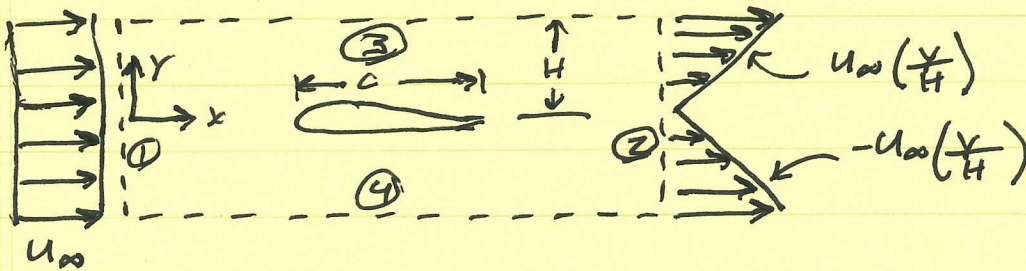


Part of Problem 1:



Assume: incompressible, steady, 2-D.

Q: What is the total Q across ③ and ④
 \uparrow volume flow rate

Use the conservation of mass equation (continuity):

$$\frac{\partial}{\partial t} \int_{cv} \rho dV + \int_{cs} \rho (\vec{v} \cdot \hat{n}) dA = 0$$

steady

$$\rho \int_{cs} (\vec{v} \cdot \hat{n}) dA = 0 \quad (\text{incompressible really const. } \rho)$$

$$\int_{\text{①}} \vec{v} \cdot \hat{n} dA + \int_{\text{②}} \vec{v} \cdot \hat{n} dA + \underbrace{\int_{\text{③}} \vec{v} \cdot \hat{n} dA + \int_{\text{④}} \vec{v} \cdot \hat{n} dA}_{Q_{34}} = 0$$

$$Q_{34} = - \int_{\text{①}} \vec{v} \cdot \hat{n} dA - \int_{\text{②}} \vec{v} \cdot \hat{n} dA$$

$$= -(-U_{\infty}) \int_{-H}^H dA - U_{\infty} \int_0^H \frac{y}{H} dA - U_{\infty} \int_{-H}^0 -\frac{y}{H} dA$$

$$dA = w dy$$

\uparrow unit width $\rightarrow w=1$

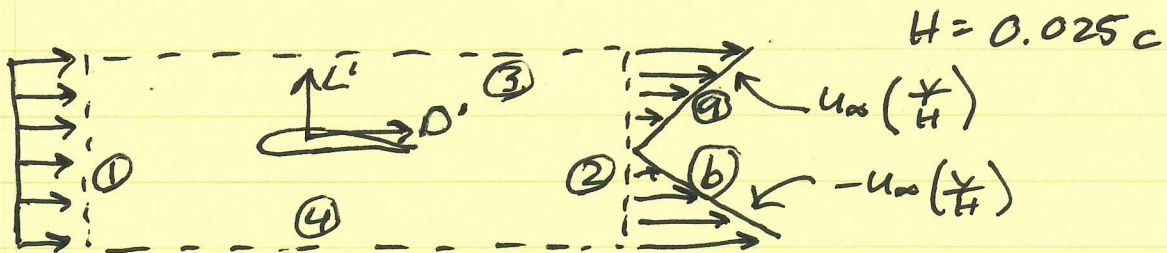
$$Q_{34} = 2U_{\infty} H - \frac{U_{\infty}}{H} \frac{y^2}{2} \Big|_0^H + \frac{U_{\infty}}{H} \frac{y^2}{2} \Big|_{-H}^0$$

$$Q_{34} = 2U_{\infty}H - \frac{U_{\infty}H}{2} - \frac{U_{\infty}H}{2}$$

$$\boxed{Q_{34} = U_{\infty}H} \text{ (per unit width)}$$

Problem 1:

Same Setup as Problem 5



Assume: Incompressible, steady, 2-D, pressure is constant on CS.
Inviscid

Q: What is the drag coefficient?

$$C_d = \frac{D'}{\frac{1}{2} \rho_{\infty} u_{\infty}^2 c}$$

Momentum equation:

$$\frac{d}{dt} \int_{CS} \vec{v} \rho dV + \int_{CS} \vec{v} \rho (\vec{v} \cdot \hat{n}) dA = \int_{CS} -p \hat{n} dA + \vec{F}_{visc} + \vec{F}_{ext}$$

Steady
const. p on CS
Inviscid

$\vec{F}_{ext} = -D' \hat{i} - L' \hat{j}$ ← ~~The fluid~~ The fluid acts on the airfoil to produce lift and drag.

The airfoil acts on the fluid with equal and opposite force by Newton's third law.
i.e. The force acting on the fluid is $= -D' \hat{i} - L' \hat{j}$

On (3) and (4) we assume: $\vec{v} = u_{\infty} \hat{i} + v(x) \hat{j}$

↑ small flow perturbations to the \hat{i} component of velocity have little effect on the boundary.

x-momentum equation:

$$\int_{\textcircled{1}} U_{\infty} \rho_{\infty} (-U_{\infty}) dy + \int_{\textcircled{2a}} \left(U_{\infty} \frac{y}{H} \right) \rho_{\infty} \left(U_{\infty} \frac{y}{H} \right) dy \quad (0 \leq y \leq H)$$

$$+ \int_{\textcircled{2b}} \left(-U_{\infty} \frac{y}{H} \right) \rho_{\infty} \left(-U_{\infty} \frac{y}{H} \right) dy \quad (-H \leq y \leq 0)$$

$$+ \int_{\textcircled{3} + \textcircled{4}} U_{\infty} \rho_{\infty} (\vec{v} \cdot \hat{n}) dy$$

$$\int_{-H}^H \textcircled{1}: -\rho_{\infty} U_{\infty}^2 (2H) = U_{\infty} \rho_{\infty} \int_{\textcircled{3} + \textcircled{4}} (\vec{v} \cdot \hat{n}) dy$$

$$\int_0^H \textcircled{2a}: \frac{\rho_{\infty} U_{\infty}^2}{H^2} \frac{y^3}{3} \Big|_0^H = \frac{\rho_{\infty} U_{\infty}^2 H}{3} \quad \downarrow \text{Q34 from Problem 5!}$$

$$\int_{-H}^0 \textcircled{2b}: \frac{\rho_{\infty} U_{\infty}^2}{H^2} \frac{y^3}{3} \Big|_{-H}^0 = \frac{\rho_{\infty} U_{\infty}^2 H}{3} = U_{\infty} \rho_{\infty} (U_{\infty} H) = \rho_{\infty} U_{\infty}^2 H$$

Put these together with the RHS

$$-D' = -\rho_{\infty} U_{\infty}^2 (2H) + \frac{\rho_{\infty} U_{\infty}^2 H}{3} + \frac{\rho_{\infty} U_{\infty}^2 H}{3} + \rho_{\infty} U_{\infty}^2 H$$

$$-D' = -\frac{\rho_{\infty} H U_{\infty}^2}{3} \Rightarrow \frac{D'}{\frac{1}{2} \rho_{\infty} U_{\infty}^2 c} = \frac{2H}{3c} \quad H = 0.025c$$

$$C_d = +\frac{2}{3} (0.025)$$

$$\boxed{C_d = 0.0167}$$

Problem 2 (3.7.1 Kuo & He)

Derive: $\rho \frac{D\vec{v}}{Dt} = \rho \vec{g} - \vec{\nabla} P = \rho \left[\frac{\partial \vec{v}}{\partial t} + (\vec{v} \cdot \vec{\nabla}) \vec{v} \right] = \rho \vec{g} - \vec{\nabla} P$

From: $\frac{d}{dt} \int_{\hat{R}} \rho \vec{v} d\hat{R} + \int_{\hat{S}} \rho \vec{v} (\vec{v} \cdot \hat{n}) d\hat{S} = -\vec{F}_e - \int_{\hat{S}} P \hat{n} d\hat{S} + \int_{\hat{R}} \rho \vec{g} d\hat{R}$

Neglecting external forces $\rightarrow \vec{F}_e = 0$

Divergence theorem: $\int_{\hat{S}} \hat{n} \cdot \vec{G} d\hat{S} = \int_{\hat{R}} \vec{\nabla} \cdot \vec{G} d\hat{R}$

Momentum flux term: $\int_{\hat{S}} \rho \vec{v} (\vec{v} \cdot \hat{n}) d\hat{S} = \int_{\hat{S}} (\hat{n} \cdot \rho \vec{v}) \vec{v} d\hat{S}$

Applying divergence $\rightarrow \int_{\hat{S}} (\hat{n} \cdot \rho \vec{v}) \vec{v} d\hat{S} = \int_{\hat{R}} \vec{\nabla} \cdot (\rho \vec{v} \vec{v}) d\hat{R}$

Expanding $\rightarrow \int_{\hat{R}} \vec{\nabla} \cdot (\rho \vec{v} \vec{v}) d\hat{R} = \int_{\hat{R}} (\vec{\nabla} \cdot \rho \vec{v}) \vec{v} d\hat{R} + \int_{\hat{R}} \rho (\vec{v} \cdot \vec{\nabla}) \vec{v} d\hat{R}$

Pressure term: $\int_{\hat{S}} \hat{n} P d\hat{S} = \int_{\hat{R}} \vec{\nabla} P d\hat{R}$

easy to see using index notation.

Collecting terms into a single integral:

$$\int_{\hat{R}} \left[\frac{\partial}{\partial t} (\rho \vec{v}) + (\vec{\nabla} \cdot \rho \vec{v}) \vec{v} + \rho (\vec{v} \cdot \vec{\nabla}) \vec{v} + \vec{\nabla} P - \rho \vec{g} \right] d\hat{R} = 0$$

Integrand must be zero since \hat{R} is an arbitrary volume:

$$\frac{\partial}{\partial t} (\rho \vec{v}) + (\vec{\nabla} \cdot \rho \vec{v}) \vec{v} + \rho (\vec{v} \cdot \vec{\nabla}) \vec{v} + \vec{\nabla} P - \rho \vec{g} = 0$$

Expanding terms: $\rho \frac{\partial \vec{v}}{\partial t} + \vec{v} \frac{\partial \rho}{\partial t} + (\vec{\nabla} \cdot \rho \vec{v}) \vec{v} + \rho (\vec{v} \cdot \vec{\nabla}) \vec{v} + \vec{\nabla} P - \rho \vec{g} = 0$

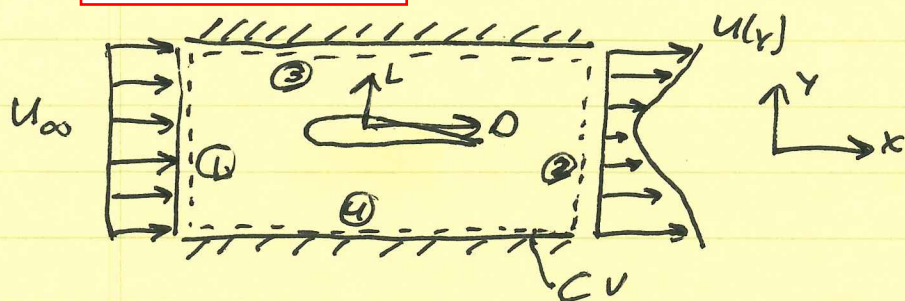
$$\left[\frac{\partial \rho}{\partial t} + (\vec{\nabla} \cdot \rho \vec{v}) \right] \vec{v}$$

continuity equation = 0

Now: $\rho \frac{\partial \vec{v}}{\partial t} + \rho (\vec{v} \cdot \vec{\nabla}) \vec{v} = \rho \vec{g} - \vec{\nabla} P$

and $\rightarrow \boxed{\rho \frac{D\vec{v}}{Dt} = \rho \vec{g} - \vec{\nabla} P}$

Problem 3:



To solve this problem we need to consider a control volume, CV. We can write down the momentum equation:

$$\frac{\partial}{\partial t} \int_{CV} \vec{v} \rho dV + \int_{CS} \vec{v} \rho (\vec{v} \cdot \hat{n}) dA = \sum \left(\text{Forces acting on the fluid} \right)$$

* Remember the forces on the right hand side (RHS) are forces acting on the fluid. These forces are equal and opposite to the forces acting on the airfoil!

The forces are:

$$\sum (\text{Forces}) = \int_{CS} -P \hat{n} dA + \vec{F}_{\text{viscous}} + \vec{F}_{\text{external}}$$

$$\vec{F}_{\text{external}} = -D \hat{i} - L \hat{j} \leftarrow \text{equal and opposite to } L \text{ and } D. \text{ Acting on the fluid.}$$

Since lift is in the \hat{j} direction we only need to consider the \hat{j} component of the momentum equation.

$$\frac{\partial}{\partial t} \int_{CV} (u_y \hat{j}) \rho dV + \int_{CS} (u_y \hat{j}) \rho (\vec{v} \cdot \hat{n}) dA = \int_{CS} -P \hat{j} dA + F_{\text{viscous}} \hat{j}$$

Steady

$-L \hat{j}$

only significant viscous forces on top and bottom wall in \hat{j} direction

$$\begin{aligned}
 & \int_{\textcircled{1}} (\cancel{V_y \hat{j}}) \rho (\vec{v} \cdot \hat{n}) dA + \int_{\textcircled{2}} (\cancel{V_y \hat{j}}) \rho (\vec{v} \cdot \hat{n}) dA \\
 & \quad \text{zero on } \textcircled{1} \quad \text{far enough downstream that } V_y \text{ on } \textcircled{2} \text{ is zero} \\
 & + \int_{\textcircled{3}} (\cancel{V_y \hat{j}}) \rho (\vec{v} \cdot \hat{n}) dA + \int_{\textcircled{4}} (\cancel{V_y \hat{j}}) \rho (\vec{v} \cdot \hat{n}) dA = 0 \\
 & \quad \text{zero on } \textcircled{3} \quad \left. \begin{array}{l} \text{no mass} \\ \text{flux through} \\ \text{a wall} \end{array} \right\} \quad \text{Same as } \textcircled{3}
 \end{aligned}$$

$$\begin{aligned}
 \Rightarrow 0 &= \int_{cs} -P \hat{j} dA - L \hat{j} \\
 L &= \int_{\textcircled{4}} P_L dA - \int_{\textcircled{3}} P_u dA
 \end{aligned}$$

\swarrow pressure distribution on lower wall \nwarrow pressure on upper wall

$$\boxed{L = \int_A (P_L - P_u) dA}$$

* Assumptions:

- 1.) Inviscid on entrance and exit
- 2.) Steady flow
- 3.) Neglecting gravity
- 4.) CV is long enough for $V_y = 0$ on $\textcircled{3}$

Problem 4

$$\psi = U_{\infty} r \sin \theta + \frac{\sigma}{2\pi} \theta \Rightarrow V_r = \frac{1}{r} \frac{\partial \psi}{\partial \theta} = U_{\infty} \cos \theta + \frac{\sigma}{2\pi r}$$

$$\frac{\sigma}{2\pi U_{\infty}} \equiv 1$$

$$V_{\theta} = -\frac{\partial \psi}{\partial r} = -U_{\infty} \sin \theta$$

Stagnation occurs when $V_r = V_{\theta} = 0$

$$\text{Stag pt: } (r, \theta) = \left(\frac{\sigma}{2\pi U_{\infty}}, \pi \right)$$

Stagnation streamline occurs for the constant: $\psi = \frac{\sigma}{2}$

$$\text{So, } \frac{\sigma}{2} = U_{\infty} r \sin \theta + \frac{\sigma}{2\pi} \theta$$

$$\left(\frac{2\pi}{\sigma} \right) \left(\frac{\sigma}{2} \right) = U_{\infty} r \sin \theta \left(\frac{2\pi}{\sigma} \right) + \left(\frac{2\pi}{\sigma} \right) \frac{\sigma}{2\pi} \theta$$

$$\pi = \underbrace{\left(\frac{2\pi U_{\infty}}{\sigma} \right)}_1 r \sin \theta + \theta$$

$$\pi = r \sin \theta + \theta$$

$$\boxed{r = \frac{\pi - \theta}{\sin \theta}}$$

← give a range of θ
and solve for r

↖ not valid for $\theta = 0, \pi$

$$\text{If } \theta = 0: \pi = r \sin(0) + (0) \rightarrow \boxed{r = \text{undefined!}}$$

$$\text{If } \theta = \pi: \pi = r \sin(\pi) + (\pi)$$

$$\pi = \pi$$

→ r can be any value
from $(0 \leq r \leq \infty)$

Plot!


```

clear
close all
clc

% Create r and theta vectors
r = linspace(1/2,10,50)';
theta = linspace(0,2*pi,51)';

% Create a grid of (r,theta) combinations
[R, Theta] = meshgrid(r,theta);

% Calculate V_r/Uinf and V_theta/Uinf
V_r_Uinf = cos(Theta) + 1./R;
V_theta_Uinf = -sin(Theta);

% Convert polar (r,theta) to Cartesian (x,y)
X = R.*cos(Theta);
Y = R.*sin(Theta);

% Find the matrix size of R and use for indices
[rows, cols] = size(R);

% Step through each (r,theta) combination
for i = 1:rows
    for j = 1:cols

        % Calculate the Cartesian velocity components using a
        % transformation matrix, Q:
        %
        % Q = [cos(theta) -sin(theta)
        %      sin(theta)  cos(theta)]
        %
        % then,
        %
        % [U; V] = Q*[V_r; V_theta]

        U(i,j) = V_r_Uinf(i,j)*cos(Theta(i,j)) - V_theta_Uinf(i,j)*sin(Theta(i,j));
        V(i,j) = V_r_Uinf(i,j)*sin(Theta(i,j)) + V_theta_Uinf(i,j)*cos(Theta(i,j));
    end
end

% Streamline calculations
theta_top = linspace(0.01,pi-0.01,1000)';
theta_bot = linspace(pi+0.01,2*pi - 0.01,1000)';
r_top = (pi - theta_top)./sin(theta_top);
r_bot = (pi - theta_bot)./sin(theta_bot);
r_pi = linspace(0,10,1000)';

% Convert to Cartesian coordinates
[x_top,y_top] = pol2cart(theta_top,r_top);
[x_bot,y_bot] = pol2cart(theta_bot,r_bot);
[x_pi,y_pi] = pol2cart(pi*ones(1000,1),r_pi);

% Plot the velocity field and streamlines

```

```

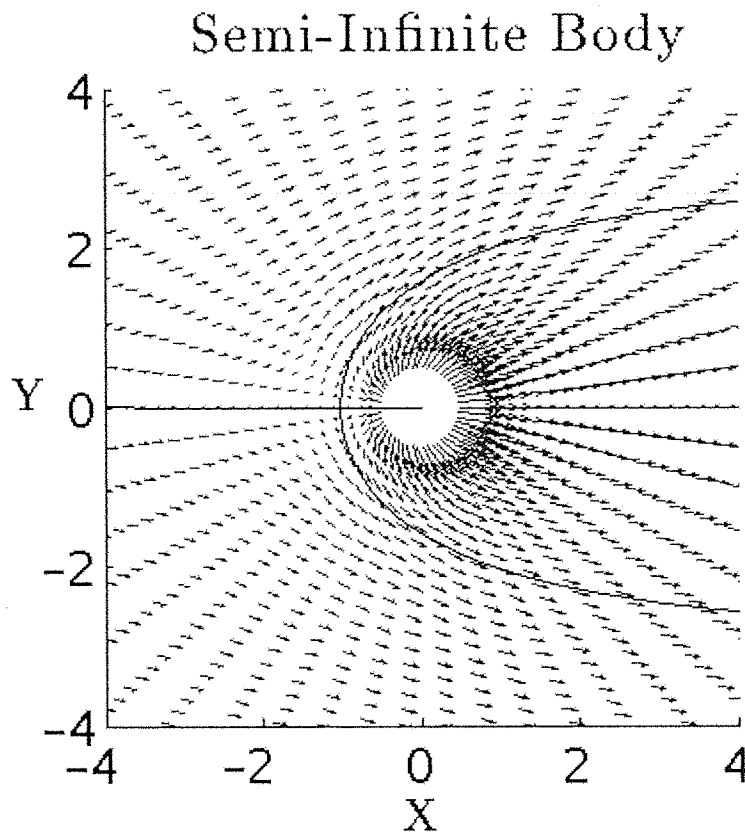
figure
set(gcf,'DefaultAxesfontsize',24,'DefaultAxesfontname','TimesNewRoman','DefaultAxesGridLineStyle','-.')

hold on
    quiver(X,Y,U,V)
    plot(x_top,y_top,'-k','LineWidth',1.5)
    plot(x_bot,y_bot,'-k','LineWidth',1.5)
    plot(x_pi,y_pi,'-k','LineWidth',1.5)
hold off
axis equal
axis([-4 4 -4 4])

title('Semi-Infinite Body','Interpreter','Latex','FontName','TimesNewRoman','FontSize',28)
xlabel('X','Interpreter','Latex','FontName','TimesNewRoman','FontSize',24)
ylabel('Y','Interpreter','Latex','FontName','TimesNewRoman','FontSize',24,'Rotation',0)

print('-depsc', '-r600','Semi_Inf_Body.eps');

```



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Problem 5:

$$\psi = -\frac{\mu}{2\pi} \frac{\sin\theta}{r}$$

$$V_r = \frac{1}{r} \frac{\partial \psi}{\partial \theta} \Rightarrow -\frac{\mu}{2\pi} \frac{\cos\theta}{r^2}$$

$$V_\theta = -\frac{\partial \psi}{\partial r} \Rightarrow \frac{\mu}{2\pi} \frac{\sin\theta}{r^2}$$

$$\textcircled{1} V_r = \frac{\partial \phi}{\partial r} \quad V_\theta = \frac{1}{r} \frac{\partial \phi}{\partial \theta} \textcircled{2}$$

$$\textcircled{1} -\frac{\mu}{2\pi} \frac{\cos\theta}{r^2} = \frac{\partial \phi}{\partial r} \Rightarrow \phi = \int -\frac{\mu}{2\pi} \frac{\cos\theta}{r^2} dr + C$$

$$\phi = -\frac{\mu}{2\pi} \left(\frac{1}{-1}\right) \frac{\cos\theta}{r} + f(\theta) + C$$

$$\phi = \frac{\mu}{2\pi} \frac{\cos\theta}{r} + f(\theta) + C \textcircled{3}$$

$$\textcircled{2} -\frac{\mu}{2\pi} \frac{\sin\theta}{r^2} = \frac{1}{r} \frac{\partial \phi}{\partial \theta} \Rightarrow \psi =$$

$$\textcircled{2} -\frac{\mu}{2\pi} \frac{\sin\theta}{r^2} = \frac{1}{r} \frac{\partial \phi}{\partial \theta} \Rightarrow \phi = \int -\frac{\mu}{2\pi} \frac{\sin\theta}{r} d\theta + C$$

$$\phi = \frac{\mu}{2\pi} \frac{\cos\theta}{r} + g(r) + C \textcircled{4}$$

Compare (3) to (4):

$$f(\theta) = g(r) = 0$$

C is an arbitrary constant that we can set to zero.

$$\boxed{\phi = \frac{\mu}{2\pi} \frac{\cos\theta}{r}}$$

Plot!

Equipotential lines \rightarrow

$$\boxed{r = \frac{\mu}{2\pi C} \cos\theta}$$

Contents

- Plot the velocity field and equipotential lines: positive μ .
- Plot the velocity field and equipotential lines: negative μ .

```

clear
close all
clc

% Create r and theta vectors
r = linspace(1/2,10,50)';
theta = linspace(0,2*pi,51)';

% Create a grid of (r,theta) combinations
[R, Theta] = meshgrid(r,theta);

% Calculate V_r and V_theta
V_r = -cos(Theta)./R.^2;
V_theta = -sin(Theta)./R.^2;
V_r_neg = cos(Theta)./R.^2;
V_theta_neg = sin(Theta)./R.^2;

% Convert polar (r,theta) to Cartesian (x,y)
X = R.*cos(Theta);
Y = R.*sin(Theta);

% Find the matrix size of R and use for indices
[rows, cols] = size(R);

% Step through each (r,theta) combination
for i = 1:rows
    for j = 1:cols

        % Calculate the Cartesian velocity components using a
        % transformation matrix, Q:
        %
        % Q = [cos(theta) -sin(theta)
        %      sin(theta)  cos(theta)]
        %
        % then,
        %
        % [U; V] = Q*[V_r; V_theta]

        U(i,j) = V_r(i,j)*cos(Theta(i,j)) - V_theta(i,j)*sin(Theta(i,j));
        V(i,j) = V_r(i,j)*sin(Theta(i,j)) + V_theta(i,j)*cos(Theta(i,j));
        U_neg(i,j) = V_r_neg(i,j)*cos(Theta(i,j)) - V_theta_neg(i,j)*sin(Theta(i,j));
        V_neg(i,j) = V_r_neg(i,j)*sin(Theta(i,j)) + V_theta_neg(i,j)*cos(Theta(i,j));
    end
end

% Equipotential line calculations
theta = linspace(0,2*pi,1000)';
C_1 = 1/2;
C_2 = 1;
C_3 = 2;

```



```

C_4 = -1/2;
C_5 = -1;
C_6 = -2;

r_1 = 1/C_1*cos(theta);
r_2 = 1/C_2*cos(theta);
r_3 = 1/C_3*cos(theta);
r_4 = 1/C_4*cos(theta);
r_5 = 1/C_5*cos(theta);
r_6 = 1/C_6*cos(theta);

% Convert to Cartesian coordinates
[x_1,y_1] = pol2cart(theta,r_1);
[x_2,y_2] = pol2cart(theta,r_2);
[x_3,y_3] = pol2cart(theta,r_3);
[x_4,y_4] = pol2cart(theta,r_4);
[x_5,y_5] = pol2cart(theta,r_5);
[x_6,y_6] = pol2cart(theta,r_6);

```

Plot the velocity field and equipotential lines: positive mu.

```

figure
set(gcf,'DefaultAxesfontsize',24,'DefaultAxesfontname','TimesNewRoman','DefaultAxesGridLineStyle','-.')
num = 1;

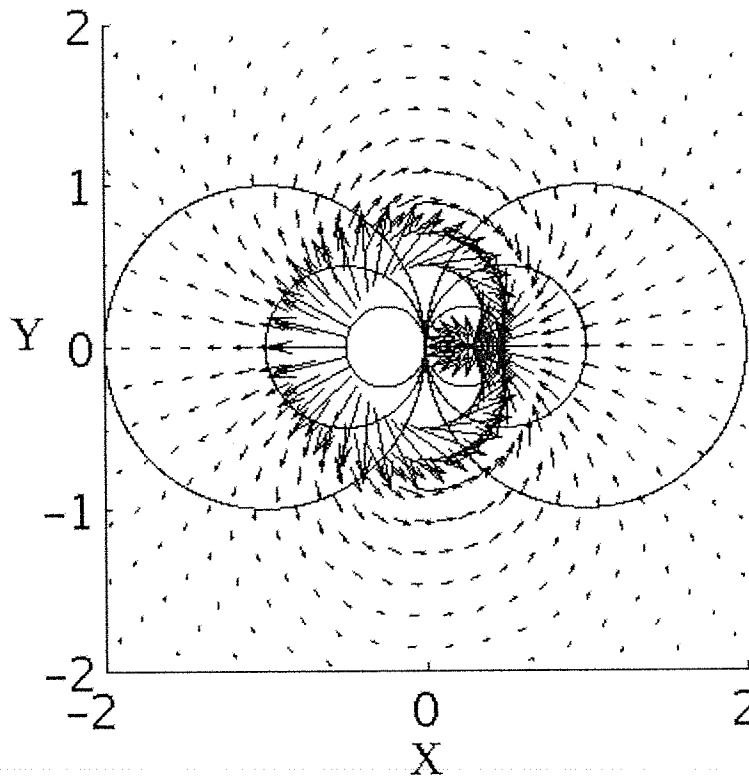
hold on
plot(x_1,y_1,'-k','LineWidth',num)
plot(x_2,y_2,'-k','LineWidth',num)
plot(x_3,y_3,'-k','LineWidth',num)
plot(x_4,y_4,'-k','LineWidth',num)
plot(x_5,y_5,'-k','LineWidth',num)
plot(x_6,y_6,'-k','LineWidth',num)
quiver(X,Y,U,V,'k')
hold off
axis equal
axis([-2 2 -2 2])

title('Doublet Flow','Interpreter','Latex','FontName','TimesNewRoman','FontSize',28)
xlabel('X','Interpreter','Latex','FontName','TimesNewRoman','FontSize',24)
ylabel('Y','Interpreter','Latex','FontName','TimesNewRoman','FontSize',24,'Rotation',0)

print('-depsc', '-r600','Doub_Flow_Pos.eps');

```

Doublet Flow



Plot the velocity field and equipotential lines: negative μ .

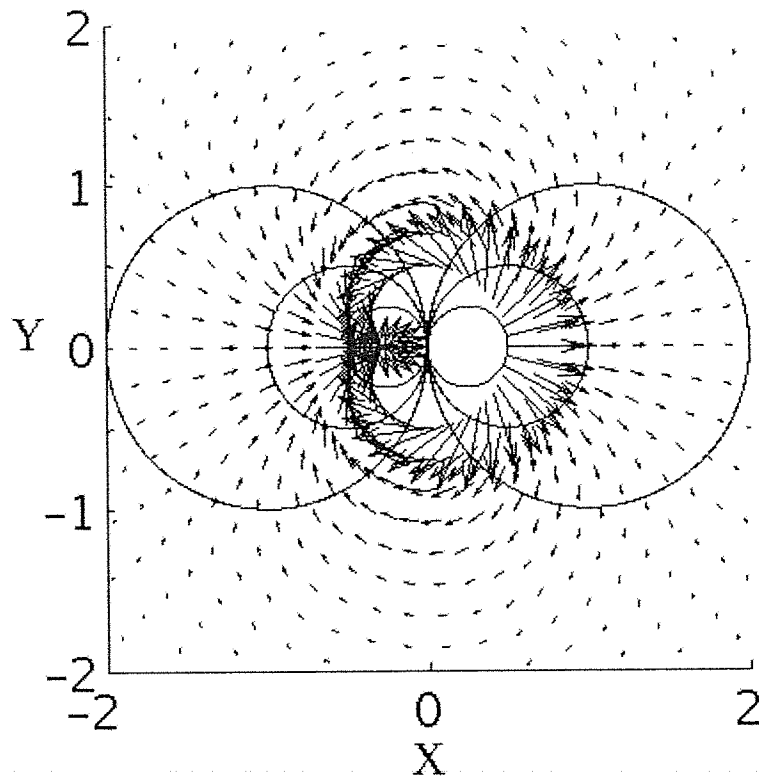
```
figure
set(gcf,'DefaultAxesfontsize',24,'DefaultAxesfontname','TimesNewRoman','DefaultAxesGridLineStyle','-.')
num = 1;

hold on
plot(x_1,y_1,'-k','LineWidth',num)
plot(x_2,y_2,'-k','LineWidth',num)
plot(x_3,y_3,'-k','LineWidth',num)
plot(x_4,y_4,'-k','LineWidth',num)
plot(x_5,y_5,'-k','LineWidth',num)
plot(x_6,y_6,'-k','LineWidth',num)
quiver(X,Y,U_neg,V_neg,'k')
hold off
axis equal
axis([-2 2 -2 2])

title('Doublet Flow','Interpreter','Latex','FontName','TimesNewRoman','FontSize',28)
xlabel('X','Interpreter','Latex','FontName','TimesNewRoman','FontSize',24)
ylabel('Y','Interpreter','Latex','FontName','TimesNewRoman','FontSize',24,'Rotation',0)

print('-depsc', '-r600','Doub_Flow_Neg.eps');
```


Doublet Flow



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