

# Cosmology Homework #3

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## ① Dynamics of a Universe with a Cosmological Const.

$$\text{Given } \dot{a}^2 = H_0^2 a^2 \left( \Omega_{m,0} \frac{a_0^3}{a^3} + \Omega_{k,0} \frac{a_0^2}{a^2} + \Omega_{\Lambda,0} \right) \quad \text{--- } ①$$

Let's first change this into something easier to work with

$$\text{multiply } a^2 \text{ in } \rightarrow \dot{a}^2 = H_0^2 \left( \Omega_{m,0} \frac{a_0^3}{a} + \Omega_{k,0} a_0^2 + a^2 \Omega_{\Lambda,0} \right)$$

$$\text{differentiate with respect to time. } \rightarrow 2\ddot{a} = H_0^2 \left[ -\Omega_{m,0} \frac{a_0^3}{a^2} + 2\Omega_{\Lambda,0} a \right]$$

a) When at a minimum,  $\ddot{a} = 0$

$$\text{Thus, } 0 = H_0^2 \left[ -\Omega_{m,0} \frac{a_0^3}{a^2} + 2\Omega_{\Lambda,0} a \right]$$

But, given in the problem,  $H_0 \neq 0$  and must be a positive quantity. So, letting  $a = a_{\min}$

$$-\Omega_{m,0} \frac{a_0^3}{a_{\min}^2} + 2\Omega_{\Lambda,0} a_{\min} = 0 \rightarrow 2\Omega_{\Lambda,0} a_{\min} = \Omega_{m,0} \frac{a_0^3}{a_{\min}^2}$$

$$a_{\min}^3 = \frac{\Omega_{m,0}}{2\Omega_{\Lambda,0}} (a_0^3) \rightarrow a_{\min} = \left( \frac{\Omega_{m,0} a_0^3}{2\Omega_{\Lambda,0}} \right)^{\frac{1}{3}} \rightarrow \boxed{a_{\min} = \left( \frac{\Omega_{m,0}}{2\Omega_{\Lambda,0}} \right)^{\frac{1}{3}} a_0}$$

## ① Dynamics of a Universe with a Cosmological Const. Cont.

b) If we let  $a = a_{\min}$  in ①, we can say

$$\dot{a}_{\min}^2 = H_0^2 a_{\min}^2 \left[ \cancel{\Omega_{M,0}} a_0^3 \left( \frac{2\Omega_{A,0}}{\cancel{\Omega_{M,0}}} a_0^3 \right) + \cancel{\Omega_{k,0}} a_0^2 \left( \left( \frac{2\Omega_{A,0}}{\cancel{\Omega_{M,0}}} \right)^{\frac{2}{3}} a_0^2 \right) + \Omega_{A,0} \right]$$

$$\dot{a}_{\min}^2 = H_0^2 a_{\min}^2 \left[ 2\Omega_{A,0} + \Omega_{k,0} \left( \frac{2\Omega_{A,0}}{\Omega_{M,0}} \right)^{\frac{2}{3}} + \Omega_{A,0} \right]$$

$$\boxed{\dot{a}_{\min}^2 = H_0^2 a_{\min}^2 \left[ 3\Omega_{A,0} + \left( \frac{2\Omega_{A,0}}{\Omega_{M,0}} \right)^{\frac{2}{3}} \right]}$$

c) From our last homework, we know that

$$H(z) = H_0 (1+z) \left[ \Omega_{M,0} (1+z) + \Omega_{rel,0} (1+z)^2 + \frac{\Omega_{A,0}}{(1+z)^2} + 1 - \Omega_0 \right]^{\frac{1}{2}}$$

Though, since our universe is flat,  $\Omega_0 = 1$ . Using our equation that relates  $z$  to  $a$ , we can say

$$a = \frac{a_0}{(1+z)} = a_{\min} \rightarrow \left( \frac{\Omega_{M,0}}{2\Omega_{A,0}} \right)^{\frac{1}{3}} a_0 = \frac{a_0}{(1+z)}$$

$$(1+z) = \left( \frac{2\Omega_{A,0}}{\Omega_{M,0}} \right)^{\frac{1}{3}} \rightarrow z = \left( \frac{2\Omega_{A,0}}{\Omega_{M,0}} \right)^{\frac{1}{3}} - 1$$

# ① Dynamics of a Universe with a Cosmological Const. Cont.

Plugging this in, along with other things we know, we get

$$H(z) = H_0 \left( 1 + \left[ \left( \frac{2\Omega_{\Lambda,0}}{\Omega_{m,0}} \right)^{1/3} - 1 \right] \right) \left[ \Omega_{m,0} \left( 1 + \left[ \left( \frac{2\Omega_{\Lambda,0}}{\Omega_{m,0}} \right)^{1/3} - 1 \right] \right) + \right]$$

$\curvearrowleft + \Omega_{rel,0} \left( 1 + \left[ \left( \frac{2\Omega_{\Lambda,0}}{\Omega_{m,0}} \right)^{1/3} - 1 \right] \right)^2 + \Omega_{\Lambda,0} \left( 1 + \left[ \left( \frac{2\Omega_{\Lambda,0}}{\Omega_{m,0}} \right)^{1/3} - 1 \right] \right)^2 \right]^{1/2}$

for flat Uni.  $\Omega_{rel,0} = 0$

$$H(z) = H_0 \left( \frac{2\Omega_{\Lambda,0}}{\Omega_{m,0}} \right)^{1/3} \left[ \Omega_{m,0} \left( \frac{2\Omega_{\Lambda,0}}{\Omega_{m,0}} \right)^{1/3} + \Omega_{\Lambda,0} \left( \frac{\Omega_{m,0}}{2\Omega_{\Lambda,0}} \right)^{2/3} \right]$$

$$H(z) = H_0 \Omega_{m,0} \left( \frac{2\Omega_{\Lambda,0}}{\Omega_{m,0}} \right)^{2/3} + H_0 \Omega_{\Lambda,0} \left( \frac{\Omega_{m,0}}{2\Omega_{\Lambda,0}} \right)^{1/3}$$

- ∴ If the bracketted part of the equation is negative, it means that the hubble constant is negative. If  $H_0$  is (-), the universal rate of expansion is decreasing.

## ② Blackbody Radiation and the CMB Dipole

a) Given that  $B_{v_{\max}}(T)$  is a maximum, we know that  $\frac{dB_v(T)}{dT} = 0$ . So, let's take the derivative and solve.

$$\frac{dB_v(T)}{dT} = \frac{d}{dT} \left( \frac{2hv^3}{c^2(e^{\frac{hv}{kT}} - 1)} \right) \rightarrow 0 = \frac{2h^2v^4}{c^2kT^2(e^{\frac{hv}{kT}} - 1)} + \frac{2h^2v^4}{c^2kT^2(e^{\frac{hv}{kT}} - 1)^2}$$

$$\frac{-2h^2v^4}{c^2kT^2(e^{\frac{hv}{kT}} - 1)} = \frac{2h^2v^4}{c^2kT^2(e^{\frac{hv}{kT}} - 1)^2} \rightarrow -1 = \frac{1}{(e^{\frac{hv}{kT}} - 1)}$$

$$-e^{\frac{hv}{kT}} + 1 = 1 \rightarrow e^{\frac{hv}{kT}} = 0 \rightarrow \ln(e^{\frac{hv}{kT}}) = \ln(0)$$

God fears one thing, and it's  $\ln(0)$ . So let's set that to a constant and let that be someone else's problem to deal with.

$$\frac{hv}{kT} = C \rightarrow h \text{ and } k \text{ are const. absorb them.} \rightarrow \frac{v}{T} = C \rightarrow v = v_{\max}$$

and Bam! we get Wein's Law

$$\frac{v_{\max}}{T} = C$$

b) Given the Sun's peculiar velocity is  $v_p \approx 300 \frac{\text{km}}{\text{s}}$ , we can use the microwave background radiation Temperature equation to solve.

$$T_{\text{sun}} = T_0 (1+z) \quad \text{where } T_0 = T_{\text{CMB}}$$

Using values we know, we can solve ( $z \approx \frac{v_p}{c}$ )

$$T_{\text{sun}} = (2.725 \text{ K}) \left( 1 + \frac{(300 \times 10^3 \frac{\text{m}}{\text{s}})}{(3 \times 10^8 \frac{\text{m}}{\text{s}})} \right) \rightarrow T_{\text{sun}} = 2.727 \text{ K}$$

$$\Delta T = T_{\text{sun}} - T_{\text{CMB}} = (2.728 \text{ K}) - (2.725 \text{ K}) \rightarrow \Delta T = 0.00273 \text{ K}$$

### ③ The age of the Universe

a) We are given that  $\Omega_{k,0}=0$ ,  $\Omega_{m,0}+\Omega_{\Lambda,0}=1$ , and  $\Omega_0=1$ .

Plugging that into our Hubble Parameter we get

$$H(z) = H_0(1+z) \left[ \Omega_{m,0}(1+z) + \Omega_{k,0}(1+z)^2 + \frac{\Omega_{\Lambda,0}}{(1+z)^2} + 1 - \Omega_0 \right]^{\frac{1}{2}}$$

$$H(z) = H_0(1+z) \left[ \Omega_{m,0}(1+z) + 0 + \frac{\Omega_{\Lambda,0}}{(1+z)^2} + 1 - (1) \right]^{\frac{1}{2}}$$

Letting  $z = \frac{a_0}{a} - 1$ , we can say

$$H\left(\frac{a_0}{a}-1\right) = H_0\left(1+\left(\frac{a_0}{a}-1\right)\right) \left[ \Omega_{m,0}\left(1+\left(\frac{a_0}{a}-1\right)\right) + \Omega_{\Lambda,0}\left(\frac{1}{1+\left(\frac{a_0}{a}-1\right)}\right)^2 \right]^{\frac{1}{2}}$$

$$H\left(\frac{a_0}{a}-1\right) = H_0\left(\frac{a_0}{a}\right) \left[ \Omega_{m,0}\left(\frac{a_0}{a}\right) + \Omega_{\Lambda,0}\left(\frac{a}{a_0}\right)^2 \right]^{\frac{1}{2}}$$

$$H\left(\frac{a_0}{a}-1\right) = H_0 \left[ \Omega_{m,0}\left(\frac{a_0^3}{a^3}\right) + \Omega_{\Lambda,0} \right]^{\frac{1}{2}}$$

If we set this equal to  $\frac{\dot{a}}{a}$ , we can integrate and simplify.

$$\frac{\dot{a}}{a} = H_0 \left[ \Omega_{m,0}\left(\frac{a_0^3}{a^3}\right) + \Omega_{\Lambda,0} \right]^{\frac{1}{2}} \rightarrow H_0 = \frac{\dot{a}}{a \left[ \Omega_{m,0}\left(\frac{a_0^3}{a^3}\right) + \Omega_{\Lambda,0} \right]^{\frac{1}{2}}}$$

### ③ The age of the Universe Cont.

$$\int H_0 dt = \int \frac{da}{a \left[ \Omega_{m,0} \left( \frac{a_0^3}{a^3} \right) + \Omega_{\Lambda,0} \right]^{1/2}}$$

Since  $\frac{a_0}{a}$  is super gross, let's make a u substitution for  $u = a/a_0 \rightarrow a = u(a_0) \Rightarrow da = du(a_0)$ . We get

$$H_0 t = \int \frac{du(a_0)}{\left( u a_0 \right) \left[ \Omega_{m,0} \left( \frac{a_0^3}{(ua_0)^3} \right) + \Omega_{\Lambda,0} \right]^{1/2}} \rightarrow H_0 t = \int \frac{du}{u \left[ \Omega_{m,0} \left( \frac{1}{u^3} \right) + \Omega_{\Lambda,0} \right]^{1/2}}$$

Similarly, we can make this substitution in our  $\tan(\theta)$  expression to help us will simplifying our integral.

$$\tan(\theta) = \left( \frac{\Omega_{m,0}}{\Omega_{\Lambda,0}} \right)^{1/2} (1+z)^{3/2} \rightarrow \text{let } z = \frac{a}{a_0} - 1 \rightarrow z = u - 1$$

$$\tan(\theta) = \left( \frac{\Omega_{m,0}}{\Omega_{\Lambda,0}} \right)^{1/2} (1+(u-1))^{3/2} \rightarrow \tan(\theta) = \left( \frac{\Omega_{m,0}}{\Omega_{\Lambda,0}} \right)^{1/2} (u)^{3/2}$$

$$u = (\tan \theta)^{-2/3} \left( \frac{\Omega_{m,0}}{\Omega_{\Lambda,0}} \right)^{1/3} \Rightarrow du = -\frac{2 \sec^2(\theta)}{3 \tan^{5/3}(\theta)} \left( \frac{\Omega_{m,0}}{\Omega_{\Lambda,0}} \right)^{1/3} d\theta$$

### ③ The age of the Universe Cont.

Plug our new expression in.

$$H_0 t = \int \frac{\left(\frac{2}{3}\right) \sec^2 \theta \left(\frac{\Omega_{m,0}}{\Omega_{\Lambda,0}}\right)^{1/3} d\theta}{\left[\left(\tan \theta\right)^{-2/3} \left(\frac{\Omega_{m,0}}{\Omega_{\Lambda,0}}\right)^{1/3}\right] \left[\Omega_{m,0} \left(\frac{1}{\left(\tan \theta\right)^{2/3} \left(\frac{\Omega_{m,0}}{\Omega_{\Lambda,0}}\right)^{1/3}}\right)^3 + \Omega_{\Lambda,0}\right]^{1/2}}$$

$$H_0 t = -\frac{2}{3} \int \frac{\sec^2 \theta \tan^{2/3} \theta d\theta}{\tan^{5/3} \theta \left[\Omega_{m,0} \left(\frac{\Omega_{\Lambda,0}}{\tan^2 \theta \Omega_{m,0}}\right) + \Omega_{\Lambda,0}\right]^{1/2}}$$

$$H_0 t = -\frac{2}{3} \int \frac{\sec^2 \theta}{\tan \theta} \left(\Omega_{\Lambda,0} (\tan^2 \theta + 1)\right)^{1/2} d\theta$$

$$H_0 t = -\frac{2}{3} \int \frac{\sec^2 \theta}{\tan \theta} \left(\Omega_{\Lambda,0} (\sec^2 \theta)\right)^{1/2} d\theta$$

$$H_0 t = -\frac{2}{3} \int \frac{\sec \theta}{\tan \theta} \frac{1}{(\Omega_{\Lambda,0})^{1/2}} d\theta$$

### ③ The age of the Universe Cont.

Let's rewrite  $\tan\theta$  and  $\sec\theta$  into  $\sin\theta$  and  $\cos\theta$ .

$$H_0 t = \frac{2}{3} \left( \frac{1}{\Omega_{m,0}} \right)^{\frac{1}{2}} \int \frac{\left( \frac{1}{\cos\theta} \right)}{\left( \frac{\sin\theta}{\cos\theta} \right)} d\theta \rightarrow H_0 t = \frac{2}{3} \left( \frac{1}{\Omega_{m,0}} \right)^{\frac{1}{2}} \int \frac{d\theta}{\sin\theta}$$

And finally Integrate.

$$H_0 t = \frac{2}{3 \Omega_{m,0}^{\frac{1}{2}}} \ln \left[ \frac{1 + \cos\theta}{\sin\theta} \right] \rightarrow t = \frac{2}{3 H_0 \Omega_{m,0}^{\frac{1}{2}}} \ln \left[ \frac{1 + \cos\theta}{\sin\theta} \right]$$

b) So, at  $t=0$  (Present Day) there is no redshift and  $z=0$ . Thus, we can say from given that

$$\tan\theta = \left( \frac{\Omega_{m,0}}{\Omega_{r,0}} \right)^{\frac{1}{2}} (1+0)^{\frac{3}{2}} \rightarrow \tan\theta = \left( \frac{\Omega_{m,0}}{\Omega_{r,0}} \right)^{\frac{1}{2}} (1) \rightarrow \tan^2\theta = \frac{\Omega_{m,0}}{\Omega_{r,0}}$$

We need to find  $\cos\theta$  and  $\sin\theta$ , so

$$\cos\theta = \frac{1}{\sec\theta} = \frac{1}{\sqrt{1 + \tan^2\theta}} \rightarrow \cos\theta = \left( 1 + \frac{\Omega_{m,0}}{\Omega_{r,0}} \right)^{-\frac{1}{2}} = \left( \frac{1}{\Omega_{r,0}} \right)^{-\frac{1}{2}}$$

### ③ The age of the Universe Cont.

$$\sin \theta = \tan \theta \cdot \cos \theta \rightarrow \sin \theta = \left( \frac{\Omega_{m,0}}{\Omega_{\Lambda,0}} \right)^{-1/2} \left( \frac{1}{\Omega_{\Lambda,0}} \right)^{-1/2} = \left( \frac{1}{\Omega_{m,0}} \right)^{-1/2}$$

Plug it all in

$$t_0 = \frac{2}{3H_0 \Omega_{\Lambda,0}^{1/2}} \ln \left[ \frac{1 + \cos \theta}{\sin \theta} \right] \rightarrow t_0 = \frac{2}{3H_0 \Omega_{\Lambda,0}^{1/2}} \ln \left[ \frac{1 + (\Omega_{\Lambda,0})^{-1/2}}{(1 - \Omega_{\Lambda,0})^{-1/2}} \right]$$

$$t_0 = \frac{2}{3H_0 \Omega_{\Lambda,0}^{1/2}} \ln \left[ \frac{(1 + \Omega_{\Lambda,0}^{-1/2})}{(1 - \Omega_{\Lambda,0})^{-1/2}} \right]$$

## ④ Inflation

From class, we know that  $t_{\text{GUT}} \sim 10^{-36}$  s. Since we are given  $m_M = m_{\text{GUT}}$ , we can make the statement.

$$n_M(t_{\text{GUT}}) = \frac{1 \text{ monopole}}{(2ct_{\text{GUT}})^3} = \frac{1 \text{ monopole}}{(2(3 \times 10^8)(10^{-36}))^3} \sim 10^{82} \text{ m}^{-3}$$

This shows us that monopoles will decrease in density proportional to  $a^{-3}$ . And, knowing that the temperature of radiation decreases with universal expansion, we can say

$$a(t_{\text{GUT}})^3 n_{\text{GUT}} \approx \frac{T_{\text{CMB}}}{T_{\text{GUT}}} = \frac{(2.73 \text{ K})}{(10^{28} \text{ K})} \rightarrow 3 \times 10^{-28}$$

$$n_{M,0} = a(t_{\text{GUT}})^3 n_{\text{GUT}} \rightarrow 10^{-82} \times 10^{82} \text{ m}^3 \rightarrow 1 \text{ m}^{-3}$$

Thus, without expansion, we expect to see a magnetic monopole about every cubic meter of space. Now, knowing  $E_{M,0} = (10^{-6})(4870 \frac{\text{MeV}}{\text{m}^3}) = 4.87 \times 10^{-3} \frac{\text{MeV}}{\text{m}^3}$  and assuming expansion is responsible for reducing density, we can say

## ④ Inflation Cont.

$$n_{M,0} = \frac{\epsilon_{M,0}}{m_{GUT} C^2} \rightarrow (m_{GUT} C^2)(1 \text{ m}^{-3}) = \epsilon_{M,0}$$

Let  $m_{GUT} C^2 = a \rightarrow (a)^3 (1 \text{ m}^{-3}) = 4.87 \times 10^{-3} \frac{\text{MeV}}{\text{m}^3}$

NOW, we know a inflation will look like  $a = e^{-N}$ .  
 And, if we use the book value for  $m_{GUT} C^2$ , we can solve for N.

$$(4.87 \times 10^{-3} \text{ MeV})^{\frac{1}{3}} = e^{-N} \rightarrow -N = \ln(4.87 \times 10^{-3} \text{ MeV})^{\frac{1}{3}}$$

$$N = -\ln(1.775 \times 10^{-7}) \rightarrow \boxed{N \approx 15.54}$$