

# Chapter 1

## Introduction

### 1.1 Vibronic Models Background

Since the breakthrough of quantum physics, it was claimed that the whole of chemistry was in essence solved. The Schrödinger equation had it all, as long as you were able to solve it. Despite the rapid growth in computing, for the majority of meaningful and sizeable chemical systems, the calculation of its exact solutions remains intractable. Therefore, modern quantum chemistry relies on clever mathematical approximations to calculate energies sufficient for predictive and chemical accuracy. In 1927, Born and Oppenheimer proposed the Born-Oppenheimer Approximation (BOA), a method that separates the nuclear and electronic regimes and allows one to calculate the energy of adiabatic potential energy surfaces (PES). The transformed electronic Hamiltonian was more manageable. However, many realistic chemical processes are innately non-adiabatic, and have electronic states that are not well separated and are nearly degenerate. Prime examples where non-adiabatic effects show an important role are the photo-isomerization of retinal and transition metal complex spectra. Vibrational-electronic (**vib-ronic**) coupling models aim to simulate the mechanics revolving around non-adiabatic dynamics. Maturation in both post-Hartree-Fock and computing methods have since enabled quantum and computational chemists alike to explore sophisticated non-adiabatic phenomena such as conical intersections and intersystem crossings.

Research in non-adiabatic chemistry has since gained momentum in recent decades leading to the consequent construction of vibronic models that can go beyond the BOA and can capture a proper non-adiabatic picture. Vibronic models in a diabatic basis is theorized to be one viable approach to non-adiabatic treatments. This procedure is referred to as a diabatization. Diabatic states (diabats) have the advantage of diminishing the tricky non-adiabatic couplings (NACs) that arise from the BOA adiabatic states (adiabats), and maintaining smooth character along the PES. To solve for dynamics, many proceed to use the classic Multiconfiguration Time-dependent Hartree (MCTDH) algorithm. A long term goal of the Nooijen group is to formulate novel alternative operator techniques like Vibrational Electronic Coupled Cluster (VECC) to calculate vibrationally-resolved electronic spectra and statistical mechanics based on a vibronic model framework [1].

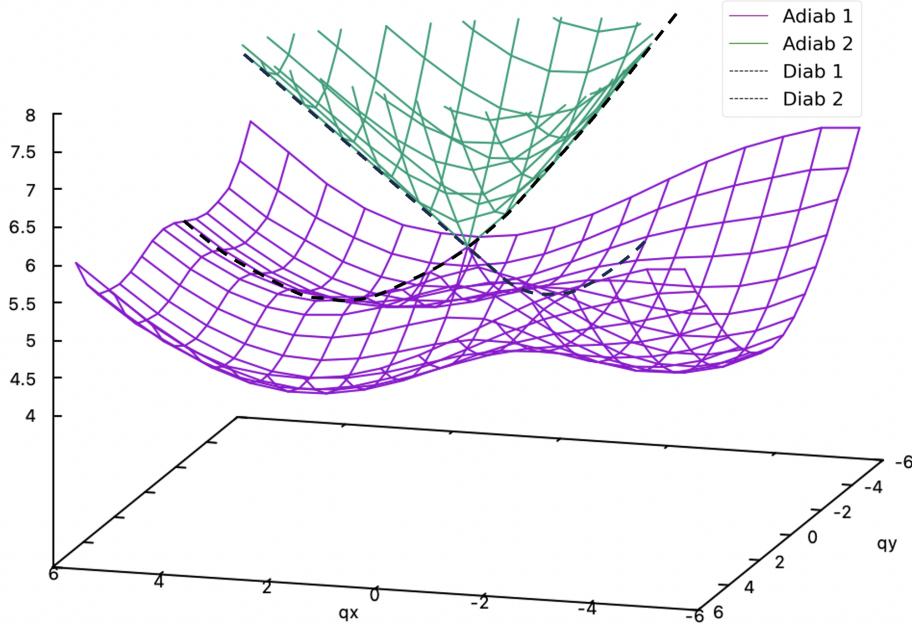


Fig 1. Stereotypical depiction of a conical intersection (CoIn) involving both adiabatic states (adiabats) and diabatic states (diabats) commonly seen in vibronic models. Source: Zeng.

The central goal in this project is to produce an accessible diabatization protocol that can supply MCTDH and VECC with vibronic model Hamiltonians suitable for a wide range of applications such as spin-orbit coupling (SOC) effects and spectroscopy. For the past two years at the Nooijen and Zeng group, we have jointly endeavoured to pursue this goal.

## 1.2 Modern Vibronic Models - Timeline

The pioneering work of Cederbaum, Köppel, and Domcke gave rise to the original idea of vibronic models in their 1981 seminal paper [2]. Determination of a vibronic model Hamiltonian that simultaneously coupled vibrational modes with an arbitrary set of electronic states was given in an *ab initio* procedure. Their study proceeded to simulate spectroscopy and demonstrated the need for the inclusion of vibronic coupling effects to match with findings in experimental spectra. Neugebauer and Nooijen found this approach useful and further developed a diabatization scheme suitable for Time-dependent Density Functional Theory (TD-DFT) absorption spectrum [3, 4]. Vibronic models were used at this point to model multiple potential energy surfaces in a limited region of configuration space, akin to a harmonic oscillator approximation applied to a single surface. Moreover, this established a formalized and standard method for extracting (from quantum chemistry calculations) vibronic model Hamiltonians in a routine fashion. This scheme inspired Santoro and coworkers, who in recent years recasted Nooijen’s scheme into a single-reference linear order Gaussian-MCTDH TD-DFT diabatization scheme purposed for studying the excited state dynamics of hexahelicene [5]. However, hexahelicene is merely an organic molecule. Ongoing frontier research by Domcke and Mondal aims to apply vibronic models for transition metal complex spectra, a notoriously difficult task due to the dilemma of low-lying degenerate electronic states. It is difficult to resolve the states and PES character for such systems. Therefore, a more ideal approach is to pivot towards a diabatization scheme that is multi-reference in nature, and can capture relativistic effects like spin-orbit coupling.

Such existence of a diabatization protocol that includes spin-orbit coupling was first devised by Toby Zeng, an expert who specializes in vibronics and has massively contributed to the field of arbitrarily large Hamiltonian formalisms and Jahn-Teller (JT) interactions [6]. The foundation of the Zeng SOC protocol derives from previous work done by Atchity

and Rudenberg on configurational uniformity, and specifically uses GAMESS-US package Generalized Multi-Configuration Quasidegenerate Perturbation Theory (GMC-QDPT) calculations. Diabatic states are directly calculated using Nakamura and Truhlar's version of configurational uniformity in GAMESS, whereby diabatic molecular orbitals (DMOs) are first prepared, and then computation of diabatic state functions follows. Calculations involved here are typed as multi-state multi-reference perturbation theory calculations in a given active space reference as generalized and implemented by Nakano. GMC-QDPT addresses the requisite Multi-configuration Self Consistent Field (MCSCF) component of the model by performing Occupation Restricted Multiple Active Space (ORMAS) SCF. At this point, Zeng assembles all these pieces and opens up the possibility of adding spin orbit coupling terms to the model Hamiltonian. The basis of the vibronic models in this project will use the diabatic energies and couplings yielded by this process.

Familiarity of transition metal vibronic models with SOC and spectroscopy is not recognized as new. In 1999 Glaesemann, Gordon (of GAMESS), and Nakano scrutinized the behaviour of  $FeCO^+$  in a wavefunction study. First author Mondal has in particular discussed  $E \times e$  Jahn-Teller and SOC effects in transition-metal fluorides, such as  $CoF_3$  with a 17 orbital and 30 electron active space - CAS(30,17). Their analysis included generation of  $D_{3h}$  trigonal planar symmetry vibronic model Hamiltonians and spectra. Previous applications of the Zeng SOC protocol were towards pnictogen hydride cations with trigonal symmetry:  $PH_3^+$ ,  $AsH_3^+$ , and  $SbH_3^+$  for the purposes of investigating vibronic coupling and SOC characteristics. No spectroscopy components were included. So far, this context and state of research converges on a case for which the project's diabatization protocol should be applied to: a large  $D_{3h}$  iron-carbonyl complex with low-lying degenerate electronic states. The candidate molecule is  $Fe(CO)_5$ . If successful, the appealing challenge of  $Fe(CO)_5$  should exhibit demonstrable JT, SOC, and photochemical effects as observed in the literature.

## 1.3 Project Scope: Pivot to Zeng Protocol and Limitations of Single-reference Linear Vibronic Models

To capitalize on all the work presented so far, I have created a novel automated diabatization code in Python to be incorporated into the Zeng SOC protocol. This is my main contribution. In my previous work, the Santoro protocol was used to generate vibronic models in Gaussian 16 for MCTDH testing. Naturally, they did not suffice for transition metal complexes. As shown in Fig 2, the vibronic models were found to not resolve any fine structure in metallic vibronic spectra, and for cases such as caffeine, the spectra was still inaccurate despite meticulous checking of the model. Santoro has likewise acknowledged the inaccuracies with using the linear protocol due to complications in finding correct parameters for calculations. The most pressing concern of using the Santoro protocol at the Nooijen lab was that the output vibronic model Hamiltonian (operator file) format was incompatible for VECC input.

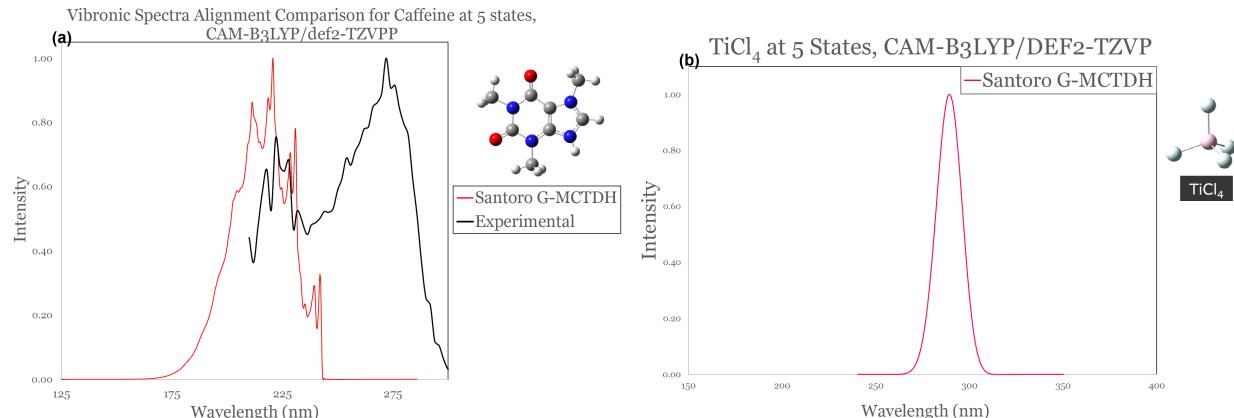


Fig 2. (a) Caffeine. Slightly matching peaks in the sub 240nm region observed, appears to need a linear red-shift for better agreement. (b) Titanium tetrachloride. The predictive accuracy here of a smooth single peak is recognized to be nowhere near reality, due to limitations of single reference TD-DFT methods applied to transition metal complexes.

To improve on the capabilities of the vibronic models available for VECC, such limitations of the Santoro G-MCTDH method prompted for a pivoting towards the Zeng protocol. Prior to this project, only the top half path shown in Fig 3. (Santoro-MCTDH-Spectra) was explored. Here, VECC was also able to perform excited state simulations like MCTDH,

however VECC's statistical mechanics ideally necessitate Zeng models for proper functionality. Hence to reiterate the declared task: we shall construct a composite approach where the Zeng protocol is first used to manufacture vibronic model Hamiltonians to be plugged in as input, and the latter half of non-adiabatic simulation shall be taken care of by VECC. A Zeng-VECC-Spectra pathway. Once the Zeng protocol can be successfully connected to both MCTDH and VECC, then it will allow for a variety of models and dataset to be used for practical comparison of the gold standard MCTDH method with the new VECC method.

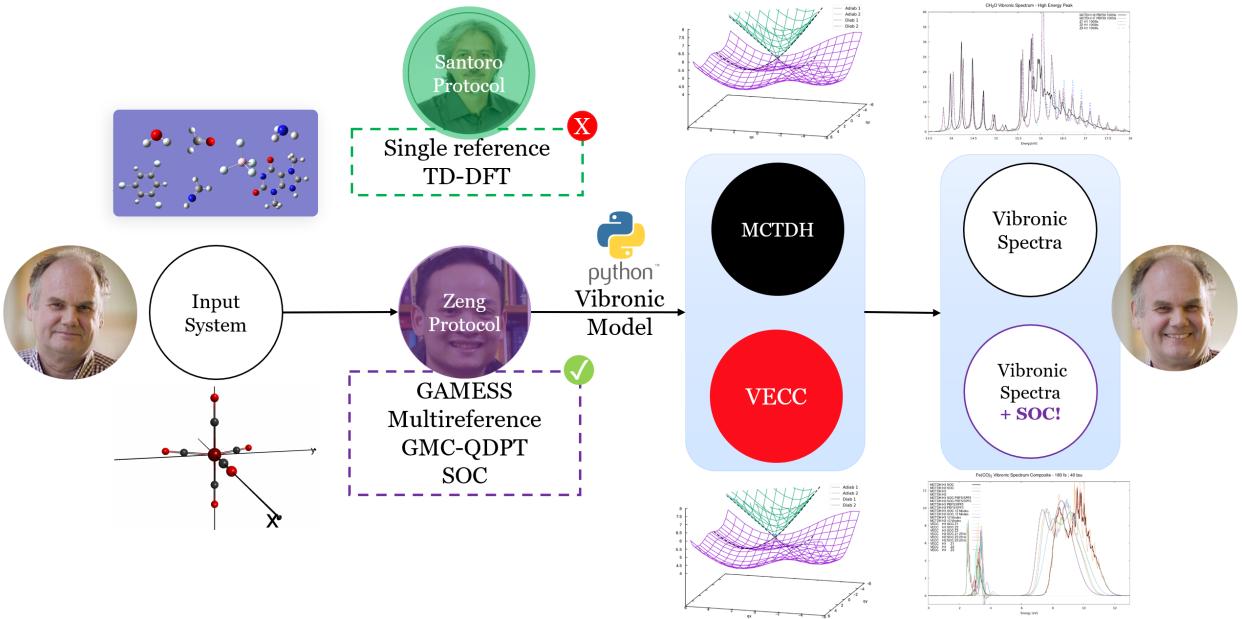


Fig 3. Original proposed research plan, the right half motivates the left half.

Formerly, the Zeng protocol diabatization was not automated. Each calculation in the series needed for full diabatization was handcrafted. In early 2023, Zeng completed the first step towards accomplishing the task: a prototype 'machine gun' Bash code that automates diabatization. This version could effectively produce up to quadratic order Zeng non-SOC models, and was implemented to be extracted from GAMESS calculations. Neil Raymond and I have iterated on the prototype and designed a seamless operating protocol. The code was first translated into Python, and then received a wide range of upgrades to ensure that it is SOC-compatible and readily plugged into MCTDH/VECC for spectroscopy.

## **1.4 Outline of Thesis**

In this thesis:

1. The first chapter gives an overview of vibronics and the research project objectives.
2. Fundamental quantum chemistry concepts and mathematics is presented.
3. Full elucidation of Python code to automate the Zeng SOC protocol.
4. Applications of the protocol to realistic transition metal systems and vibronic spectroscopy.
5. Conclusion and future outlook.

# Chapter 2

## Theoretical Overview and Methodology

### 2.1 Electronic structure theory

In an ideal world, the Schrödinger equation would be readily solved by computers and the exact solution for energy would be known.

Let the full molecular Schrödinger's Equation be defined as:

$$\hat{H}(\vec{R})\Psi_n(\vec{r}; \vec{R}) = E_n\Psi_n(\vec{r}; \vec{R}) \quad (2.1)$$

The full molecular Hamiltonian:

$$\hat{H} = -\sum_{\alpha}^M \frac{1}{2M_{\alpha}} \nabla_{\alpha} \cdot \nabla_{\alpha} - \sum_i^N \frac{1}{2m_e} \nabla_i \cdot \nabla_i + \sum_{i,j>i}^N \frac{1}{r_{ij}} + \sum_{i,\alpha}^{N,M} \frac{-Z_A}{r_{i\alpha}} + \sum_{\alpha,\beta \neq \alpha}^M \frac{Z_A Z_B}{R_{\alpha\beta}} \quad (2.2)$$

Which can concisely be represented by kinetic and potential energy operators  $\hat{T}$  and  $\hat{V}$ :

$$\hat{H} = \hat{T}_N + \hat{T}_e + \hat{V}_{ee} + \hat{V}_{Ne} + \hat{V}_{NN} \quad (2.3)$$

Note the subscripts  $e, N$  are for electronic and nuclear correspondingly.

$$\hat{H} = \hat{T}_N + \hat{H}_{elec} \quad (2.4)$$

$$\hat{H}_{elec} = \hat{T}_e + \hat{V}_{ee} + \hat{V}_{Ne} + \hat{V}_{NN} \quad (2.5)$$

Let the wavefunction,  $\Psi$ , be the sum over the electronic eigenstates  $\phi_\lambda(\vec{r}; \vec{R})$  and nuclear eigenstates  $\chi_\lambda(\vec{R})$ :

$$\Psi(\vec{r}; \vec{R}) = \sum_{\lambda} \phi_{\lambda}(\vec{r}; \vec{R}) \chi_{\lambda}(\vec{R}) \quad (2.6)$$

Here,  $\vec{r}$  indicates electron coordinates, and  $\vec{R}$  refers to nuclear configuration. The subscripts  $\lambda$  and  $\mu$  denote electronic states. To describe a parametric dependence of  $\vec{r}$  on  $\vec{R}$ ,  $(\vec{r}; \vec{R})$  is used. Signifying that although an explicit dependence on  $\vec{R}$  is not present in the function,  $\vec{r}$  values are predicated at a single  $\vec{R}$  parameter [7]. Therefore, if  $\vec{R}$  changes, so does  $\phi_{\lambda}(\vec{r}; \vec{R})$ .

Proceed with full expansion of Hamiltonian acting on the wavefunction:

$$\begin{aligned} & \hat{H}(\vec{R}) \Psi(\vec{r}; \vec{R}) \\ &= \left[ - \sum_{\alpha} \frac{1}{2M_{\alpha}} \vec{\nabla}_{\alpha} \cdot \vec{\nabla}_{\alpha} + \hat{H}_{elec} \right] \left[ \sum_{\lambda} \phi_{\lambda}(\vec{r}; \vec{R}) \chi_{\lambda}(\vec{R}) \right] \end{aligned} \quad (2.7a)$$

$$\begin{aligned} &= \left( - \sum_{\alpha} \frac{1}{2M_{\alpha}} \right) \sum_{\lambda} [\chi_{\lambda}(\vec{R}) \vec{\nabla}_{\alpha}^2 \phi_{\lambda}(\vec{r}; \vec{R}) + \vec{\nabla}_{\alpha} \phi_{\lambda}(\vec{r}; \vec{R}) \vec{\nabla}_{\alpha} \chi_{\lambda} \right. \\ &\quad \left. + \phi_{\lambda}(\vec{r}; \vec{R}) \vec{\nabla}_{\alpha}^2 \chi_{\lambda}(\vec{R}) + \vec{\nabla}_{\alpha} \chi_{\lambda} \vec{\nabla}_{\alpha} \phi_{\lambda}(\vec{r}; \vec{R}) + E_{\lambda}(\vec{R}) \phi_{\lambda}(\vec{r}; \vec{R}) \chi_{\lambda}(\vec{R}) ] \right) \end{aligned} \quad (2.7b)$$

$$\begin{aligned} &= \left( - \sum_{\alpha} \frac{1}{2M_{\alpha}} \right) \sum_{\lambda} [\chi_{\lambda}(\vec{R}) \vec{\nabla}_{\alpha}^2 \phi_{\lambda}(\vec{r}; \vec{R}) \\ &\quad + \phi_{\lambda}(\vec{r}; \vec{R}) \vec{\nabla}_{\alpha}^2 \chi_{\lambda}(\vec{R}) + 2 \vec{\nabla}_{\alpha} \phi_{\lambda}(\vec{r}; \vec{R}) \vec{\nabla}_{\alpha} \chi_{\lambda} + E_{\lambda}(\vec{R}) \phi_{\lambda}(\vec{r}; \vec{R}) \chi_{\lambda}(\vec{R}) ] \end{aligned} \quad (2.7c)$$

Integrating against  $\phi_{\mu}(\vec{r}, \vec{R})$  over all electronic coordinates,  $\vec{r}$ , and asserting that these states

are orthonormal at a single fixed  $\vec{R}$ , produces a sum of four terms:

$$= \sum_{\lambda} \int \phi_{\mu}^*(\vec{r}; \vec{R}) \sum_{\alpha} -\frac{1}{2M_{\alpha}} \vec{\nabla}_{\alpha}^2 \phi_{\lambda}(\vec{r}; \vec{R}) dr \chi_{\lambda}(\vec{R}) \quad (\mathbf{A})$$

$$+ \sum_{\lambda} \sum_{\alpha} \delta_{\lambda\mu} \frac{-1}{2M_{\alpha}} \vec{\nabla}_{\alpha}^2 \chi_{\lambda}(\vec{R}) \quad (\mathbf{B})$$

$$+ \sum_{\lambda} \sum_{\alpha} \int \phi_{\mu}^*(\vec{r}; \vec{R}) \vec{\nabla}_{\alpha} \phi_{\lambda}(\vec{r}; \vec{R}) dr \cdot -\frac{1}{2M_{\alpha}} \vec{\nabla}_{\alpha} \chi_{\lambda}(\vec{R}) \quad (\mathbf{C})$$

$$+ \sum_{\lambda} \delta_{\lambda\mu} E_{\lambda}(\vec{R}) \chi_{\lambda}(\vec{R}) \quad (\mathbf{D})$$

$$= \sum_{\lambda} E_{\lambda}(\vec{R}) \delta_{\lambda\mu} \chi_{\lambda}(\vec{R}) \quad (2.8)$$

The BOA facilitates setting the expansion in terms of a single electronic state, hence  $\lambda = \mu$ . Terms **A** and **C** in the BOA are also ignored. The **A** term represents a diagonal Born-Oppenheimer correction to potential, useful in specialized instances—as in accounting for different  $M_{\alpha}$  (e.g. isotopes) [8]. It is typically proximate enough to zero for it to be neglected. Likewise, **C** is zero for real normalized wavefunctions. **C** is referred to as the off-diagonal non-adiabatic coupling (NAC) terms. In other words, only **B** and **D** terms remain because terms with the presence of  $\int dr \vec{\nabla}_{\alpha} \phi_{\lambda}(\vec{r}; \vec{R})$  are effectively zero if there exists little change in  $\phi_{\lambda}$  with respect to  $\vec{R}$ , thus implying no coupling between electronic wavefunctions and nuclear configuration [8].

As for a small example, so far the BOA shows:

$$\begin{bmatrix} \hat{T}_N + E_1(\vec{R}) + A_1(\vec{R}) & N.A.C. \\ N.A.C. & \hat{T}_N + E_2(\vec{R}) + A_2(\vec{R}) \end{bmatrix} \quad (2.9)$$

Ultimately, the Schrödinger equation with the BOA applied appears as:

$$\left[ \hat{T}_N + E_{\lambda}(\vec{R}) \right] \Psi_{\lambda,n}(\vec{R}) = E_n \Psi_{\lambda,n}(\vec{R}) \quad (2.10)$$

Where label  $n$  is for rotational-vibrational levels and recalling  $\lambda$  is for electronic states. Attention to the term  $E_\lambda(\vec{R})$ , which is a function of nuclear configuration.  $E_\lambda(\vec{R})$  is defined to be a single point sufficiently able to define a potential energy surface (PES),  $E_\lambda(\{\vec{R}\})$ . Now the electronic wavefunction  $\phi_\lambda(\vec{r}; \vec{R})$ , wavefunction  $\chi_n(\vec{R})$  (detailing nuclear motion), and potential energy surface  $E_\lambda(\vec{R})$  in Equation (2.10) can be solved for at a single electronic state  $\lambda$  with any given fixed-geometry  $\vec{R}$ . In practice of using BOA concepts, Figure 2.1 demonstrates how the electronic ground state PES can be theoretically constructed. Understanding surfaces in this fashion gives rise to adiabatic states or true Born-Oppenheimer surfaces [9].

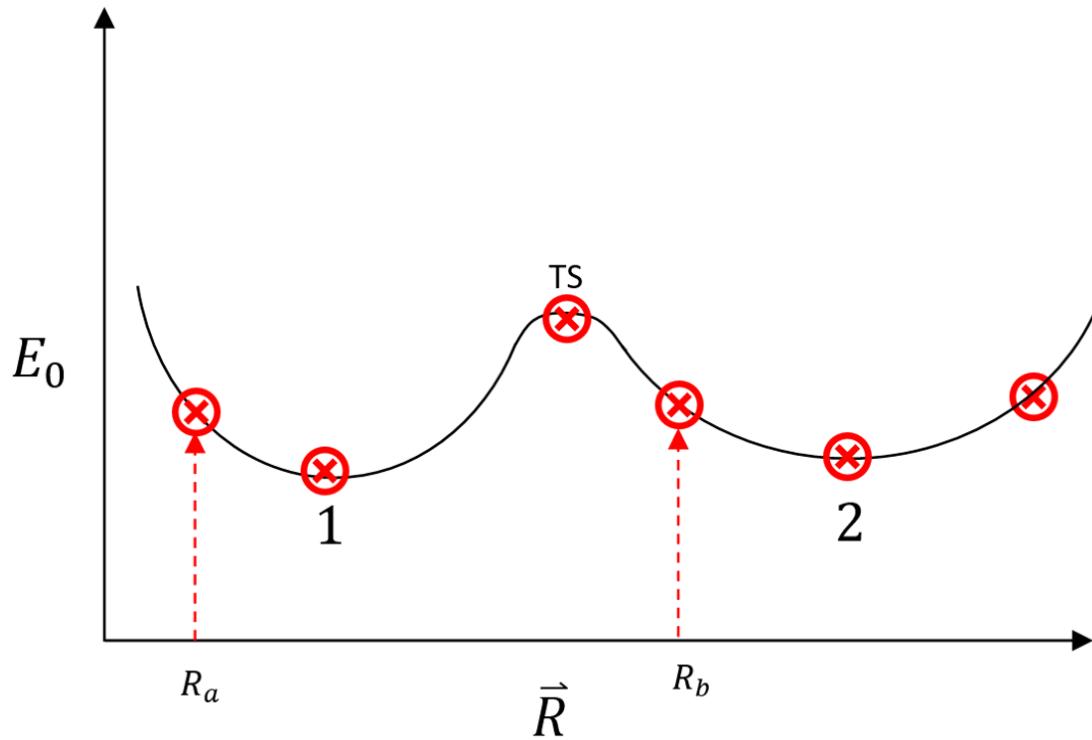


Figure 2.1: A single electronic ground state surface within the BOA leads to sampling and subsequently calculating points at fixed nuclear configurations  $(\vec{R}_a, \dots, \vec{R}_b, \dots)$ .

### 2.1.1 Adiabatic and Diabatic Basis Connection

Earlier discussions on the importance of non-adiabatic dynamics to vibronic models and spectra can now deepen. Although the BOA was highlighted to reduce calculation complexity (i.e. avoiding evaluation of integral terms), it fails when the NACs are relevant—??(a), occurring mainly in two situations:

- i. The adiabatic state character changes rapidly as a function of geometry.
- ii. Degenerate electronic states, separations on the order of vibrational energy level values.

When these conditions exist, Yarkony has described how conical intersections and avoided crossings between separate single excited states develop, and in order to simulate spectra successfully, the model needs to account for neglected terms and be governed by multiple PES to describe nuclear motion [10].

Recall, term **C** contains an integral, which is inversely proportional to energy separation:

$$\int \phi_\mu^*(\vec{r}; \vec{R}) \nabla_\alpha \phi_\lambda(\vec{r}; \vec{R}) \ dr \propto \frac{1}{E_\mu - E_\lambda} \quad (2.11)$$

Ultimately, the Schrödinger equation with the BOA applied appears as:

$$[\hat{T}_N + E_\lambda(\vec{R})] \Psi_{\lambda,n}(\vec{R}) = E_n \Psi_{\lambda,n}(\vec{R}) \quad (2.12)$$

$$\begin{bmatrix} \hat{T}_N + E_1(\vec{R}) + A_1(\vec{R}) & N.A.C. \\ N.A.C. & \hat{T}_N + E_2(\vec{R}) + A_2(\vec{R}) \end{bmatrix} \Rightarrow \begin{bmatrix} \hat{T}_N + E_{11}(\vec{R}) & E_{12}(\vec{R}) \\ E_{21}(\vec{R}) & \hat{T}_N + E_{22}(\vec{R}) \end{bmatrix} \quad (2.13)$$

Obtaining this matrix of diabatic form, it can be utilized to describe potentials and a coupled

set of vibrational wavefunctions  $\chi_a(\vec{R})$ :

$$\begin{pmatrix} \hat{T}_N + E_{11}(\vec{R}) & E_{12}(\vec{R}) \\ E_{21}(\vec{R}) & \hat{T}_N + E_{22}(\vec{R}) \end{pmatrix} \begin{pmatrix} \chi_1(\vec{R}) \\ \chi_2(\vec{R}) \end{pmatrix} = E \begin{pmatrix} \chi_1(\vec{R}) \\ \chi_2(\vec{R}) \end{pmatrix} \quad (2.14)$$

$$\Psi_a(\vec{r}; \vec{R}) = \sum_{\lambda} U_{\lambda a}(\vec{R}) \phi_{\lambda}(\vec{r}; \vec{R}) \quad (2.15)$$

$$S_{\lambda\mu}(\vec{R}) = \int \phi_{\mu}^*(\vec{r}; \vec{R}_{q_0}) \phi_{\lambda}(\vec{r}; \vec{R}_{q_0 \pm dq_i}) \quad (2.16)$$

$$\sum_{\lambda} \int \phi_{\mu}^*(\vec{r}; \vec{R}) \hat{H}_{elec} \Psi_a(\vec{r}; \vec{R}) \phi_{\lambda}(\vec{r}; \vec{R}) dr \chi_{\lambda}(\vec{R}) = \sum_{\lambda} E_{\mu\lambda} \chi_{\lambda}(\vec{R}) \quad (2.17)$$

$$E_{\mu\lambda}(q) = E_{\mu\lambda}(0) + \sum_i^M E_{\mu\lambda}^i q_i + \frac{1}{2} \sum_{i>j}^M E_{\mu\lambda}^{ij} q_i q_j + \dots \quad (2.18)$$

$$E_{\mu\lambda}^i = \frac{E_{\mu\lambda}(\vec{R}_{q_0} + dq_i) - E_{\mu\lambda}(\vec{R}_{q_0} - dq_i)}{2dq_i} \quad (2.19)$$

## **Vibronic models - short summary?**

Possibly have some equations here if you need to refer to them later? Optional

## 2.2 Building Wavefunctions

### 2.2.1 Multi-configuration Self Consistent Field (MCSCF)

To improve on the traditional SCF method, when a linear combination of Slater determinants (instead of single) are used to construct the reference wave function as seen in MCSCF, this enables static correlation to be expressed. MCSCF wavefunctions are useful in cases where electrons show high correlation, such as excited states.

### CASSCF & ORMASSCF

Complete Active Space Self Consistent Field (CASSCF) is one special type of MCSCF. The characteristic feature of CASSCF is that one must proceed to take the entire chosen active space and consider all electronic configurations within it. If there is limitation in the active space, this is Restricted Active Space (RAS). ORMASSCF on the other hand, is a subset of CASSCF, but provides comparable accuracy at a cheaper cost. In order to use ORMASSCF reference wavefunctions in GAMESS, the REFTYP=ORMAS flag must be set within the \$GMCPT group.

## 2.2.2 Active Space Design

Designing the active space is a pedantic problem. To elucidate the active space without prior intuition is a challenging task. Many iterations on assessing the orbitals and adjusting the dimensions of the active space are expected to be performed in order to nail down the active space. This human in the loop part. What does ORMAS address? Cool paper: <https://arxiv.org/pdf/2111.15587#:~:text=ORMAS%20is%20a%20construction%20of, and%20minimum%20number%20of%20electrons.>

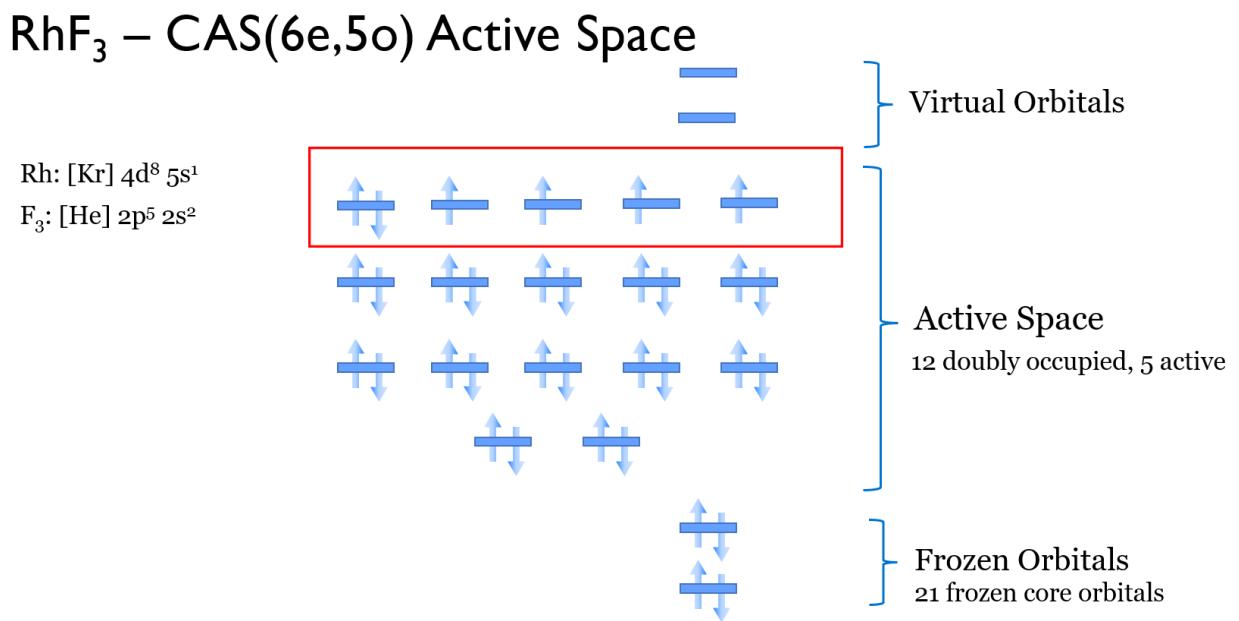


Figure 2.2: Quantized energy level diagram showing orbitals and electrons in a  $\text{CAS}(6,5)$  active space of  $\text{RhF}_3$ .

## GMC-QDPT

Nakano and coworkers have devised the Generalized Multiconfiguration - Quasi Degenerate Perturbation Theory (GMC-QDPT) scheme. GMC-QDPT itself is classified as a composite method that aims to recover static and dynamical correlation alike. The 'generalized' label refers to zero-restrictions in variational space. This means beyond CASSCF and RASSCF,

so general MCSCF wavefunctions are applicable for submission as input. The QDPT latter half indicates a subsequent multireference calculation that proceeds. This is needed because multireference calculations will save the dynamic correlation aspect. Perturbation theory (PT) can "correct the many MCSCF states simultaneously". In short, it is intended for multistate calculations, and optimal for heavy/relativistic calculation purposes. Interestingly, when benchmarked with CAS-QDPT, there is only a 0.1 eV and 1 kcal/mol  $\Delta$ . Nakano has prepared an excellent student reference guide for performing GAMESS GMCPT and GMC-QDPT calculations: <https://ccl.scc.kyushu-u.ac.jp/~nakano/gmcpt.html>. As per the GAMESS input documentation manual, to invoke GMC-QDPT, one needs to use \$GMCPT group in GAMESS (6 letters due to Fortran limit). To achieve the second order quasidegenerate PT correction, one needs to use the MPLEVL=2 and MRPT=GMCPT flags.

An important consideration is that FCI is often impossible for large systems due to the computational cost growing factorially in terms of reference space determinants, so this efficient scheme is more feasible instead. As Rahm remarked, the full level of theory considered here is therefore "a ridiculously high GMC-QDPT ORMAS SOC" <sup>1</sup>.

## Diabatic Molecular Orbitals

Configurational uniformity's implementation in GAMESS is a concept devised by Nakamura and Truhlar to manufacture diabats from adiabats. In their framework, which expands on previous work by Atchity and Rudenberg,a diabatization-adapted molecular orbitals (DMOs) representation is obtained, and then subsequent DSFs (diabatic state functiosn) are constructed. The idea is to optimize for uniformity in electronic configurations that originate from DMOs. Model space

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<sup>1</sup>Prof. Martin Rahm presentation at York University

## 2.3 Jahn-Teller and Relativistic Effects

### 2.3.1 Group Theory and Symmetry

### 2.3.2 E x e JT problems

### 2.3.3 Spin Orbit Coupling (SOC)

SOC scales with Z, heavier the atom, more effect. Since it is a relativistic effect and a small correction to the H, we call this 'perturbation'. Goeppert Mayer's analogy of twirling dancers, clockwise/counter-clockwise

## 2.4 Calculation of Spectra

### 2.4.1 Traditional Franck-Condon

### 2.4.2 Time-correlation functions

### 2.4.3 MCTDH

Consultation with Dieter Meyer of MCTDH

### 2.4.4 VECC

VECC Paper: <https://arxiv.org/pdf/2312.14164.pdf> Our model comparisons have the following assumptions: although we can submit more information to MCTDH than VECC (e.g. off-diagonal coupling of state energies and excited state-to-state transition dipole moments), our comparisons are limited to whatever VECC accepts. Therefore, the difficulty arises from

finding a hybrid format compatible to both dynamics-simulation-schemes that can capture the model. A technical challenge, mainly.

Focusing on VECC, it is potentially advantageous due to its generation of results in a theoretical  $O(n^2)$  time scaling, where  $n$  is the number of normal modes. In comparison, the modern multi-layer MCTDH approach is computationally uneconomical with its exponential scaling and requires extensive user input to dictate efficient calculation. Despite VECC being less systematic, it has the critical advantage of demanding very few user choices; limited to simple on/off switches. VECC's present implementation requires construction of vibronic model Hamiltonians to be plugged in as input, which is likewise an expensive and tedious task.

#### **2.4.5 Thermal Effects: Thermal Field Vibrational Electronic CC (TF-VECC)**

### **2.5 Computational Hardware and Software**

#### **2.5.1 GAMESS-US**

GAMESS is an ab initio quantum chemistry package. The general pipeline of how calculations are run: HF/DFT/GVB/MCSCF  $\Rightarrow$  CI/MP2/CC (these are correlation corrections)  $\Rightarrow$  SOC/DK (these are relativistic effect corrections). One interesting thing about GAMESS is that recently there is possibility for GPU-acceleration.

Fancy method, otherwise why not just DFT B3LYP 6-31G\* @ everything ...

## 2.5.2 nlogin, ComputeCanada (Computational Resources)

On nlogin, the available compute is as follows: cpu[005-006] intel,westmere, X5670 12 48000  
io003 intel,ivybridge,E5-2 16 128000 io[001-002] intel,westmere,X5670 12 48000

On ComputeCanada Graham cluster, the available compute per node is 32 CPUs with 125GB of memory.

The GAMESS memory requirements for parallelized calculations is as follows:  $((\text{mwords} * \text{p} + \text{memddi}) * 8) / 1024 = \text{TOTAL NUMBER OF GB}$ . Where p is the number of processors, mwords is the amount of memory (in Words) for each processor, memddi is the distributed memory for optimization runs, and 8/1024 to account for the conversion of words to gigabytes. Therefore, on nlogin, the max amount of compute per run would be 16 CPUs at 4 gigabytes. For reference, see <https://docs.alliancecan.ca/wiki/GAMESS-US>.

To compile Zeng's GMC-QDPT compatible GAMESS version on ComputeCanada clusters, it is complicated. First, the *gmcpt.src* module must be obtained from Zeng. The next step is to *module load intel/2023.2.1* and run *./config* in GMS' source directory; this will prompt a set of questions in which the answers are as follows: linux64, enter, enter, enter, ifort, 10, none. After you must enter the ddi/ directory and run *./compddi* and move the ddikick executable one directory above. Finally, you must run *./lked* to link GAMESS to your path.

# Chapter 3

## Diabatization Code Development

The newly developed diabatization code will be deconstructed, along with how it is used to produce vibronic models and compute spectra. As previously described in the project scope (1.3), this is my main contribution, and the centerpiece of the project. The process of diabatization can be considered to have two broad stages: first, we use Zeng Protocol MCSCF to obtain a ‘template’<sup>1</sup>, and then we proceed to diabatize the system’s Hamiltonian. I briefly outline the individual steps of the MCSCF calculations in Section 3.1. Explanation presented in Section 3.2 shows how the machinery of the code performs the diabatization. In addition, I describe in detail the practical challenges encountered, as well as their solutions, in automating this extensive and unwieldy process. Spectroscopy related workflows will be discussed in Section 3.3, with the final Section 3.4 briefly touching on extra details.

### 3.1 MCSCF Pipeline

The MCSCF process can be broken down into four individual steps, with some steps requiring very specialized knowledge of GMC-QDPT and ORMASSCF intricacies<sup>2</sup>.

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<sup>1</sup>A reference geometry model and some assorted parameters

<sup>2</sup>Consultation of Nakano’s student guide [11] is highly advised.

The approach taken here is a slightly modified version as seen in Zeng's work [cite](#).

1. Appendix [5.1.1](#) Geometry optimization + frequency (Hessian)
2. Appendix [5.1.2](#) GMCPT
3. Appendix [5.1.3](#) DMO calculation
4. Appendix [5.1.4](#) REFDET calculation

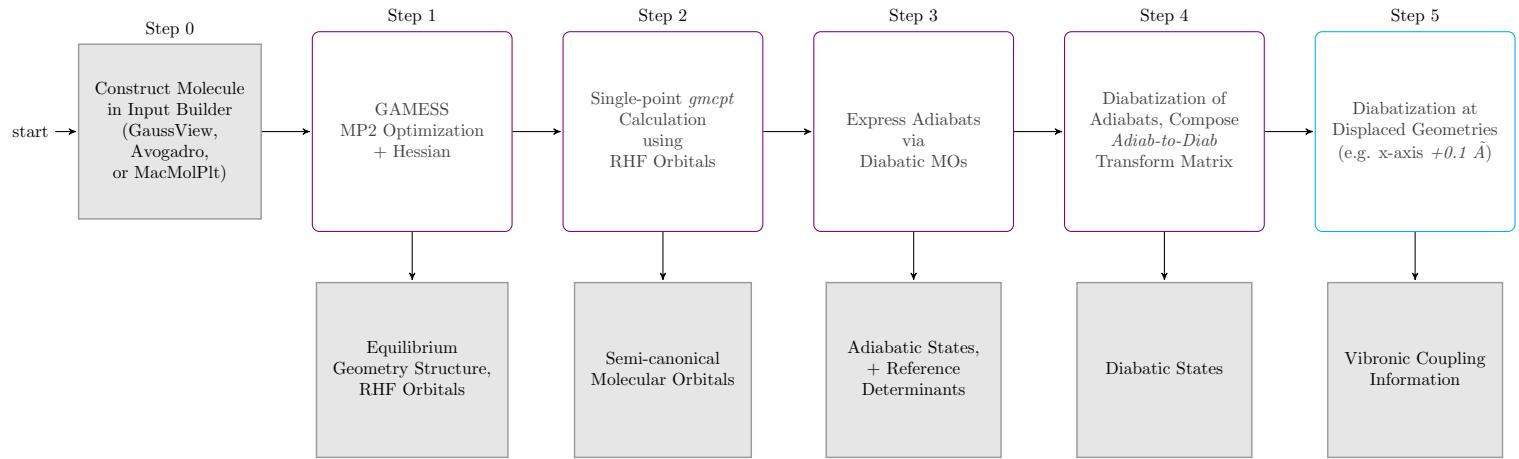


Figure 3.1: GMC-QDPT ORMASSCF vibronic coupling model scheme formulated by Zeng.

From routine usage of the Zeng protocol, there were inconvenient parts such as Step 3 and 4, where it can be tedious to compose the input needed. It typically requires extracting out information from the preceding calculation's punched out orbitals or energies. We have thus created a package of collected Zeng protocol scripts where users can select from a list of 8 options to assist in running through the protocol. These scripts have saved an incredible amount of time and frustration. A walk through of each step for an example  $\text{RhF}_3$  SOC system is provided in the appendix. Once Step 4 is completed, vibronic coupling information at reference geometry is directly available, and the Step 4's input file can serve as a template for the diabatization to proceed.

### 3.1.1 Spin Orbit Coupling: Imaginary Changes Everything

SOC enters into the described procedure at Step 2 by setting the individual `$transt` flag. Another relevant flag to include here would be `knospn` in the `$GMCPT` group, which can turn off the system's intended spin states. As soon as these two flags are plugged in, the SOC calculation will be on its way and provide Level 1 and 2 coupling values between off-diagonal state pairs. Beware that these SOC values are printed in a dual real + imaginary component, which is the first such occurrence of imaginary in the protocol (besides a pesky imaginary frequency in some symmetries).

## 3.2 Automated Machine Gun Diabatization Code: *'dist\_allmodes\_pm.py'*

The current code can calculate up to quadratic order coupling constants. Majority of the code implementation was done in collaboration with Neil Raymond, and is publicly available on GitHub: <https://github.com/bjb2chen/VEmodel>. At this point, the intended use of the diabatization code is to take the template constructed by the Zeng protocol and generate the vibronic model Hamiltonians in a suitable manner for both VECC and MCTDH. Much of the output format for *dist* is setup to be connected to Raymond's project scripts and in a mutable style based on model parameters. The user must carefully set the pivotal parameters for diabatization, such as a selected mode list and the multiplication order at each mode. For example, `x2` is square on displacement size (`qsize`) at a mode, `x3` would be cubed `qsize`. Triple checking of parameters is recommended before executing the code.

The code now takes care of all these subsequent steps automatically:

1. Section 3.2.1 Extraction of frequencies and reference structure (refG) coordinates, storing them inside dictionaries.
2. Submit refG calculation.
3. Precomputation of distortions and subsequent submission of calculations in 'machine gun' fashion.
4. Collection of diabatic vibronic model data.
5. Fitting PES plots for higher order modes (if necessary).

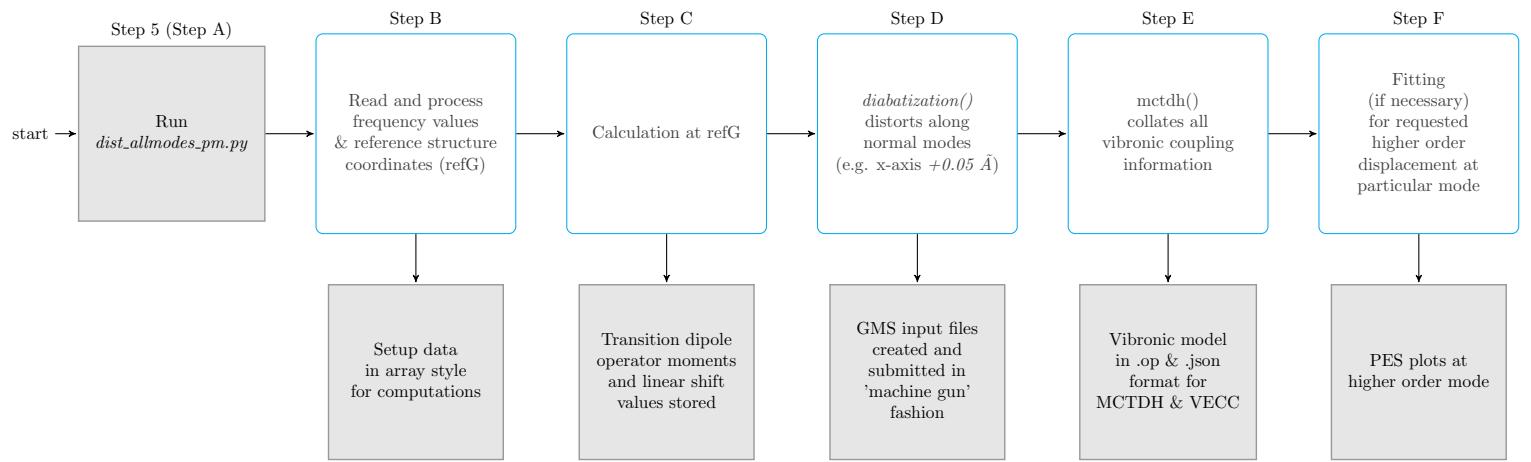


Figure 3.2: '*dist\_allmodes\_pm.py*'

### 3.2.1 Frequencies and Reference Structure

Frequencies must be extracted from the reference structure geometry and stored in arrays.

```
1      to do.
```

Listing 3.1: t.b.d.

1	##### SAMPLE REF STRUCT #####
2	N 7.0 0.0000000000 -0.0000000000 -0.1693806842
3	H 1.0 -0.4653267700 0.8059696078 0.2564602281
4	H 1.0 -0.4653267700 -0.8059696078 0.2564602281
5	H 1.0 0.9306535400 0.0000000000 0.2564602281

Listing 3.2: Sample reference structure

### 3.2.2 Step 2: Submit diabatization calculations Displacements

Special details about +1/+2/+3/+4

Special details about selecting certain modes to have +8 etc

Loops in python are notoriously slow. We refactored the excessive looping that distorted the modes into array operations. By vectorizing the calculation, we have advanced an efficient scheme.

```
1 # -----
2 # NEIL MAGIC CODE (^o^)
3 displacements = {
4     "+1": reference[:, NEW] + 1.0 * R_array[NEW, :] * mode_array[:, :],
5     "-1": reference[:, NEW] - 1.0 * R_array[NEW, :] * mode_array[:, :],
6     "+2": reference[:, NEW] + 2.0 * R_array[NEW, :] * mode_array[:, :],
7     "-2": reference[:, NEW] - 2.0 * R_array[NEW, :] * mode_array[:, :],
8 }
9 #           (Z*3, 1)           (1,   N)           (Z*3, N)
10 #          (9, 1)            (1,   3)            (9,  3)
11 #          (12, 1)           (1,   6)           (12,  6)
12
13 # -----
```

Listing 3.3: Precomputing linear displacements

### 3.2.3 Step 3: Harvest results

SOC flag

1. (1) OPTIMIZED RHF
- (2) OPTIMIZED ROHF
- (3) Semi-canonical MOs
- (4) MCSCF Natural Orbitals
- (5) OPTIMIZED MCSCF
- (6) DMO group
- (7) REFDET GROUP
- (8) EQUILIBRIUM GEOMETRY

Loops in python are notoriously slow. So in collaboration with Neil Raymond, we refactored the excessive looping that distorted the modes into array operations. This drastically improved the performance. Furthermore, it also made the entire program more dynamic in a sense that it can adapt and be extended for a variety of purposes. An example here is fitting. This situation arose where distortion beyond-quadratic-order at a single mode was needed, and the recasted arrays handled this trivially as compared to looping. Another key aspect of this change was the fact that the arrays forced us to understand exactly what dimensions are required in the diabatization code's various calculation modules. We have stitched together various functionalities such as: text processing, coordinate distortion, input file generation, GAMESS job submission, and operator file generation.

PPT idea: use opacity overlay to highlight and gray out parts of the screen for emphasis.

We originally opted to extract data using the slow pythonic 'read' method. However, this proved to be debilitating when operating on large output files to a point where the bottleneck could be hours when running the script. Therefore we opted to use grep instead. When it came down to the full computational study on FeCO, surprisingly even grep was deemed too slow. The average size of the FeCO output file (there are ~1000 outputs) is 70 megabytes or around 2 million lines of text. When we grepped the SOC data, it took around 4 hours. So finally we pivoted to using memory-mapping. This not only made the program more performant, reducing the runtime of the script processing all 28 modes of FeCO from 10300 seconds (with grep) to a mere hour.

If diabatic energy, we need to extract along the diagonal of the states array (a,a). If diabatic coupling, we need to extract along the off-diagonal of the upper triangular states array (a,b). This needs to be performed for linear, quadratic, and bilinear.

Update on May 16/17th work: Neil helped make it read JSON. We want to do SOC for VECC. We found incredible difficulty in making sure the JSON reads the .op model correctly due to a load of issues.

Profusely

Toby Zeng created a custom DMO using symmetrized reference geometry orbitals. 39  
(a1'), 45,46 (e") 47,48 (e') 51 (a1'\*) 53,54 (e'\*) became the SOC's 44, 45,46, 47,48, 49, 50,51.

We have discussed the prospects for ... We have advanced the prospects for an efficient scheme ... A sign of things to come ...

FeCO was unstable. To stabilize the converger, we had to build up the molecule's active space bunch by bunch.

### 3.2.4 diabatize() and mctdh() functions

### 3.2.5 Distorting coordinates and calculating up to quadratic coupling constants

### 3.2.6 Fitting on n-th order models

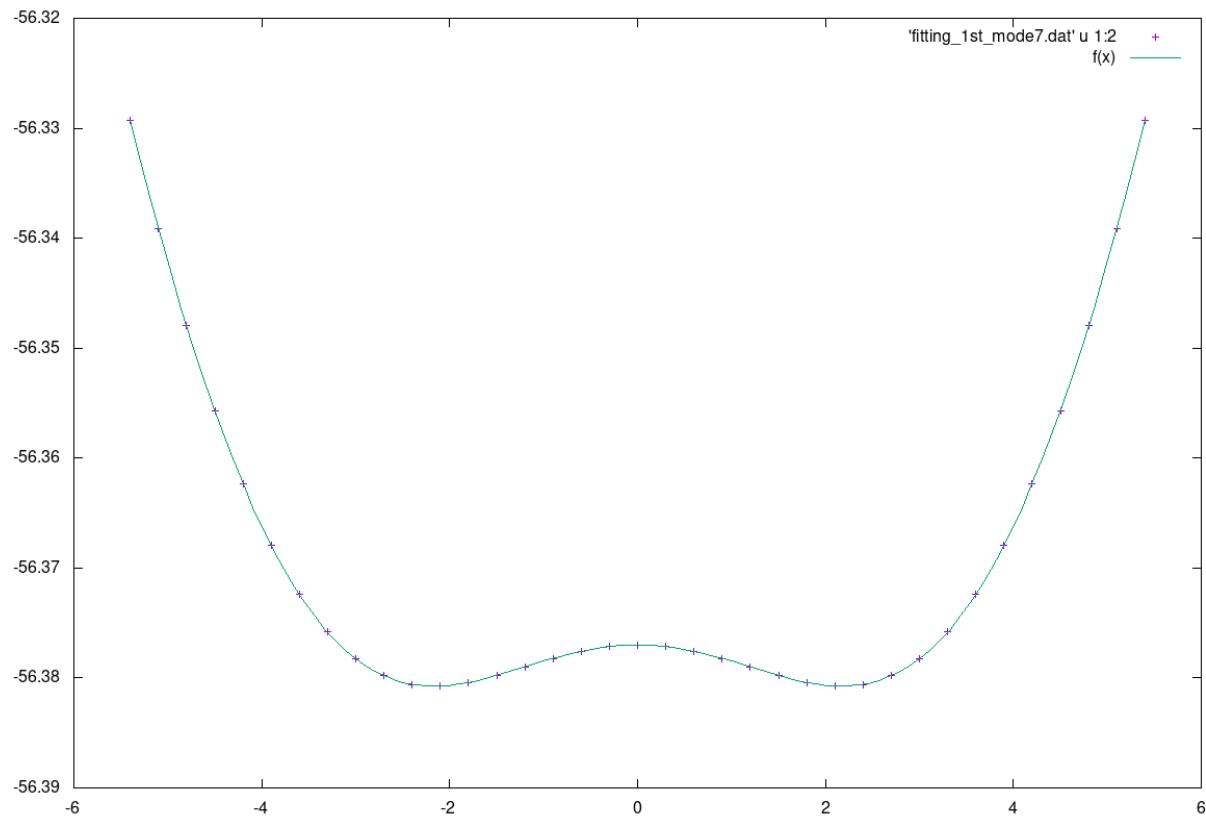


Figure 3.3: PES of  $D_{3h}$  NH<sub>3</sub>, Mode 7, 10 Displacements, qsize = 0.30

Bells and whistles

### **3.3 Generating spectra via MCTDH and VECC**

To crack this problem, rewritten program with Neil. The program is not perfect, but we can still use it to do relatively functional calculations and build models.

Negative magnitude in spectra = too high of a tau Useful for MCTDH and VECC alike.

### **3.4 Assorted comments and considerations**

Human custom input is needed to setup the first four steps in the MCSCF component, which revolves around dictating parameters such as active space and reference geometry; step five is where the Python automated code enters to fulfill diabatization. Most of the pertinent information relating to the machinery will be covered. The code can fulfill up to quadratic order coupling constants. After the diabatization calculations are complete, vibronic Hamiltonians can be harvested, wavepackets propagated, and spectra generated. Stumbles, hurdles, obstacles...

### **3.5 In Summary**

# Chapter 4

## Realistic Applications to Systems of Interest

### 4.1 Details about running calculations

The spectra simulations shown here have many parameters that go into influencing their overall shape and appearance. That is to say, they are highly influenced by whichever input model parameters are given and tuned. The interplay of the parameters sometimes even causes unphysical results (negative intensity) or incompatibility. The most important factors include symmetry of the original system, ionic or neutral character, and inclusion of quadratic and higher order terms. However, an important reminder here is that the objective of the research is to not produce beautiful (i.e. resolved, coherent) spectra, nor accurate to experimental. The sole focus is to demonstrate the vibronic model process and its ability to produce demonstrably different results when extra parameters are inputted into the Hamiltonian. A prime example is a case where say a combination of lower SPF/PBF yields a more resolved spectra. Although this might be the case, it is not a good spectra due to its lacking of rdgcheck and rdgpop results. Additionally, in that case we are only sampling the boundary and it is not correct. FINETUNING IS EVERYTHING

## **4.2 Details about producing spectra**

A 3–5 sentence paragraph that briefly describes (cite wikipedia or a textbook or something) the basic idea of producing spectra. As previously described (ref to introduction/background somewhere) spectra are produced by fourier transforming auto correlation function (ACF). I used MCTDH's `autospec86` program to do this. Explain more.

Highlight the parameters and how they're connected to the equation (show equation from `autospec` documentation) (borrow stuff from Neil's thesis)

JULY 7th: THE WHOLE PROBLEM FOR DISCREPANCY BETWEEN MCTDH AND VECC WITH SOC IS DUE TO cross (NOT auto) IN INPUT FILE!!!

## **4.3 any other random stuff you need to explain before the systems?**

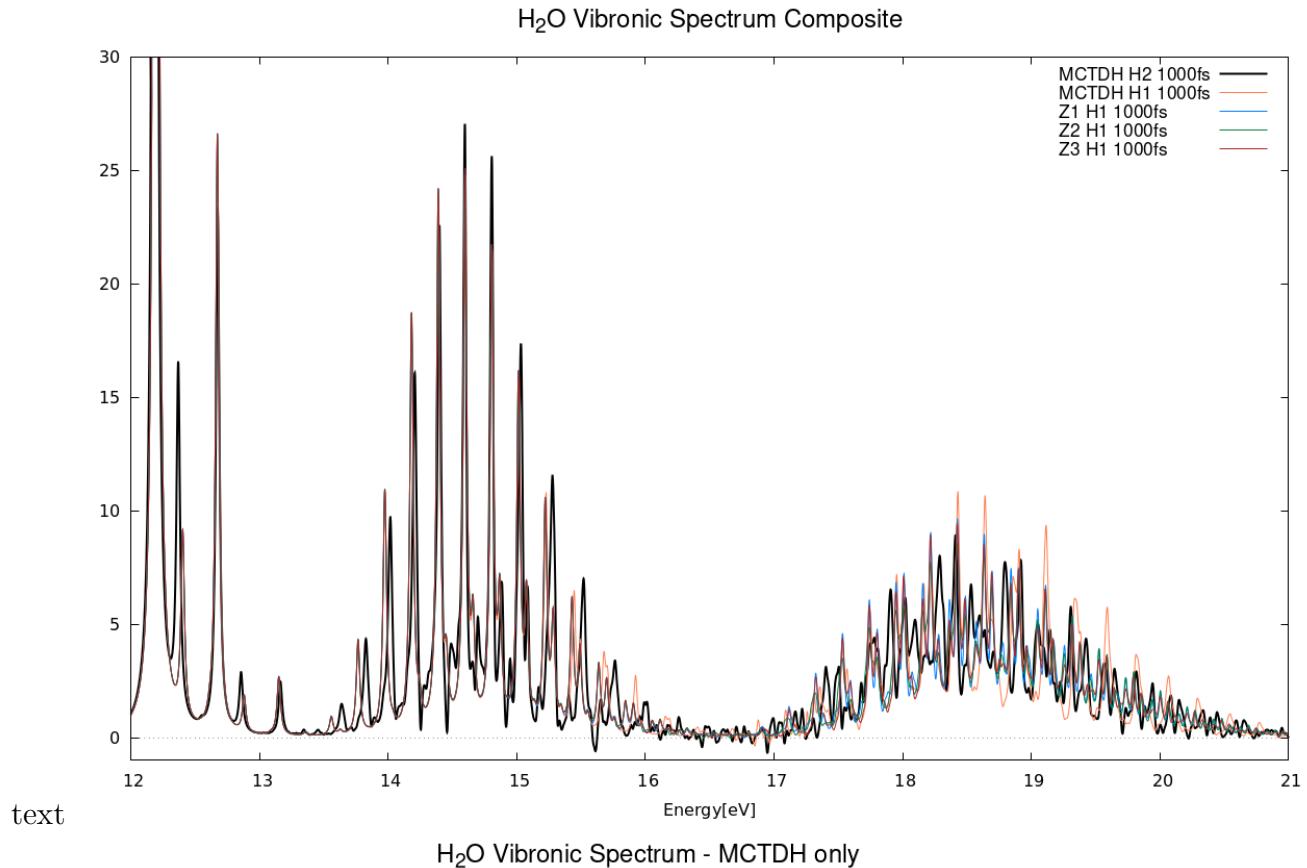
## **4.4 Preliminary systems**

brief description of each system and results you expect and this is to prove that it has no issues in easy cases.

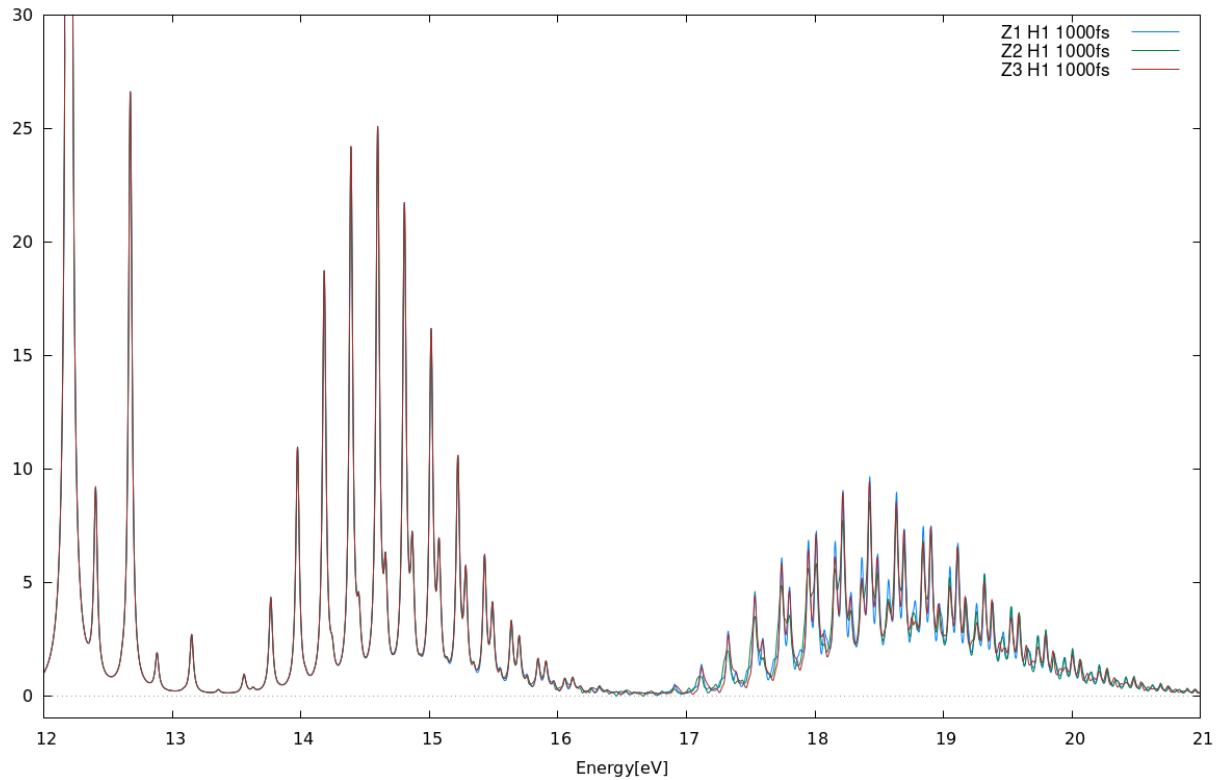
Here I show results for  $n$  systems which are all relatively small. The purpose of these systems is to show  $|x\rangle$ . In particular Water is interesting because ...



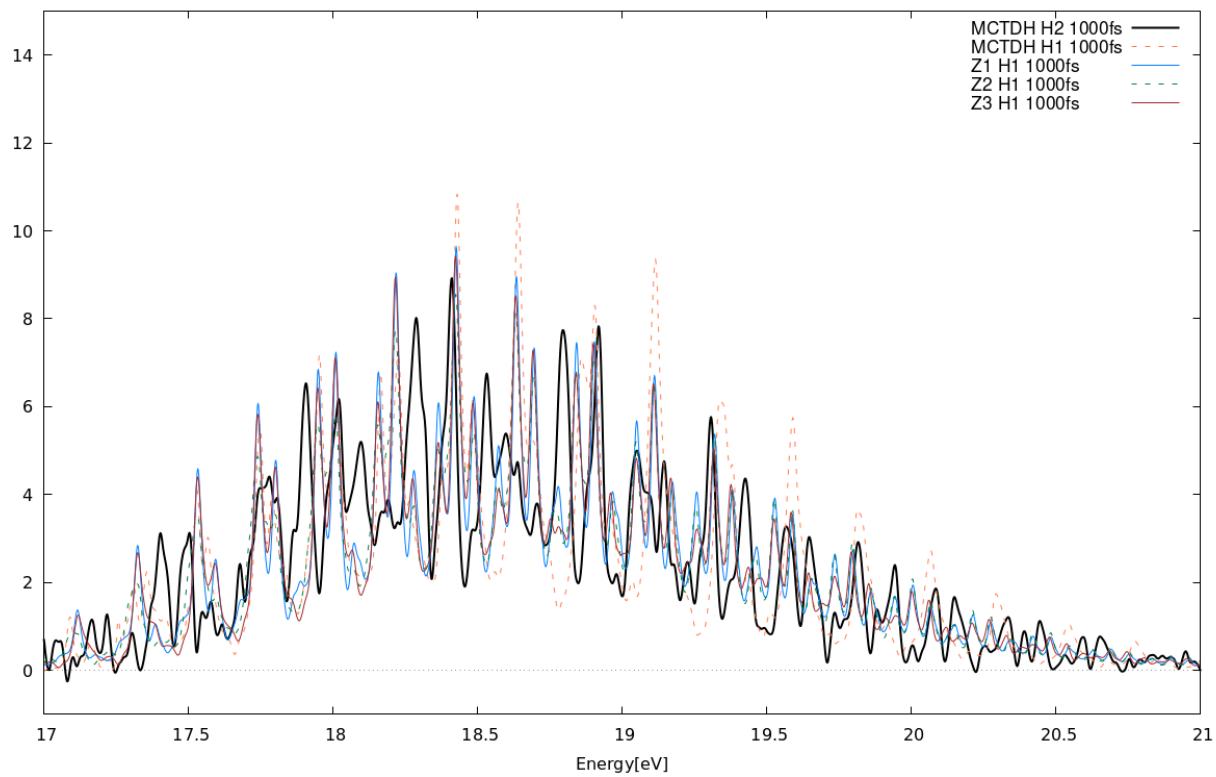
#### 4.4.1 H<sub>2</sub>O



H<sub>2</sub>O Vibronic Spectrum - VECC only



H<sub>2</sub>O Vibronic Spectrum - 17 to 21 eV

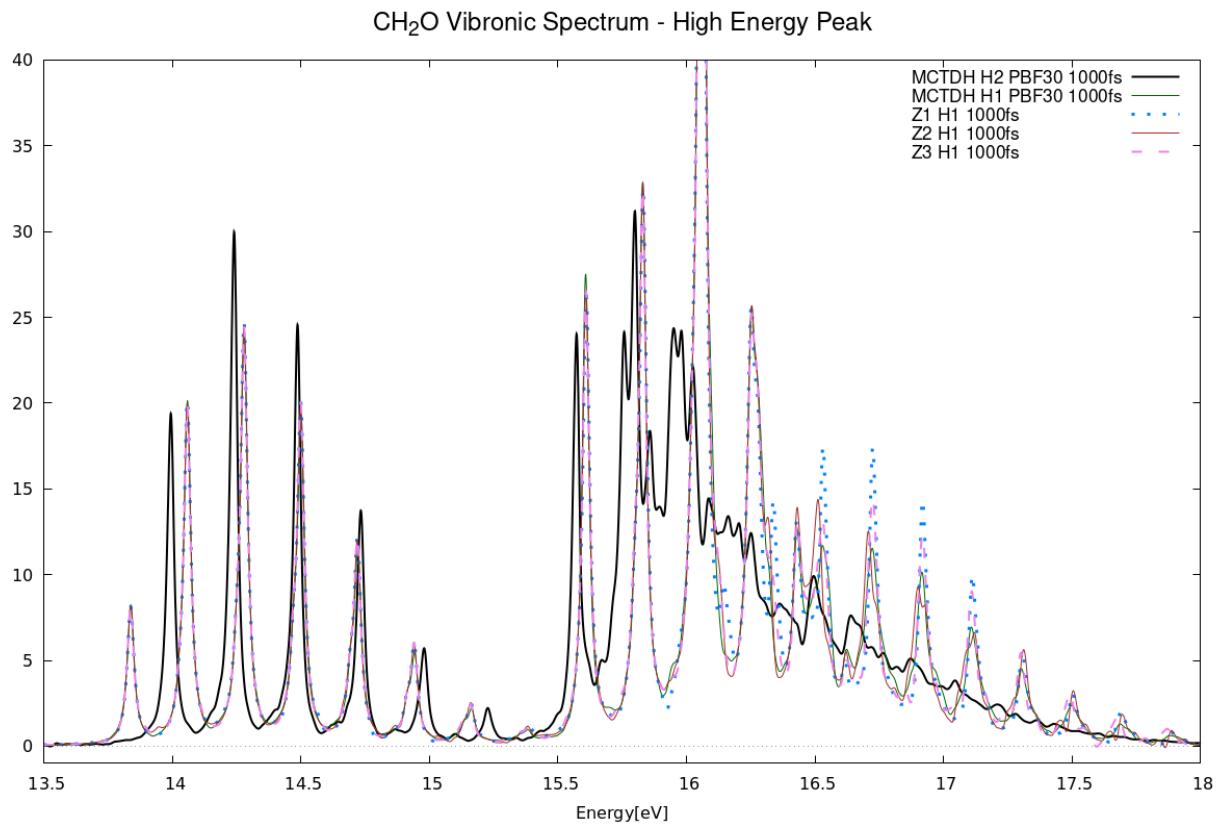
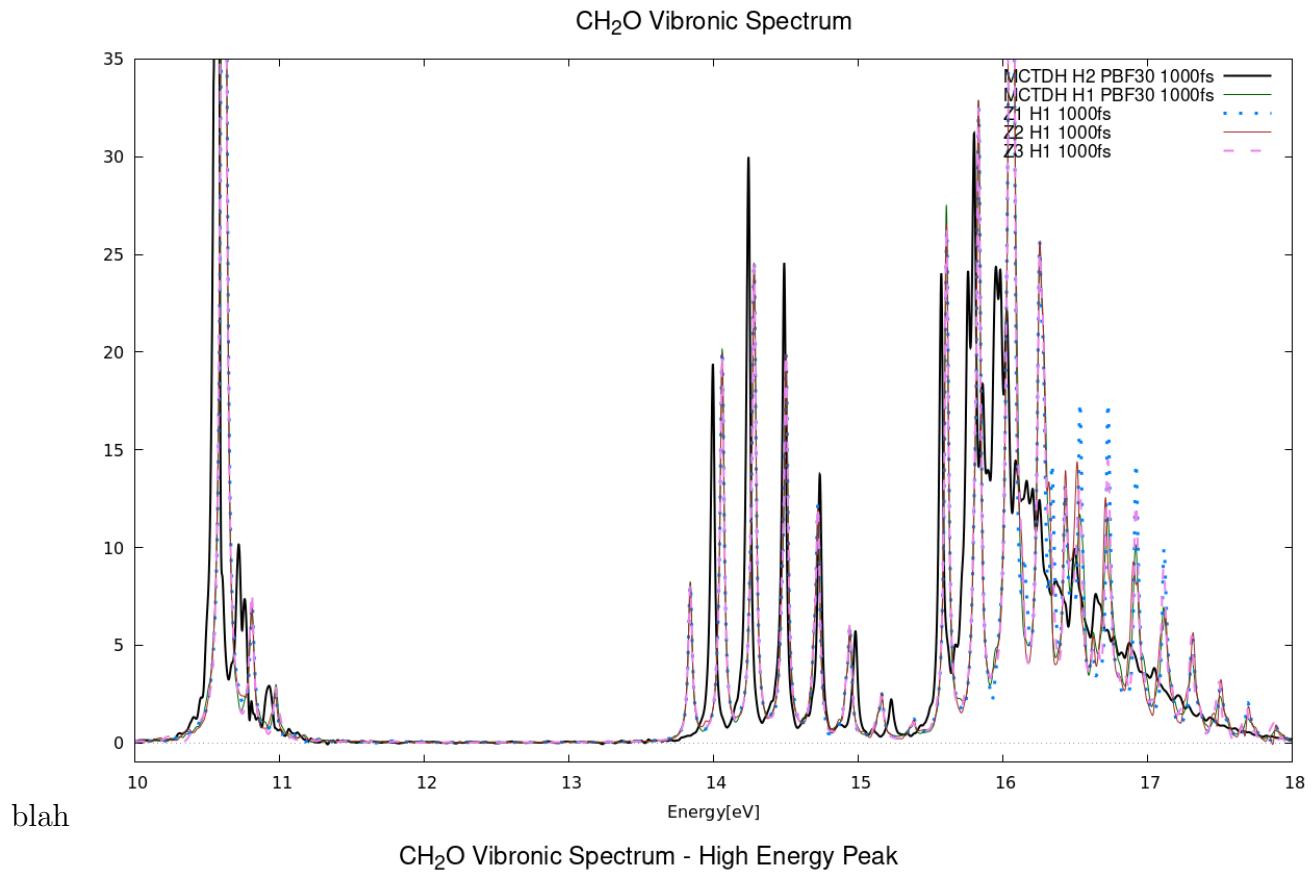


Water was the first molecular species investigated in this research project and was the original test case for running GAMESS calculations. The active space of  $H_2O^+$  is 6o7e.

```
1 nmofzc=0 nmodoc=1 nmoact=6 nelact=7 mstart(1)=2 icharg = 1 mult = 2  
gbasis = ccd state(1) = 4
```



#### 4.4.2 CH<sub>2</sub>O

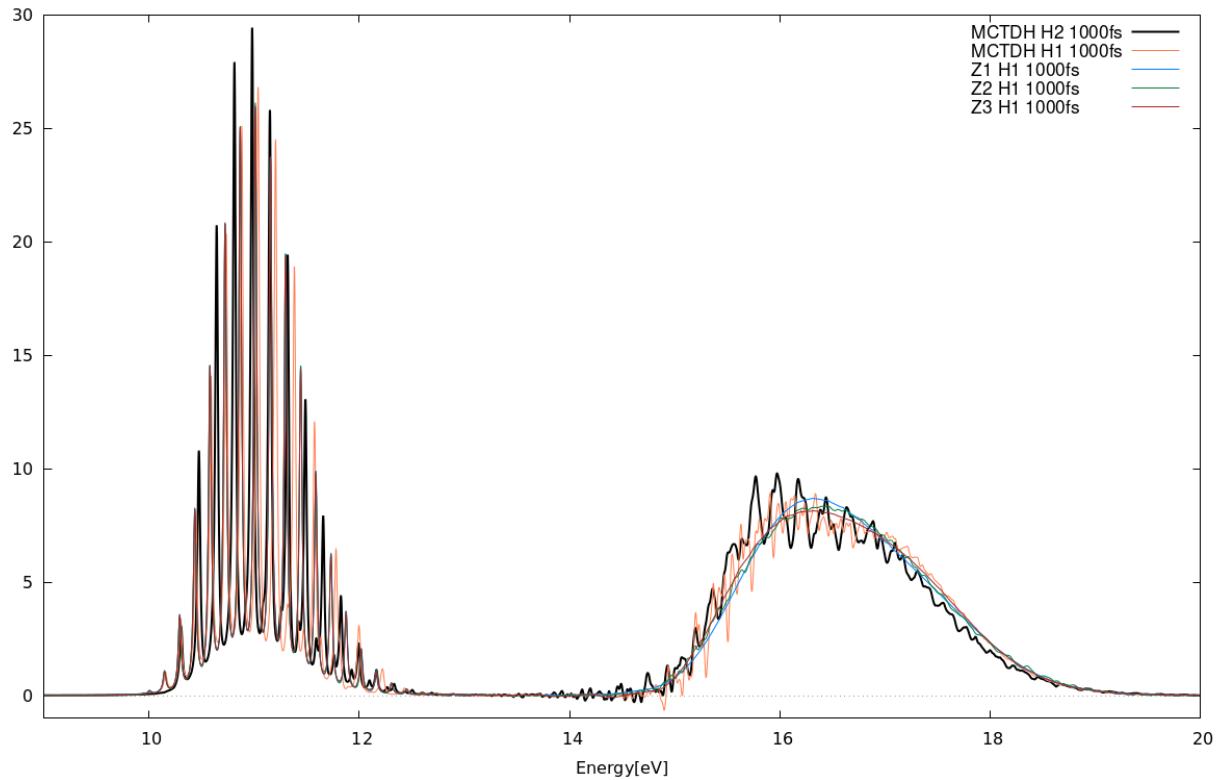


add short description of file show table with Hamiltonian parameters show image of molecule?  
(probably unnnecessary) show spectra from MCTDH and VECC (maybe show constant, linear, quadratic, SOC?) discuss the results? are they identical, good this system was just top test xyz, and it behaved as we expected.

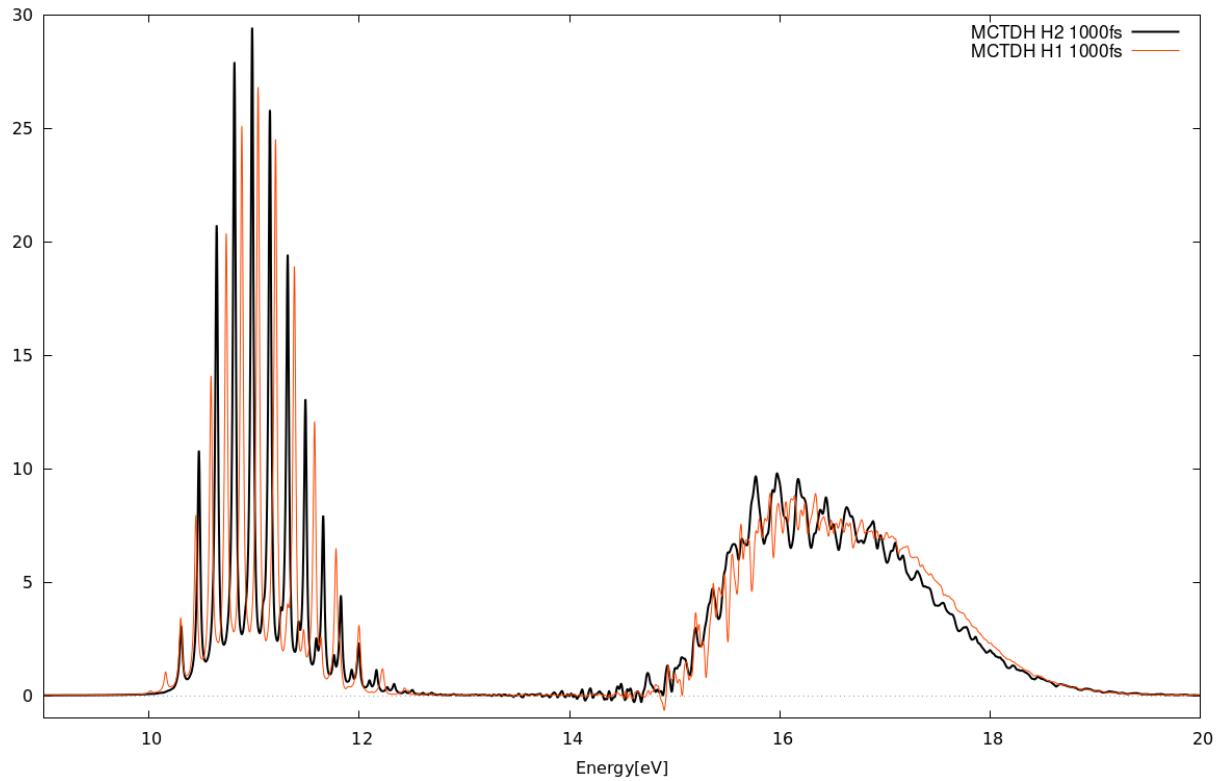
#### 4.4.3 NH3

blahj

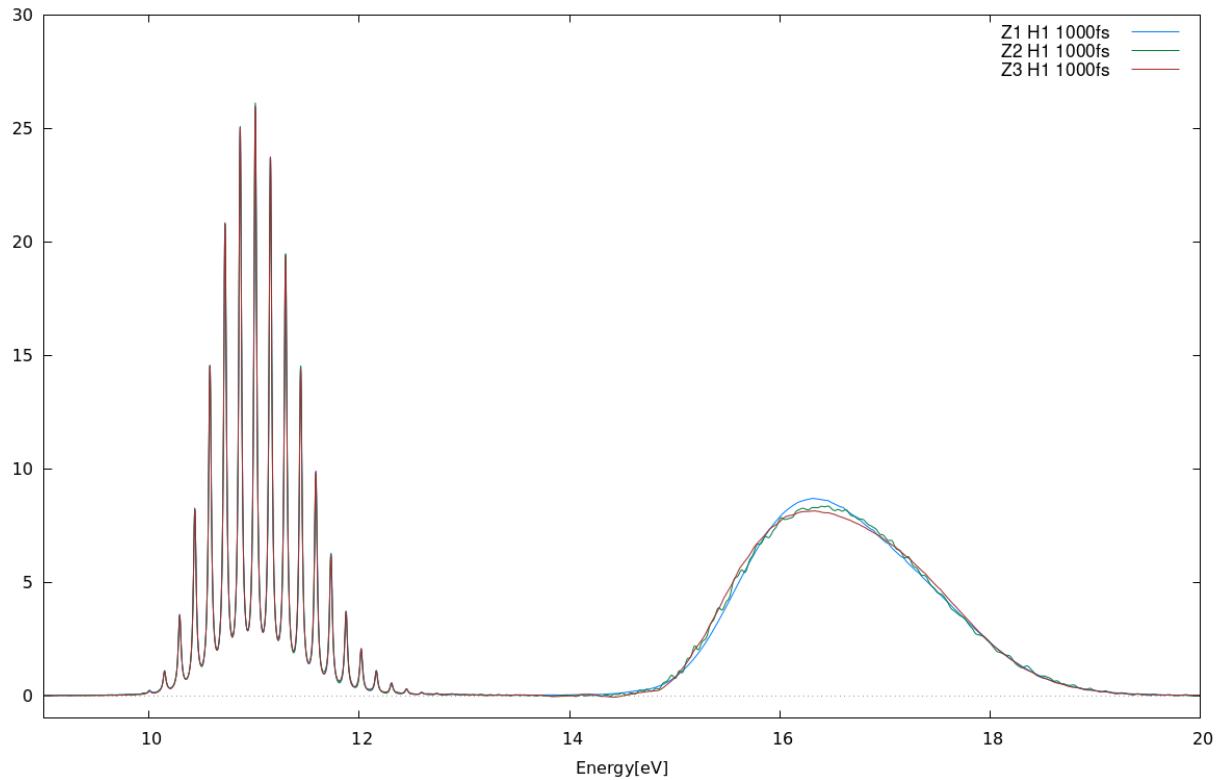
NH<sub>3</sub> Vibronic Composite Spectrum



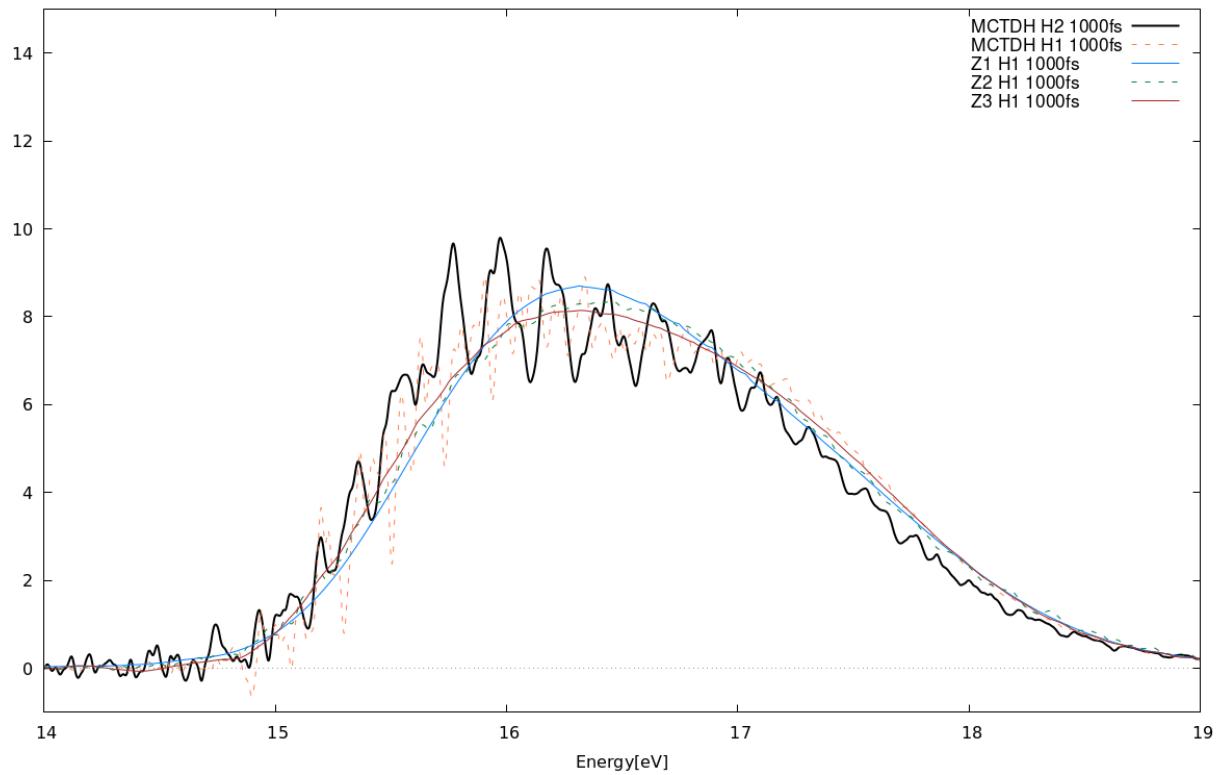
NH<sub>3</sub> Vibronic Spectrum - MCTDH only comparison



### NH<sub>3</sub> Vibronic Spectrum - VECC only comparison



### NH<sub>3</sub> Vibronic Spectrum Comparison - 14 to 19 eV

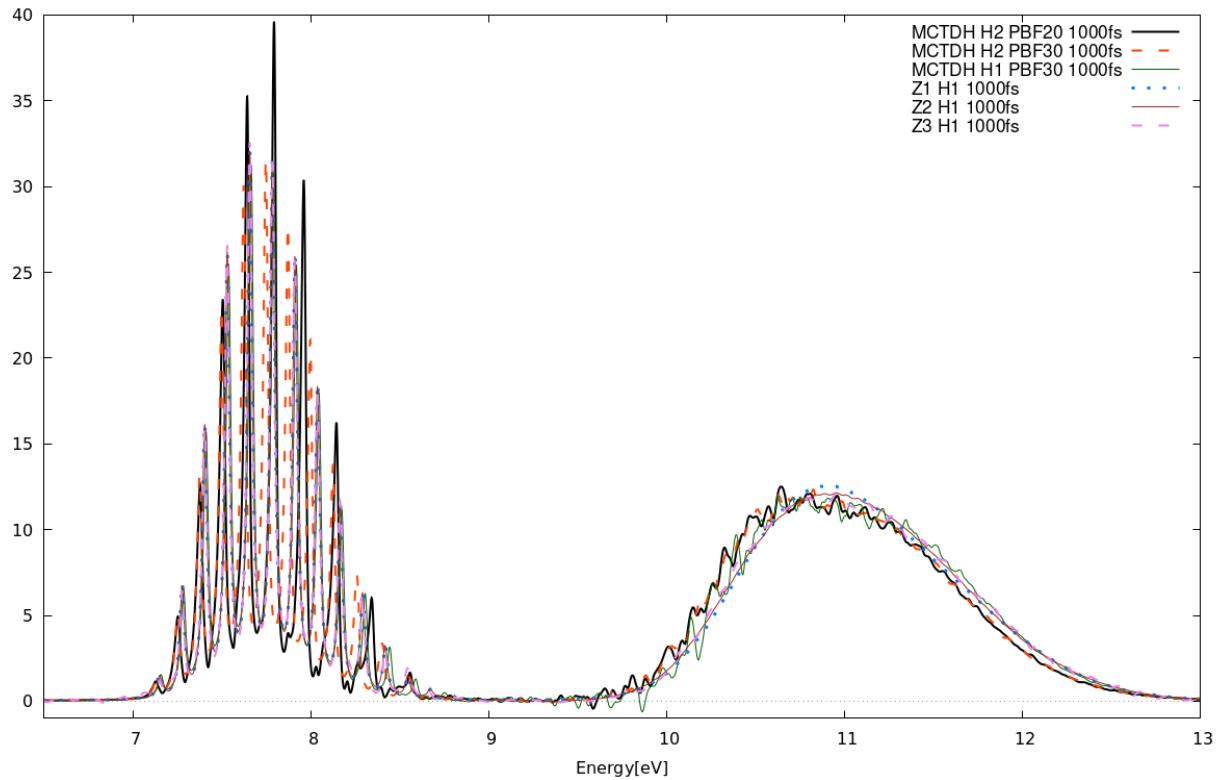


add short description of file show table with Hamiltonian parameters show image of molecule? (probably unnecessary) show spectra from MCTDH and VECC (maybe show constant, linear, quadratic, SOC?) discuss the results? are they identical, good this system was just top test xyz, and it behaved as we expected.

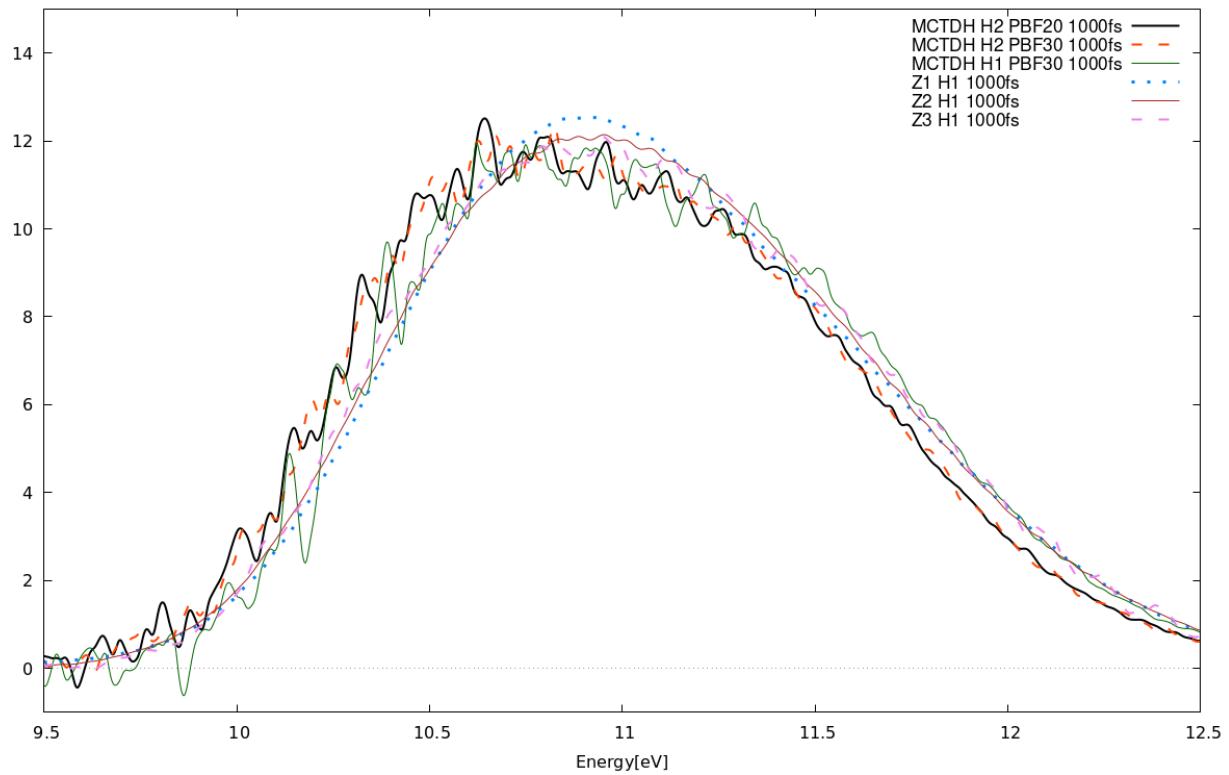
#### 4.4.4 PH3

blah

### PH<sub>3</sub> Vibronic Spectrum



### PH<sub>3</sub> Vibronic Spectrum - High Energy Peak



## 4.5 Complex systems

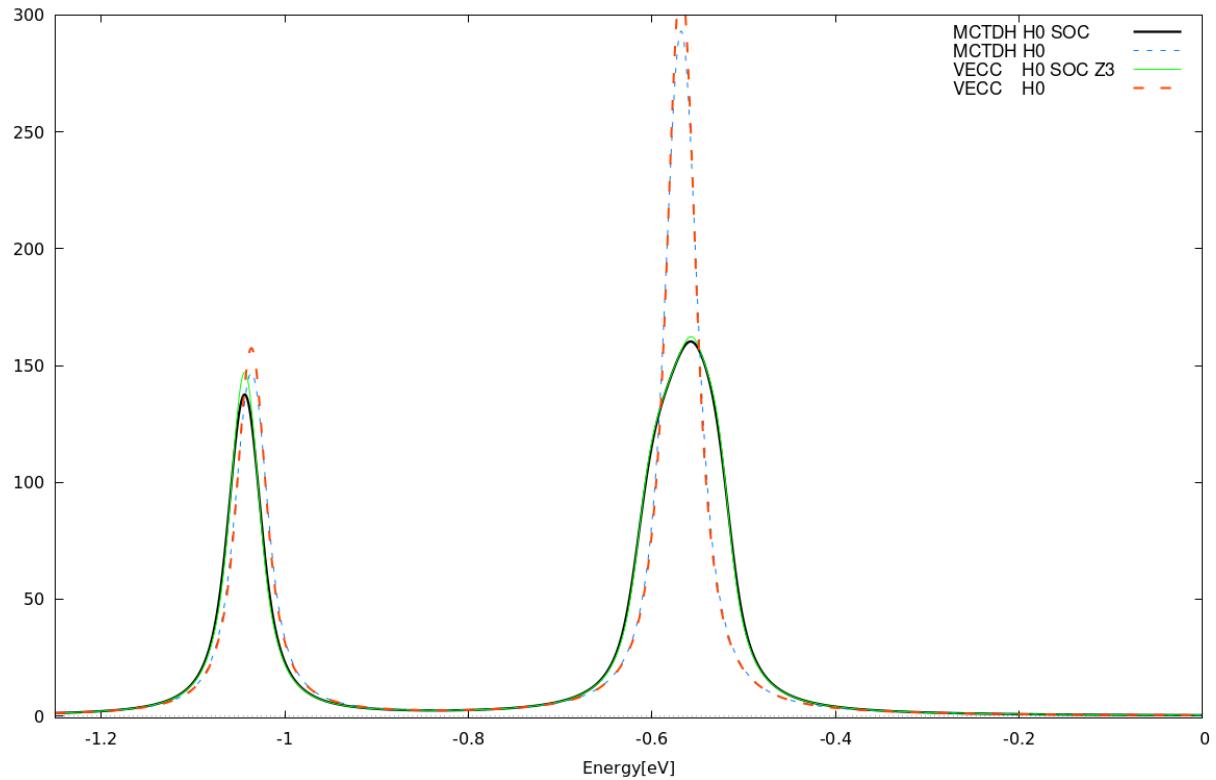
As we explained in (section intro) transition metal complexes are of high importance because *some reason*. We primarily focused on CoF<sub>3</sub>, RhF<sub>3</sub> and Fe(CO)<sub>5</sub> because of *some reason*. The two smaller systems (CoF<sub>3</sub>, RhF<sub>3</sub>) are good for benchmarking the model generation as well as fine-tuning the calculation parameters for MCTDH and VECC. The large system (Fe(CO)<sub>5</sub>) is expected to be very difficult and has only some citation of a paper and demonstrate how its real tough and scary

To align VECC and MCTDH, it was a monumental task to get things right. We had to consult with Dieter Meyer! Finetuning. Hurdles as Marcel would say... Once they did align, it was a wonderful moment, almost breathtaking. IMAGINARY CHANGES EVERYTHING! Concerted effort to figure it all out. Full capacity to do SOC

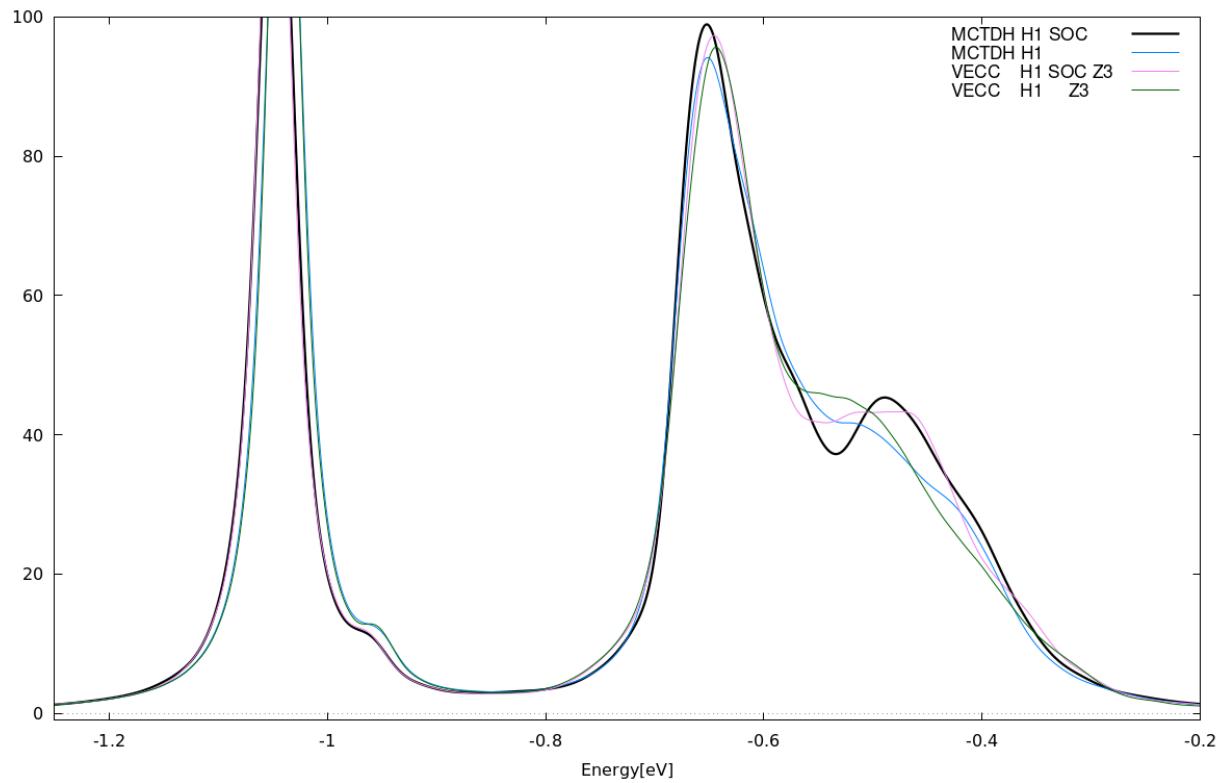
### 4.5.1 CoF<sub>3</sub>

blah, no empty page

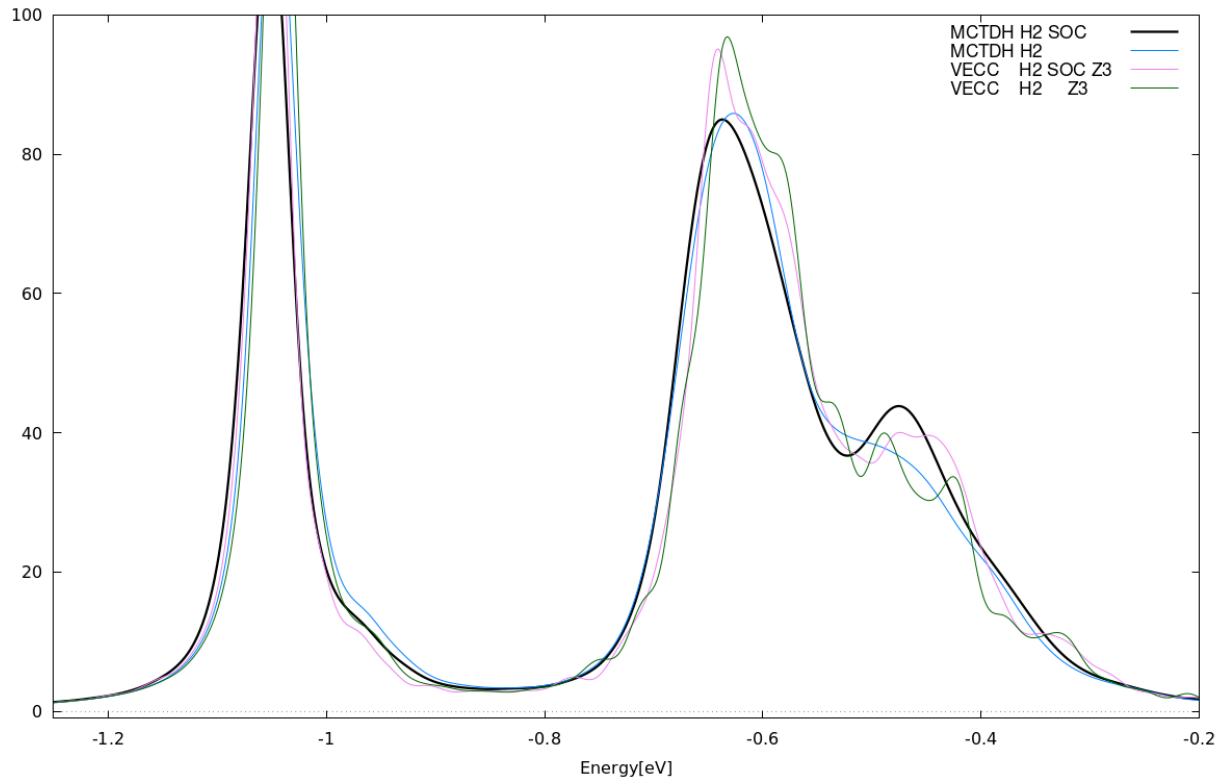
CoF<sub>3</sub> Vibronic Spectrum Composite - Constants



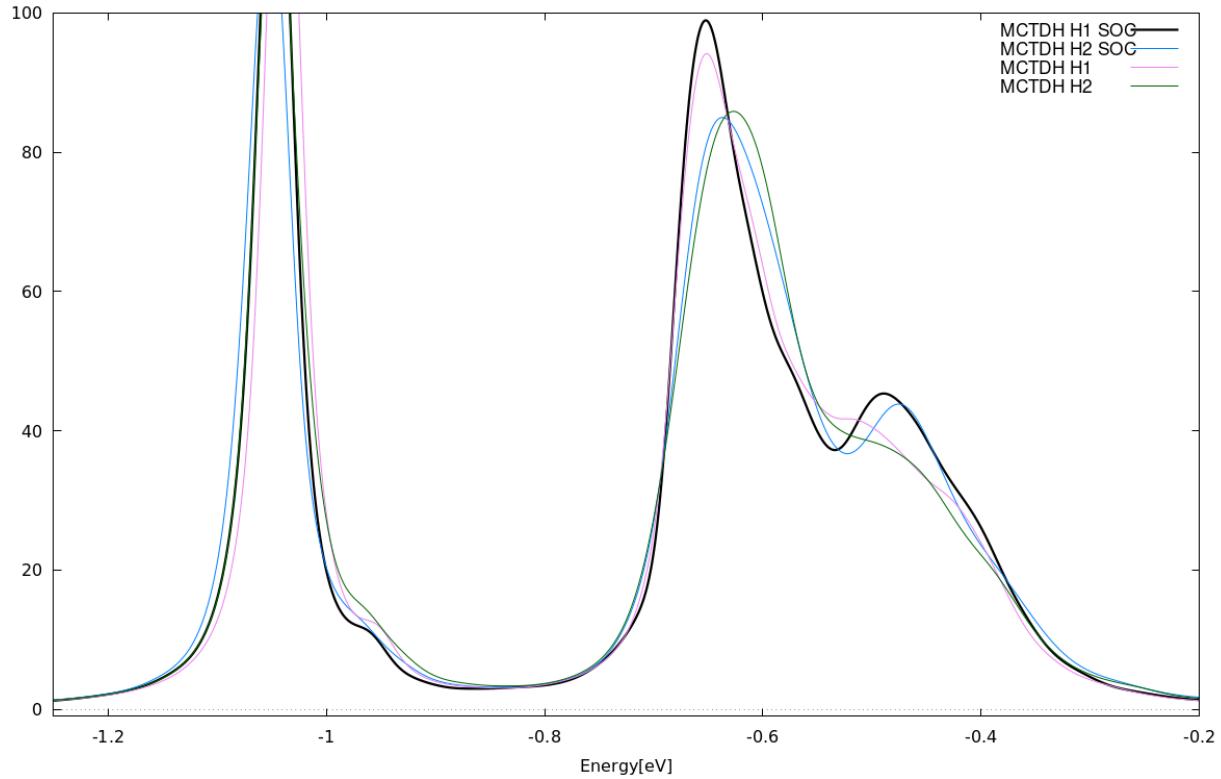
CoF<sub>3</sub> Vibronic Spectrum - Linear Model Comparison



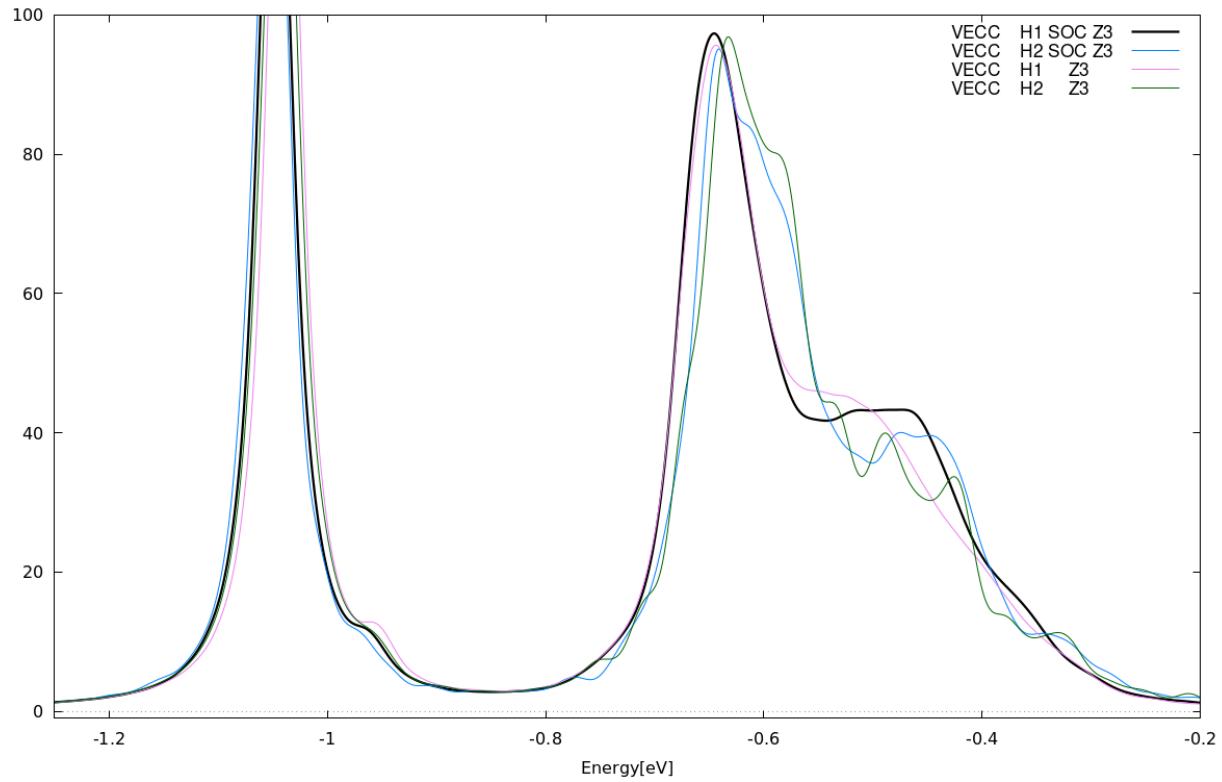
CoF<sub>3</sub> Vibronic Spectrum - Quadratic Model Comparison



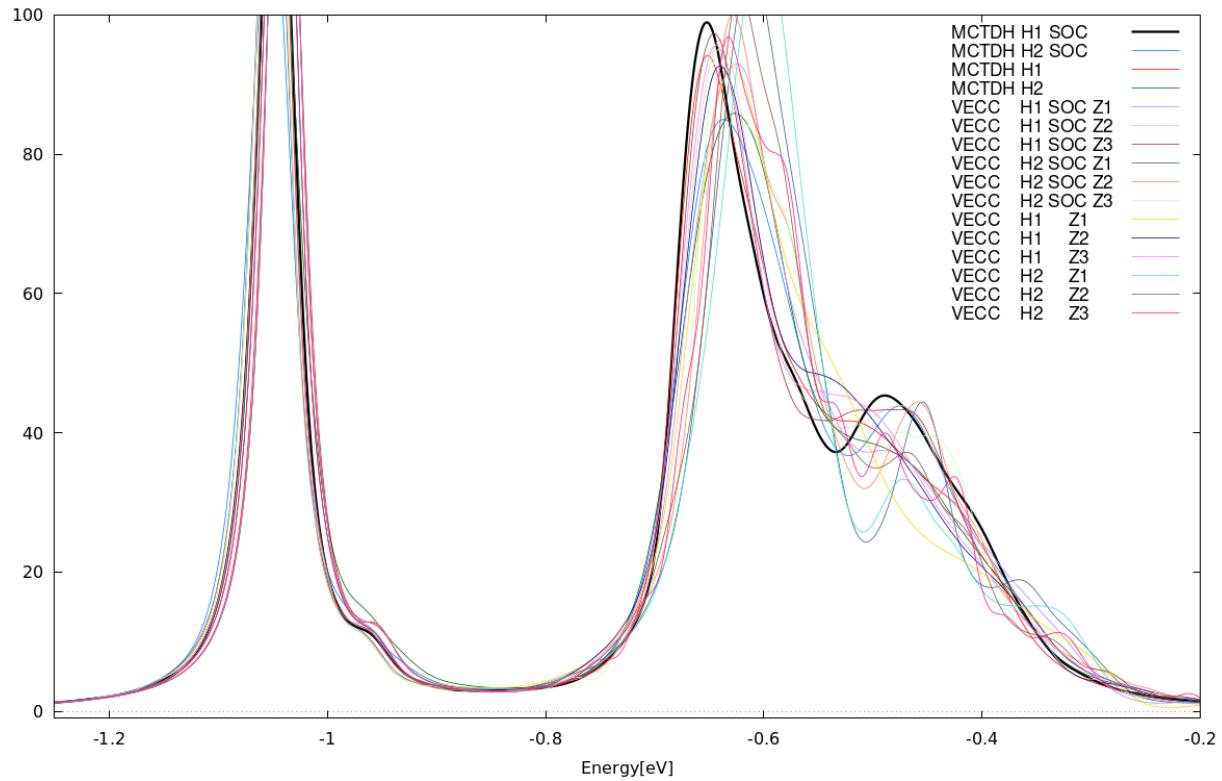
CoF<sub>3</sub> Vibronic Spectrum - MCTDH SOC Effect Comparison



CoF<sub>3</sub> Vibronic Spectrum - VECC SOC Effect Comparison



CoF<sub>3</sub> Vibronic Spectrum Composite



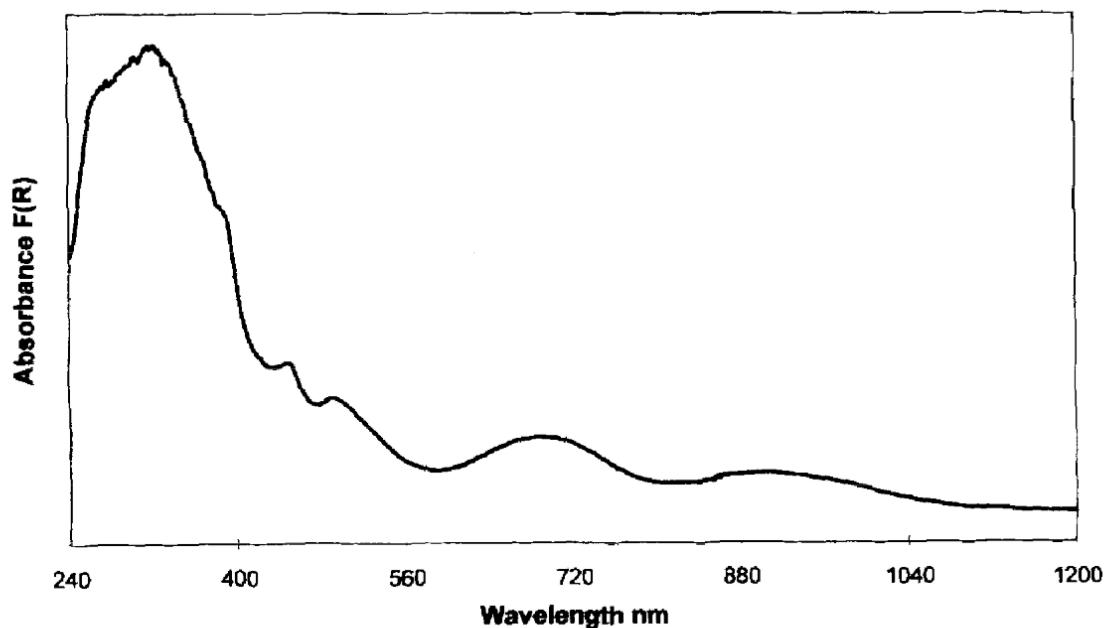


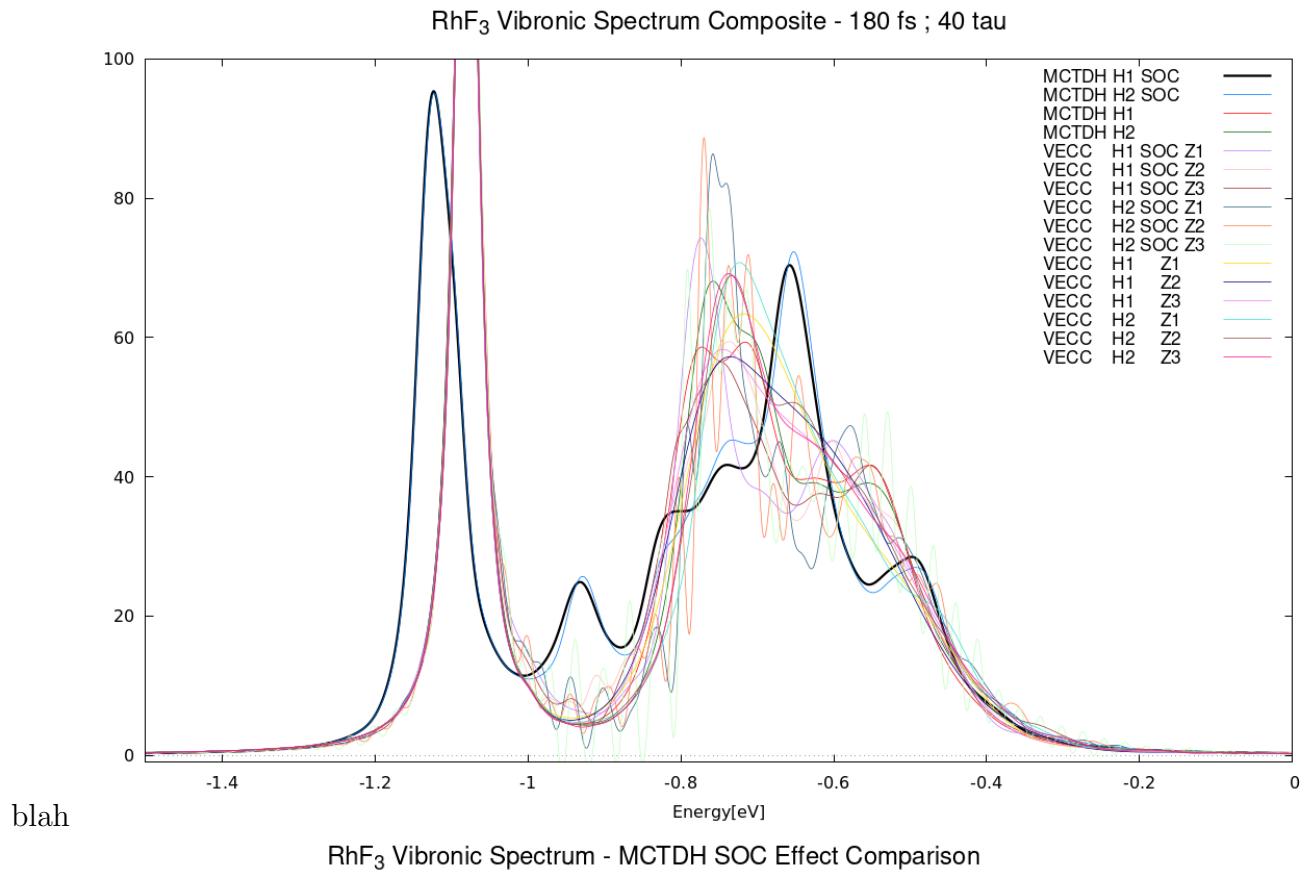
Fig. 3. Diffuse reflectance spectrum of  $\text{CoF}_3$ .

taken from "UV-visible spectroscopic studies of group 8-10 metal trifluorides" //

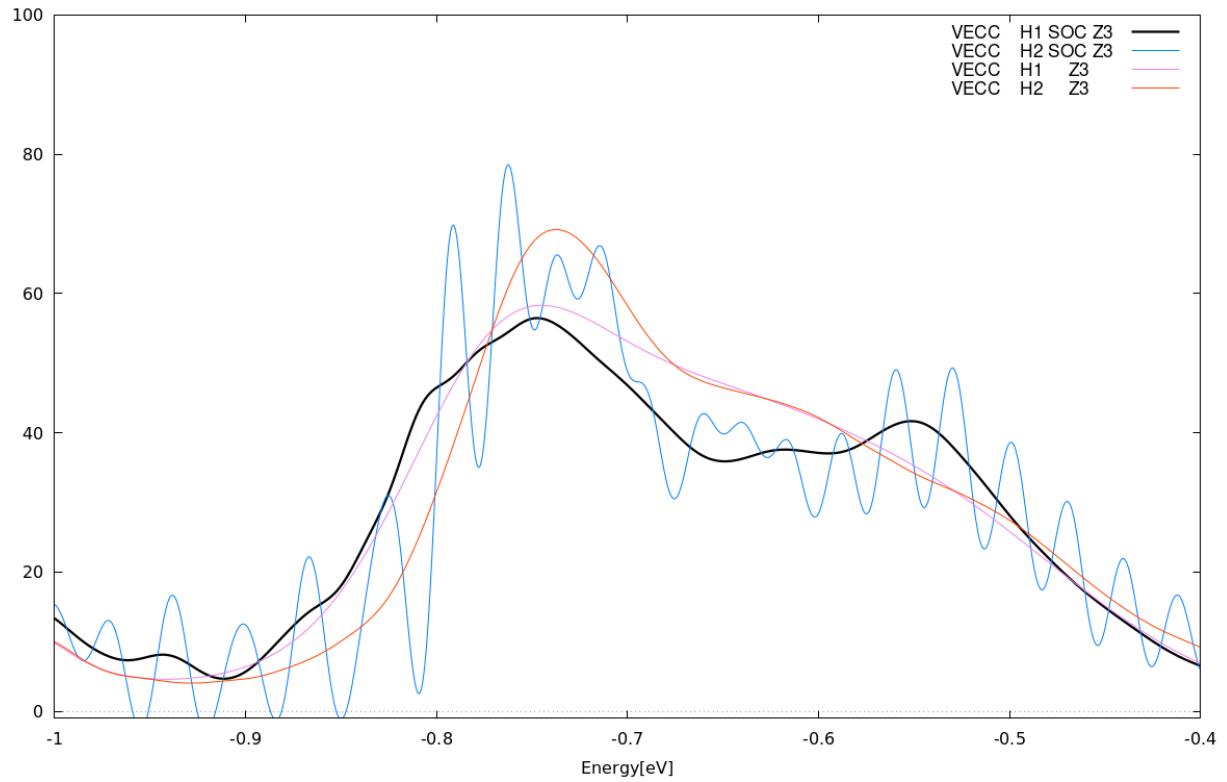
add short description of file show table with Hamiltonian parameters show image of molecule? (probably unnecessary) show spectra from MCTDH and VECC (maybe show constant, linear, quadratic, SOC?) discuss the results? are they identical, good this system was just top test xyz, and it behaved as we expected.



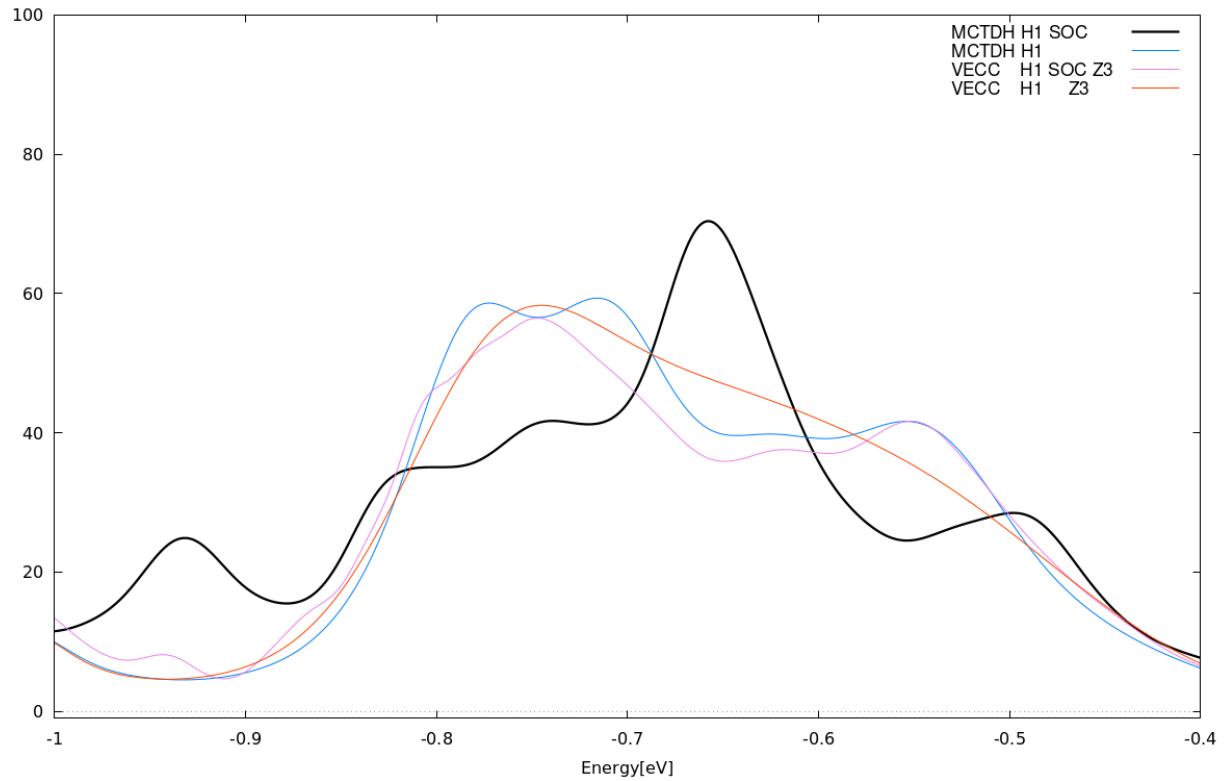
### 4.5.2 RhF<sub>3</sub>



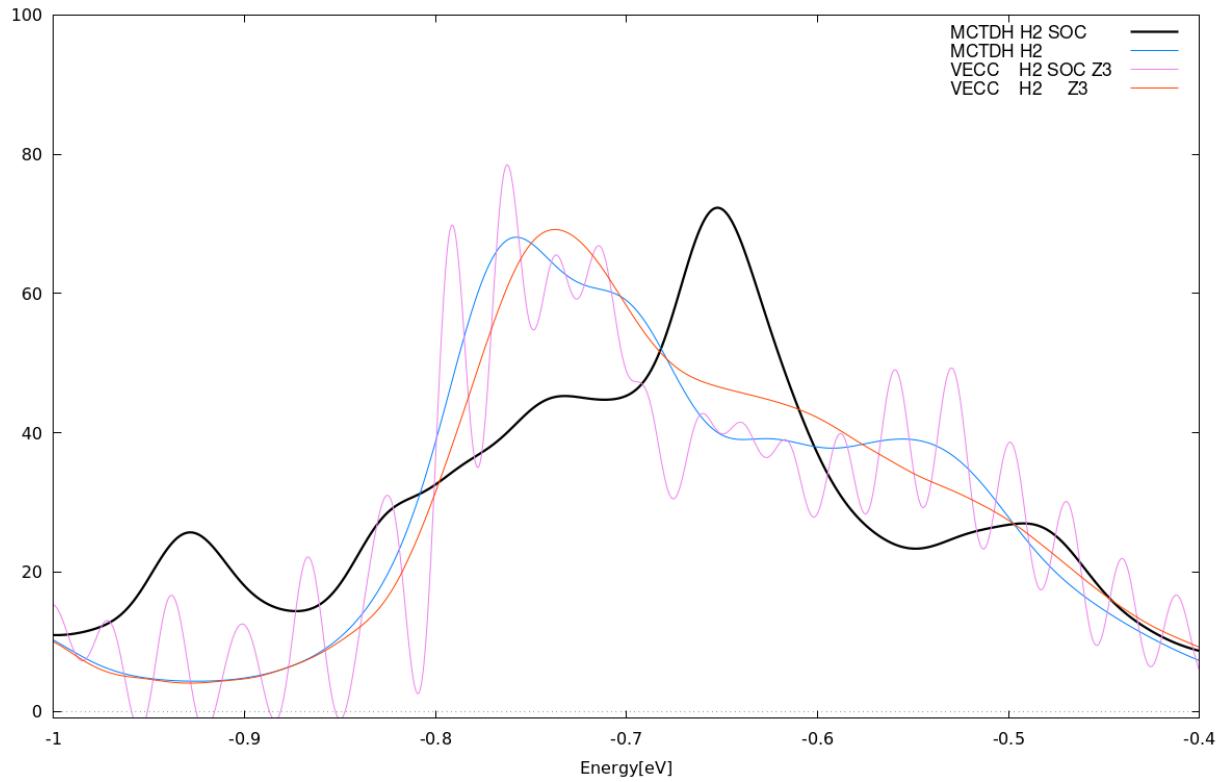
RhF<sub>3</sub> Vibronic Spectrum - VECC SOC Effect Comparison



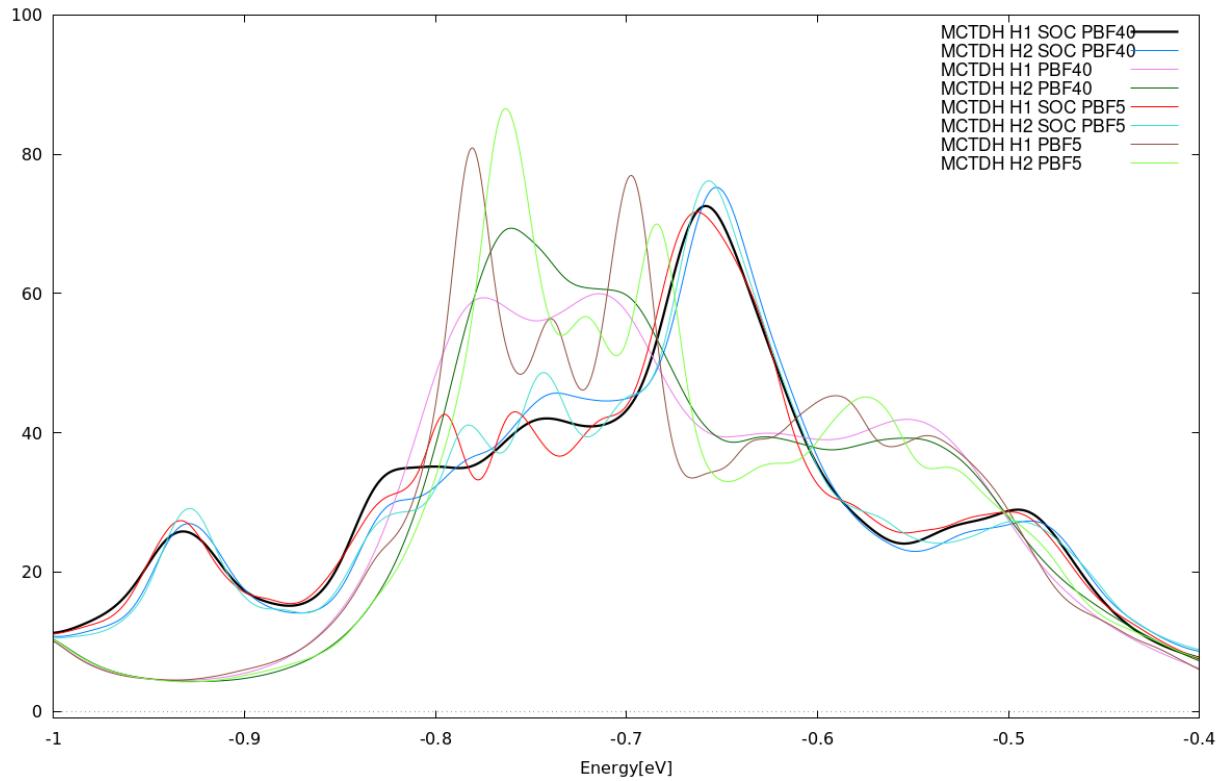
RhF<sub>3</sub> Vibronic Spectrum - Linear Model Comparison



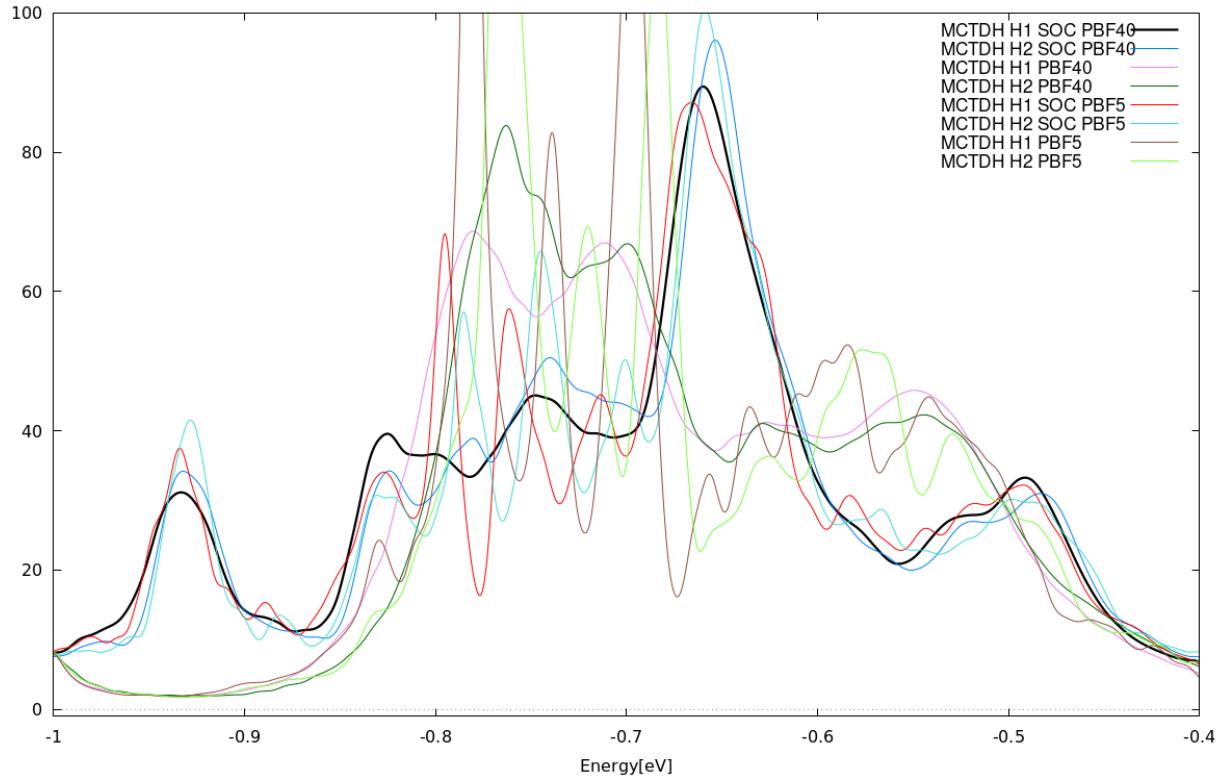
### RhF<sub>3</sub> Vibronic Spectrum - Quadratic Model Comparison



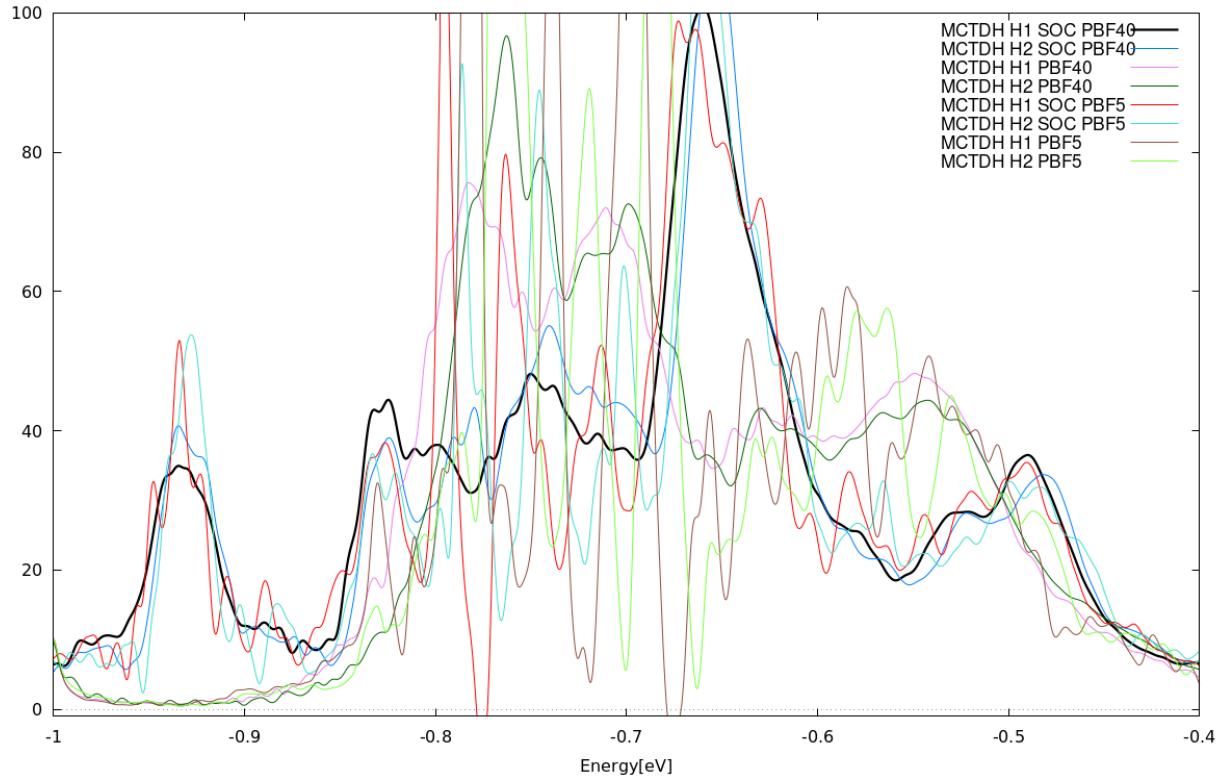
### RhF<sub>3</sub> Vibronic Spectrum Composite - 1000 fs ; 40 tau



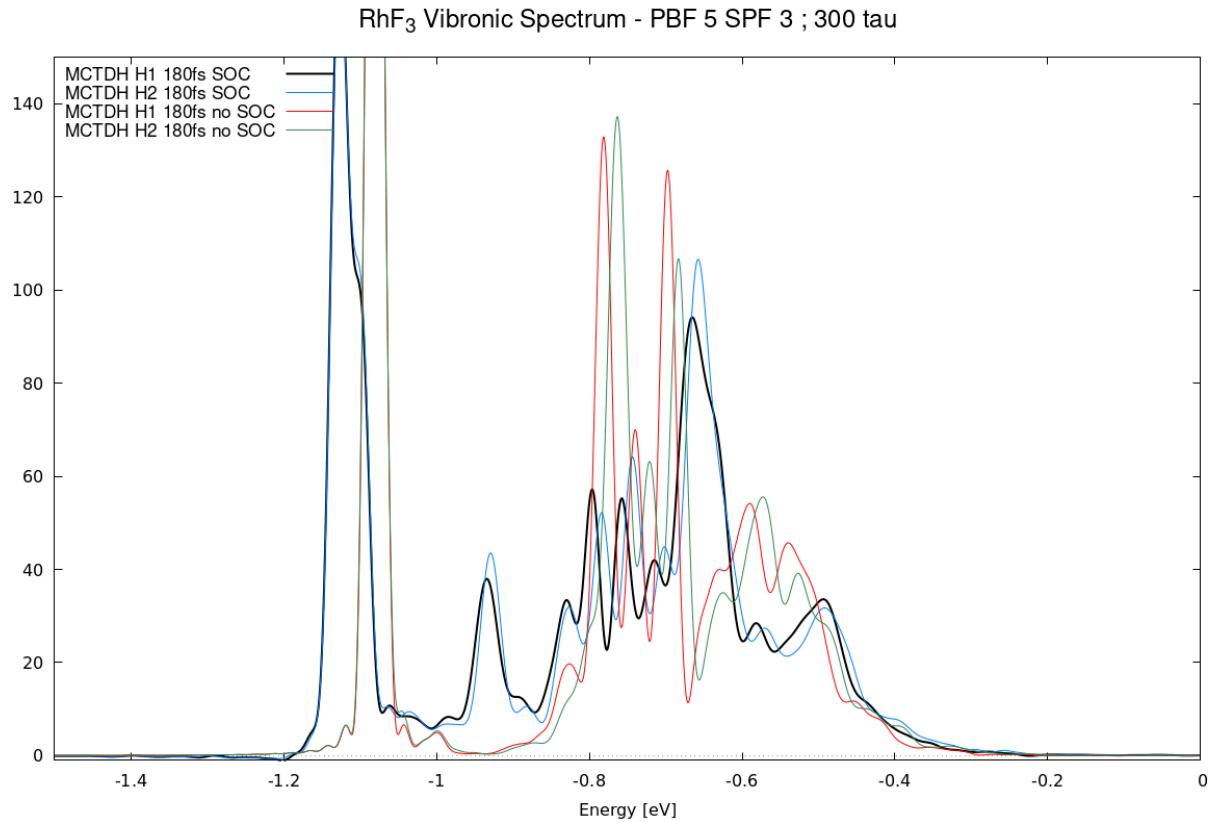
RhF<sub>3</sub> Vibronic Spectrum Composite - 1000 fs ; 90 tau



RhF<sub>3</sub> Vibronic Spectrum Composite - 1000 fs ; 200 tau



TAU MAKES A HUGE DIFFERENCE FOR RHF3! ESPECIALLY IF IT IS LONG PROPAGATION MCTDH.



HOWEVER THE PROBLEM IS THAT VECC CANNOT HANDLE LARGE TAU AND CANNOT HANDLE LARGE FS PROPAGATION ELSE DIVERGES SO 180FS 40 TAU IS IDEAL FOR COMPARISON!!

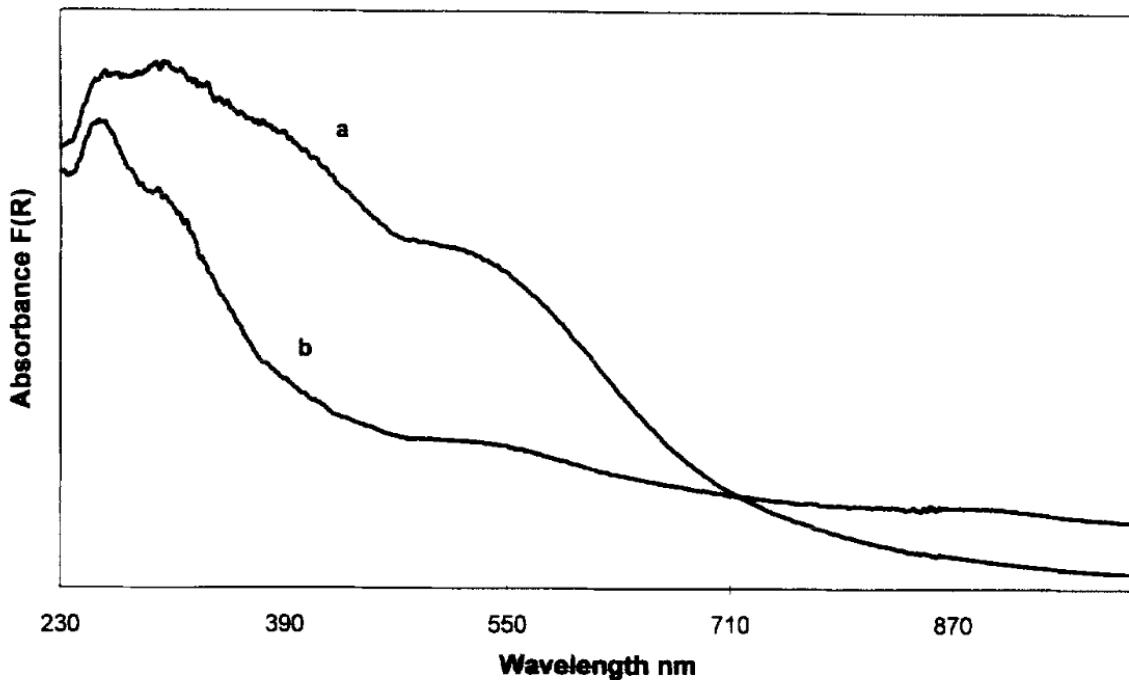


Fig. 4. Diffuse reflectance spectra of (a) RhF<sub>3</sub> and (b) IrF<sub>3</sub>.

taken from "UV-visible spectroscopic studies of group 8-10 metal trifluorides" //

add short description of file show table with Hamiltonian parameters show image of molecule? (probably unnecessary) show spectra from MCTDH and VECC (maybe show constant, linear, quadratic, SOC?) discuss the results? are they identical, good this system was just top test xyz, and it behaved as we expected.

#### 4.5.3 Fe(CO)<sub>5</sub>

Postulated, a priori

characterized by

add short description of file show table with Hamiltonian parameters show image of molecule? (probably unnecessary) show spectra from MCTDH and VECC (maybe show constant, linear, quadratic, SOC?) discuss the results? are they identical, good this system was just top test xyz, and it behaved as we expected.

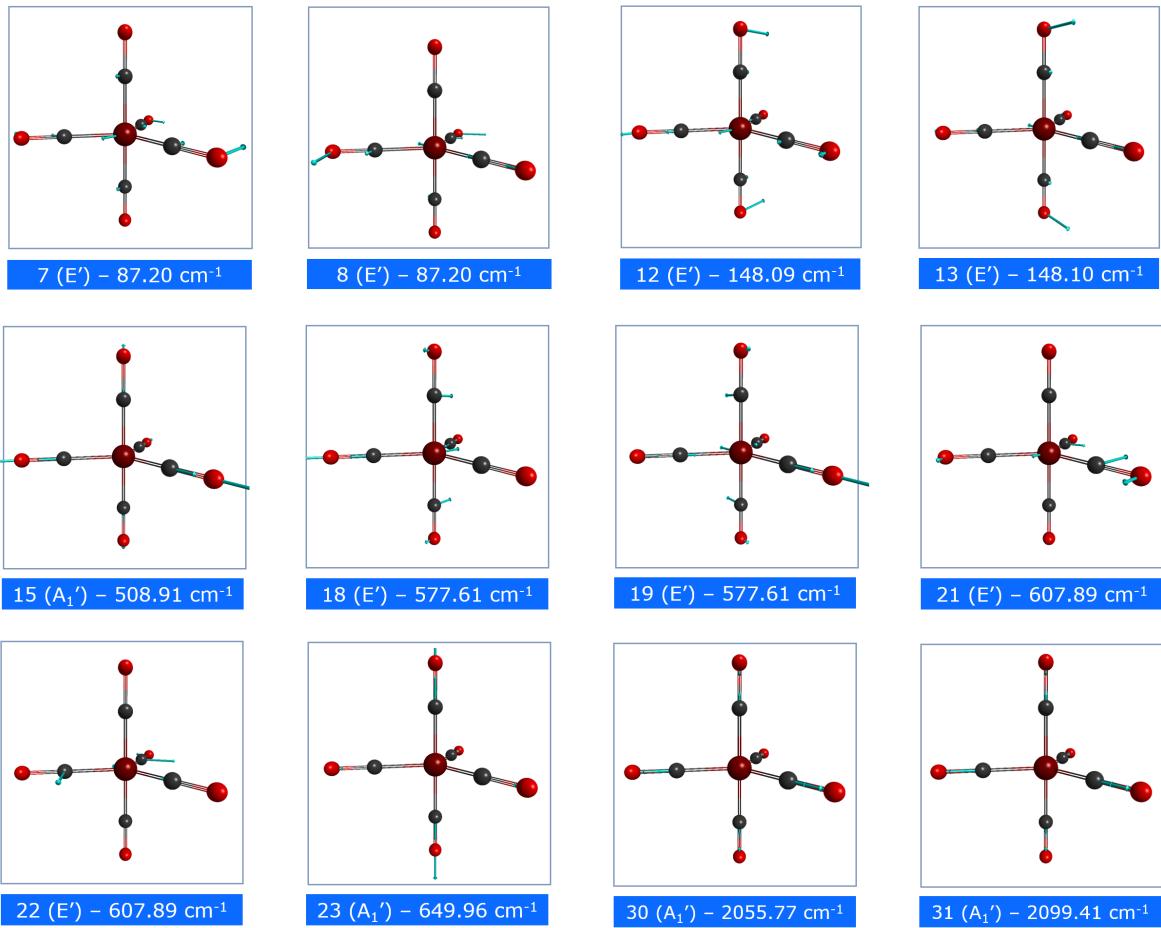


Fig . Screened 12 out of 27 modes.

#### 4.5.4 Vibronic spectra of test case pnictogen hydrides

#### 4.5.5 Spectra of $D_{3h}$ SOC transition metal complexes ( $\text{CoF}_3$ and $\text{RhF}_3$ )

### 4.6 Iron pentacarbonyl, $\text{Fe}(\text{CO})_5$

carbonyl ligand, due to carbonyls' known stabilizing effect on metals

Using  $\text{Fe}(\text{CO})_5$  as an example: we first recognized that the full valence active space was 11o18e. However, doing a oneshot for even the ground state equilibrium geometry calculation

did not result in anywhere near convergence. Instead, we had to slowly build up the transition metal compound's ground state geometry by filling in bunches of electrons and orbitals at a time. The initial calculation was a 4+ cation, with a multiplicity of 1. Next, a calculation of 1+ charge and multiplicity of 4, due to three unpaired electrons between a set of  $a''$  and  $e'$  of orbitals 46 through 48. This calculation resolved the orbital ranking:  $a''$  was lower than  $e'$ . Finally, a third calculation was ultimately performed on the fully electron-filled molecule and this neutral model yielded the full equilibrium geometry. CHEMICAL INTUITION NEEDED!

We first started by doing an optimization of  $\text{Fe}(\text{CO})_5^{4+}$ . Then, we introduced 2 electrons (multiplicity of 3), and then made it neutral. Observed that the irreducible representation orbital ranking of HOMO was  $e'$  and LUMO was  $a'_1$ . After obtaining the equilibrium geometry, we proceeded to fill the active space. 4o4e, 5o8e, 6o8e, 6o10e, 8o10e, 9o12e, and finally the full valent 11o18e. However, 11o18e did not seem to converge no matter what we tried. So we opted for 8o10e to tide us over and achieve a reasonable spectrum.

Curse of dimensionality when doing 33 modes ... gigantic ... ML-MCTDH is needed according to Prateek: <https://mattermodeling.stackexchange.com/questions/535/what-is-the-largest-system-for-which-vibrationally-resolved-electronic-spectra-h>

So we screened modes using effective vibronic coupling magnitude above a threshold. Linear couplings only. See JCTC singlet and triplet fission paper. Heart-wrenching that MCTDH could only compile up to 220k parameters when we needed 260k for our full model. Emailed Prateek and H.D.-Meyer. We can assert that only these 12 modes belong to the model. Incredible struggle, but nah, we'd win. Selected screen modes: [7, 8, 12, 13, 15, 18, 19, 21, 22, 23, 30, 31]. 12 out of 27. TZ: show a spreadsheet of screenshots here, displaying the ones we selected, does it make sense according to our analysis?

(see email with Dieter) We have good news. Prof. Meyer advised that the MCTDH-84 version I was running was outdated and that upgrading to the newest 86 version would do the trick. Fortunately Toby had a copy of 86 too on hand. Not only has the parameter limitation seemed to be resolved, but it similarly resolves a secondary issue in 84 where parallelized runs were erroring out (but works now in 86!).

Stitching together various functionalities, tuning

May 29th: we have done plotting of diabatic energies today and saw that mode 15 was equatorial dissociation and mode 23 was axial dissociation. Equatorial dissociation leads to one of the COs on the horizontal plane leaving, resulting in the axial COs to bend down and resulting in Td symmetry. Whereas, axial dissociation results in one CO on principal axis leaving and resulting in C3v symmetry.

An optimal and desirable system for exhibiting both JT and SOC effects is a large iron complex with low-lying degenerate electronic states and carbonyl ligand, due to carbonyls' known stabilizing effect on metals (Nooijen discussion reference)

<https://stackoverflow.com/questions/19412382/gnuplot-line-types#19420678>

We are using ORMAS+GMCPT! <https://ccl.scc.kyushu-u.ac.jp/~nakano/gmcpt.html>

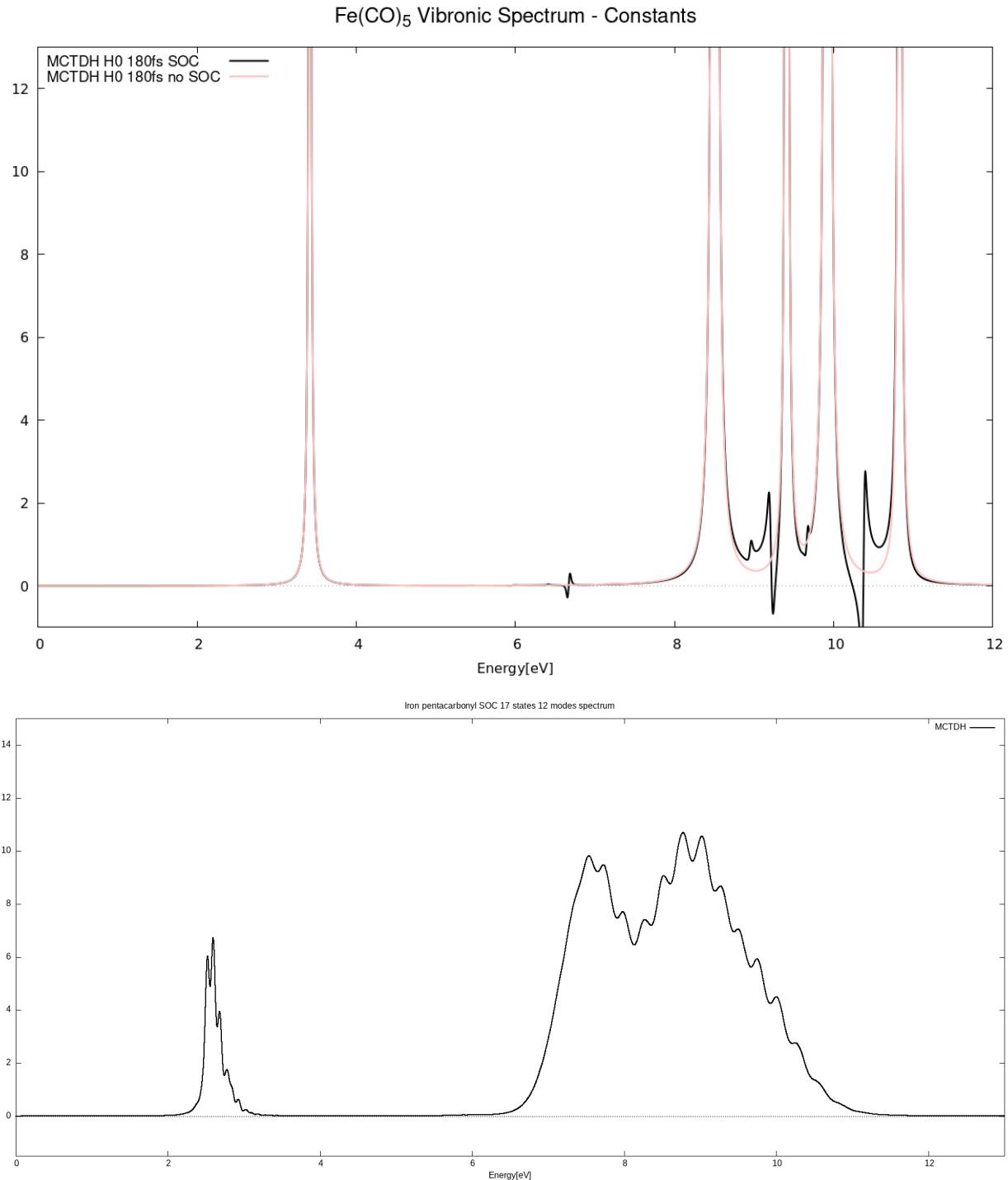
Normally we use Douglas-Kroll (DK) in the `$relfwn` group, however Sapporo basis sets like RELWFN=LUT-IOTC instead.

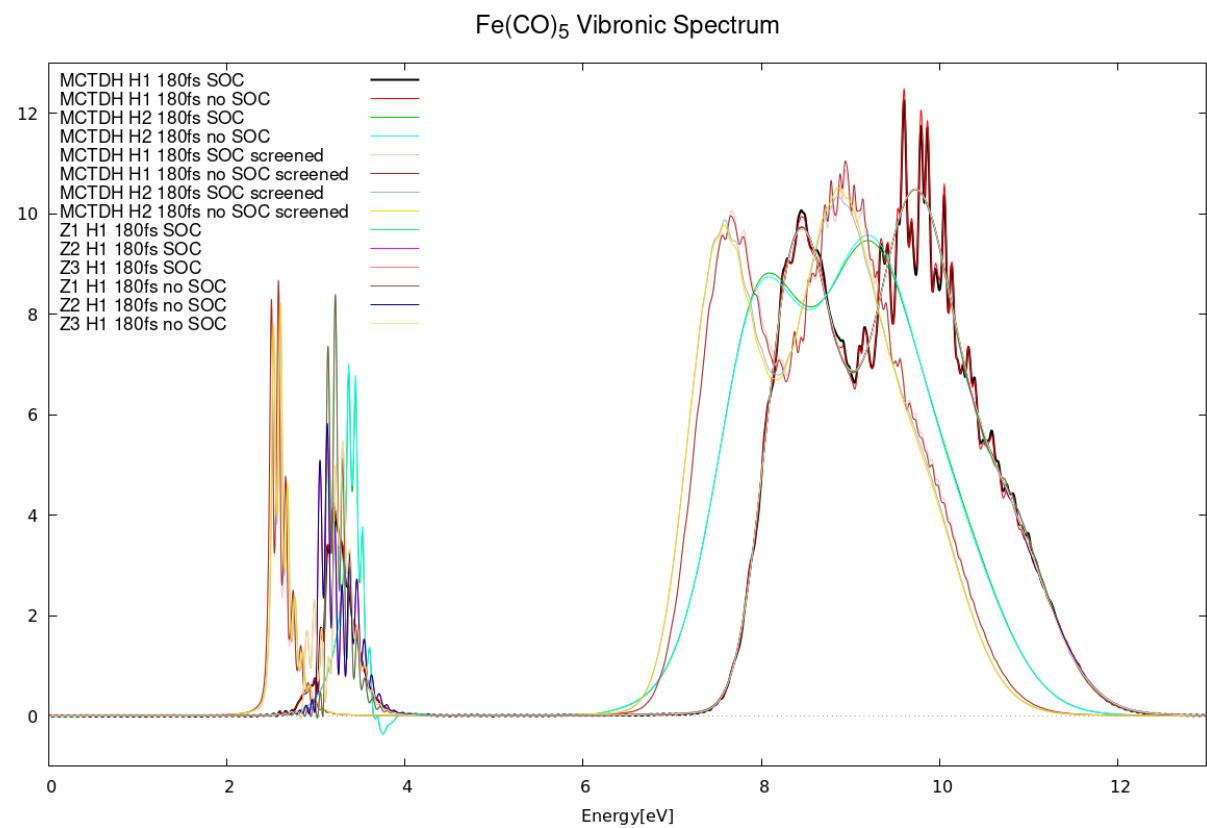
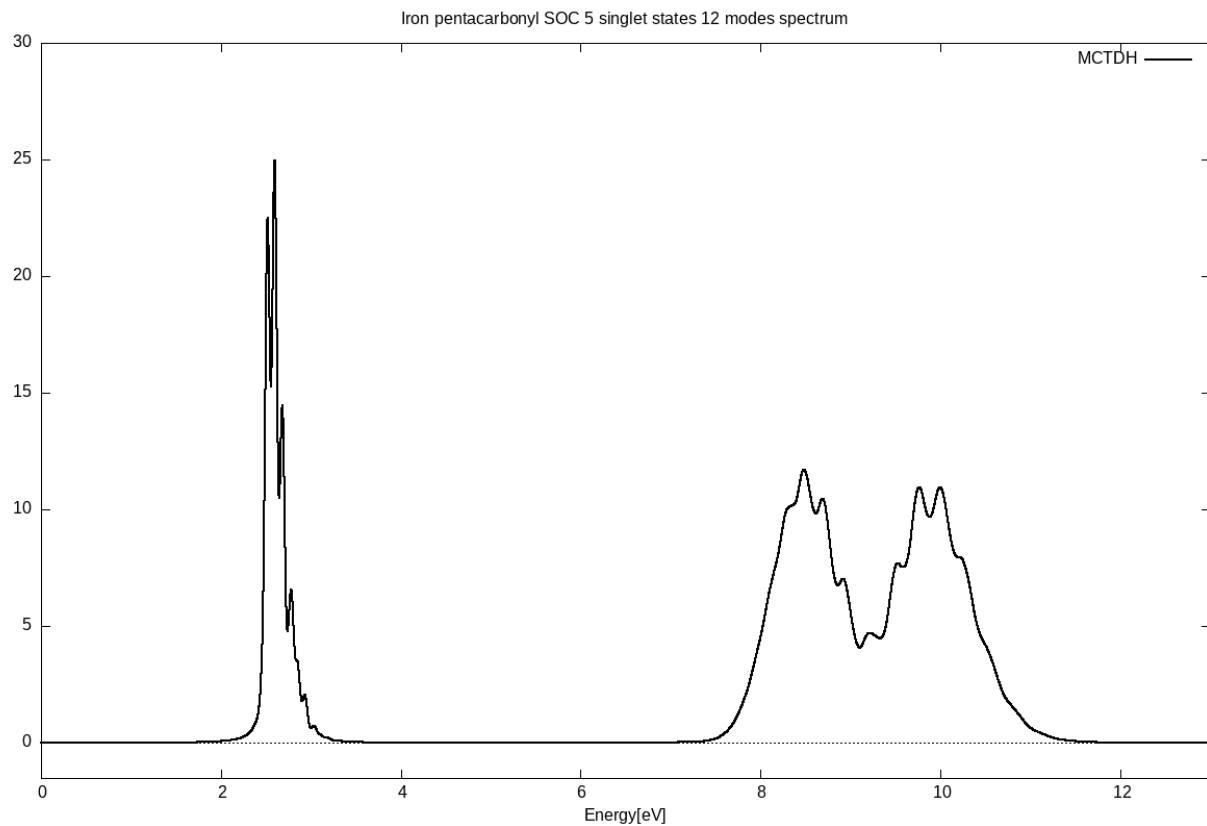
Scale breaks things. We have to fix things, redesign, change this, change that... Emphasize the technical hiccups, airplane dashboard ... It is fickle. On a macro level VECC gets it right, however does not get the fine structure peaks.

We cannot simplify everything down to a one-size fits all simple model. The sheer amount of parameters and caveats each spectra run takes translates into it being incomplete in some way if we do not alter the parameters specific for the case we are looking at. Overall you

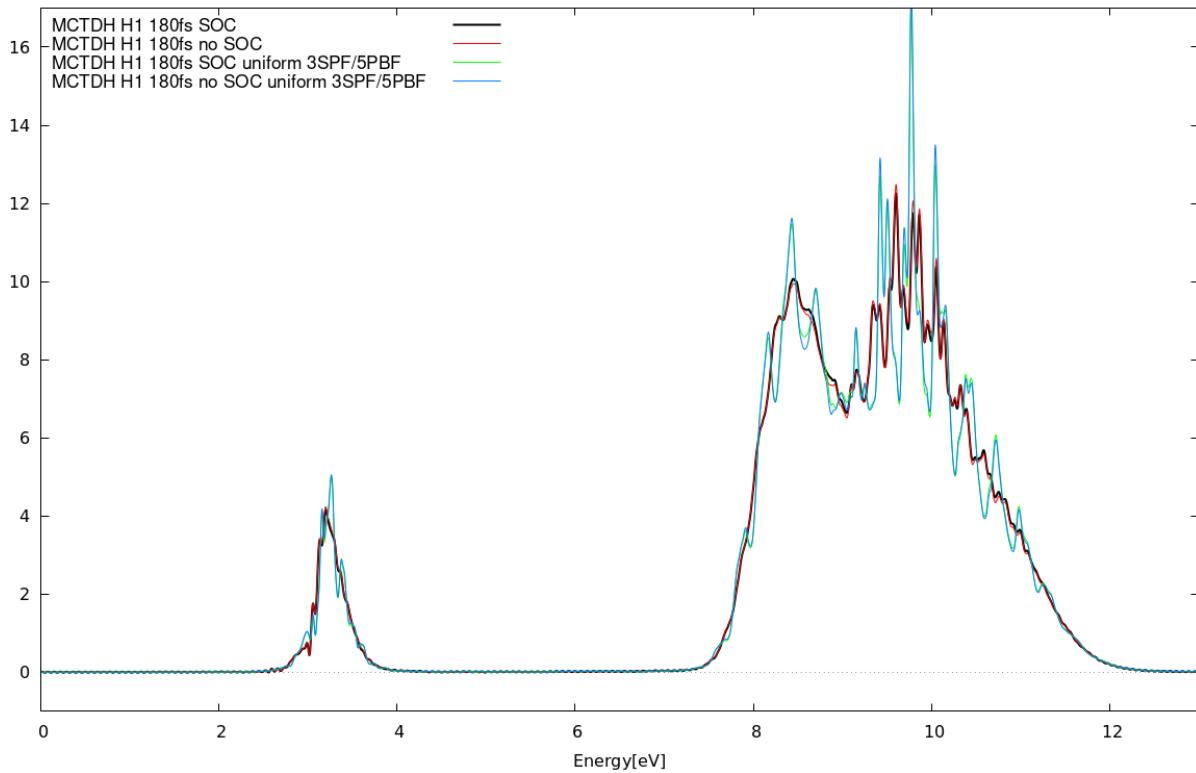
have to make decisions and cleverly choose an array of compatible parameters. Having reservations about a job is needed, oftentimes a double or even triple check is prudent.

FeCO SOC effect not observed probably because it is too subtle and weak.

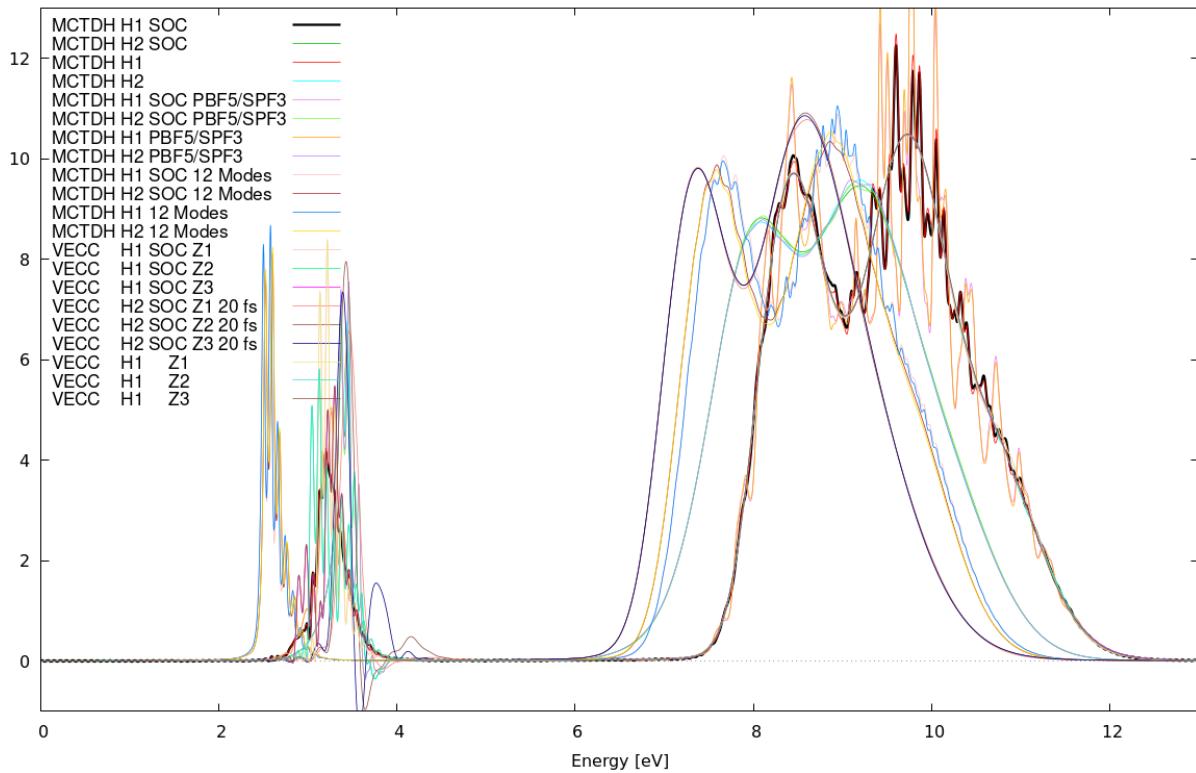




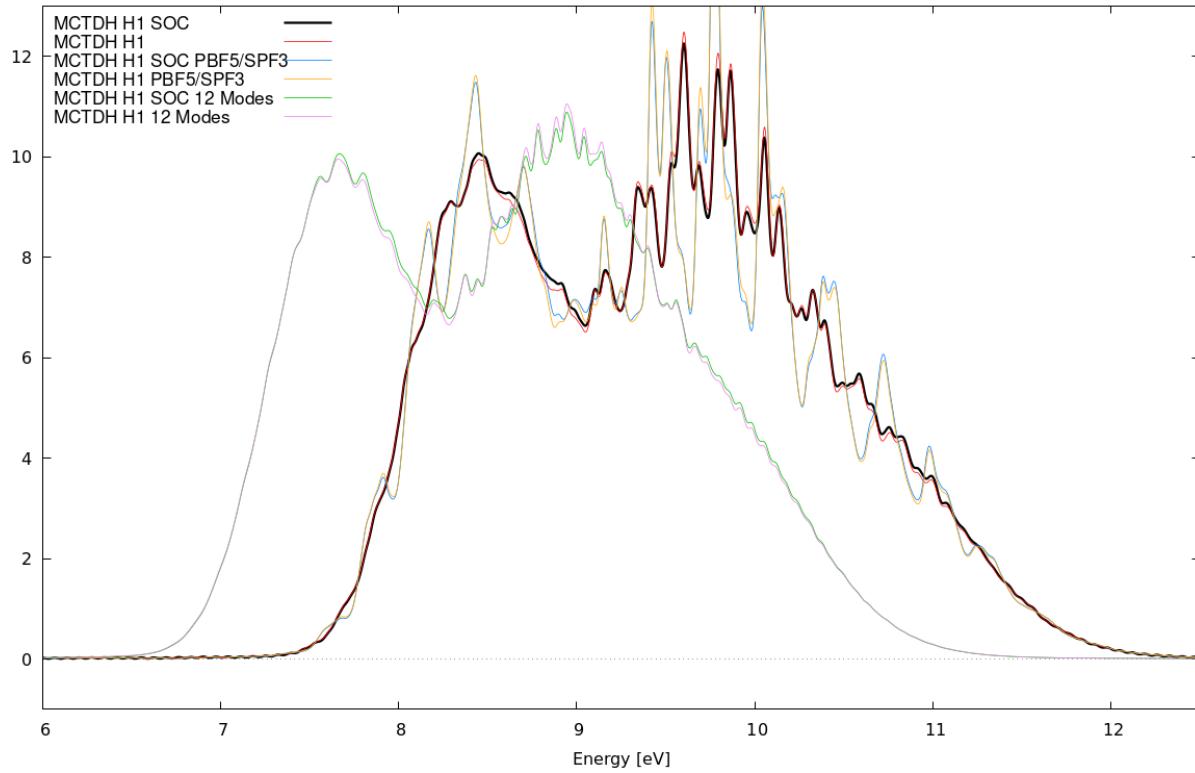
Fe(CO)<sub>5</sub> Vibronic Spectrum



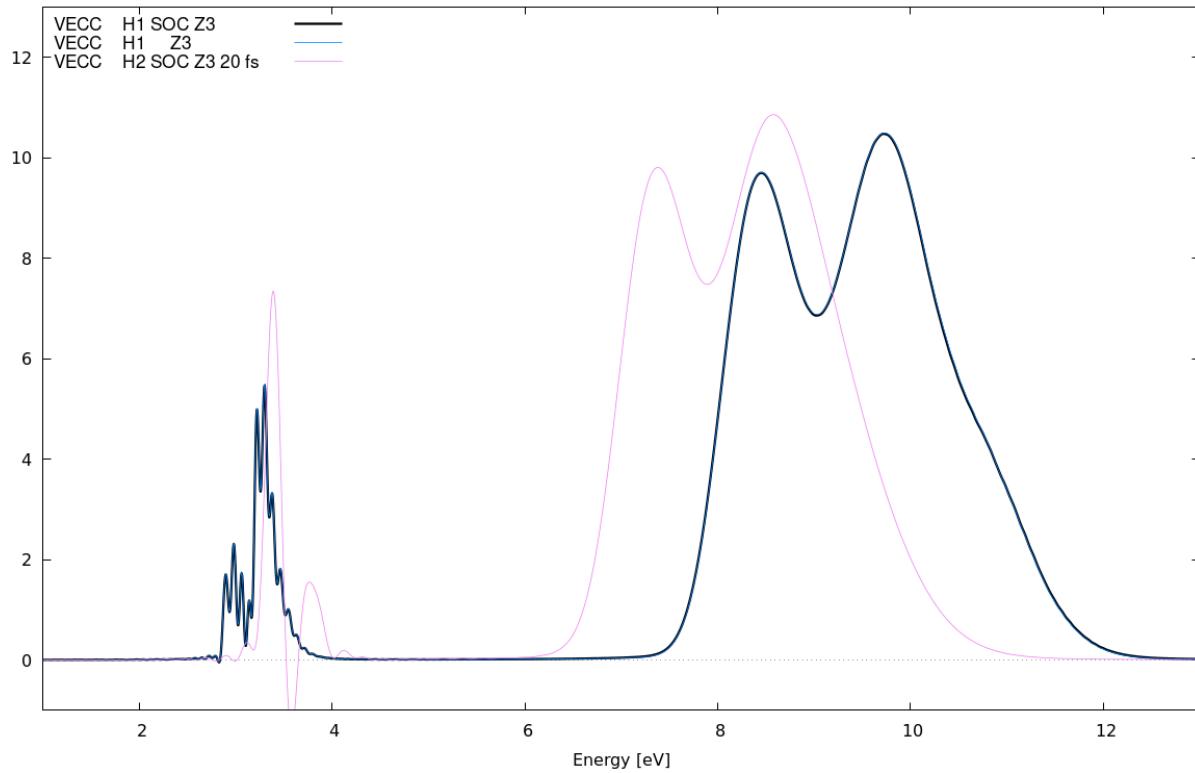
Fe(CO)<sub>5</sub> Vibronic Spectrum Composite - 180 fs ; 40 tau



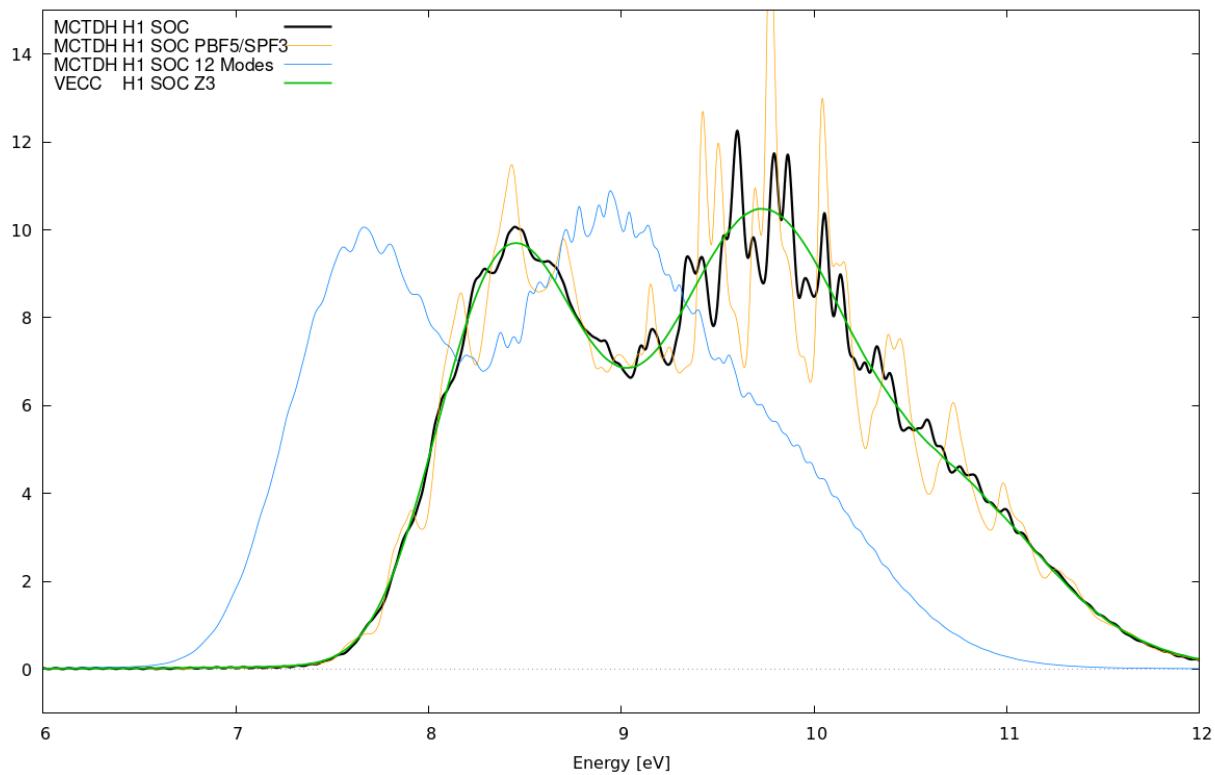
Fe(CO)<sub>5</sub> Vibronic Spectrum - MCTDH SOC Effect Comparison



Fe(CO)<sub>5</sub> Vibronic Spectrum - VECC SOC Effect Comparison

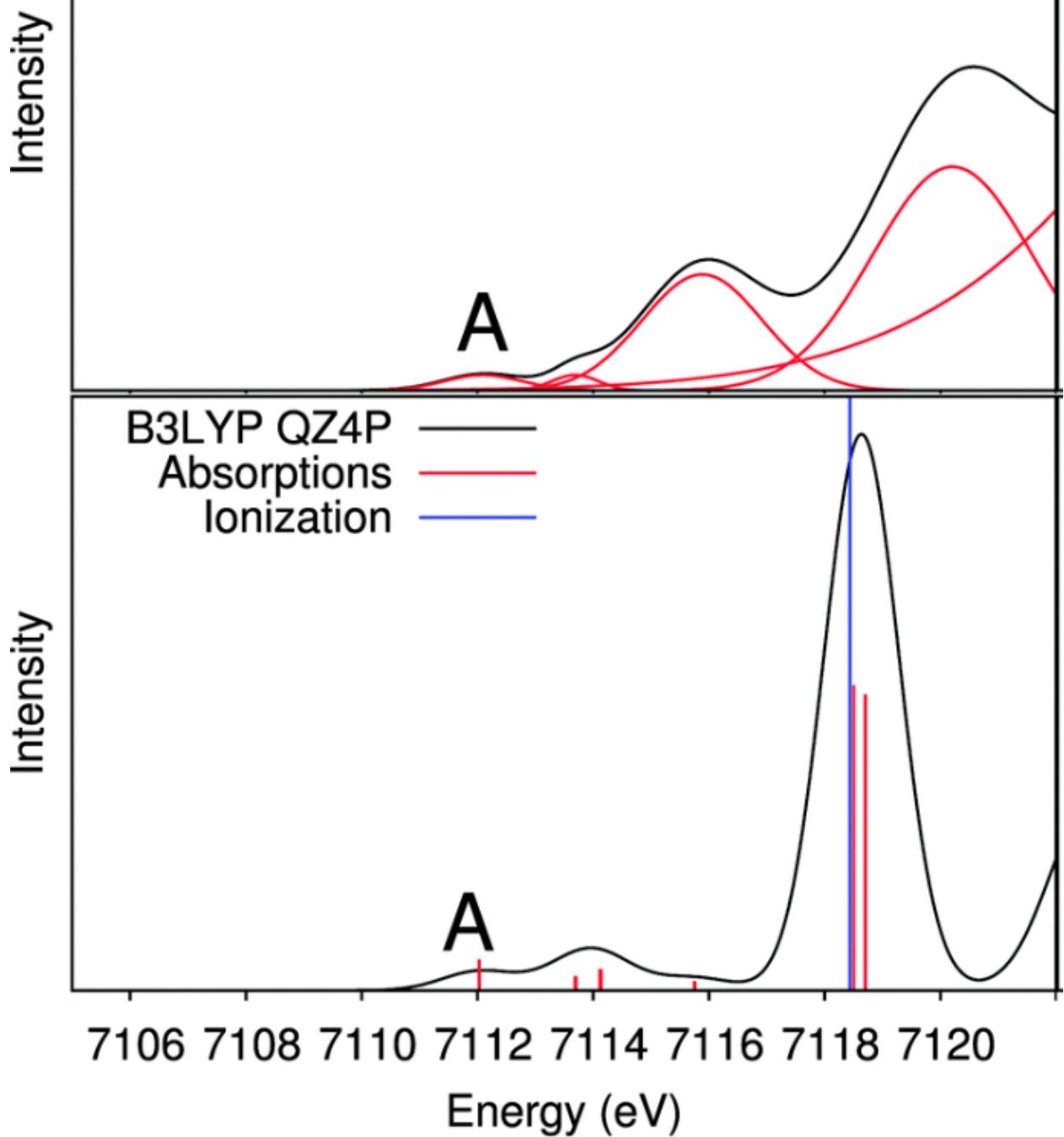


### Fe(CO)<sub>5</sub> Vibronic Spectrum - 6 to 12 eV



Lack of tremble and jitters?

Experiment  
 $\text{Fe}(\text{CO})_5$



taken from "High-resolution X-ray absorption spectroscopy of iron carbonyl complexes" //

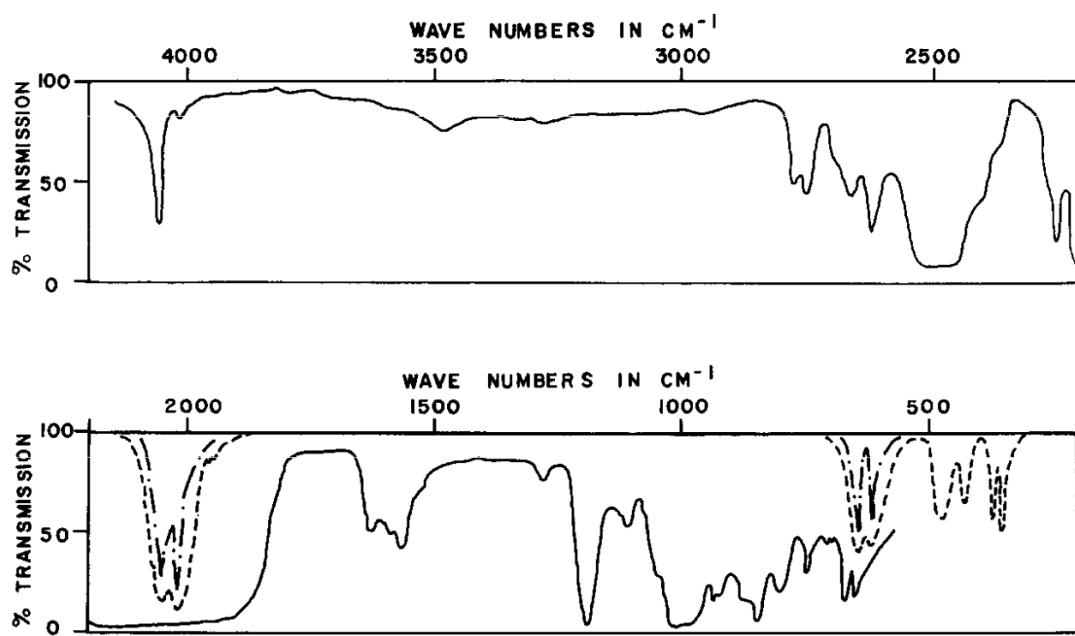
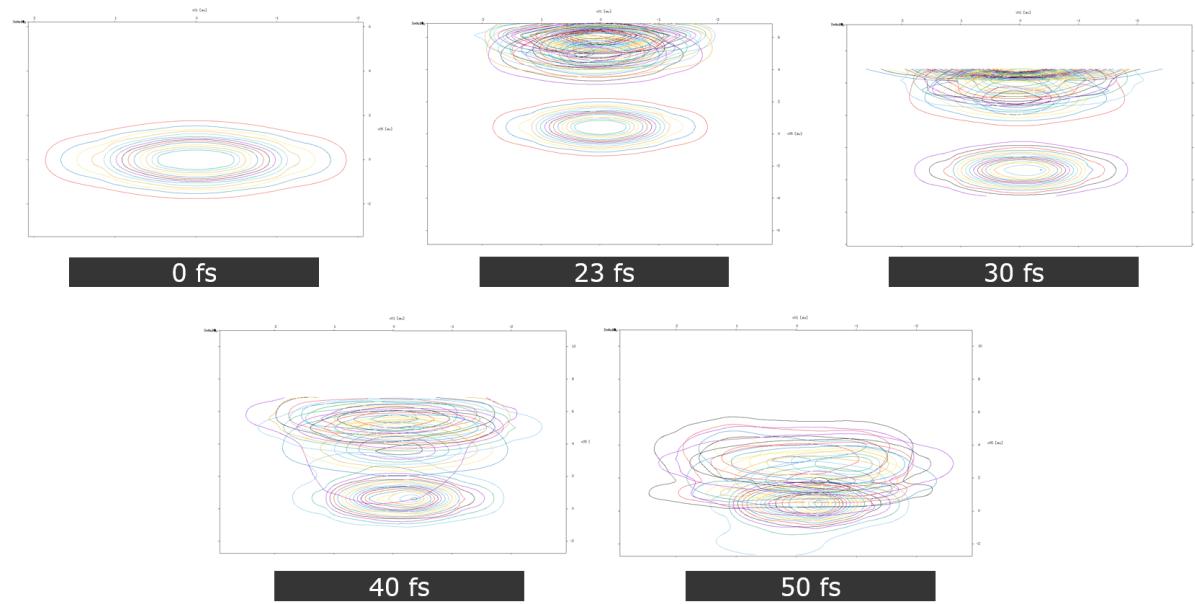


FIG. 1. Infrared absorption spectrum of iron pentacarbonyl. — Liquid at 25°C in 0.5-mm cell. - - Gas at 25°C in 10-cm cell (25-mm partial pressure of iron pentacarbonyl in 735 mm of nitrogen). - · - Gas at 25°C in 10-cm cell (5-mm partial pressure of iron pentacarbonyl in 755 mm of nitrogen).

taken from "Infrared Spectra and Normal Coordinate Analysis of Iron Pentacarbonyl" //



$D_{3h}$  Character Table

Table 4.1: Character Table for the  $D_{3h}$  Point Group

$D_{3h}$	$E$	$2C_3$	$3C'_2$	$\sigma_h$	$2S_3$	$3\sigma_v$
$A'_1$	1	1	1	1	1	1
$A'_2$	1	1	-1	1	1	-1
$A''_1$	1	1	1	-1	-1	-1
$A''_2$	1	1	-1	-1	-1	1
$E'$	2	-1	0	2	-1	0
$E''$	2	-1	0	-2	1	0

No matter how much we tried to increase the primitive basis grid size, the end state is always highly populated, meaning that the potentials are unbound.

AVOID DOING TDH AT ALL COSTS!!! IT EXPLODED MY SCRATCH

```
1 [bjb2chen@narval3 ~]$ diskusage_report
2
3          Description           Space      #
4          of files
5
6          /home (project bjb2chen)    636M/50G
7          12k/500k
8
9          /scratch (project bjb2chen)   19T/20T
10         18k/1000k
11
12         /project (project def-mnooijen) 138G/1000G
13         14k/500k
14
15         /nearline (project def-mnooijen) 98k/1000G
16         4/5000
17
18
19 [bjb2chen@narval3 Apr26_redo_SOC]$ du -ah SOC_CORRECTION/ | sort -hr |
20   head -n 20
21 19T      SOC_CORRECTION/TDH/mctdh/op_FeC027Q_17st_fullmodes
```

Haruyuki Nakano. Hisao Nakamura. See conversation with Neil on June 3rd.

SOC: <https://pubs.aip.org/aip/jcp/article/146/14/144103/195065> Metal trifluorides: <https://www.sciencedirect.com/scientific-topics/metal-trifluorides>

Diabatic states: G. J. Atchity and K. Ruedenberg, Theor. Chem. Acc. 97, 47 (1997).<https://doi.org/10.1007/s001020050001>

<https://pubs.aip.org/aip/jcp/article/115/22/10353/946057/The-direct-calculation-of-diabatic-states-based-on-UV-spectra>: <https://onlinelibrary.wiley.com/doi/10.1002/qua.10724> Linear vibronic models: <https://pubs.acs.org/doi/10.1021/acs.jctc.1c00022> (Santoro)

We are using ORMAS+GMCPT! Student guide that really crystallized my knowledge of GMC-QDPT: <https://ccl.scc.kyushu-u.ac.jp/~nakano/gmcpt.html>

Iron pentacarbonyl, haruyuki nakano and mark s gordon <https://pubs.rsc.org/en/content/articlepdf/199>

This is a very good diagram showing interdependencies: [https://en.wikipedia.org/wiki/Complete\\_active\\_immutability](https://en.wikipedia.org/wiki/Complete_active_immutability)

This is a good consideration for future work, running the calculations via GPU acceleration: <https://www.nvidia.com/es-la/data-center/gpu-accelerated-applications/gamess/>

MCTDH troubleshooting

[https://en.wikipedia.org/wiki/Dung\\_beetle](https://en.wikipedia.org/wiki/Dung_beetle)

```
1 # This is a comment  
2 def hello_world():  
3     print("Hello, world!")
```

Listing 4.1: Example Python Code

## 4.7 Various things i need to remember / 12 billion caveats

1. VECC doesn't support off-diagonal E matrix terms; when reading from a .op file it only fills in the diagonal elements of the electronic matrix (E)

```
# save the reference Hamiltonian into the energies array
energies = np.zeros((A, A))
for a in range(A):
    list_of_words = lines[a].split()
    assert list_of_words[0] == f"EH_s{a+1:02}_s{a+1:02}"
    assert list_of_words[-1] == "ev"
    energies[a, a] = list_of_words[2]
```

2. VECC doesn't support more than 1D electron TDM
3. VECC diverges on quadratic after 20fs. You can salvage the run with high autospec damping, but it is fickle.
4. 4 thing
5. 5 thing

# Chapter 5

## Conclusion

Diabatization is at the heart of current vibronics research. Although it provides a numerical result, it can be applied to generalized situations. Perhaps one day there will be analytical models that can solve vibronics entirely. We cannot make sweeping conclusions about ...

<https://www.sci.uwaterloo.ca/~nooijen/website/research.html>

<https://yorkspace.library.yorku.ca/server/api/core/bitstreams/302f290d-535e-462e-8ca5-1e>  
content

<https://pubs.acs.org/doi/10.1021/jp0456990>

<https://www.sci.uwaterloo.ca/~nooijen/website/vibron/VC/Manual.pdf>

[https://www.sci.uwaterloo.ca/~nooijen/website/summary\\_papers.pdf](https://www.sci.uwaterloo.ca/~nooijen/website/summary_papers.pdf)

<https://onlinelibrary.wiley.com/doi/full/10.1002/qua.10724>

<https://ccl.scc.kyushu-u.ac.jp/~nakano/gmcpt.html>

## 5.1 MCSCF Pipeline

Inside the repository are other accompanying videos with commentary running through the full fledged protocol from input geometry to spectra. RhF<sub>3</sub> molecule will be used as an example here.

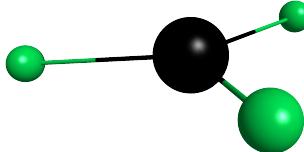


Figure 5.1: RhF<sub>3</sub>

### 5.1.1 Step 1: Geometry optimization + frequency (Hessian)

Readers can follow along on: <https://youtu.be/KJBxNBxqhHA>.

Following standard practice for GAMESS users, input geometry coordinates are first set using the wxMacMolPlt GUI builder tool. Typically the \$data group appears as follows, and GAMESS will duplicate the single fluorine atom given by the appropriate  $D_{3h}$  geometry. In the appendices, I have provided a standardized template that can be used as a starting point for all GAMESS calculations in this protocol.

```
1 $data
2 DNH 3
3
4 Rh     45.0      0.0   0.0   0.0
5 F      9.0       2.07  0.0   0.0
6 $end
```

Listing 5.1: Example  $D_{3h}$  RhF<sub>3</sub> Input Geometry

Important additional flags that are needed to be modified from the default template:

```
1 $contrl runtyp=optimize                      # 'geometry optimization'
2 $contrl scftyp=rohf    mult=5 icharg=0      # ROHF, multiplicity of 5, neutral
```

```

3 $basis gbasis=SPKrTZP          # Sapporo triple zeta basis set
4 $contrl relwfn=LUT-IOTC       # relativistic effects for heavy
      atom
5 $scf diis=.t                 # turn on DIIS SCF converger

```

Listing 5.2: Additional flags for Step 1.

Table 5.1: Step 1 information.

Input	Output
rungms RhF3_SPK_mp2_rohf_D3h_mult5_diis_gh.inp ncpu ngb nhour	RhF3_*_gh.out

After the input file is successfully setup, one can submit the job to the cluster via the following input command. The resulting output file should locate the important 'EQUILIBRIUM GEOMETRY', and list out the  $3N-6 = 6$  vibrational normal modes at the end of calculation.

### 5.1.2 Step 2: GMCPT

This is the crucial MCSCF (ORMAS) active space setup step.

```

1 $contrl runtyp=energy nosym=1 scftyp=mcscf relwfn=LUT-IOTC
2 $guess norb=168
3 $mcscf cistep=gmcci fullnr=.t. finci=mos
      accuracy=1.d-6 maxit=200 $end
4 $mrmp mrpt=gmcpt rdvecs=.f. $end
5 $gmcpt nmofzc=21 nmodoc=12 nmoact=5 nelact=6 stsym(1)=A
      reftyp=ormas nspace=1 mstart(1)=34 mine(1)=6 maxe(1)=6
6 kstate(1)=1,1,1,1,1,1,1,1,1,1,1,1,1,1,1,1,1,1,1,1,1,1,1,1,1,1,1
7 wstate(1)=1,1,1,1,1,1,1,1,1,1,1,1,1,1,1,1,1,1,1,1,1,1,1,1,1,1,1,1
8 iwgt=0 thrwgt=-1.0 krot=.f. ksdoe=.f. knospn=.t.
9      thrde=-1.0 edshft=0.02 $end
10
11

```

```

12 $transt $end
13 $data 'EQUILIBRIUM GEOMETRY' coordinates
14 $VEC group of ROHF orbitals

```

Listing 5.3: Additional flags for Step 2.

Table 5.2: Step 2 information.

Command Line	Input	Output
python3 toolkit.py python3 toolkit.py	RhF3_gh.out RhF3_gh.dat	ref_structure vec.dat (ROHF orbs)
rungms .inp ncpu ngb nhour	RhF3_gmcpt.inp	RhF3_gmcpt.out

### Note: Talk about GMCPT characteristics

GMCPT is tedious to get right. The recommendation is to consult the students' guide for GMCPT/GMC-QDPT calculations prepared by Nakano and frequently run many experimental jobs. One must ensure that they get the combo of frozen core, doubly occupied, and active orbitals/electrons right. Likewise, if the orbital ranking you input is unphysical, it can potentially result in a poor calculation. If you desire to introduce SOC to your system, attention to add the *\$transt \$end* flag. As it turns out, if you add SOC, and you are running on *nlogn*, the memory mapping fails! This is due to a bug on compiling GMCPT module on *nlogn*, causing memory corrupted bytes to appear near the termination of the calculation. However, do not be alarmed, as all the meaningful SOC information will be printed out before the memory corruption happens in the output file. If you run the exact same calculations on *ComputeCanada* servers, there will be no problem.

### 5.1.3 Step 3: DMO Calculation

#### Discuss semi canonical orbitals to DMO

Here, you will simply use the toolkit script to extract out the Semi-canonical orbitals and \$DMO group from the second step's *.dat* file. Be sure you know the active space's *nmoact* and *mstart* values, as the toolkit script will require these inputs to extract out the correct portion of diabatic molecular orbitals from the full Semi-canonical group. Step 3's input will be identical to Step 2, however the prior orbitals shall be deleted and in its place concatenate the Semi-canonical orbitals along with \$DMO group to the end of the file.

```
1 ~ Step 2 Input ~  
2 $VEC group of Semi-canonical orbitals (replacing old ROHF orbitals)  
3 $DMO group
```

Listing 5.4: Additional flags for Step 3.

Table 5.3: Step 3 information.

Command Line	Input	Output
python3 toolkit.py	RhF3_gmcpt.dat	active_space_orbs (Semi-canonical MOs)
python3 toolkit.py	RhF3_gmcpt.dat	dmo.dat (DMO group)
rungms .inp ncpu ngb nhour	RhF3_dmo.inp	RhF3_dmo.out

### 5.1.4 Step 4:

Rerun the previous Step 3's process on the new \$DMO outputs to obtain the Semi-canonical MOs and \$DMO group again. After, you will run toolkit on the *.out* file from Step 3, and from it can extract the \$REFDET group automatically. Note that one can actually handcraft a \$REFDET group for better accuracy, using knowledge of dominant Slater determinants and

state configurations by reviewing Step 3's output. Once more, concatenate it to the exact same Step 3 input, and this will be the Step 4 input file.

```

1 ~ Step 3 Input ~
2 $VEC group of Semi-canonical orbitals (fresh from Step 3 outputs)
3 $DMO group (fresh from Step 3 outputs)
4 $REFDET group

```

Listing 5.5: Additional flags for Step 4.

Table 5.4: Step 4 information.

Command Line		Input	Output
python3 toolkit.py		RhF3_dmo.dat	active_space_orbs (Fresh Semi-canonical MOs)
python3 toolkit.py		RhF3_dmo.dat	dmo.dat (DMO group)
rungms .inp ncpu ngb nhour		RhF3_refdet.inp	RhF3_refdet.out

### Cleanup: temp.inp

After Step 4 terminates successfully, there should be full vibronic coupling results. If this is the case, then the diabatization template `temp.inp` shall be prepared from Step 4's input. The only difference is that the `$data` group must be deleted, and a new `$data` section appended to the end of the template.

```

1 ~ Step 4 Input ~
2 ~ early $data group deleted ~
3 $data
4 title of diabatization
5 C1

```

Listing 5.6: Additional flags for Step 5.