau \\ u_r = V_2^T W^{-1} \tau \end{cases}$$

$$\bar{\Gamma}_{ij}\triangleq V_{i}^{T}\big(\Gamma-\Omega\big)V_{j}, i,j=1,2, \Gamma_{r}\triangleq\bar{\Gamma}_{22}, \Gamma_{o}\triangleq\bar{\Gamma}_{12}$$

Decoupled equation of motion/constrained force:

$$\dot{v}_r + \Gamma_r v_r = u_r$$

$$\Gamma_o v_r = u_o + \Gamma_r^T \lambda$$

Lagrange's Equations with Constraints

3 Gauge invariant formulation (Aghili, F.)

□ Gauge-invariant Formulation (Aghili):

$$\bar{u} = \begin{bmatrix} u_o \\ u_r \end{bmatrix}, \begin{cases} u_o = V_1^T W^{-1} \tau \\ u_r = V_2^T W^{-1} \tau \end{cases}$$

$$\bar{\Gamma}_{ij} \triangleq V_i^T (\Gamma - \Omega) V_j, i, j = 1, 2, \Gamma_r \triangleq \bar{\Gamma}_{22}, \Gamma_o \triangleq \bar{\Gamma}_{12}$$

♦ Decoupled equation of motion/constrained force:

$$\dot{v}_r + \Gamma_r v_r = u_r$$

$$\Gamma_o v_r = u_o + \Gamma_r^T \lambda$$

♦ Combined equation of motion (Kane's equation):

$$\boxed{ \begin{array}{c} \frac{\mathrm{d}}{\mathrm{d}t} \left[\begin{array}{c} q \\ v_r \end{array} \right] = \left[\begin{array}{c} W^{-T} V_2 \\ -\Gamma_r \end{array} \right] v_r + \left[\begin{array}{c} 0 \\ I \end{array} \right] u_r }$$

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3 Gauge invariant formulation (Aghili, F.)

□ Gauge-invariant Formulation (Aghili):

 \diamond Composite error vector ε

 $q = q(\theta), \theta \in \mathbb{R}^{n-r}$: generalized coordinate (of $Q \subset \mathbb{R}^n$)

$$\nu_r = B(\theta)\dot{\theta}, B(\theta) \triangleq V_2^T W^T J, J \triangleq \frac{\partial q}{\partial \theta}, \tilde{\theta} \triangleq \theta - \theta_d, \tilde{\nu}_r \triangleq \nu_r - \nu_{r_d} \Rightarrow$$

$$\varepsilon \triangleq \tilde{v}_r + BK_p\tilde{\theta} = B(\dot{\tilde{\theta}} + K_p\tilde{\theta})$$
: composite error

$$s \triangleq v_{r_d} - BK_p \tilde{\theta} = v_r - \varepsilon = B(\dot{\theta}_d - K_p \tilde{\theta}), \tilde{\lambda} \triangleq \lambda - \lambda_d$$

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3 Gauge invariant formulation (Aghili, F.)

☐ Gauge-invariant Formulation (Aghili):

 \diamond Composite error vector ε

$$q = q(\theta), \theta \in \mathbb{R}^{n-r}$$
: generalized coordinate (of $Q \subset \mathbb{R}^n$)

$$v_r = B(\theta)\dot{\theta}, B(\theta) \triangleq V_2^T W^T J, J \triangleq \frac{\partial q}{\partial \theta}, \tilde{\theta} \triangleq \theta - \theta_d, \tilde{v}_r \triangleq v_r - v_{r_d} \Rightarrow$$

$$\varepsilon \triangleq \tilde{v}_r + BK_p\tilde{\theta} = B(\dot{\tilde{\theta}} + K_p\tilde{\theta})$$
: composite error

$$s \triangleq v_{r_d} - BK_p \tilde{\theta} = v_r - \varepsilon = B(\dot{\theta}_d - K_p \tilde{\theta}), \tilde{\lambda} \triangleq \lambda - \lambda_d$$

♦ Hybrid position/force control

$$u_r = \dot{s} + \Gamma_r s - K_d \varepsilon$$

$$u_o = -\Lambda_r^T \lambda_d + \Gamma_o v_r$$

Note: integration term $K_I \int (\lambda - \lambda_d)$ is missing from u_o . (Continues next slide)

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3 Gauge invariant formulation (Aghili, F.)

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□ Gauge-invariant Formulation (Aghili):

$$\dot{v}_r + \Gamma_r v_r = u_r = \frac{\mathrm{d}}{\mathrm{d}t} (v_r - \varepsilon) + \Gamma_r (v_r - \varepsilon) - K_d \varepsilon \Rightarrow \dot{\varepsilon} = -(\Gamma_r + K_d) \varepsilon$$
$$\Lambda_r^T \lambda + \Gamma_o v_r = u_o = -\Lambda_r^T \lambda_d + \Gamma_o v_r \Rightarrow \Lambda_r^T \tilde{\lambda} + K_I \int \tilde{\lambda} = 0$$

3 Gauge invariant formulation (Aghili, F.)

☐ Gauge-invariant Formulation (Aghili):

$$\begin{split} \dot{v}_r + \Gamma_r v_r &= u_r = \frac{\mathrm{d}}{\mathrm{d}t} (v_r - \varepsilon) + \Gamma_r (v_r - \varepsilon) - K_d \varepsilon \Rightarrow \dot{\varepsilon} = - (\Gamma_r + K_d) \varepsilon \\ \Lambda_r^T \lambda + \Gamma_o v_r &= u_o = - \Lambda_r^T \lambda_d + \Gamma_o v_r \Rightarrow \Lambda_r^T \tilde{\lambda} + K_I \int \tilde{\lambda} = 0 \end{split}$$

♦ Equivalence to the Geometric approach

$$\begin{bmatrix} P_{\omega} = WV_2V_2^TW^{-1} \\ I - P_{\omega} = WV_1V_1^TW^{-1} \end{bmatrix} \Rightarrow V^TW^{-1}P_{\omega}\tau = \begin{bmatrix} 0 \\ V_2^TW^{-1}\tau \end{bmatrix} = \begin{bmatrix} 0 \\ u_r \end{bmatrix}$$
$$V^TW^{-1}(I - P_{\omega})\tau = \begin{bmatrix} V_1^TW^{-1}\tau \\ 0 \end{bmatrix} = \begin{bmatrix} u_o \\ 0 \end{bmatrix}$$

(Continues next slide)

Chapter 4 Manipulator Dynamics

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3 Gauge invariant formulation (Aghili, F.)

□ Gauge-invariant Formulation (Aghili):

1. Geometric approach (K_p, K_I, K_d) :

$$u = \underbrace{\left[\begin{array}{c} V_1^T W^{-1} A^T \left(-\lambda_d + K_I \int \left(\lambda - \lambda_d \right) \right) \\ V_2^T W^T J \left(\ddot{\theta}_d - K_d \tilde{\theta} - K_p \tilde{\theta} \right) \end{array} \right]}_{fb} + \underbrace{V^T W^{-1} C_1}_{ff}$$

2. Gauge-invariant formulation $(\tilde{K}_p, \tilde{K}_d)$:

$$u = \begin{bmatrix} u_o \\ u_r \end{bmatrix} = \underbrace{\begin{bmatrix} V_1^T W^{-1} A^T (-\lambda_d) \\ V_2^T W^T J (\ddot{\theta}_d - K_d' \dot{\tilde{\theta}} - K_p' \tilde{\theta}) \end{bmatrix}}_{fb} + \underbrace{C_1'}_{ff} \Rightarrow$$

$$K_d' \triangleq \tilde{K}_p + (V_2^T W^T J)^{-1} \tilde{K}_d (V_2^T W^T J)$$

$$K_p' \triangleq (V_2^T W^T J)^{-1} \tilde{K}_d (V_2^T W^T J) \tilde{K}_p$$

Constraints
References

Lagrange's Equations with

† End of Section †

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4.6 References

Section 4.6

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